# Fractional Jet Multiplicity

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# Outline

- New observable: fractional jet multiplicity
- This talk: discuss analytic predictions and unique features [DB, J.Thaler, and J.Walsh 1501.01965]

#### Fractional Jet Multiplicity *definition*  $E$ rootional Jot $N$ Since we will be at each will be at each of the collision of the coll

$$
\widetilde{N}_{\text{jet}} = \sum_{i \in \text{event}} \frac{E_i}{E_{iR}} \Theta(E_{iR} > E_{\text{cut}})
$$
  

$$
E_{iR} = \sum_j E_j \Theta(\Delta \theta_{ij} < R)
$$
  
= energy in a cone around particle i

#### Fractional Jet Multiplicity *definition*  $E$ rootional Jot $N$ Since we will be at each will be at each of the collision of the coll



Example of distributions in  $e^+e^- \rightarrow$  jets



Example of distributions in  $e^+e^- \rightarrow$  jets



I will describe analytically the near-integer distribution

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Collimated particles give integer number of jets



Collimated particles give integer number of jets gi ve integer number of jets



 $\widetilde{N}_{jet} = \sum_i E_i / E_{iR}$  $=$   $1/3$ esterne<br>Experience  $\mathbf{E}_{\mathbf{i}}$ <sub>jet</sub> =  $\mathbf{\Sigma}_{\mathbf{i}}$  **E**<sub>i</sub> $\mathbf{E}_{\mathbf{i}}$  / **E**<sub>iR</sub>

Collimated particles give integer number of jets gi ve integer number of jets



$$
\widetilde{N}_{\rm jet} = \Sigma_{\rm i} E_{\rm i} / E_{\rm iR}
$$

 $=$   $1/3+1/3$ 

Collimated particles give integer number of jets gi ve integer number of jets



$$
\widetilde{N}_{\rm jet} = \Sigma_{\rm i} E_{\rm i} / E_{\rm iR}
$$

 $= 1/3 + 1/3 + 1/3 = 1$ 

Less collimated particles give fractional number of jets



Less collimated particles give fractional number of jets **c** es give fractional number of jets



$$
\widetilde{N}_{\rm jet} = \Sigma_{\rm i} E_{\rm i} / E_{\rm iR}
$$

 $= 1/2 + 1/3 + 1/2 = 4/3$ 

Soft limit drives the approach to the integer value ap proach to the integer value

$$
\widetilde{N}_{\text{jet}} = \Sigma_{\text{i}} E_{\text{i}} / E_{\text{iR}}
$$
  
\n
$$
\begin{pmatrix}\n1 \\
2 \\
z_3\n\end{pmatrix}\n\begin{pmatrix}\n3 \\
z_3\n\end{pmatrix}\n= E_3 / E_{\text{tot}} = 1/3
$$
  
\n
$$
= 4/3 \approx 1.3
$$

Consider now one particle going soft  $E_1=E_2>E_{\text{cut}}$  and  $E_3\rightarrow 0$ 

Soft limit drives the approach to the integer value ap proach to the integer value

$$
\widetilde{N}_{\text{jet}} = \Sigma_{\text{i}} E_{\text{i}} / E_{\text{IR}}
$$
\n
$$
\widetilde{Z}_{\text{2}} = E_{\text{3}} / E_{\text{tot}} = 1/9
$$
\n
$$
\widetilde{Z}_{\text{1}} = 1.1
$$

Consider now one particle going soft  $E_1=E_2>E_{\text{cut}}$  and  $E_3\rightarrow 0$ 

Soft limit drives the approach to the integer value ap proach to the integer value

$$
\widetilde{N}_{\text{jet}} = \Sigma_{\text{i}} E_{\text{i}} / E_{\text{iR}}
$$
\n
$$
\begin{pmatrix}\n1 & 2 \\
2 & 3\n\end{pmatrix}\n_{\text{z3}} = E_{3}/E_{\text{tot}} = 0
$$
\n
$$
\widetilde{N}_{\text{jet}} = \Sigma_{\text{i}} E_{\text{i}} / E_{\text{iR}}
$$

Consider now one particle going soft  $E_1=E_2>E_{\text{cut}}$  and  $E_3\rightarrow 0$ 

<u>The contract of the contract </u>

Well separated clusters of particles give integer number of jets

Well separated clusters of particles give integer number of jets er s of particles give integer number of jets



$$
\widetilde{N}_{\rm jet} = \Sigma_{\rm i} E_{\rm i} / E_{\rm iR}
$$

 $= 1/2+1/2+1=2$ 

## Fractional Jet Multiplicity *properties*

- To get non-integer behavior we need at least three particles.  $LO = O(\alpha_s^2)$  in pQCD
- They have to be collimated ( $\sim$  within 2R) and occupy special regions of phase space
- Near-integer behavior is driven by soft logarithms

**Consider**  $e^+e^- \rightarrow \text{jets}$ 



 $\Delta_{2-} = 2 - \widetilde{N}_{\text{jet}}, \quad \Delta_{2+} = \widetilde{N}_{\text{jet}} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{\text{jet}}$  $d\sigma$  $d\Delta_{2\pm}$ = z<br>Z  $d\Phi_4 \mathcal{T} (e^+e^- \to 4 \text{ partons}) \mathcal{F}(\Delta_{2\pm}, \Phi_4)$ 



 $\Delta_{2-} = 2 - \widetilde{N}_{\text{jet}}, \quad \Delta_{2+} = \widetilde{N}_{\text{jet}} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{\text{jet}}$ 

 $d\sigma$  $d\Delta_{2\pm}$ = z<br>Z  $d\Phi_4 \mathcal{T} (e^+e^- \to 4 \text{ partons}) \mathcal{F}(\Delta_{2\pm}, \Phi_4)$  $\mathcal{T}(e^+e^- \to 4 \text{ partons}) \simeq \mathcal{T}(e^+e^- \to q\bar{q}) \cdot \sum \mathcal{T}_k^{\text{coll}}(1 \to 3)$ *k*  $k \in \{q \rightarrow g g q \,, \bar{q} \rightarrow g g \bar{q} \,, q \rightarrow q' \bar{q}' q \,, \bar{q} \rightarrow q' \bar{q}' \bar{q} \}$ 



 $\Delta_{2-} = 2 - \widetilde{N}_{\text{jet}}, \quad \Delta_{2+} = \widetilde{N}_{\text{jet}} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{\text{jet}}$ 

$$
\frac{d\sigma}{d\Delta_{2\pm}} = \int d\Phi_4 \, \mathcal{T}(e^+e^- \to 4 \text{ partons}) \mathcal{F}(\Delta_{2\pm}, \Phi_4)
$$
\n
$$
\mathcal{T}(e^+e^- \to 4 \text{ partons}) \simeq \mathcal{T}(e^+e^- \to q\bar{q}) \cdot \sum_k \mathcal{T}_k^{\text{coll}}(1 \to 3) \qquad \text{measurement}
$$
\n
$$
k \in \{q \to ggg, \bar{q} \to gg\bar{q}, q \to q'\bar{q}'q, \bar{q} \to q'\bar{q}'\bar{q}\} \qquad \text{function}
$$

# Near-integer distribution

*soft logarithms and rapidity divergences*



### Near-integer distribution *soft logarithms and rapidity divergences*

Modes controlling near-integer behavior



Fixed order result

$$
\frac{d\sigma}{d\Delta_{2\pm}} = \sigma_0 \delta(\Delta_{2\pm})
$$
  
+  $\kappa_1 \left(\frac{\alpha_s}{\pi}\right)^2$   $\mathcal{L}_1(\Delta_{2\pm}) + \kappa_0 \left(\frac{\alpha_s}{\pi}\right)^2$   $\mathcal{L}_0(\Delta_{2\pm})$  + non-singular terms

- We calculated  $k_1$  and  $k_0$  (which include different color structures and leading dependence on  $z_{\text{cut}} = E_{\text{cut}} / E_{\text{tot}}$ )
- Get non singular from Event2

#### Event2 comparison:



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Fractional jet multiplicity

- Near-integer driven by soft-logs only
- Hybrid event-shape / jet-algorithm behavior
- Non-additive / non-factorizable / non-global

### Near-integer distribution *Hybrid event-shape / jet algorithm*



# Near-integer distribution

*additivity / factorizability / non-global logs*



### Near-integer distribution *beyond FO, collinear functions*



$$
C_{q,\bar{q}}(\Delta_{2\pm}) = \delta(\Delta_{2\pm}) + \sum_{n=2}^{\infty} \left(\frac{\alpha_s}{\pi}\right)^n \sum_{k=-1}^{n-1} \left[ \kappa_{k,+}^{(n)} \mathcal{L}_k(\Delta_{2\pm}) + \kappa_{k,-}^{(n)} \mathcal{L}_k(\Delta_{2\pm}) \right]
$$

### Near-integer distribution beyond FO, collinear functions a candidate factorization theorem for 2*<sup>±</sup>* in Sec. 5, which includes a renormalization-groupinder "collinear". In Sec. 6, we arrive the construction". In Sec. 6, we have a second to extend the second our

 $d\sigma$  $d\Delta_{2\pm}$  $\simeq$  $\simeq \sigma_0 \, C_q(\Delta_{2\pm}) \otimes C_{\overline{q}}(\Delta_{2\pm})$ 2 Aspects of Fractional Jets



Wide angle soft emissions give enhanced logarithmic contributions to  $\dot{N}_{jet}$ =2 cross-section rapidity distances<br>
ra e iet<sup>=2</sup> cross-section

#### Near-integer distribution *beyond FO, collinear functions* ar intograp dictribution UGH THE YY FRONT ON WATER ON ANALYTIC CALCULATIONS TO PARAMETERS TO PARAMETERS TO PARAMETERS TO PHONE IN SECTION TO PARAMETERS TO PARAMETE y change, commodition in Sec. 8. The appendices contained results and details.

 $d\sigma$  $d\Delta_{2\pm}$  $\simeq$  $\simeq \sigma_0 \, C_q(\Delta_{2\pm}) \otimes C_{\overline{q}}(\Delta_{2\pm})$  $\overline{\mathrm{d} \Delta_{2\pm}}$  -  $\sigma_0 C_q(\Delta_{2\pm}) \otimes C_{\bar{q}}(\Delta_{2\pm})$ 



At higher orders they also enhance fractional N<sub>jet</sub><br>cross section cross-section e jet(*E*cut*, R*) = <sup>X</sup>

### Near-integer distribution *beyond FO, collinear functions*

 $d\sigma$  $d\Delta_{2\pm}$  $\simeq$  $\simeq \sigma_0 \, C_q(\Delta_{2\pm}) \otimes C_{\overline{q}}(\Delta_{2\pm})$ 



Note that this contribution is contained in C

### Near-integer distribution *a candidate factorization theorem*

$$
\frac{d\sigma}{d\Delta_{2\pm}} \simeq \sigma(\widetilde{N}_{\rm jet} = 2) \left[ C_q(\Delta_{2\pm}) \otimes C_{\overline{q}}(\Delta_{2\pm}) \right]
$$

### Near-integer distribution *a candidate factorization theorem*

$$
\frac{d\sigma}{d\Delta_{2\pm}} \simeq \sigma(\widetilde{N}_{\rm jet} = 2) \left[ C_q(\Delta_{2\pm}) \otimes C_{\overline{q}}(\Delta_{2\pm}) \right]
$$

$$
\sigma(\widetilde{N}_{\rm jet} = 2) = \sigma_0 H_{q\overline{q}}(Q, \mu) J_q(Q, R, z_{\rm cut}, \mu) J_{\overline{q}}(Q, R, z_{\rm cut}, \mu) S_{q\overline{q}}(R, z_{\rm cut}, \mu)
$$

$$
+ \sigma_2^{\rm non-fac}(Q, R, z_{\rm cut}, \mu)
$$

Improved distributions

- Include  $O(\alpha_s^4)$  terms from convolutions
- Running coupling  $\alpha_s(\mu), \mu = Q\sqrt{\Delta_{2\pm}}$







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# Conclusions

- Fractional jet multiplicity can be used as a novel and more powerful probe of jet formation. E.g. quark/gluon discrimination?
- It has peculiar analytic properties. However, I showed we still have very good analytic control, which in principle is improvable. E.g. generalize to LHC case?
- Wide dynamic range, potential test of matching/merging matrix-element/parton shower. Potential phenomenological applications e.g. in multijet final states?

# Backup

*near-integer phase space configurations*

$$
\Delta_{2-} = 2 - \widetilde{N}_{\rm jet}, \quad \Delta_{2+} = \widetilde{N}_{\rm jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{\rm jet}
$$



#### Backup  $\Omega$  $\Omega$ plicities <sup>2</sup>, 2+, and <sup>3</sup>. Compared to Fig. 4, the value of *N*ejet is 1 unit higher, because the event

*near-integer phase space configurations* contains an additional isolated parton (not shown). For each observable, we show the corresponding contributions from di↵erent angular regions and soft limits. Circles represent cones of radius *R*, large blue dots represent energetic partons, small red dots soft partons with *z<z*cut. The angular regions

$$
\Delta_{2-} = 2 - \widetilde{N}_{\rm jet}, \quad \Delta_{2+} = \widetilde{N}_{\rm jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{\rm jet}
$$



## Backup *rapidity divergences*

 $\Delta = z_1z_2$ 

1 soft: 
$$
I_{sc}(\Delta) = \int_0^\infty \frac{dz_1}{z_1} \int_0^1 \frac{dz_2}{z_2} (z_1 z_2)^{-2\epsilon} \delta(\Delta - z_1 z_2)
$$
  
\n2 soft:  $I_{cs}(\Delta) = \int_0^1 \frac{dz_1}{z_1} \int_0^\infty \frac{dz_2}{z_2} (z_1 z_2)^{-2\epsilon} \delta(\Delta - z_1 z_2)$   
\n1, 2 soft:  $I_{ss}(\Delta) = \int_0^\infty \frac{dz_1}{z_1} \int_0^\infty \frac{dz_2}{z_2} (z_1 z_2)^{-2\epsilon} \delta(\Delta - z_1 z_2)$ 

### Backup *rapidity divergences*

$$
\Delta = z_1 z_2
$$
energy-sharing "rapidity"  

$$
s = z_1 z_2, \quad y = 1/2 \log(z_1/z_2)
$$

1 soft: 
$$
I_{sc}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \,\Theta\left(-\frac{1}{2}\ln(1/\Delta) < y\right)
$$
  
2 soft:  $I_{cs}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \,\Theta\left(y < \frac{1}{2}\ln(1/\Delta)\right)$   
1, 2 soft:  $I_{ss}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy$ 

### Backup *rapidity divergences*

$$
\Delta = z_1 z_2
$$
energy-sharing "rapidity"  

$$
s = z_1 z_2, \quad y = 1/2 \log(z_1/z_2)
$$

1 soft: 
$$
I_{sc}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \,\Theta\left(-\frac{1}{2}\ln(1/\Delta) < y\right) \left(\frac{\nu}{E_J}\right)^{\eta} s^{-\eta/2} e^{y\eta}
$$
  
\n2 soft:  $I_{cs}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \,\Theta\left(y < \frac{1}{2}\ln(1/\Delta)\right) \left(\frac{\nu}{E_J}\right)^{\eta} s^{-\eta/2} e^{-y\eta}$   
\n1, 2 soft:  $I_{ss}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \left(\frac{\nu}{E_J}\right)^{\eta} s^{-\eta/2} |2\sinh y|^{-\eta}$  rapidity regulators

$$
I_{\text{full}}(\Delta) = I_{sc}(\Delta) + I_{cs}(\Delta) + I_{ss}(\Delta)
$$