Fractional Jet Multiplicity

Daniele Bertolini UC Berkeley, Lawrence Berkeley National Laboratory



Outline

- New observable: fractional jet multiplicity
- This talk: discuss analytic predictions and unique features [DB, J.Thaler, and J.Walsh 1501.01965]

$$\widetilde{N}_{jet} = \sum_{i \in event} \frac{E_i}{E_{iR}} \Theta(E_{iR} > E_{cut})$$
$$E_{iR} = \sum_j E_j \Theta(\Delta \theta_{ij} < R)$$
$$= \text{energy in a cone around particle i}$$



Example of distributions in $e^+e^- \rightarrow \text{jets}$



Example of distributions in $e^+e^- \rightarrow \text{jets}$



I will describe analytically the near-integer distribution

Daniele Bertolini - SCET 2015

Collimated particles give integer number of jets



Consider three hard particles $E_1 = E_2 = E_3 > E_{cut}$

Collimated particles give integer number of jets



Consider three hard particles $E_1=E_2=E_3>E_{cut}$

$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$

 \mathbf{I}

Collimated particles give integer number of jets



$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$

= |/3 + |/3

Consider three hard particles $E_1 = E_2 = E_3 > E_{cut}$

Collimated particles give integer number of jets



$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$

= |/3+|/3+|/3=|

Consider three hard particles $E_1=E_2=E_3>E_{cut}$

Less collimated particles give fractional number of jets



Consider three hard particles $E_1=E_2=E_3>E_{cut}$

Less collimated particles give fractional number of jets



$$N_{jet} = \Sigma_i E_i / E_{iR}$$

= |/2+|/3+|/2=4/3

Consider three hard particles $E_1 = E_2 = E_3 > E_{cut}$

Soft limit drives the approach to the integer value

$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$
$$= 4/3 \approx 1.3$$

Consider now one particle going soft $E_1 = E_2 > E_{cut}$ and $E_3 \rightarrow 0$

Soft limit drives the approach to the integer value

$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$

$$\simeq 1.1$$

$$\simeq 1.1$$

Consider now one particle going soft $E_1 = E_2 > E_{cut}$ and $E_3 \rightarrow 0$

Soft limit drives the approach to the integer value

$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$
$$= I$$
$$z_3 = E_3/E_{tot} = 0$$

Consider now one particle going soft $E_1 = E_2 > E_{cut}$ and $E_3 \rightarrow 0$

Well separated clusters of particles give integer number of jets

Consider three hard particles $E_1 = E_2 = E_3 > E_{cut}$

Well separated clusters of particles give integer number of jets



$$\widetilde{N}_{jet} = \Sigma_i E_i / E_{iR}$$

= |/2+|/2+|=2

Consider three hard particles $E_1 = E_2 = E_3 > E_{cut}$

Fractional Jet Multiplicity properties

- To get non-integer behavior we need at least three particles. LO = $O(\alpha_s^2)$ in pQCD
- They have to be collimated (~ within 2R) and occupy special regions of phase space
- Near-integer behavior is driven by soft logarithms

Consider $e^+e^- \rightarrow \text{jets}$



 $\Delta_{2-} = 2 - \widetilde{N}_{jet}, \quad \Delta_{2+} = \widetilde{N}_{jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{jet}$ $\frac{d\sigma}{d\Delta_{2\pm}} = \int d\Phi_4 \, \mathcal{T}(e^+e^- \to 4 \text{ partons}) \, \mathcal{F}(\Delta_{2\pm}, \Phi_4)$



 $\Delta_{2-} = 2 - \widetilde{N}_{jet}, \quad \Delta_{2+} = \widetilde{N}_{jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{jet}$

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} = \int \mathrm{d}\Phi_4 \, \mathcal{T}(e^+e^- \to 4 \text{ partons}) \, \mathcal{F}(\Delta_{2\pm}, \Phi_4)$$

$$\mathcal{T}(e^+e^- \to 4 \text{ partons}) \simeq \mathcal{T}(e^+e^- \to q\bar{q}) \cdot \sum_k \mathcal{T}_k^{\mathrm{coll}}(1 \to 3)$$

$$k \in \{q \to ggq, \bar{q} \to gg\bar{q}, q \to q'\bar{q}'q, \bar{q} \to q'\bar{q}'\bar{q}\}$$



 $\Delta_{2-} = 2 - \widetilde{N}_{jet}, \quad \Delta_{2+} = \widetilde{N}_{jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{jet}$

Near-integer distribution

soft logarithms and rapidity divergences



Near-integer distribution

soft logarithms and rapidity divergences

Modes controlling near-integer behavior



Fixed order result

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} = \sigma_0 \delta(\Delta_{2\pm}) \\ + \kappa_1 \left(\frac{\alpha_s}{\pi}\right)^2 \mathcal{L}_1(\Delta_{2\pm}) + \kappa_0 \left(\frac{\alpha_s}{\pi}\right)^2 \mathcal{L}_0(\Delta_{2\pm}) + \text{non-singular terms}$$

- We calculated k_1 and k_0 (which include different color structures and leading dependence on $z_{cut} = E_{cut} / E_{tot}$)
- Get non singular from Event2

Event2 comparison:



Event2 comparison:



Daniele Bertolini - SCET 2015

Fractional jet multiplicity

- Near-integer driven by soft-logs only
- Hybrid event-shape / jet-algorithm behavior
- Non-additive / non-factorizable / non-global

Near-integer distribution *Hybrid event-shape / jet algorithm*



Near-integer distribution

additivity / factorizability / non-global logs

Additive	Factorizable	Global Logs		
Yes	Yes	Yes	Thrust	
Yes	Yes	No	Hemisphere Mass	
Yes	No	Yes	Jade Algorithm Rate	
Yes	No	No	$\sum_{i} E_i / Q \Theta(E_{iR} - E_{cut})$	
No	Yes	Yes	?	
No	Yes	No	?	
No	No	Yes	?	
No	No	No	fractional jet multiplicity	



$$C_{q,\bar{q}}(\Delta_{2\pm}) = \delta(\Delta_{2\pm}) + \sum_{n=2}^{\infty} \left(\frac{\alpha_s}{\pi}\right)^n \sum_{k=-1}^{n-1} \left[\kappa_{k,+}^{(n)} \mathcal{L}_k(\Delta_{2\pm}) + \kappa_{k,-}^{(n)} \mathcal{L}_k(\Delta_{2\pm})\right]$$

 $\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} \simeq \sigma_0 \, C_q(\Delta_{2\pm}) \otimes C_{\bar{q}}(\Delta_{2\pm})$



Wide angle soft emissions give enhanced logarithmic contributions to \widetilde{N}_{jet} =2 cross-section

 $\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} \simeq \sigma_0 \, C_q(\Delta_{2\pm}) \otimes C_{\bar{q}}(\Delta_{2\pm})$



At higher orders they also enhance fractional \widetilde{N}_{jet} cross-section

 $\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} \simeq \sigma_0 \, C_q(\Delta_{2\pm}) \otimes C_{\bar{q}}(\Delta_{2\pm})$



Note that this contribution is contained in C

Near-integer distribution *a candidate factorization theorem*

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} \simeq \sigma(\widetilde{N}_{\mathrm{jet}} = 2) \left[C_q(\Delta_{2\pm}) \otimes C_{\bar{q}}(\Delta_{2\pm}) \right]$$

Near-integer distribution *a candidate factorization theorem*

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Delta_{2\pm}} \simeq \sigma(\widetilde{N}_{\mathrm{jet}} = 2) \left[C_q(\Delta_{2\pm}) \otimes C_{\bar{q}}(\Delta_{2\pm}) \right]$$

$$\downarrow$$

$$\sigma(\widetilde{N}_{\mathrm{jet}} = 2) = \sigma_0 H_{q\bar{q}}(Q,\mu) J_q(Q,R,z_{\mathrm{cut}},\mu) J_{\bar{q}}(Q,R,z_{\mathrm{cut}},\mu) S_{q\bar{q}}(R,z_{\mathrm{cut}},\mu)$$

$$+ \sigma_2^{\mathrm{non-fac}}(Q,R,z_{\mathrm{cut}},\mu)$$

Improved distributions

- Include $O(\alpha_s^4)$ terms from convolutions
- Running coupling $\alpha_s(\mu), \mu = Q\sqrt{\Delta_{2\pm}}$







Daniele Bertolini - SCET 2015

Conclusions

- Fractional jet multiplicity can be used as a novel and more powerful probe of jet formation. E.g. quark/gluon discrimination?
- It has peculiar analytic properties. However, I showed we still have very good analytic control, which in principle is improvable. E.g. generalize to LHC case?
- Wide dynamic range, potential test of matching/merging matrix-element/parton shower.
 Potential phenomenological applications e.g. in multijet final states?

Backup

near-integer phase space configurations

$$\Delta_{2-} = 2 - \widetilde{N}_{jet}, \quad \Delta_{2+} = \widetilde{N}_{jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{jet}$$

Observable	\mathcal{R}_A	\mathcal{R}_B	\mathcal{R}_C
Δ_{2-}			_
Δ_{2+}		$ \begin{array}{c} g_2 \\ $	$ \begin{array}{c} g_1 & g_2 \\ \bullet & \bullet \\ \end{array} $ $ \begin{array}{c} g_1 & g_3 & g_2 \\ \bullet & \bullet \\ \end{array} $ $ \begin{array}{c} g_1 & g_3 & g_2 \\ \bullet & \bullet \\ \end{array} $
Δ_{3-}			_

Backup

near-integer phase space configurations

$$\Delta_{2-} = 2 - \widetilde{N}_{jet}, \quad \Delta_{2+} = \widetilde{N}_{jet} - 2, \quad \Delta_{3-} = 3 - \widetilde{N}_{jet}$$

Observable	Region	Expression	Limit	Cuts
Δ_{2-}	\mathcal{R}_A	$z_1 z_2$	1, 2 soft	$z_1 + z_2 < z_{\rm cut}$
Δ_{2-}	\mathcal{R}_B	$z_1 z_2$	1, 2 soft	$z_1 + z_2 < z_{\rm cut}$
Δ_{2+}	\mathcal{R}_A	$z_1(1-z_2^2)/z_2$	1 soft	$z_2 > z_{\rm cut}$
Δ_{2+}	\mathcal{R}_B	$z_2(1-z_1^2)/z_1$	2 soft	$z_1 > z_{\rm cut}$
Δ_{2+}	\mathcal{R}_C	$z_1 z_2 (2 - z_2) / (1 - z_2)$	1 soft	_
Δ_{2+}	\mathcal{R}_C	$z_1 z_2 (2 - z_1) / (1 - z_1)$	2 soft	_
Δ_{2+}	\mathcal{R}_C	$2z_1z_2$	1, 2 soft	_
Δ_{3-}	\mathcal{R}_A	$z_2[1-z_1(1-z_1)]/[z_1(1-z_1)]$	2 soft	$z_1 > z_{\rm cut}$
Δ_{3-}	\mathcal{R}_B	$z_1[1-z_2(1-z_2)]/[z_2(1-z_2)]$	1 soft	$z_2 > z_{ m cut}$

Backup rapidity divergences

 $\Delta = z_1 z_2$

1 soft:
$$I_{sc}(\Delta) = \int_{0}^{\infty} \frac{\mathrm{d}z_{1}}{z_{1}} \int_{0}^{1} \frac{\mathrm{d}z_{2}}{z_{2}} (z_{1}z_{2})^{-2\epsilon} \,\delta(\Delta - z_{1}z_{2})$$

2 soft: $I_{cs}(\Delta) = \int_{0}^{1} \frac{\mathrm{d}z_{1}}{z_{1}} \int_{0}^{\infty} \frac{\mathrm{d}z_{2}}{z_{2}} (z_{1}z_{2})^{-2\epsilon} \,\delta(\Delta - z_{1}z_{2})$
1, 2 soft: $I_{ss}(\Delta) = \int_{0}^{\infty} \frac{\mathrm{d}z_{1}}{z_{1}} \int_{0}^{\infty} \frac{\mathrm{d}z_{2}}{z_{2}} (z_{1}z_{2})^{-2\epsilon} \,\delta(\Delta - z_{1}z_{2})$

Backup rapidity divergences

$$\Delta = z_1 z_2$$
 energy-sharing "rapidity"
 $s = z_1 z_2, \quad y = 1/2 \log(z_1/z_2)$

1 soft:
$$I_{sc}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \,\Theta\left(-\frac{1}{2}\ln(1/\Delta) < y\right)$$

2 soft: $I_{cs}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy \,\Theta\left(y < \frac{1}{2}\ln(1/\Delta)\right)$
1, 2 soft: $I_{ss}(\Delta) = \Delta^{-1-2\epsilon} \int_{-\infty}^{\infty} dy$

Backup rapidity divergences

$$\Delta = z_1 z_2$$
 energy-sharing "rapidity"
 $s = z_1 z_2, \quad y = 1/2 \log(z_1/z_2)$

$$I_{\text{full}}(\Delta) = I_{sc}(\Delta) + I_{cs}(\Delta) + I_{ss}(\Delta)$$