

Probing TeV LR Seesaw for Neutrino Mass and Origin of matter

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Nu@Fermilab, 2015

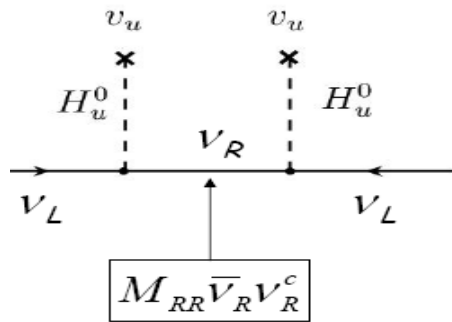
Weinberg operator and naïve lore about scale of L-violation

$$\text{SM+} \quad \lambda \frac{LHLH}{M} \rightarrow m_\nu = \lambda \frac{v_{wk}^2}{M} \rightarrow \text{nu Majorana}$$

- Neutrino osc data $\rightarrow m_\nu \ll eV$
- **Scale of L violation:** $\sim 10^{14}$ GeV for $\lambda \sim 1$
- Dimensional analysis arguments, however, can be quite misleading (*e.g. $K_L - K_S$ mass diff.*)!!
- To explore true scale, UV completion of Weinberg operator essential (*build models*) !!

Seesaw as step towards UV completion of Weinberg Op.

- Add right handed N and a Majorana mass for it:
Seesaw mechanism:



$$m_\nu \cong - \frac{h_\nu^2 v_{wk}^2}{M_R}$$

Minkowski'77, Gell-Mann, Ramond, Slansky; Yanagida; Glashow; Mohapatra, Senjanovic'79

- Majorana mass of N \rightarrow Majorana nu (SM seesaw)
- **Major bonus:** Leptogenesis as origin of matter from N-decay in combination with sphalerons.
- Small h, lower seesaw scale: (talks by: Deppisch, Dev, Lopez-Pavon, Ruiz, Molinaro@INFO2015)



TeV seesaw beyond SM

- Search for BSM UV complete seesaw models
- Guiding principle (*assume as little as possible*)
 - (i) *Existence of N should be predicted by theory*
 - (ii) *Seesaw scale be related to local symmetry*
- Two simple theories that conform to these:
 - (i) Left-right model where N is the parity partner of ν_L and seesaw scale is $SU(2)_R$ scale could be TeV
 - (ii) SO(10) GUT where $N+15$ SM fermions = 16 spinor and seesaw scale = GUT scale. (*Hard to test*)



This talk: TeV LR seesaw

- A “natural” TeV scale theory for neutrinos
- Minimal SUSY LR *requires* TeV scale L-violation
- How to probe this TeV scale theory in colliders
- Leptogenesis with TeV scale ~~L~~ and constraints



Left-Right Model Basics

- LR basics: Gauge group: $SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$

- Fermions

$$\begin{pmatrix} u_L \\ d_L \end{pmatrix} \stackrel{P}{\Leftrightarrow} \begin{pmatrix} u_R \\ d_R \end{pmatrix} \quad \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} \stackrel{P}{\Leftrightarrow} \begin{pmatrix} \nu_R \\ e_R \end{pmatrix}$$

■

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$$L = \frac{g}{2} [\vec{J}_L^\mu \cdot \vec{W}_{\mu L} + \vec{J}_R^\mu \cdot \vec{W}_{\mu R}]$$

- Parity is spontaneously broken symmetry: $M_{W_R} \gg M_{W_L}$
(Mohapatra, Pati, Senjanovic'74-75)



Other advantages of LR

- A more physical electric charge formula than SM
- Solves strong CP problem without the axion and limits $M_{WR} < 1000$ TeV.
- With supersymmetry, provides a naturally stable dark matter (automatic R-parity)

New Higgs fields and Yukawa couplings

- LR bidoublet: $\phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}$

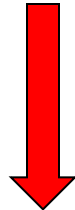
- Triplet to break B-L and generate seesaw: $\Delta = \begin{pmatrix} \frac{1}{\sqrt{2}} \Delta^+ & \Delta^{++} \\ \Delta^0 & -\frac{1}{\sqrt{2}} \Delta^+ \end{pmatrix}$

$$\mathcal{L}_Y = h \bar{L} \phi R + \tilde{h} \bar{L} \tilde{\phi} R + f R R \Delta_R + h.c.$$

$$\langle \Delta_R \rangle = \begin{pmatrix} 0 & 0 \\ v_R & 0 \end{pmatrix} \quad \phi = \begin{pmatrix} \kappa & 0 \\ 0 & \kappa' \end{pmatrix}$$

Seesaw scale is $SU(2)_R$ breaking Scale

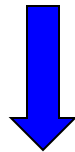
$$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$



$$\nu_R \quad (\Delta L=2)$$

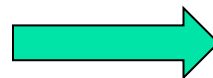
$$M_N = f \nu_R$$

$$SU(2)_L \times U(1)_Y$$



$$\kappa$$

$$U(1)_{em}$$



$$M_{\nu, N} = \begin{pmatrix} 0 & h\kappa \\ h\kappa & f\nu_R \end{pmatrix}$$

Seesaw

$$m_\nu \simeq -\frac{(h\kappa)^2}{M_N}$$

- If $\nu_R \sim \text{TeV}$, L-violation is TeV scale
- Any theoretical justification for TeV ν_R ?

SUSY as theory justification for TeV W_R

- Supersymmetrize this minimal LR model
- First consequence: Tree level global minimum violates electric charge: $\langle \Delta^{++} \rangle \neq 0$

(i) unless R-parity is broken (Kuchimanchi, R. N. M.'94, '95)

■ +

(ii) W_R mass has an upper limit:

$$M_{W_R} \leq \frac{M_{SUSY}}{f}$$

i.e. W_R is in TeV range !

Minimal SUSYLR with exact R-parity

- Extend with a singlet and **add one loop** → RP exact !

(Babu, R. N. M.'08; Babu, Patra'14; Basso, Fuks, Krauss, Porod'15)

- *Upper bound on W_R required to conserve electric charge; $M_{WR} < 7 \text{ TeV}$* (Porod et al., private communications)

- *Implies a light ($< \text{TeV}$) doubly charged Higgs*

$$M_{\Delta^{++}} \leq 2 \text{ TeV} \quad (\text{Porod et al.})$$

- *Neutrino masses from usual seesaw*

Seesaw formula in TeV LR models: Type II small

■ Generic LR models with parity down to TeV,

→ Seesaw formula $m_\nu \simeq f \frac{\kappa^2}{v_R} - m_D^T \frac{1}{f v_R} m_D$

(apriori large)

Seesaw formula in TeV LR models: Type II small

Generic LR models with parity down to TeV,

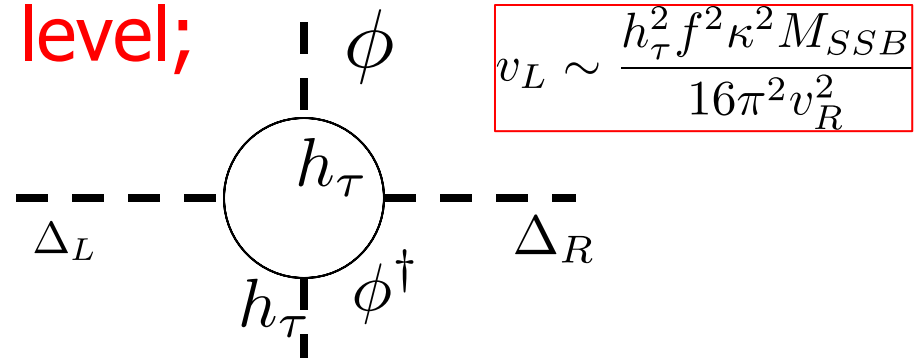
→ Seesaw formula
$$m_\nu \simeq f \frac{\kappa^2}{v_R} - m_D^T \frac{1}{f v_R} m_D$$

Two theories where first term is small:

(i) decouple P breaking from $SU(2)_R$ $g_R < g_L$

(ii) SUSYLR → zero at tree level;

1-loop small



Small Neutrino masses with TeV WR

- $\mathcal{L}_\gamma = h\bar{L}\phi R + \tilde{h}\bar{L}\tilde{\phi}R + h.c.$
- Using $\phi = \begin{pmatrix} \kappa & 0 \\ 0 & \kappa' \end{pmatrix} \rightarrow \begin{aligned} M_\ell &= h\kappa + \tilde{h}\kappa' \\ m_D &= h\kappa' + \tilde{h}\kappa \end{aligned}$
- How to get small m_ν for TeV seesaw:
 - $\kappa' = 0; \tilde{h} \sim 10^{-5.5} (\sim h_e^{SM})$
 - Cancellation with κ', κ similar
 - assume texture for Dirac mass*

$\left. \begin{array}{l} h, \tilde{h} \\ \text{much larger} \end{array} \right\}$

Making TeV scale seesaw "natural" Case (iii)

■ Neutrino Mass texture: (for SM → Kersten, Smirnov; Pilaftsis, Underwood)

$$m_D = \begin{pmatrix} m_1 & \delta_1 & \epsilon_1 \\ m_2 & \delta_2 & \epsilon_2 \\ m_3 & \delta_3 & \epsilon_3 \end{pmatrix} \quad M_N = \begin{pmatrix} 0 & M_1 & 0 \\ M_1 & 0 & 0 \\ 0 & 0 & M_2 \end{pmatrix}$$

$$m_{D_{1,2,3}} \sim \text{GeV} \rightarrow Y_\nu \sim 10^{-2} - 10^{-4}$$

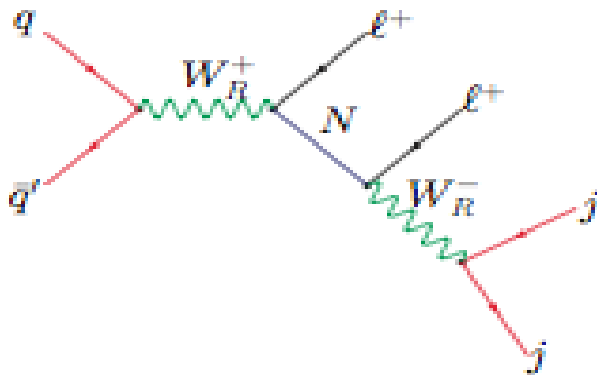
- Sym limit $\epsilon_i, \delta_i \rightarrow 0 \rightarrow m_\nu = m_D M_R^{-1} m_D^T = 0$
- sym. Br. $\delta_i, \epsilon_i \ll m_i \rightarrow$ for TeV $M_R \rightarrow$ small m_ν
- Small δ, ϵ arise from one loop SUSY breaking effects;
Good fit to neutrinos (Dev, Lee, RNM'13)

Probing TeV LR:

(i) WR Signals at LHC

$l_i l_k j j$

(a) LR Seesaw signals at LHC



$$N \rightarrow l^\pm jj$$

(Keung, Senjanovic'83;
Gunion, Kayser'84)

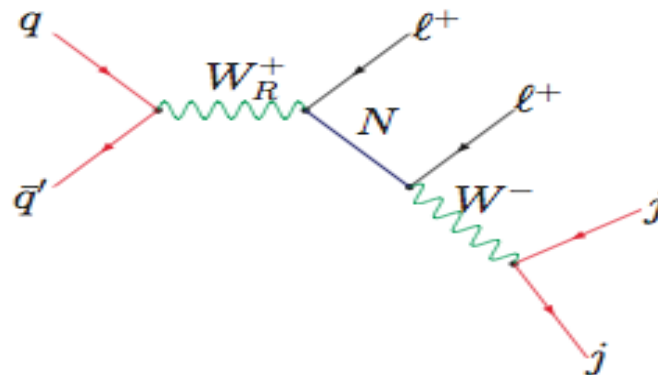
$$A_{l+l+jj} \propto M_{N,ik}^{-1}$$

(b)

(Tello, Nemevsek, Senjanovic; Chen, Dev, RNM)

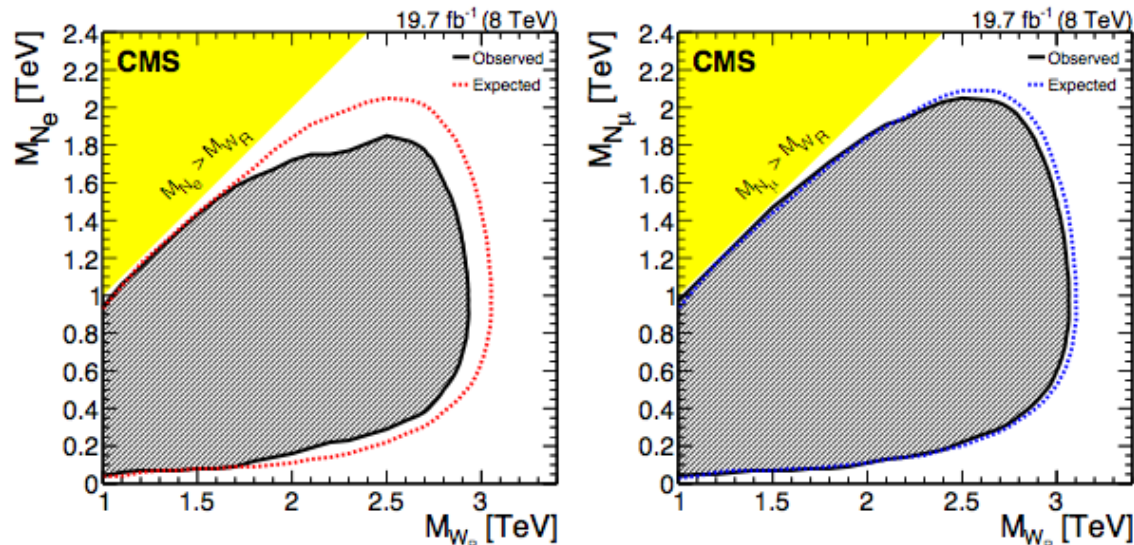
$$q\bar{q} \rightarrow W_R \rightarrow l + N;$$

$$N \rightarrow lW_L$$



Current LHC analysis: only W_R graph

■ Current W_R limits from CMS, ATLAS 2.9 TeV;



- 14-TeV LHC reach for $M_{W_R} < 6$ TeV with 300 fb⁻¹
- Higgs signals at LHC (Nemevsek@CETUP2015)

LHC anomalies ~ 2 TeV

$(W_R?)$

- 3.4σ $WZ \rightarrow JJ$ excess (ATLAS)
- CMS JJ excess 1.8σ excess
- 2.2σ Wh excess (ATLAS)
- 2.8σ $eejj$ excess (CMS)
- 2.6σ excess WW and ZZ channel (ATLAS)

W_R interpretation of LHC anomalies



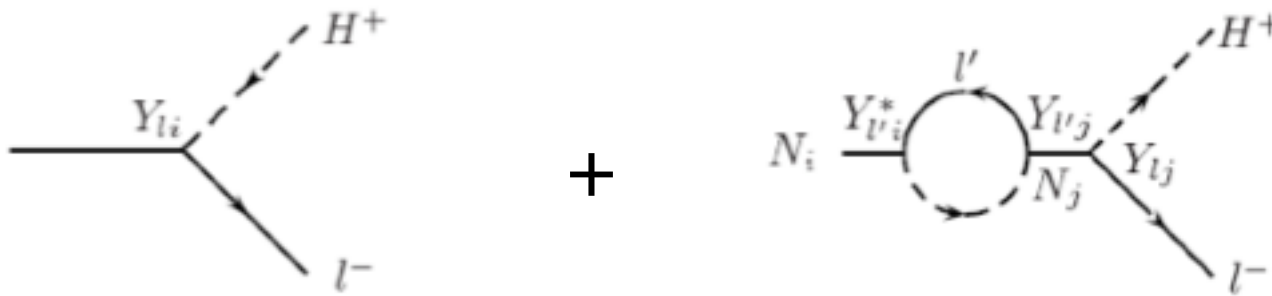
- $eejj$: (Deppisch Gonzalo, Patra, sahu, Sarkar; Heikinheimo, Raidal, Spethman; Aguilar-Saavedra, Joachim; Fowlie, Marzola; Gluza, Jelinsky)
- Diboson+...: (Hisano et al. Dobrescu, Liu; Gao, Ghosh, Sinha, Yu; Cheung, Keung et al; Cao, Dong, Zhang; Bremmer, Hewett, Kopp, Rizzo, Tattersal; Krauss, Porod)
- Telling W_R from W' : (Han, Lewis, Ruiz, Si'12)

Does Leptogenesis work in TeV W_R models

- Since $m_\nu \simeq \frac{(Y\kappa)^2}{fv_R}$, TeV v_R means $Y \leq 10^{-5.5}$
- or larger Y with Texture
- Either case $\epsilon_{CP} \sim \frac{\text{Im}(Y^\dagger Y)^2}{4\pi Y^\dagger Y} \sim 10^{-12} - 10^{-10}$
- since $\frac{n_B}{n_\gamma} \sim 10^{-2} \epsilon_{CP} \kappa_{eff}$ (κ_{eff} = wash out)
- need enhancement \rightarrow suggest resonant leptogenesis

(II): TEV SCALE RESONANT LEPTOGENESIS:

RH neutrino mass \sim TeV scale



$$\frac{n_B}{n_\gamma} \propto \frac{\text{Im} Y^4}{|Y|^2} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 + M_2 \Gamma_2)^2}$$

Generic model requires extreme degeneracy among RHNs to get enough n_B/n_γ

Final baryon asymmetry from lepton asymmetry

■ Wash out effect important: (Buchmuller, Di Bari, Plümacher)

$$\frac{n_B}{n_\gamma} \sim 10^{-2} \epsilon_{CP} \kappa_{eff}$$

■ In LR, $\kappa_{eff} \propto \frac{\Gamma_D/\Gamma_S}{1 + \Gamma_D/\Gamma_S} \ll 1$

$$\Gamma_D \propto Y^2$$
$$\Gamma_S \propto M_{WR}^{-4}$$

■ Given Y, Washout increases as M_{WR} decreases:
→ lower bound on M_{WR}

■ Two papers: small Y: $M_{WR} > 18$ TeV (Frere, Hambye, Vertongen)

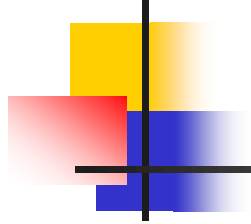
Larger Y with nu fits: $M_{WR} > 10$ TeV (Dev, Lee, RNM.'14)

■ *LHC can rule out leptogenesis idea !!*



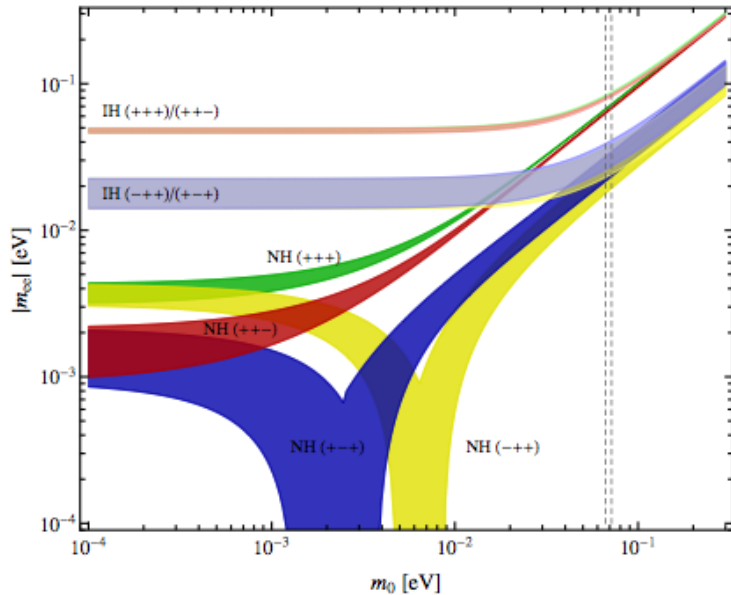
Summary

- Left-Right theories provide a simple realization of TeV scale seesaw for neutrino mass and leptogenesis with testable collider implications (W_R, Z', N_{\dots})!
- Minimal susy LR bound on $\rightarrow M_{WR} < \text{multi-TeV}$
- Leptogenesis bound on $W_R \rightarrow M_{WR} > 10 \text{ TeV}$
- If colliders find W_R with mass $< 10 \text{ TeV}$ or $M_{WR} < M_N$ leptogenesis can be ruled out.
- Another direction: Inverse seesaw in TeV LR models



Thank you for your attention !

Neutrino mass and $\beta\beta_{0\nu}$



(well-known IH bounds)

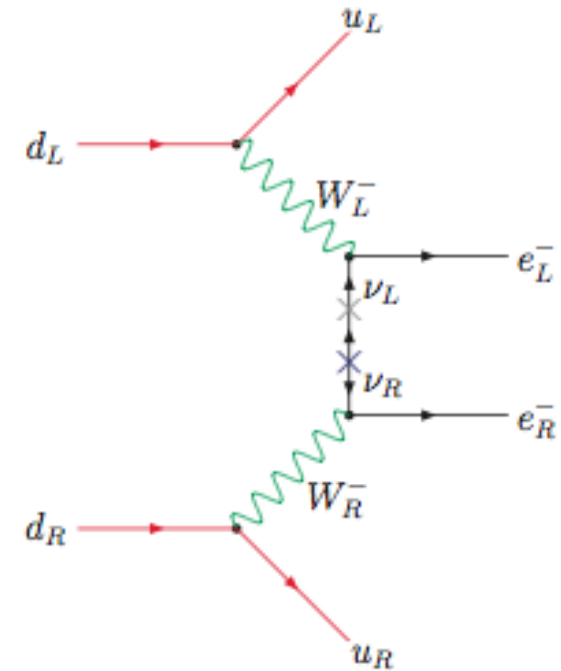
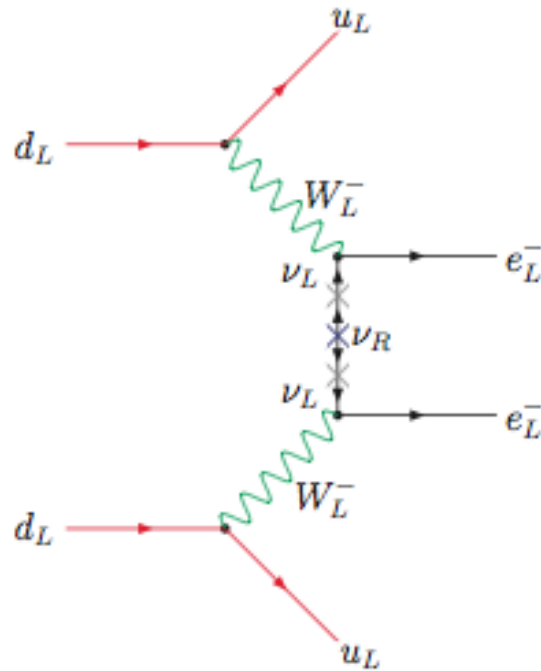
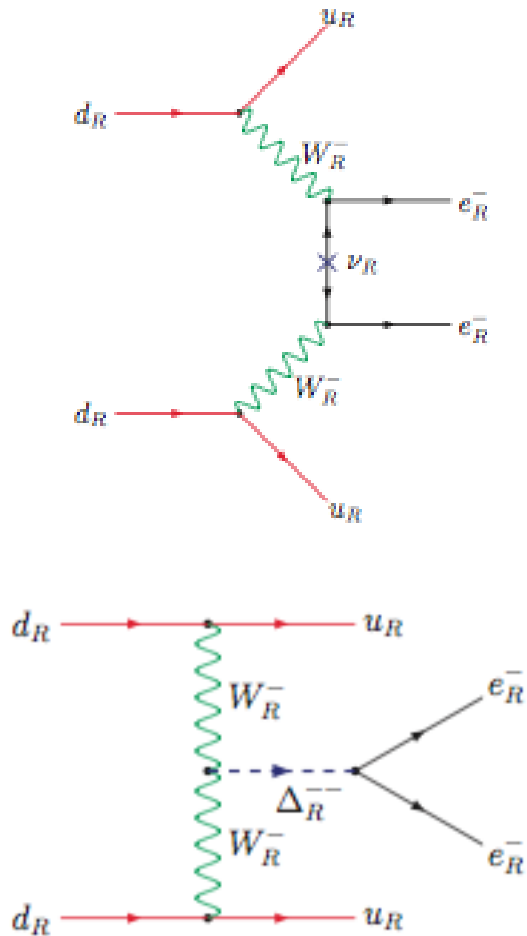
(Vissani'99; Bilenky,Pascoli,Petcov'01)

Two points to emphasize:

(i) lower bounds even for NH with sym (RNM, Nishi; 1506.)

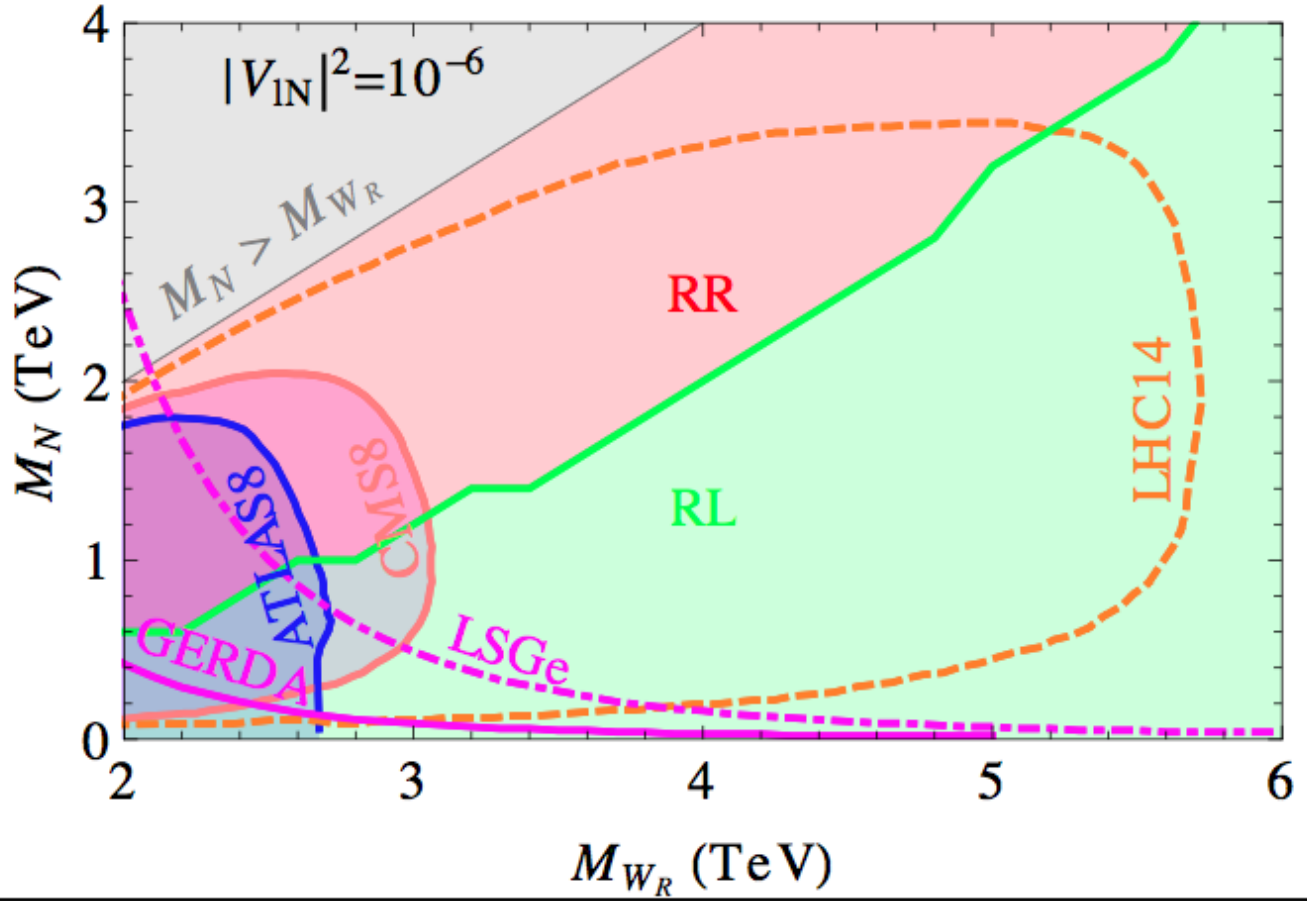
(ii) Heavy particle effects can "fake" IH (e.g. WR)

(ii) New contributions to $\beta\beta_{0\nu}$ in LR models



LHC and double beta reach for W_R

Dev, et al.



(Das, Deppisch, Kittel, Valle)

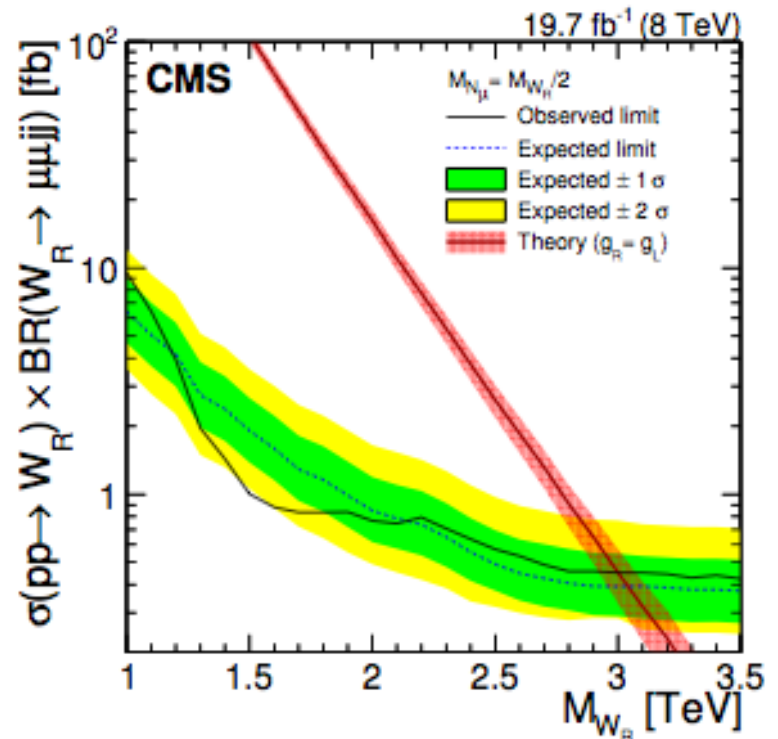
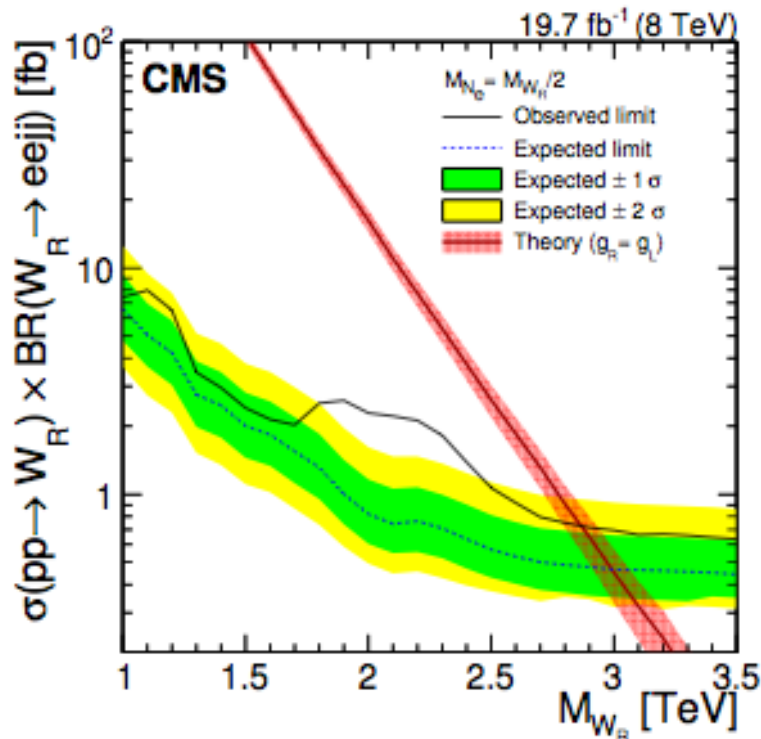


Case of $M_N > M_{WR}$

- CP conserving decay mode $N \rightarrow W_R + \ell$ dominates !
- Leptogenesis impossible (Deppisch, Harz, Hirsch'14)
- If experimentally it is found, $M_N > M_{WR}$, this by itself can rule out leptogenesis as a mechanism for origin of matter !!

Intriguing excess in CMS

■ . CMS: arXiv:1407.3683



■ Possible $M_{W_R} = 2.1$ TeV ?

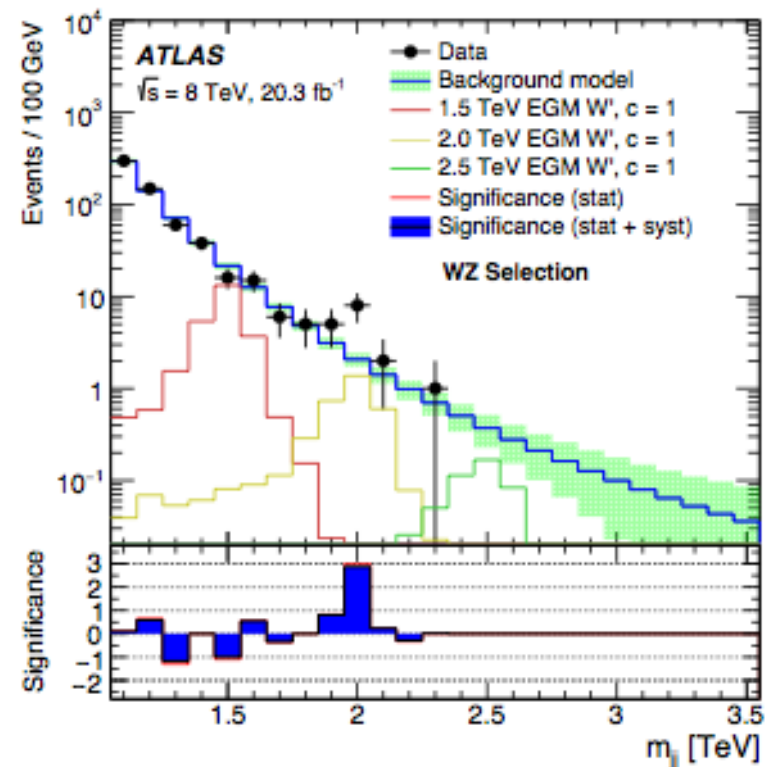
(Deppisch Gonzalo, Patra, sahu, Sarkar; Heikinheimo, Raidal, Spethman; Aguilar-Saavedra, Joachim; Fowlie, Marzola'14; Gluza, Jelinsky'15)

ATLAS Diboson anomaly

- Another W_R decay mode: $W_R \rightarrow W_L Z$ (via WL-WR mixing)
- Could it be connected to ATLAS diboson anomaly around 2-2.3 TeV?

arXiv:1506.00962

- Possibly a CMS $W_R \rightarrow Wh$ anomaly ?



LHC anomalies and LR interpretation (~ 2 TeV)

- 2 TeV W_R : $\sigma(W_R) \times B_{W_R \rightarrow WZ} \approx 600 g_R^2 B$ fb
- If no leptons $\rightarrow M_{N_{e,\mu}} \geq M_{W_R}$
- WZ channel signal at the level of 6-7 fb \rightarrow arises from $W_L - W_R$ mixing, corresponds to $\zeta_{LR} \sim 0.01$
- *Signal fits for $g_R \sim 0.5 g_L \rightarrow \sim 8$ excess events* **jjjj**
- *Predicts $\sim 2-3$ excess events in the Wh^0 channel – consistent with CMS excess for this channel.* **$b\bar{b}\ell\nu$**
- *Should not see any signal in WW and ZZ mode.*

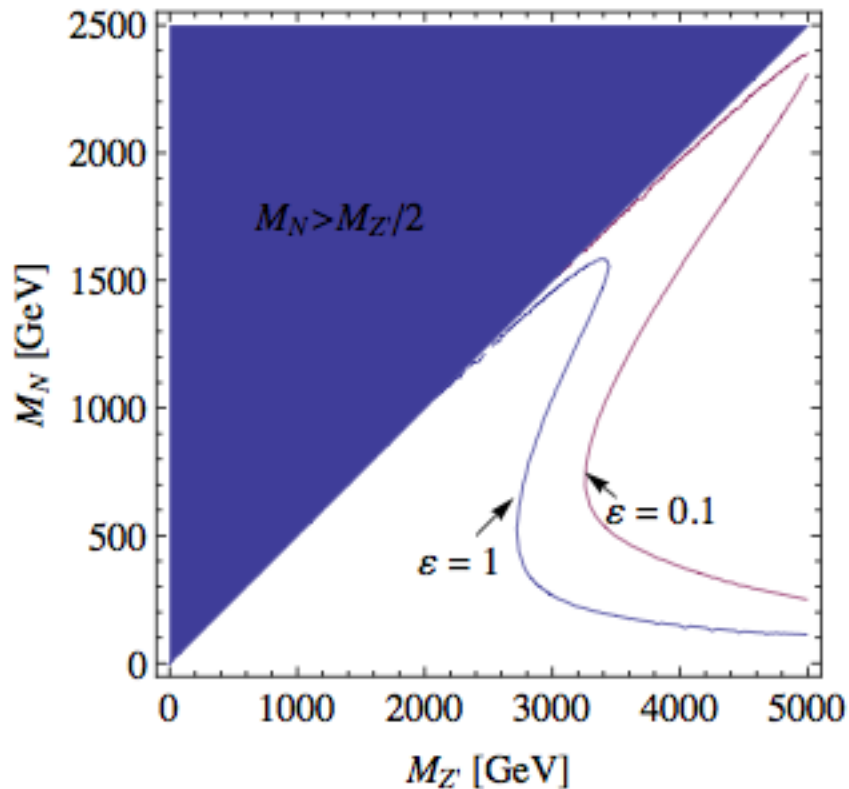


Leptogenesis with $M_{Z'} \ll M_{W_R}$

- Effective theory: $SU(2)_L \times U(1)_{I_{3R}} \times U(1)_{B-L}$
- Z' couples also to NN and effects leptogenesis
- Origin of CP asymmetry same as in WR case via resonant leptogenesis and requires deg $N_{1,2}$:
 \mathcal{E} can be as large as 1.
- Washout has no W_R contribution but only
 $NN \rightarrow Z' \rightarrow qq, ll$ type.
- Lower the Z' , more washout in generic case

Lower bound on $M_{Z'}$

(Blanchet, Chacko, Granor, RNM'2009, PRD)



$M_{Z'} > 3 \text{ TeV}$

Directly probing leptogenesis in Z' case:

- Lepton asymmetry \mathcal{E} is directly related to the following collider observable:

$$\frac{N(\ell^+\ell^+) - N(\ell^-\ell^-)}{N(\ell^+\ell^+) + N(\ell^-\ell^-)} = \frac{2 \sum_i \mathcal{E}_i}{\sum_i 1}$$

- Makes it possible to see origin of matter directly.

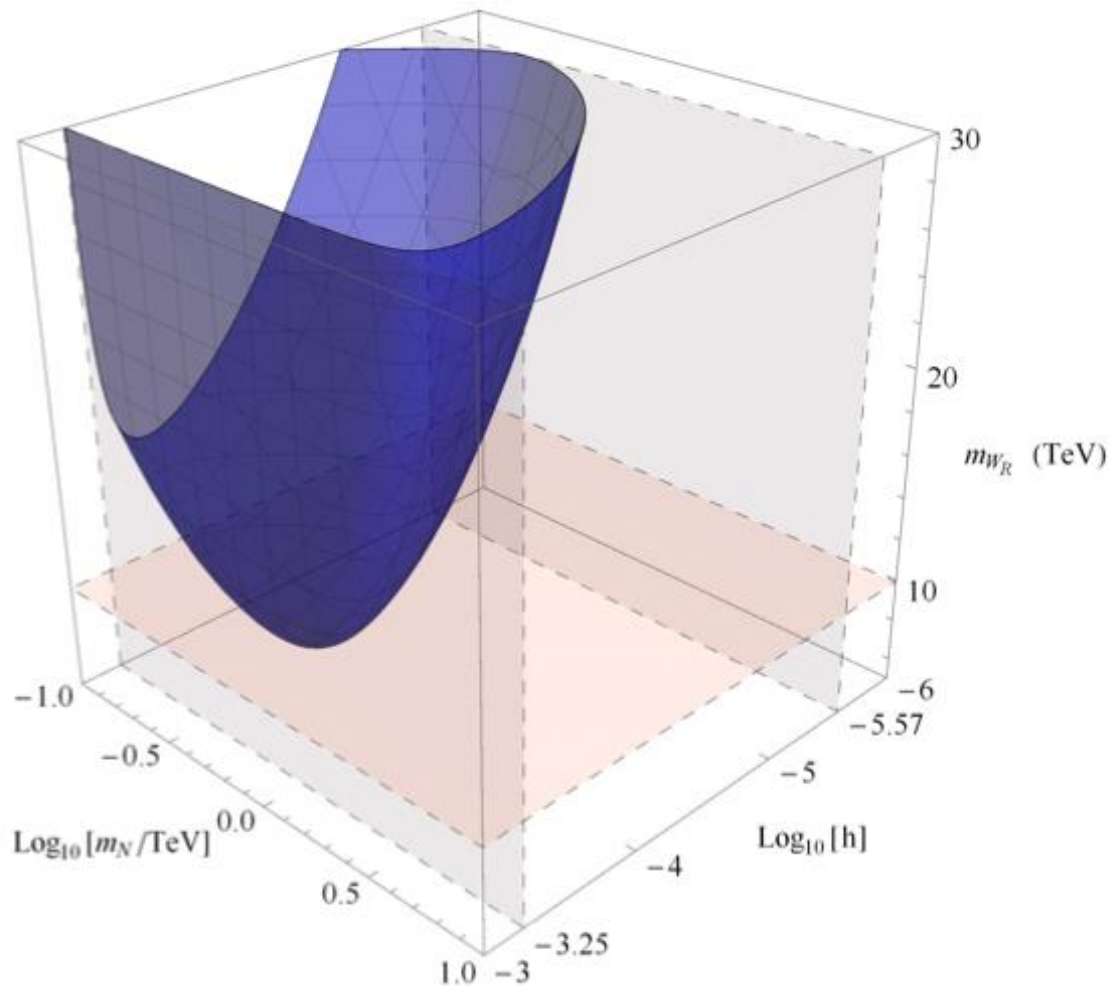
Distinguishing different mechanisms (RR vs RL)

- Look for end points in various inv. Masses:

| | <i>RL</i> | <i>LL</i> | <i>RR</i> | <i>LR</i> |
|----------------|--|--|---------------------|-------------|
| m_{ij}^2 | m_W^2 | m_W^2 | m_N^2 | m_N^2 |
| m_{ll}^2 | $\frac{(m_{W_R}^2 - m_N^2)(m_N^2 - m_W^2)}{m_N^2}$ | $\frac{(s - m_N^2)(m_N^2 - m_W^2)}{m_N^2}$ | $m_{W_R}^2 - m_N^2$ | $s - m_N^2$ |
| $m_{jl}^{>2}$ | $m_N^2 - m_W^2$ or | $m_N^2 - m_W^2$ or | m_N^2 or | m_N^2 or |
| $m_{jl}^{<2}$ | $m_{W_R}^2 - m_N^2$ | $s - m_N^2$ | $m_{W_R}^2 - m_N^2$ | $s - m_N^2$ |
| m_{llj}^2 | $m_{W_R}^2 - m_W^2$ | $s - m_W^2$ | $m_{W_R}^2$ | s |
| $m_{ljj}^{>2}$ | m_N^2 or | m_N^2 or | $m_{W_R}^2$ | s |
| $m_{ljj}^{<2}$ | $m_{W_R}^2 - m_N^2 + m_W^2$ | $s - m_N^2 + m_W^2$ | m_N^2 | m_N^2 |
| m_{uljj}^2 | $m_{W_R}^2$ | s | $m_{W_R}^2$ | s |

(Kim, Dev, RNM'15)

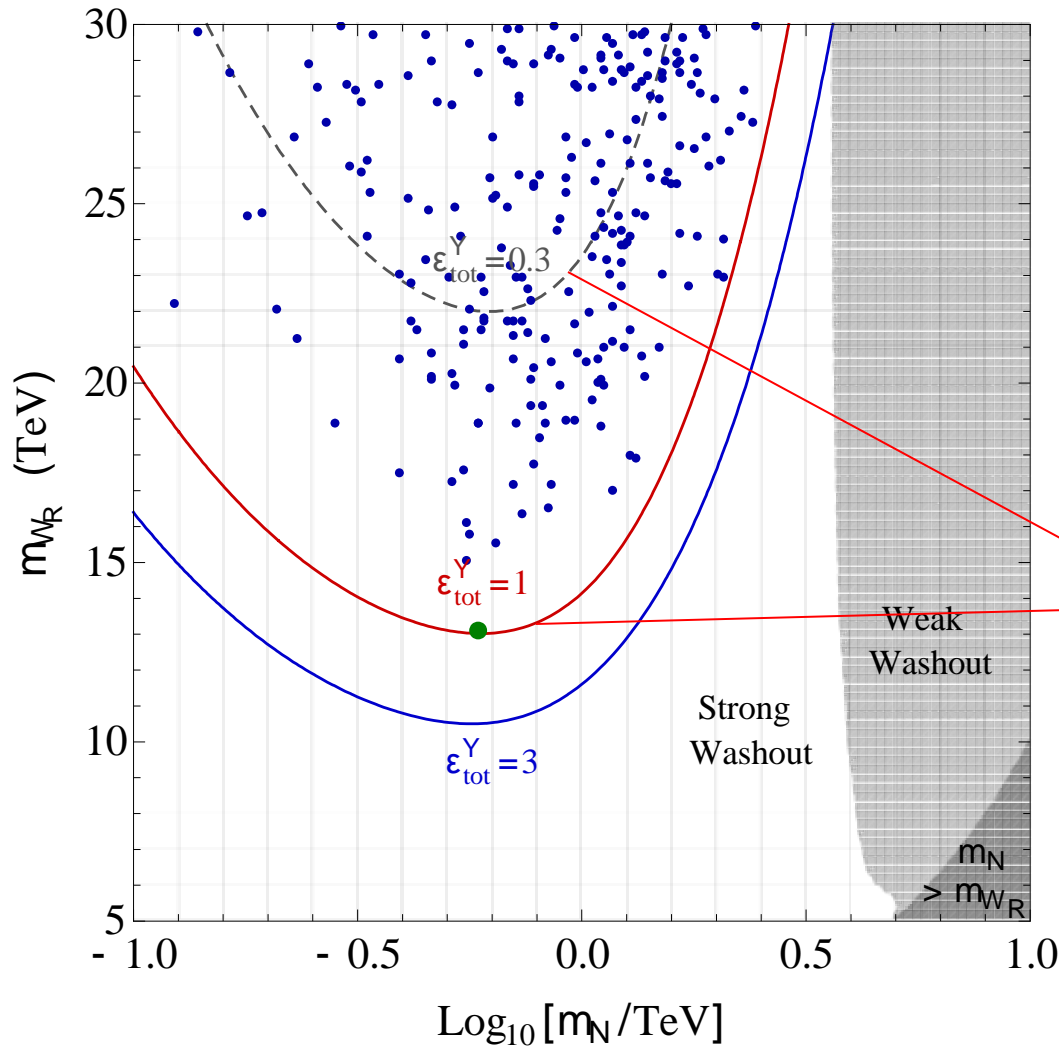
Low scale Leptogenesis Plot



$$M_{\text{WR}} > 10 \text{ TeV}$$
$$M_{\text{N}} > 585 \text{ GeV}$$

(Dev., Lee and RNM'15)

M_{WR} vs M_N Plot where leptogenesis works



(Dev, Lee, RNM'15)

$M_{WR} > 10$ TeV

$M_N > 585$ GeV

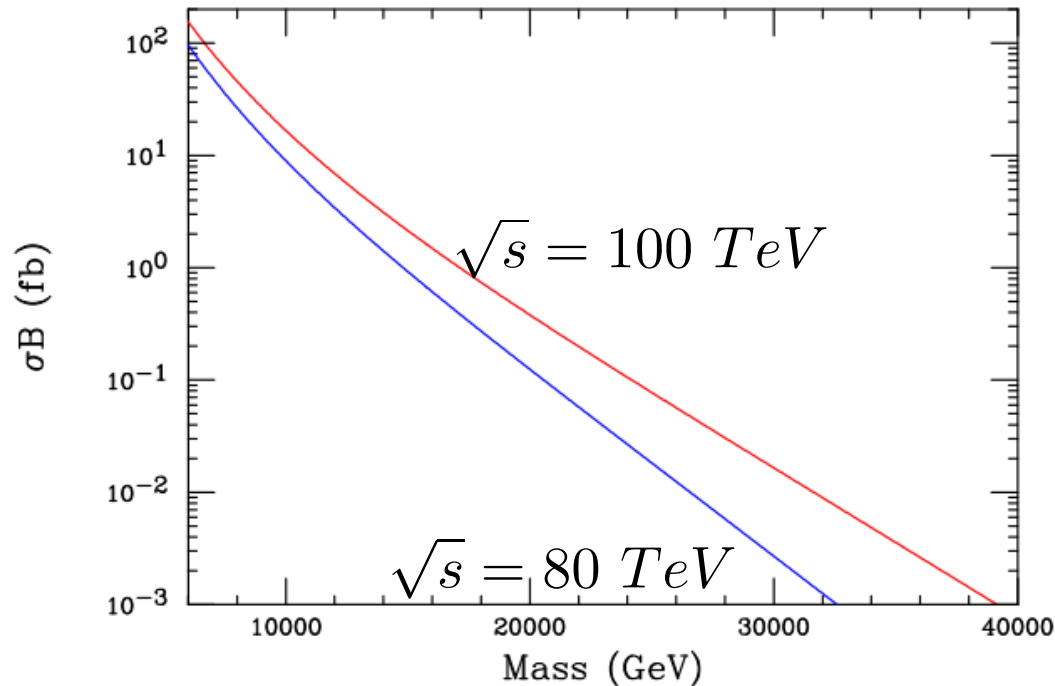
Weak
Washout

Strong
Washout

Explicit models
with nu mass fits.

Higher Mass WR probe at Future Circular colliders

- So far one study by **Rizzo**: $W_R \rightarrow \ell + \nu'$ channel



$$M_{WR} < 30 \text{ TeV}$$

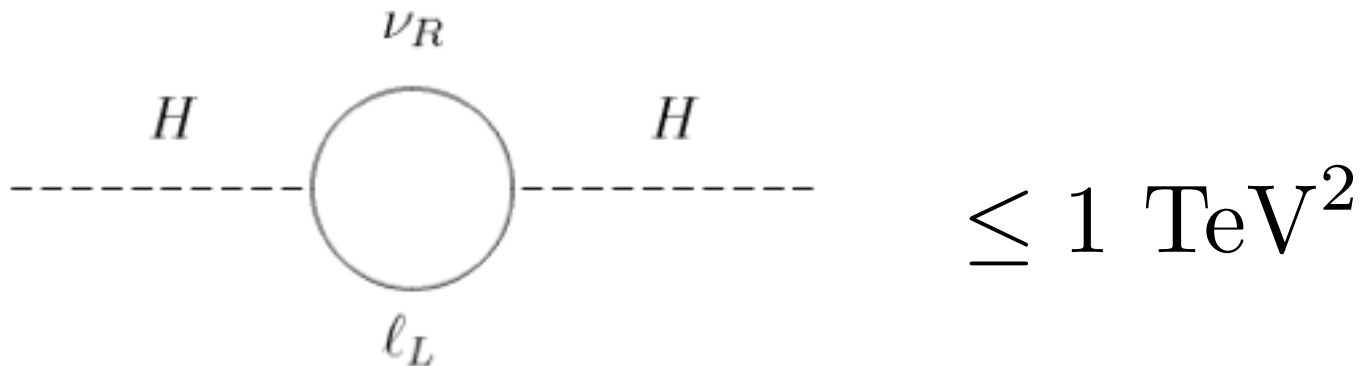
- For the $\ell^\pm \ell^\pm jj$ channel, see **Ng, Puento, Pan'15**

Right handed neutrino mass restricted by low energy obs.

- Low scale seesaw $\rightarrow \Delta_R$ masses below 10 TeV
- $\mu \rightarrow 3e, \mu \rightarrow e + \gamma, \tau \rightarrow 3e$ etc.
bounds restrict flavor structure of Δ_R coupling f and hence RHN mass texture $M_N = f v_R$!!
- One (only) allowed texture: $M_N = \begin{pmatrix} 0 & M_1 & 0 \\ M_1 & 0 & 0 \\ 0 & 0 & M_2 \end{pmatrix}$

Naturalness arguments for lower Seesaw scale

- Correction to Higgs mass from RHN Yukawa



→ $M_R < 7 \times 10^7 \text{ GeV}$ (not a GUT scale)

(Vissani'97; Clarke, Foot, Volkas'15)

- Explore TeV scale models !!



SUSY + Leptogenesis also prefer low scale seesaw

- For leptogenesis to occur, $M_N < T_{\text{reheat}}$;
- Gravitino overclosing prefers that $T_{\text{reheat}} < 10^6$ GeV (Kohri et al.)

→ Hence preference of leptogenesis for lower seesaw scale !!

Experimental searches for TeV W_R effects

- Collider searches for W_R and N : LHC

- (i) Direct W_R production

- (ii) ν - N mixing from seesaw

$$V_{\ell N} = \sqrt{\frac{m_\nu}{M_N}}$$

(Han, Ruiz et al; Senjanovic, Nemevsek, Nesti, Tello,..Deppisch, Dev, Pilaftsis;..Del Aguila et al.)

- New leptophilic Higgses: Δ^{++}, Δ^+

(Chakraborty, Gluza, Bhambaniya, Zafron,..Dutta, Goa, Ghosh, Eusebii, Kamon...)

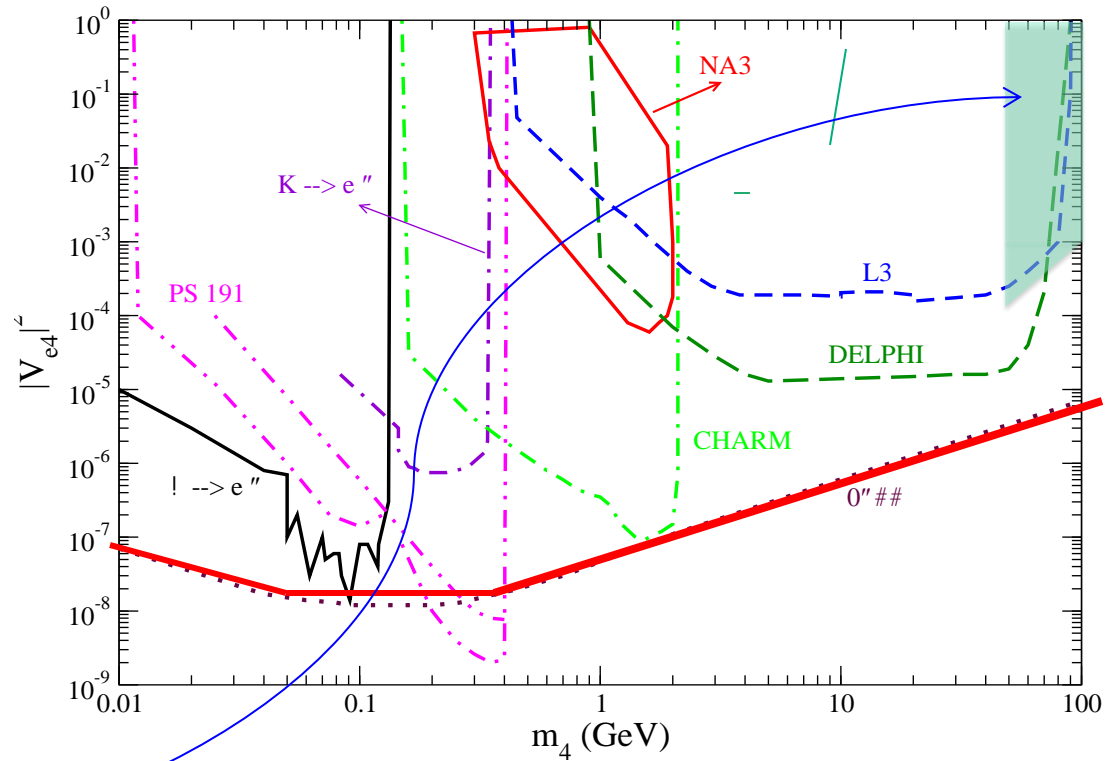
- Neutrinoless double beta decay and LFV

(RNM'86; Hirsch, Klapdor, Kovalenko'96; Das, Deppisch, Kittel, Valle; Dev, Goswami, Mitra;....)

- Light N 's and displaced vertices (Helo, Dib, Kovalenko, Ortiz,)

Lower Mass N_R

$\nu - N$ mixing
 $V_{\ell N}$



(Atre, Han, Pascoli, Zhang; Antusch, Fisher'14)

Bounds from LHC Higgs decay to $e^+ e^- E_T$ from $pp \rightarrow h \rightarrow \ell N$

$N \rightarrow W^+ \ell^-, Z + \nu$

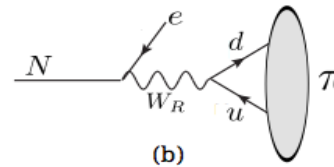
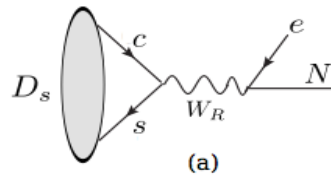
(Dev, Francischini, RNM'12 ; Gago, Hernandez, Perez, Losada, Briceno'15)

Beam Dump searches

■ Displaced vertices

(Castillo-Feliosela, Helo, Dib, Kovalenko, Ortiz'15)

■ $M_N < 1.8$ GeV



SHIP setup

■ Reach:

$M_{W_R} \leq 18$ TeV

