

# Charged Higgs production in association with a $W$ or a top

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- Charged Higgs production
- Higher-order corrections
- $tH^-$  production
- $H^-W^+$  production



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## Charged Higgs production

A charged Higgs would be sure sign of new physics

2-Higgs doublet models

LHC has good potential for discovery

I will discuss two production processes

$bg \rightarrow tH^-$  and  $b\bar{b} \rightarrow H^-W^+$

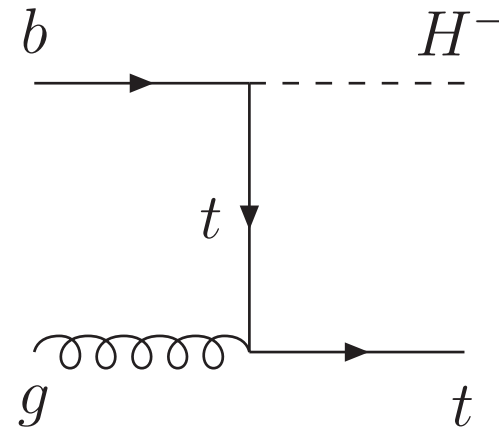
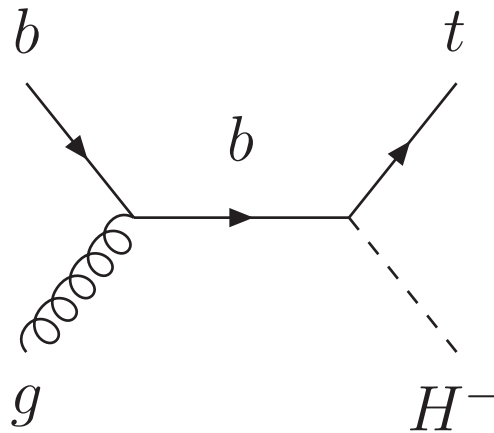
Probe into electroweak and Higgs physics

Higher-order corrections are significant

very massive final states

Soft-gluon corrections are important

## $tH^-$ production



Top is the heaviest known elementary particle

Decays before hadronization

Born cross section for  $bg \rightarrow tH^- \propto \alpha\alpha_s(m_b^2 \tan^2 \beta + m_t^2 \cot^2 \beta)$

$\tan \beta = v_2/v_1$  ratio of vevs of two Higgs doublets

## Higher-order corrections

$$b(p_b) + g(p_g) \longrightarrow t(p_t) + H^-(p_H)$$

Define  $s = (p_b + p_g)^2$ ,  $t = (p_b - p_t)^2$ ,  $u = (p_g - p_t)^2$   
and  $s_4 = s + t + u - m_t^2 - m_H^2$

At partonic threshold  $s_4 \rightarrow 0$

Soft corrections  $\left[ \frac{\ln^k(s_4/m_H^2)}{s_4} \right]_+$

For the order  $\alpha_s^n$  corrections  $k \leq 2n - 1$

Resum these soft corrections for the double-differential cross section

At NNLL accuracy we need two-loop soft anomalous dimensions

Derive approximate cross sections at NNLO

## Soft-gluon Resummation

moments of the partonic cross section with moment variable  $N$ :

$$\hat{\sigma}(N) = \int (ds_4/s) e^{-Ns_4/s} \hat{\sigma}(s_4)$$

factorized expression for the cross section in  $4 - \epsilon$  dimensions

$$\hat{\sigma}^{bg \rightarrow tH^-}(N, \epsilon) = \left( \prod_{i=b,g} J_i(N, \mu, \epsilon) \right) H^{bg \rightarrow tH^-}(\alpha_s(\mu)) S^{bg \rightarrow tH^-} \left( \frac{m_H}{N\mu}, \alpha_s(\mu) \right)$$

$H^{bg \rightarrow tH^-}$  is hard function and  $S^{bg \rightarrow tH^-}$  is soft function

Soft function  $S$  satisfies the renormalization group equation

$$\left( \mu \frac{\partial}{\partial \mu} + \beta(g_s, \epsilon) \frac{\partial}{\partial g_s} \right) S^{bg \rightarrow tH^-} = -2 S^{bg \rightarrow tH^-} \Gamma_S^{bg \rightarrow tH^-}$$

Soft anomalous dimension  $\Gamma_S^{bg \rightarrow tH^-}$  controls the evolution of  $S^{bg \rightarrow tH^-}$  which results in the exponentiation of logarithms of  $N$

$$\Gamma_S^{bg \rightarrow tH^-} = (\alpha_s/\pi) \Gamma_S^{(1)} + (\alpha_s/\pi)^2 \Gamma_S^{(2)} + \dots, \text{ with}$$

$$\Gamma_S^{(1)} = C_F \left[ \ln \left( \frac{m_t^2 - t}{m_t \sqrt{s}} \right) - \frac{1}{2} \right] + \frac{C_A}{2} \ln \left( \frac{m_t^2 - u}{m_t^2 - t} \right)$$

$$\Gamma_S^{(2)} = \left[ C_A \left( \frac{67}{36} - \frac{\zeta_2}{2} \right) - \frac{5}{18} n_f \right] \Gamma_S^{(1)} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

## Resummed cross section

$$\hat{\sigma}_{\text{resummed}}^{bg \rightarrow tH^-}(N) = \exp \left[ \sum_{i=b,g} E_i(N_i) \right] H^{bg \rightarrow tH^-}(\alpha_s(\sqrt{s}))$$

$$\times S^{bg \rightarrow tH^-} \left( \alpha_s \left( \frac{\sqrt{s}}{N'} \right) \right) \exp \left[ \int_{\sqrt{s}}^{\sqrt{s}/N'} \frac{d\mu}{\mu} 2 \Gamma_S^{bg \rightarrow tH^-}(\alpha_s(\mu)) \right]$$

The aNNLO soft-gluon corrections are:

$$\frac{d^2 \hat{\sigma}_{\text{aNNLO}}^{(2) bg \rightarrow tH^-}}{dt du} = F_{\text{LO}}^{bg \rightarrow tH^-} \frac{\alpha_s^2}{\pi^2} \sum_{k=0}^3 C_k^{(2)} \left[ \frac{\ln^k(s_4/m_H^2)}{s_4} \right]_+$$

with coefficients  $C_3^{(2)} = 2(C_F + C_A)^2$

$$C_2^{(2)} = (C_F + C_A) \left\{ 3C_F \left[ 2 \ln \left( \frac{m_t^2 - t}{m_t \sqrt{s}} \right) - 2 \ln \left( \frac{m_H^2 - u}{m_H^2} \right) - 1 \right] \right.$$

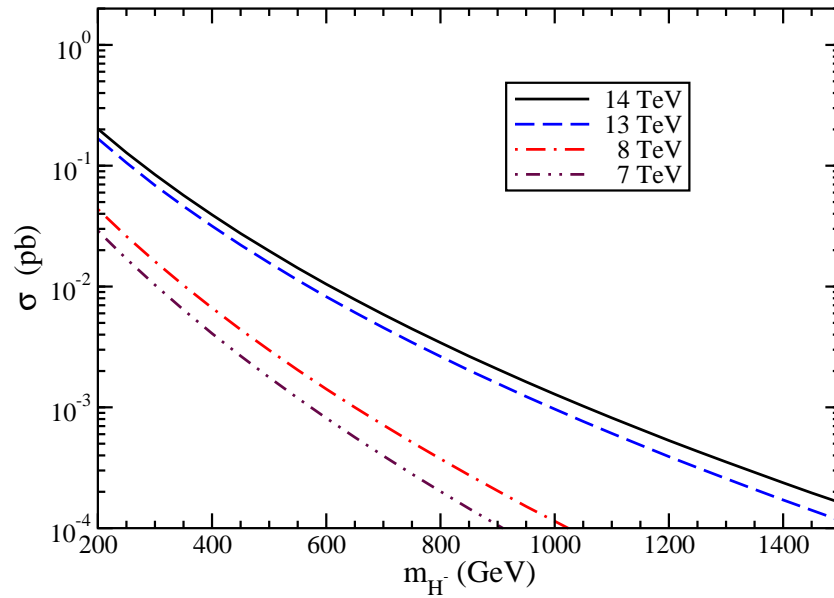
$$\left. - 3C_A \left[ \ln \left( \frac{m_t^2 - t}{m_t^2 - u} \right) + 2 \ln \left( \frac{m_H^2 - t}{m_H^2} \right) \right] - 3(C_F + C_A) \ln \left( \frac{\mu_F^2}{s} \right) - \frac{\beta_0}{2} \right\}$$

The expressions for  $C_1^{(2)}$  and  $C_0^{(2)}$  are much longer

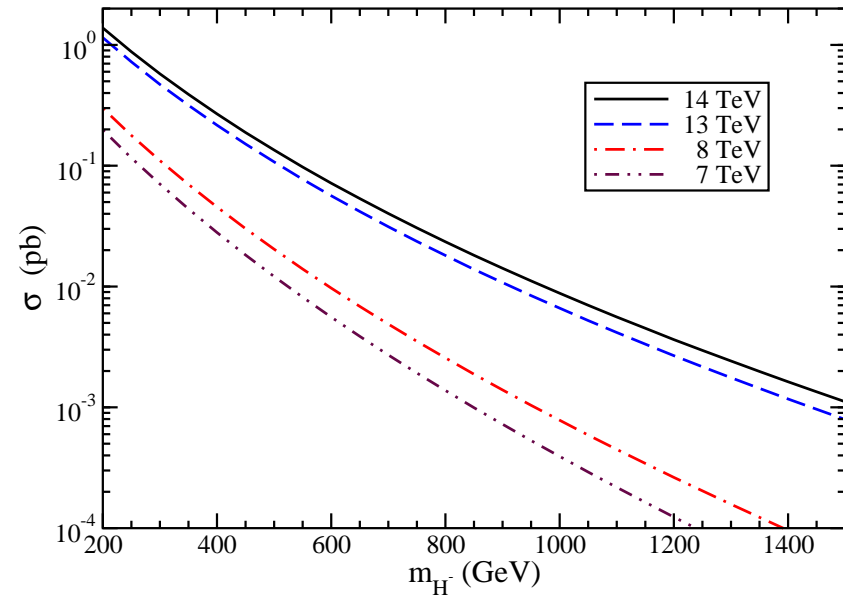
# $tH^-$ production

## Total cross sections

bg- $\rightarrow$   $tH^-$  at LHC aNNLO  $\tan\beta=10$   $\mu=m_{H^-}$



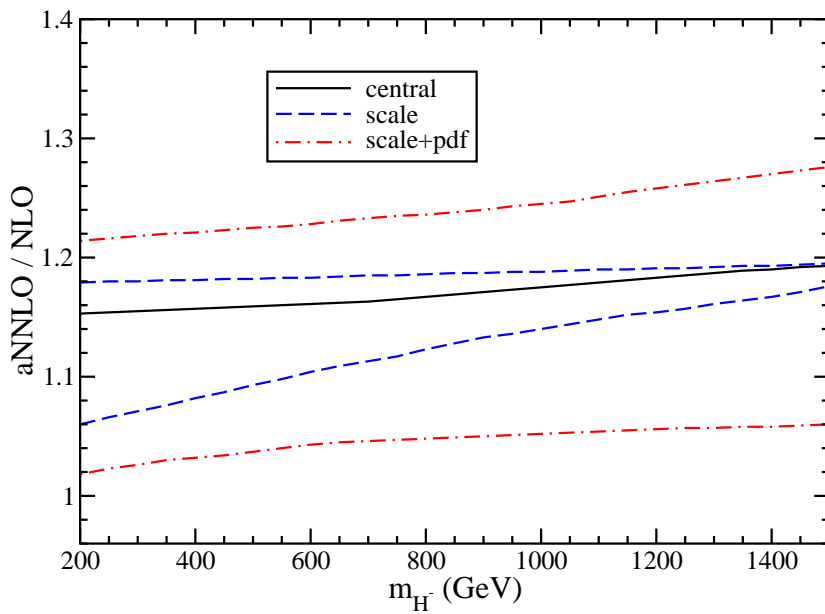
bg- $\rightarrow$   $tH^-$  at LHC aNNLO  $\tan\beta=30$   $\mu=m_{H^-}$



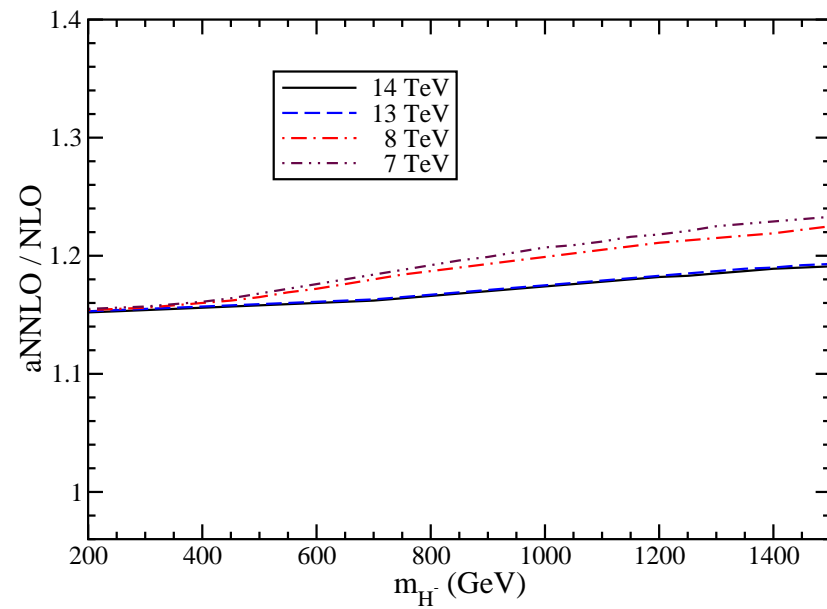
# $tH^-$ production

## $K$ -factors

bg- $\rightarrow$   $tH^-$  at LHC     $K$ -factor    13 TeV



bg- $\rightarrow$   $tH^-$  at LHC     $K$ -factor     $\mu = m_{H^-}$

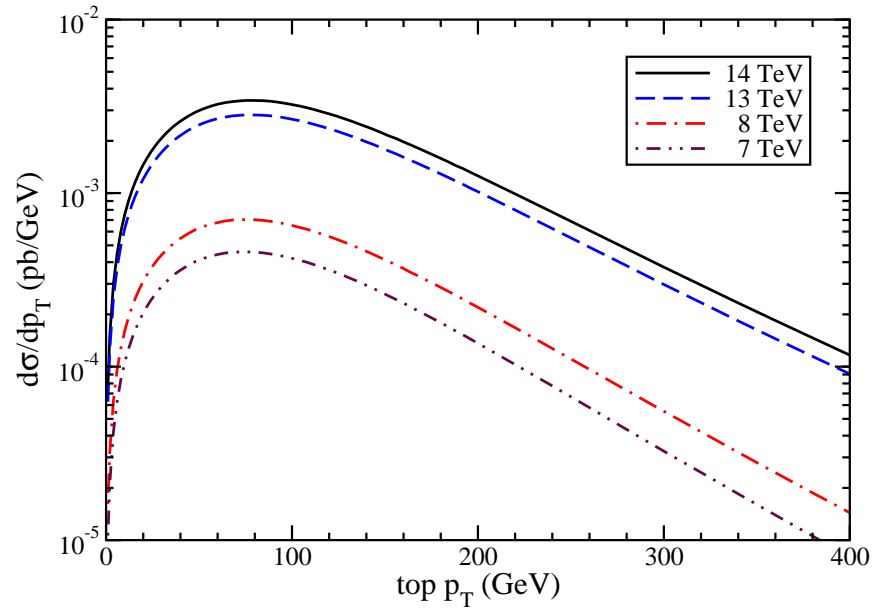




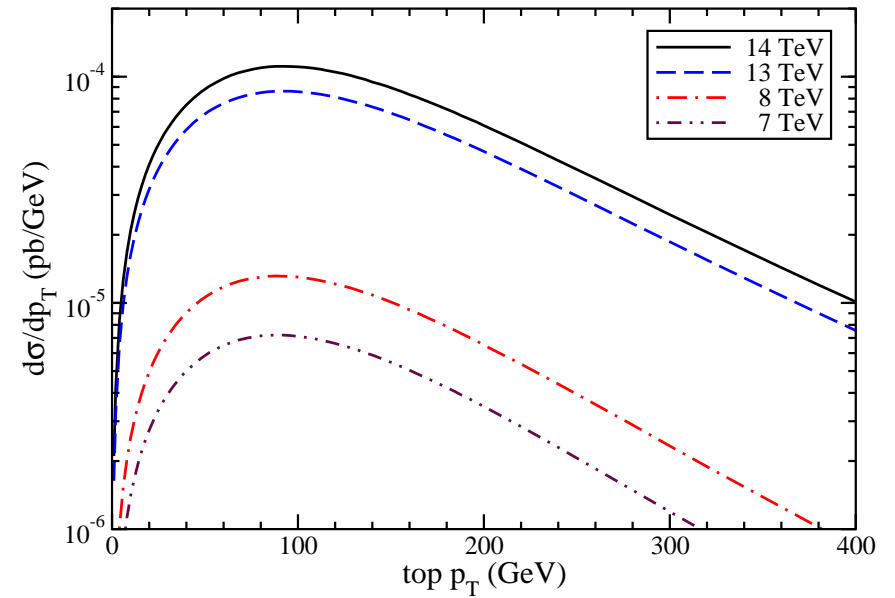
# $tH^-$ production

## Top $p_T$ distributions

bg  $\rightarrow$   $tH^-$  at LHC top  $p_T$  aNNLO  $\tan\beta=30$   $m_{H^-}=300$  GeV

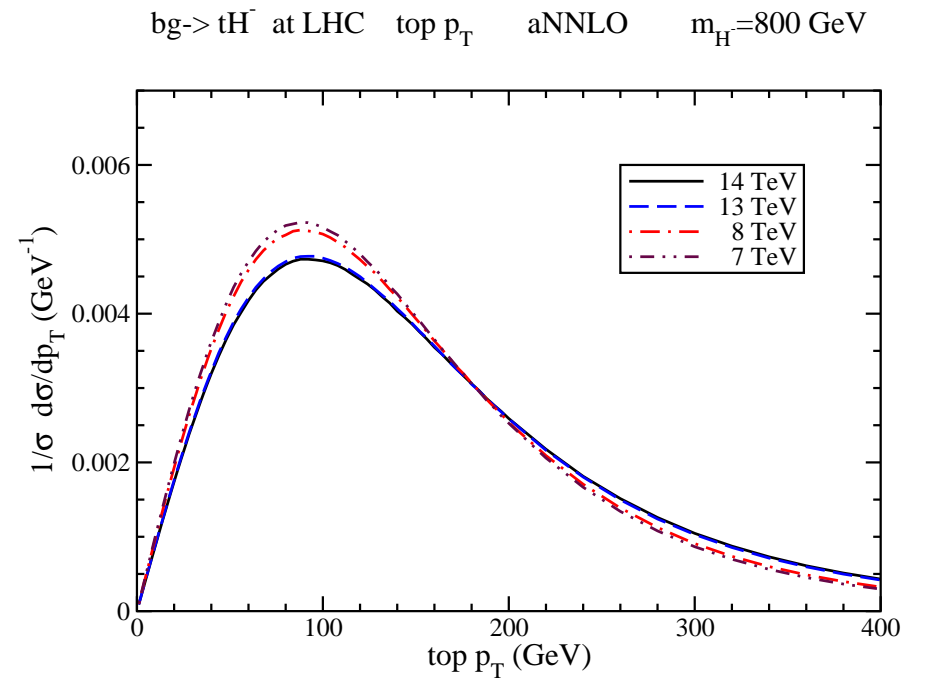
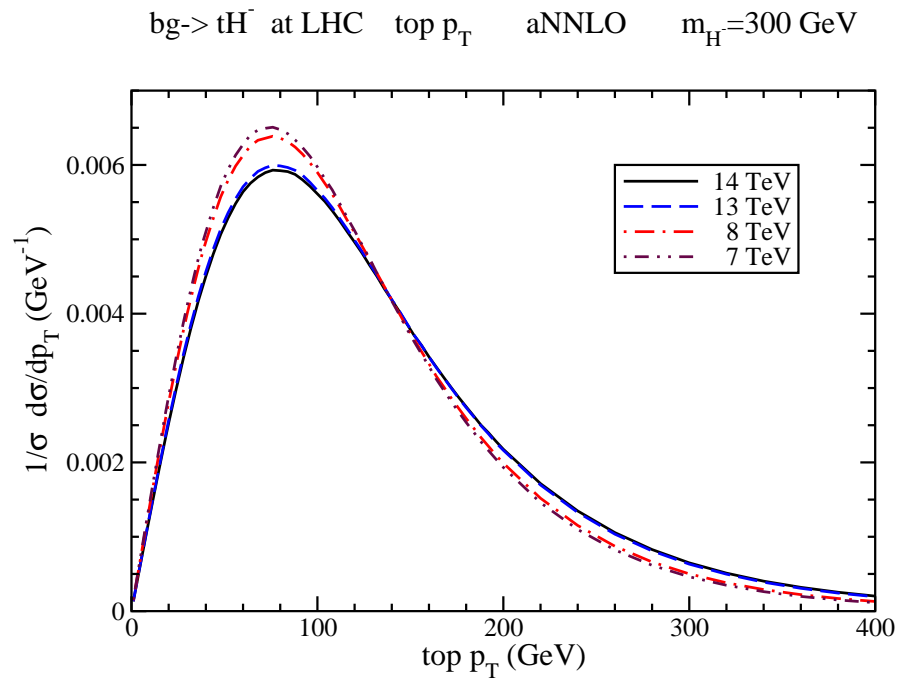


bg  $\rightarrow$   $tH^-$  at LHC top  $p_T$  aNNLO  $\tan\beta=30$   $m_{H^-}=800$  GeV



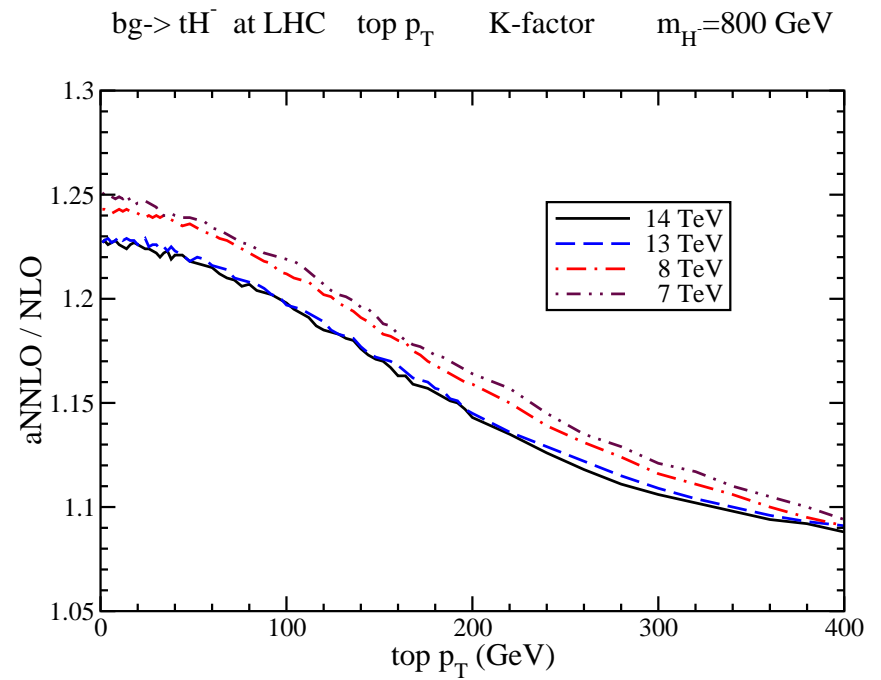
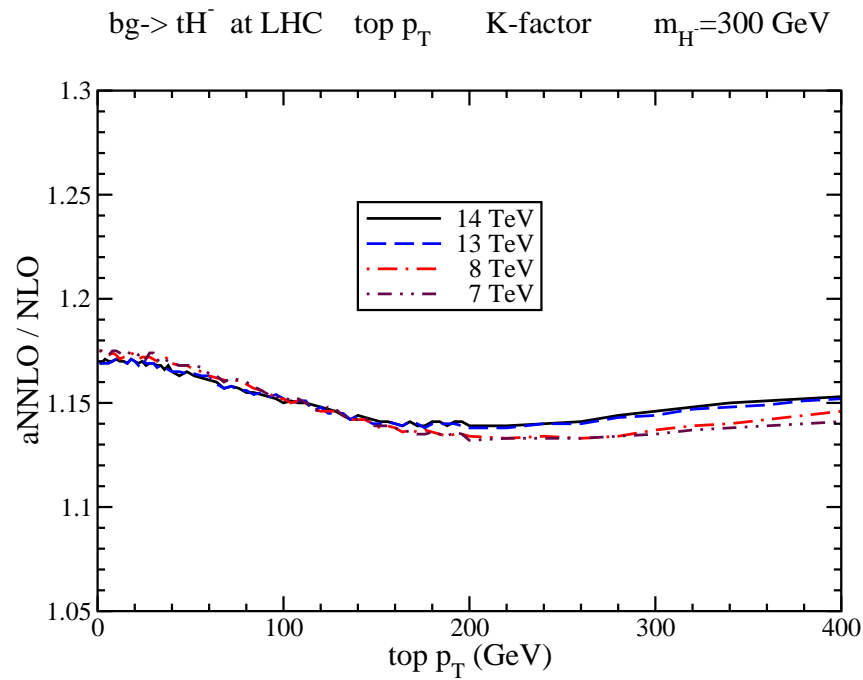
# $tH^-$ production

## Normalized top $p_T$ distributions



# $tH^-$ production

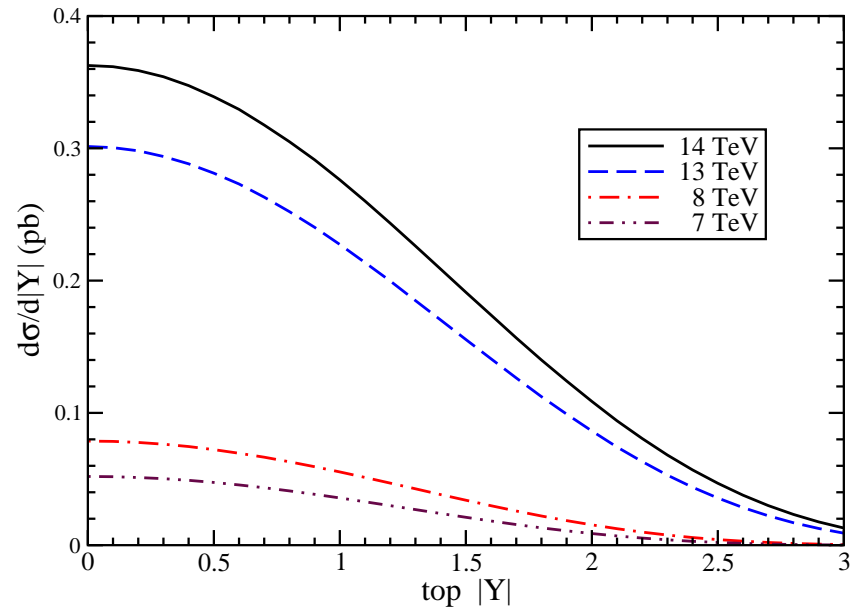
## $K$ -factors for top $p_T$ distributions



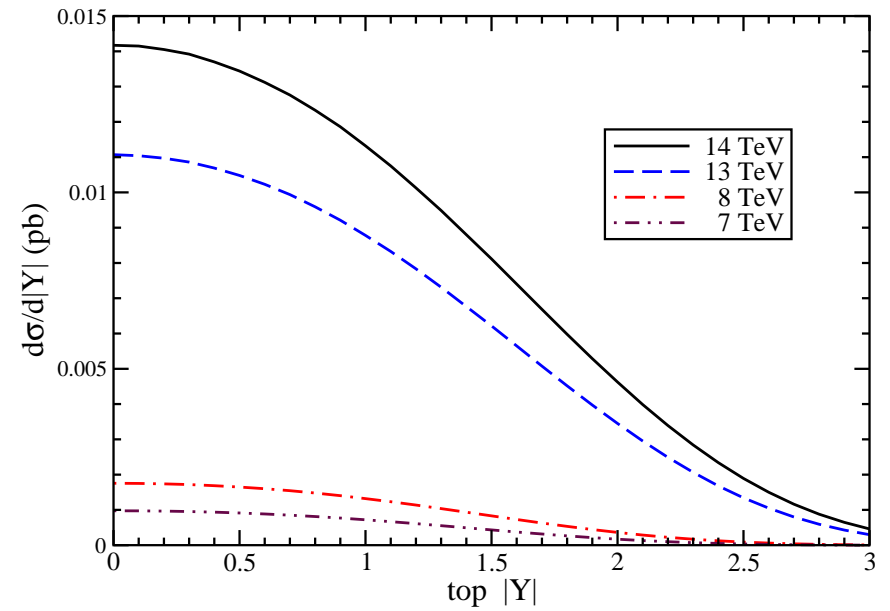
# $tH^-$ production

## Top rapidity distributions

bg- $\rightarrow$ t $H^-$  at LHC top rapidity aNNLO  $\tan\beta=30$   $m_{H^-}=300$  GeV

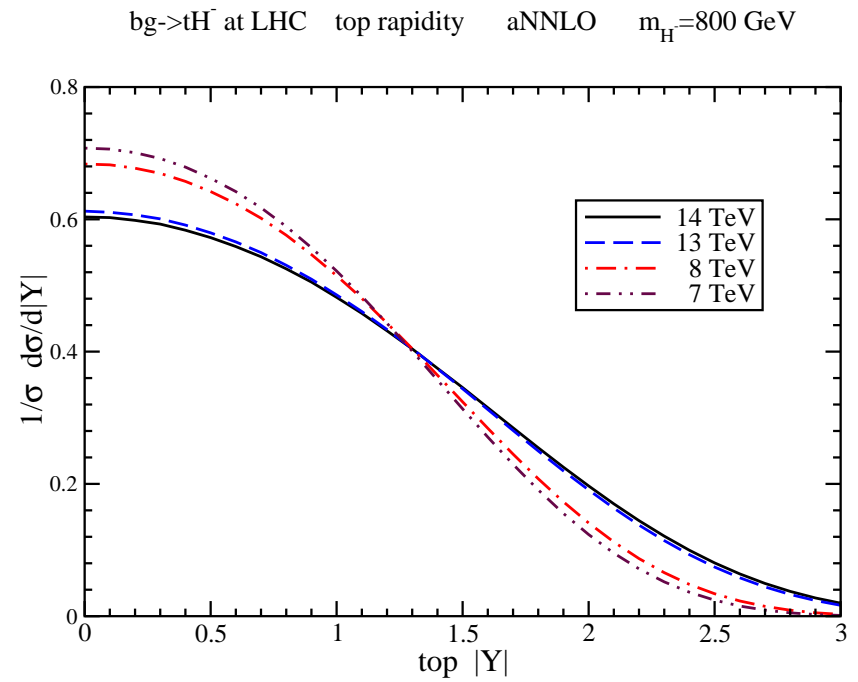
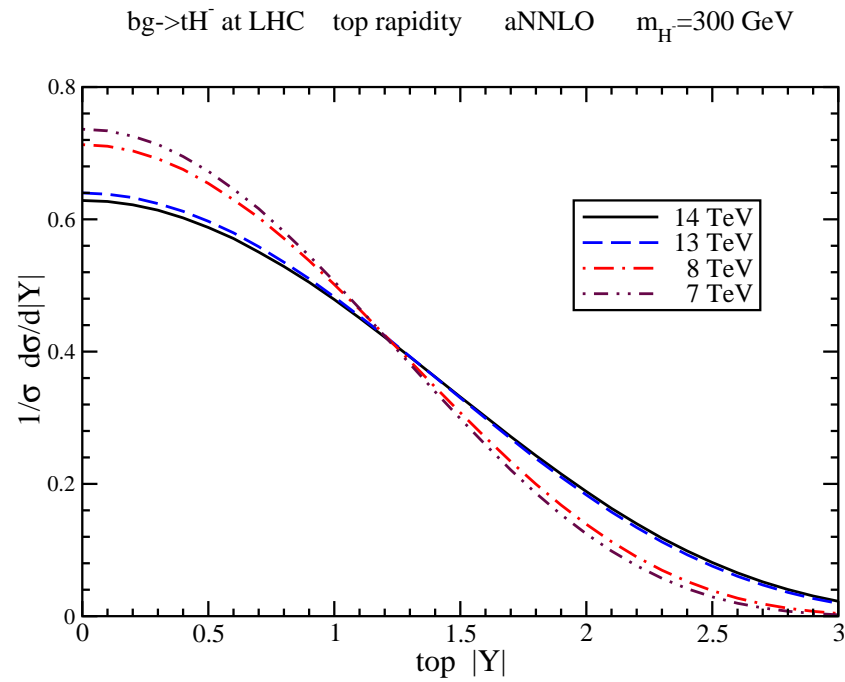


bg- $\rightarrow$ t $H^-$  at LHC top rapidity aNNLO  $\tan\beta=30$   $m_{H^-}=800$  GeV



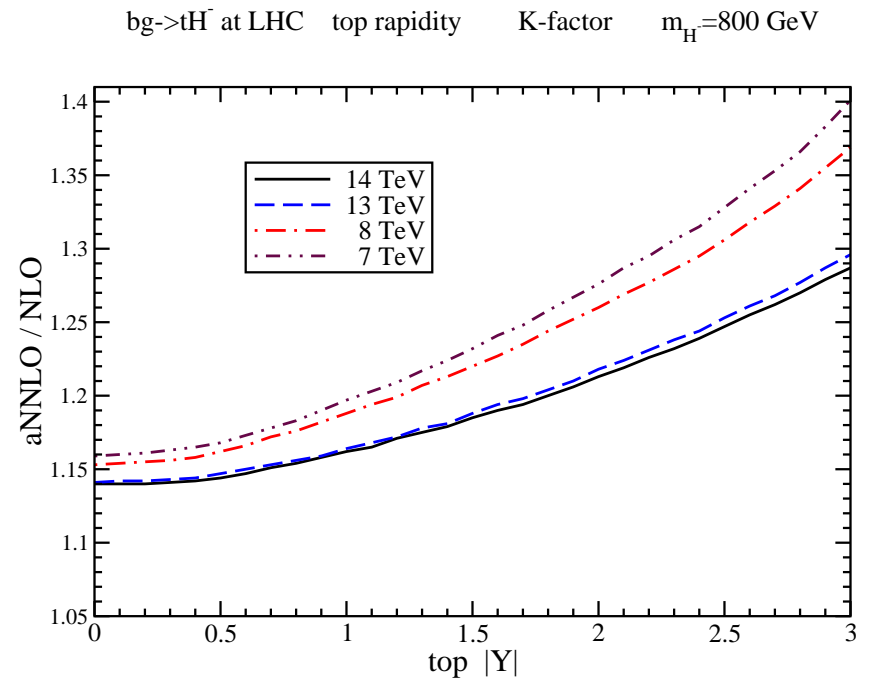
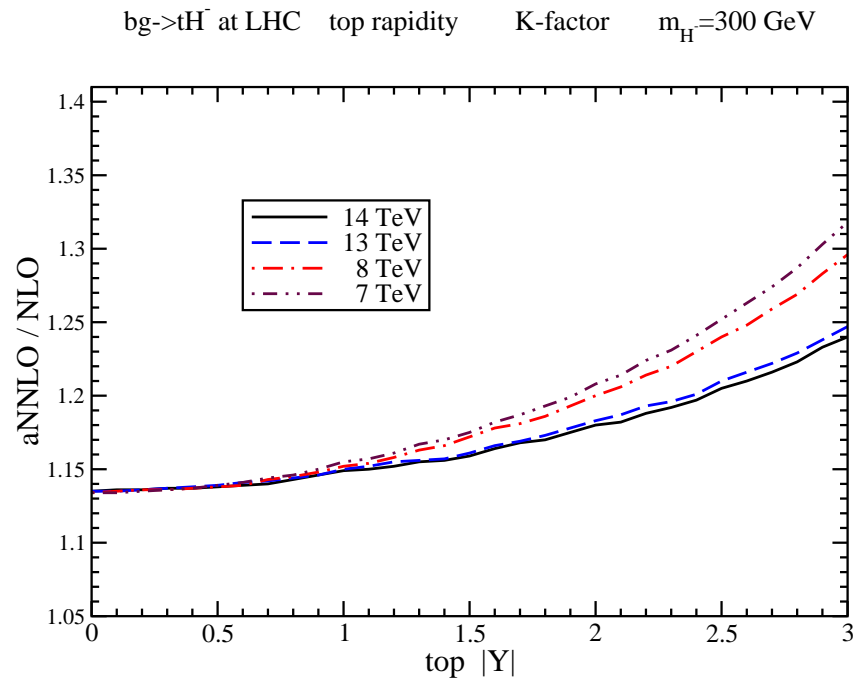
# $tH^-$ production

## Normalized top rapidity distributions

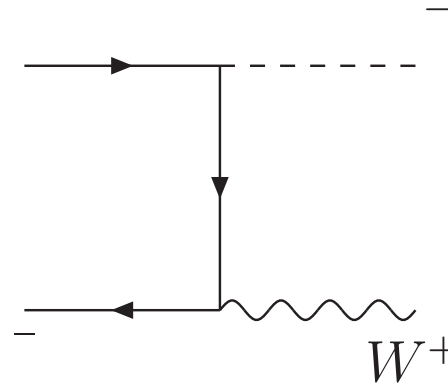


# $tH^-$ production

## $K$ -factors for top rapidity distributions



## $H^- W^+$ production



$$b(p_1) + \bar{b}(p_2) \rightarrow H^-(p_3) + W^+(p_4)$$

Define  $s = (p_1 + p_2)^2$ ,  $t = (p_1 - p_3)^2$ ,  $u = (p_2 - p_3)^2$   
and  $s_4 = s + t + u - m_H^2 - m_W^2$

At partonic threshold  $s_4 \rightarrow 0$

Soft corrections  $\left[ \frac{\ln^k(s_4/m_H^2)}{s_4} \right]_+$

factorized expression for the cross section in  $4 - \epsilon$  dimensions

$$\hat{\sigma}^{b\bar{b} \rightarrow H^- W^+}(N, \epsilon) = \left( \prod_{i=b, \bar{b}} J_i(N, \mu, \epsilon) \right) H^{b\bar{b} \rightarrow H^- W^+}(\alpha_s(\mu)) S^{b\bar{b} \rightarrow H^- W^+} \left( \frac{m_H}{N\mu}, \alpha_s(\mu) \right)$$

**Resummed cross section**

$$\begin{aligned} \hat{\sigma}_{\text{res}}^{b\bar{b} \rightarrow H^- W^+}(N) &= \exp \left[ \sum_{i=b, \bar{b}} E_i(N_i) \right] H^{b\bar{b} \rightarrow H^- W^+}(\alpha_s(\sqrt{s})) S^{b\bar{b} \rightarrow H^- W^+}(\alpha_s(\sqrt{s}/\tilde{N}')) \\ &\times \exp \left[ 2 \int_{\sqrt{s}}^{\sqrt{s}/\tilde{N}'} \frac{d\mu}{\mu} \Gamma_S^{b\bar{b} \rightarrow H^- W^+}(\alpha_s(\mu)) \right] \end{aligned}$$

The NNLO collinear and soft-gluon corrections are

$$\frac{d^2 \hat{\sigma}_{\text{aNNLO}}^{(2) b\bar{b} \rightarrow H^- W^+}}{dt du} = F_{LO}^{b\bar{b} \rightarrow H^- W^+} \frac{\alpha_s^2}{\pi^2} \left\{ -C_3^{(2)} \frac{1}{m_H^2} \ln^3 \left( \frac{s_4}{m_H^2} \right) + \sum_{k=0}^3 C_k^{(2)} \left[ \frac{\ln^k(s_4/m_H^2)}{s_4} \right]_+ \right\}$$

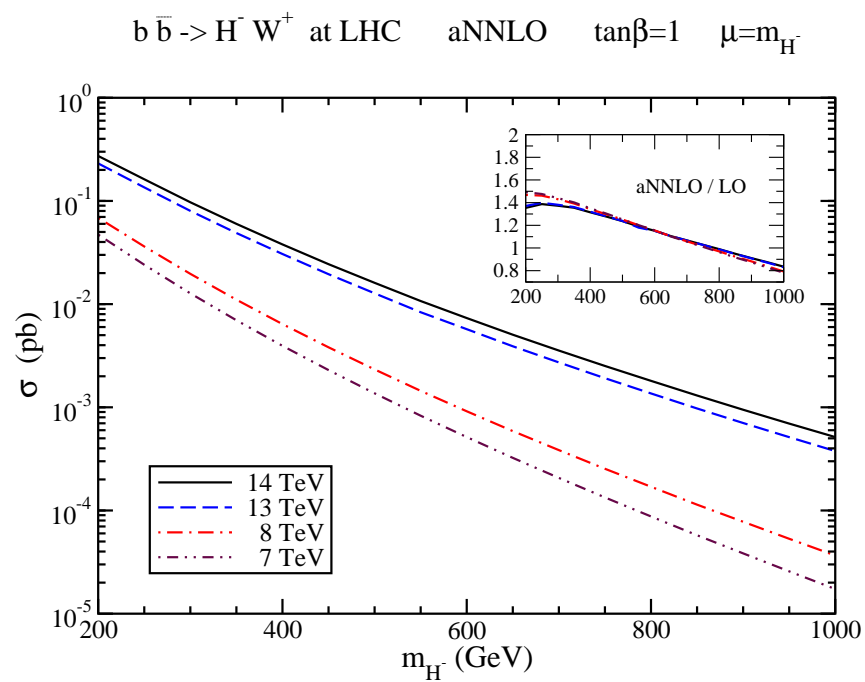
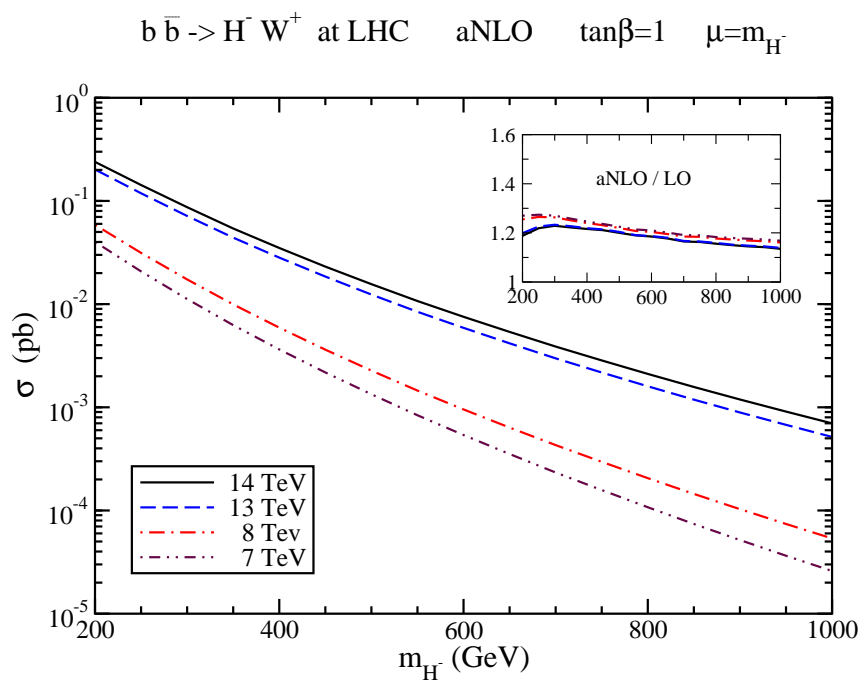
with  $C_3^{(2)} = 8C_F^2$

$$C_2^{(2)} = -12C_F^2 \left( \ln \left( \frac{(t - m_W^2)(u - m_W^2)}{m_H^4} \right) + \ln \left( \frac{\mu_F^2}{s} \right) \right) - \frac{11}{3} C_F C_A + \frac{2}{3} C_F n_f$$



# $H^-W^+$ production

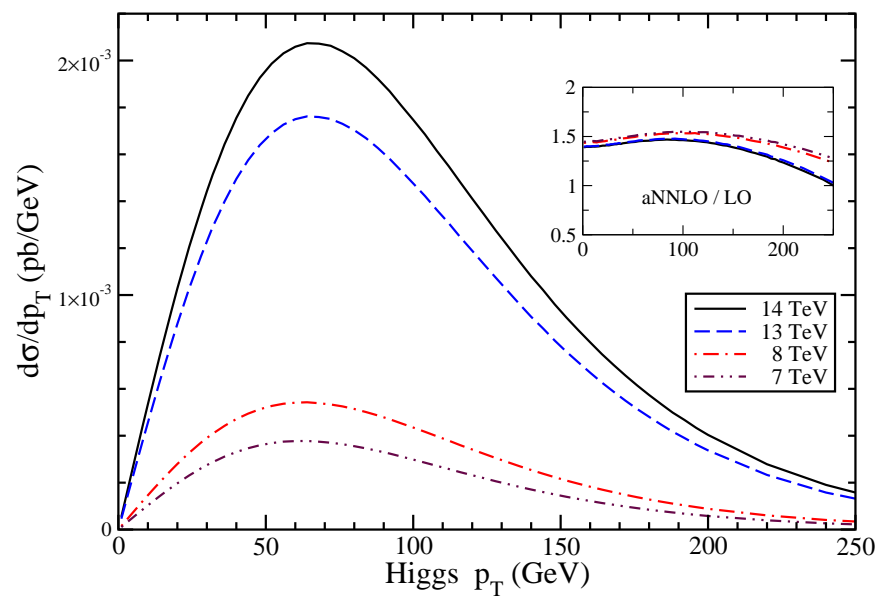
## Total cross sections



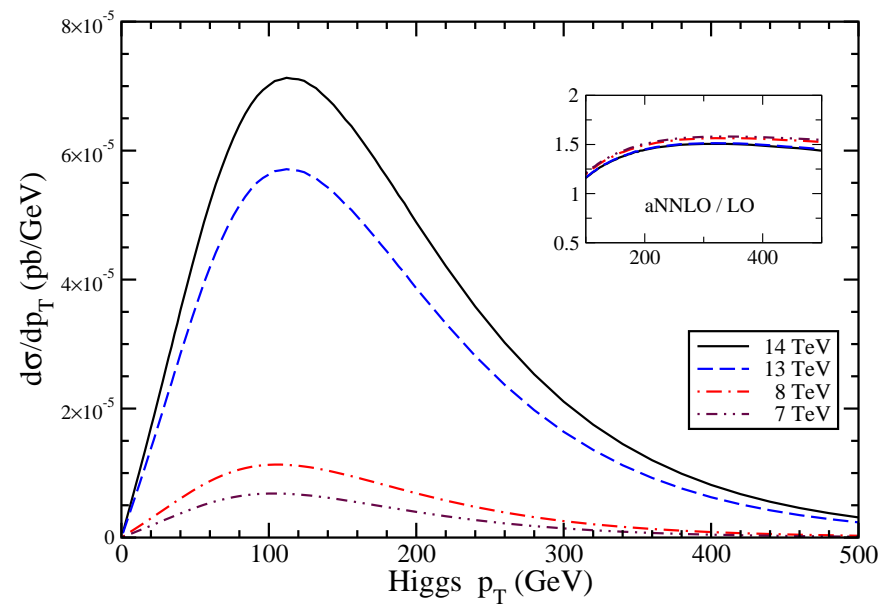
# $H^-W^+$ production

## Charged Higgs $p_T$ distributions

$b\bar{b} \rightarrow H^-W^+$  at LHC aNNLO  $\tan\beta=1$   $m_{H^-}=200$  GeV

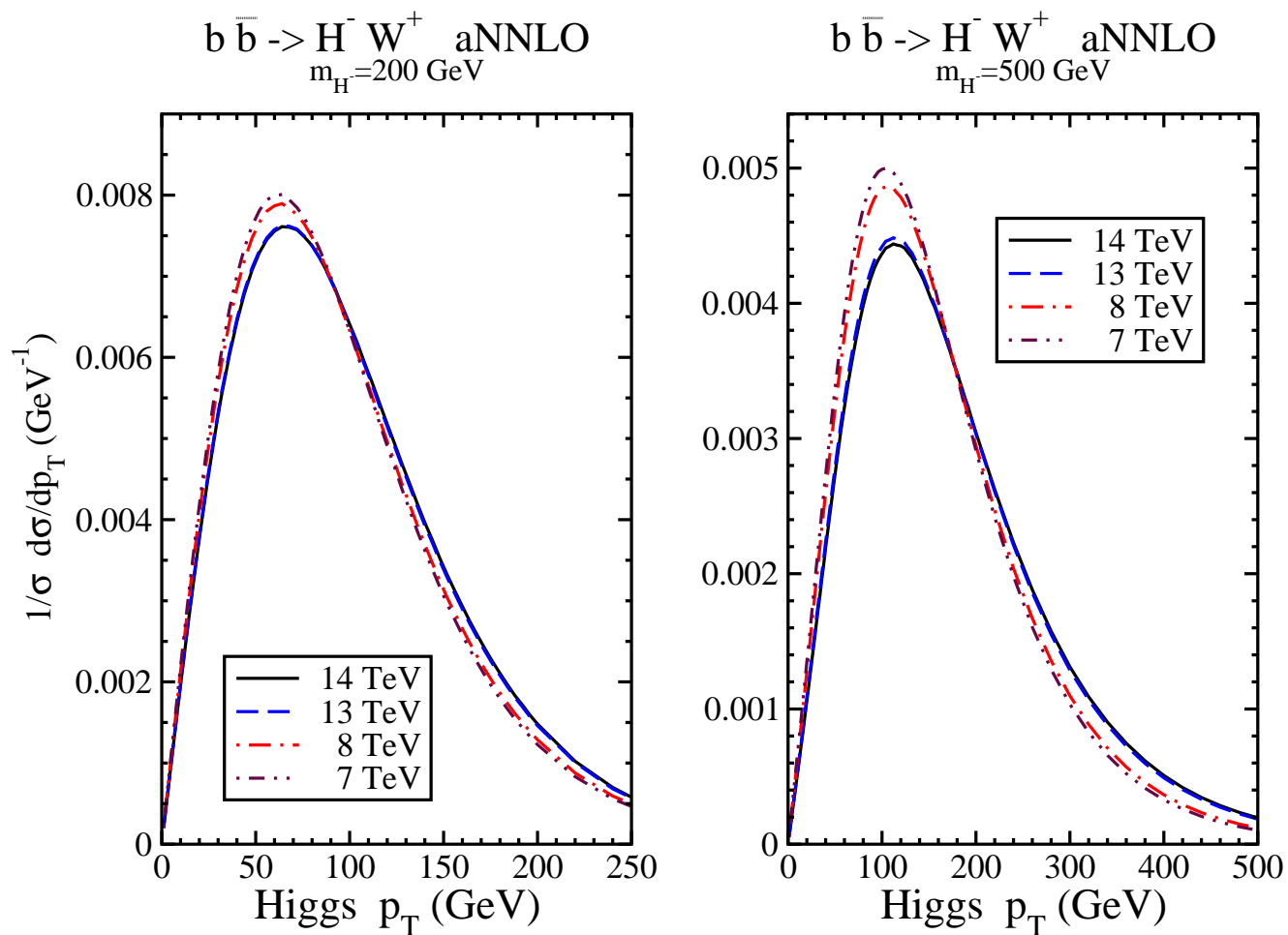


$b\bar{b} \rightarrow H^-W^+$  at LHC aNNLO  $\tan\beta=1$   $m_{H^-}=500$  GeV



# $H^-W^+$ production

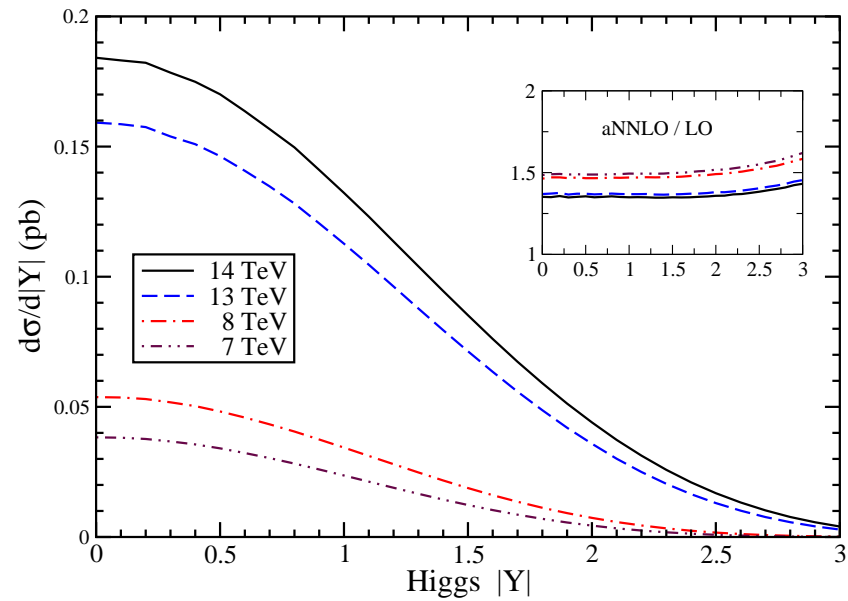
## Normalized charged Higgs $p_T$ distributions



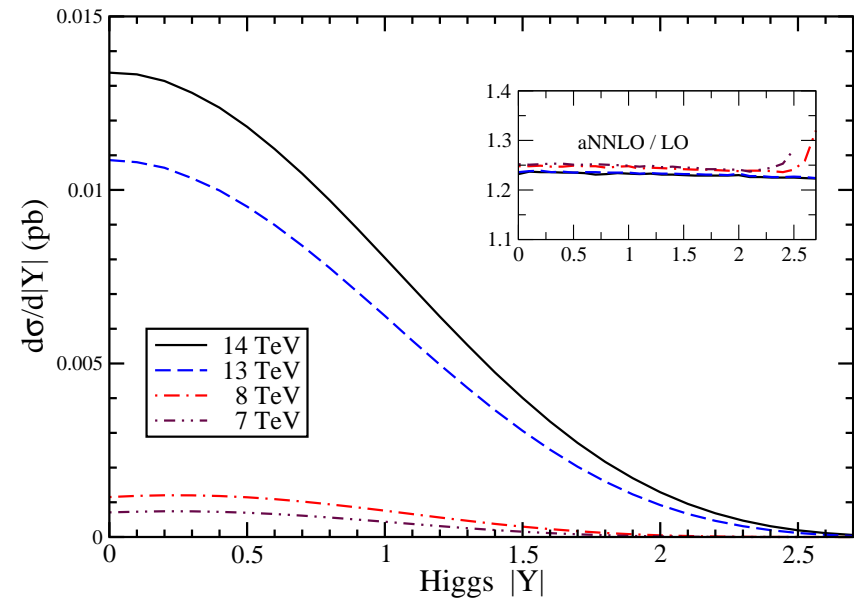
# $H^-W^+$ production

## Charged Higgs rapidity distributions

$b\bar{b} \rightarrow H^-W^+$  at LHC aNNLO  $\tan\beta=1$   $m_{H^\pm}=200$  GeV

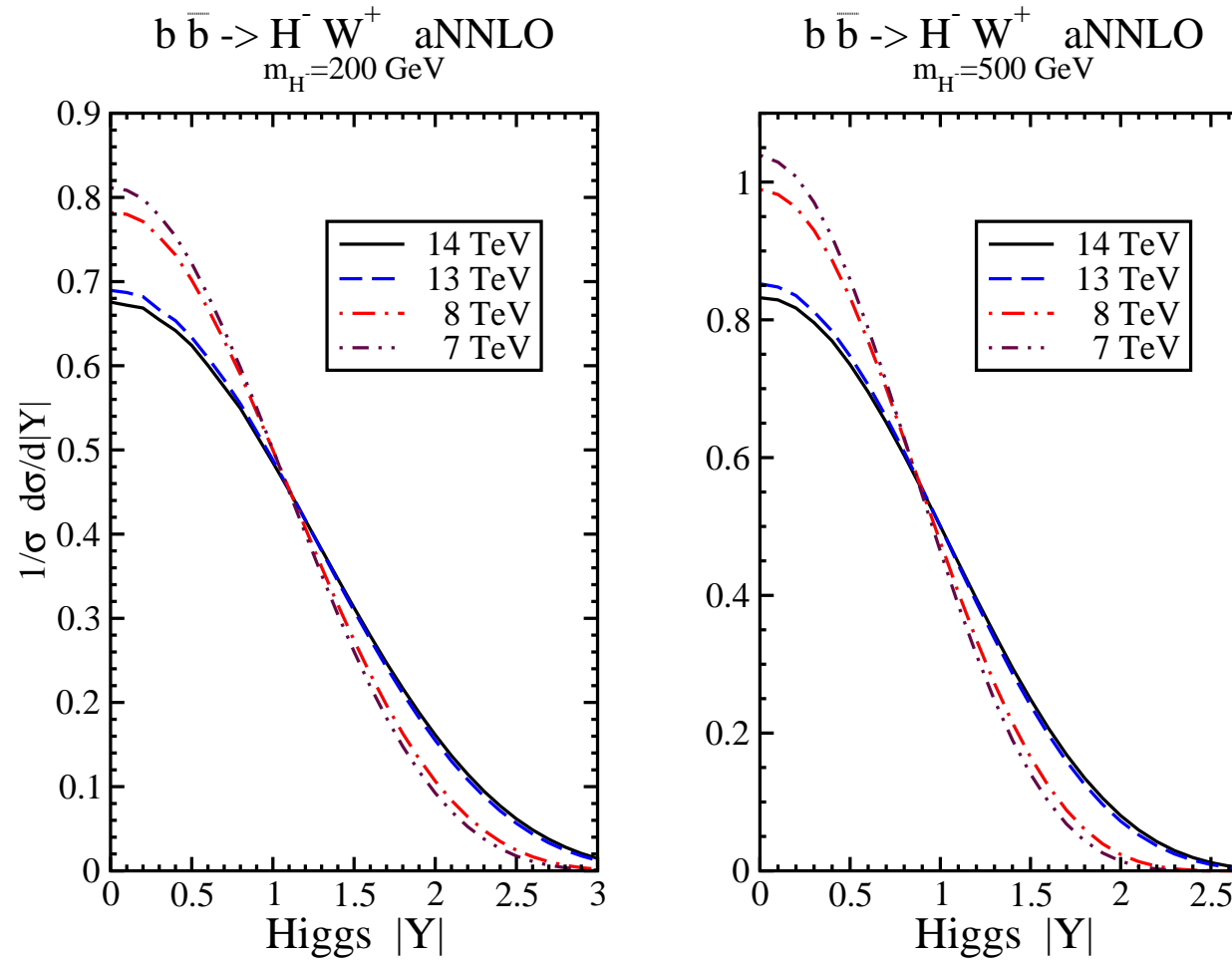


$b\bar{b} \rightarrow H^-W^+$  at LHC aNNLO  $\tan\beta=1$   $m_{H^\pm}=500$  GeV



# $H^-W^+$ production

## Normalized charged Higgs rapidity distributions



## Summary

- many new results in charged Higgs production
- total cross sections for  $tH^-$  production
- top-quark  $p_T$  and rapidity distributions in  $tH^-$  production
- total cross sections for  $H^-W^+$  production
- charged-Higgs  $p_T$  and rapidity distributions in  $H^-W^+$  production
- higher-order corrections are very significant