

Radiative corrections in a general Yukawa model

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Particles and Interactions

FWF

Der Wissenschaftsfonds.

Main question

How do radiative corrections to masses and mixing angles behave in ν -mass models?

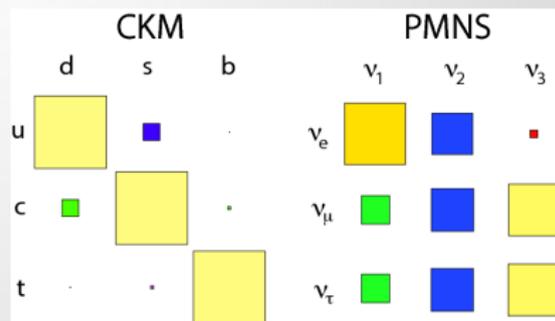
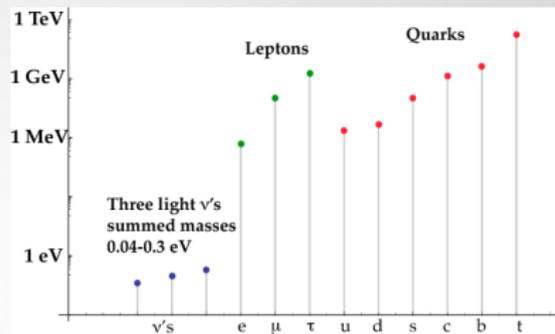


M. Fox, W. Grimus and M. Löschner,
Renormalization and radiative corrections to masses in a general Yukawa model,
arXiv:1705.09589 [hep-ph].

Motivation

Open problems:

1. Smallness of ν -masses
2. Mild hierarchy in ν -mass spectrum vs. strong hierarchy in spectra of charged leptons and quarks
3. Relation between mixing angles and masses
4. Majorana nature of neutrinos



Taken from [1]

Motivation

Flavour symmetries

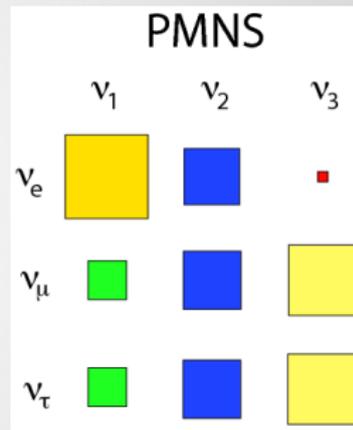
- Attempt to describe/explain structure of U_{PMNS} via symmetries of the mass matrix
- Use (combination of) discrete symmetries to approximate U_{PMNS} , e.g. μ - τ symmetry [3]

$$SM_\nu S = \mathcal{M}_\nu^*$$

$$\text{with } S = \begin{matrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{matrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$\Rightarrow |U_{\mu i}| = |U_{\tau i}| \quad \forall i$$

$$\Rightarrow \theta_{23} = 45^\circ, \quad \delta = \pm \frac{\pi}{2}.$$



Taken from [1]

- ▶ Models of this kind often show *proliferation of scalars*
- ▶ Need suitable treatment of potentially large number of scalars

Radiative Corrections in ν -mass models

Setup of general Yukawa toy model:

- n_χ Majorana fermions
- n_φ real scalars
- \mathbb{Z}_4 -Symmetry: no tree-level mass terms and (tri-)linear terms in scalar potential:

$$\mathcal{S} : \quad \chi_L \rightarrow i\chi_L \quad \varphi \rightarrow -\varphi$$

Toy model Lagrangian

$$\begin{aligned} \mathcal{L} = & i\bar{\chi}_{iL}\not{\partial}_\mu\chi_{iL} + \frac{1}{2}(\partial_\mu\varphi_a)(\partial^\mu\varphi_a) \\ & + \left(\frac{1}{2}(Y_a)_{ij}\chi_{iL}^T C^{-1}\chi_{jL}\varphi_a + \text{H.c.} \right) \\ & - \frac{1}{2}(\mu^2)_{ab}\varphi_a\varphi_b - \frac{1}{4}\lambda_{abcd}\varphi_a\varphi_b\varphi_c\varphi_d \end{aligned}$$

Radiative Corrections in ν -mass models

Generation of masses

- Masses are generated via *spontaneous symmetry breaking* (SSB)

$$\varphi_a = \mathcal{M}^{-\varepsilon/2} \bar{v}_a + h_a$$

- ▶ Majorana masses $m_0 = \sum_{a=1}^{n_\varphi} v_a Y_a$
- ▶ Scalar masses $(M_0^2)_{ab} = \mu_{ab}^2 + 3\lambda_{abcd} v_c v_d$
- Masses fully determined via model parameters, also in terms of renormalization
- Note: number of masses \ll number of model parameters
 \Rightarrow Renormalization of masses not sufficient to make theory finite

Radiative Corrections in ν -mass models

Renormalization

Want to achieve:

- ▶ Clear and straightforward approach
- ▶ Non-specific on peculiarities of the investigated models
- ▶ Take nature of masses as functions of parameters of the model serious

Radiative Corrections in ν -mass models

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Programme:

1. $\overline{\text{MS}}$ renormalization for all parameters of the model and the field strength renormalization constants (using dim. reg.)
2. Finite shifts Δv_a of vacuum expectation values such that scalar one-point functions vanish
3. Finite field strength renormalization to get on-shell self-energies

Radiative Corrections in ν -mass models

Determination of Counterterms

- (a) $\delta\hat{\lambda}_{abcd}$ determined from the quartic scalar coupling
- (b) $\delta\hat{Y}_a$ obtained from Yukawa vertex
- (c) $\delta\hat{\mu}_{ab}^2$ removes UV-divergent terms from p^2 -independent part of scalar selfenergy
- (d) $\hat{\delta}^{(\chi)}$ and $\hat{\delta}^{(\varphi)}$ determined from momentum-dependent parts of respective selfenergies.



Results

One-point functions and VEV shifts

- ▶ Unobscured treatment of tadpole contributions

$$0 = \begin{array}{c} \bullet \\ | \\ | \\ | \end{array} + \begin{array}{c} \text{⦿} \\ | \\ | \\ | \end{array} + \begin{array}{c} \times \\ | \\ | \\ | \end{array}$$
$$= \mathcal{M}^{-\varepsilon/2} \frac{i}{-M_{0a}^2} \times (-i) \left(\hat{t}_a + T_a + \Delta \hat{t}_a + \delta \hat{\mu}_{ab}^2 \hat{v}_b + \delta \hat{\lambda}_{abcd} \hat{v}_b \hat{v}_c \hat{v}_d \right)$$

- ▶ At one loop order:

$$\Delta \hat{t}_a = \hat{\mu}_{ab}^2 \Delta \hat{v}_b + 3 \hat{\lambda}_{abcd} \hat{v}_c \hat{v}_d \Delta \hat{v}_b = \left(\hat{M}_0^2 \right)_{ab} \Delta \hat{v}_b,$$

- ▶ Finite tadpole contributions are absorbed in VEV shifts

$$\Delta \hat{v}_a = - \left(\hat{M}_0^2 \right)_{ab}^{-1} (T_{\text{fin}})_b.$$

- ▶ Note: $\hat{t}_a = 0$ equivalent to finding stationary points of potential

$$\hat{t}_a = \hat{\mu}_{ab}^2 \hat{v}_b + \hat{\lambda}_{abcd} \hat{v}_b \hat{v}_c \hat{v}_d = \left. \frac{\partial V}{\partial \varphi_a} \right|_{\varphi_i = \hat{v}_i} \stackrel{!}{=} 0$$

Results

Renormalized self-energies

Majorana fermions

$$\begin{aligned}\Sigma(p) = & \Sigma^{1\text{-loop}}(p) - \not{p} \left[\hat{\delta}^{(\chi)} \gamma_L + \left(\hat{\delta}^{(\chi)} \right)^* \gamma_R \right] \\ & + \hat{v}_a \left[\delta \hat{Y}_a \gamma_L + (\delta \hat{Y}_a)^* \gamma_R \right] + \Delta \hat{v}_a \left[\hat{Y}_a \gamma_L + \hat{Y}_a^* \gamma_R \right],\end{aligned}$$

Real scalars

$$\Pi_{ab}(p^2) = \Pi_{ab}^{1\text{-loop}}(p^2) - \hat{\delta}_{ab}^{(\varphi)} p^2 + \delta \hat{\mu}_{ab}^2 + 3\delta \hat{\lambda}_{abcd} \hat{v}_c \hat{v}_d + 6\hat{\lambda}_{abcd} \hat{v}_c \Delta \hat{v}_d$$

*explicit results for self-energies and counterterms in [2]

Results

One-loop masses

Majorana fermions:

$$m_i = m_{0i} + m_{0i} \left(\Sigma_L^{(A)} \right)_{ii} (m_{0i}^2) + \text{Re} \left(\Sigma_L^{(B)} \right)_{ii} (m_{0i}^2),$$

Dirac fermions:

$$m_i = m_{0i} + \frac{1}{2} m_{0i} \left[\left(\Sigma_L^{(A)} \right)_{ii} (m_{0i}^2) + \left(\Sigma_R^{(A)} \right)_{ii} (m_{0i}^2) \right] + \text{Re} \left(\Sigma_L^{(B)} \right)_{ii} (m_{0i}^2)$$

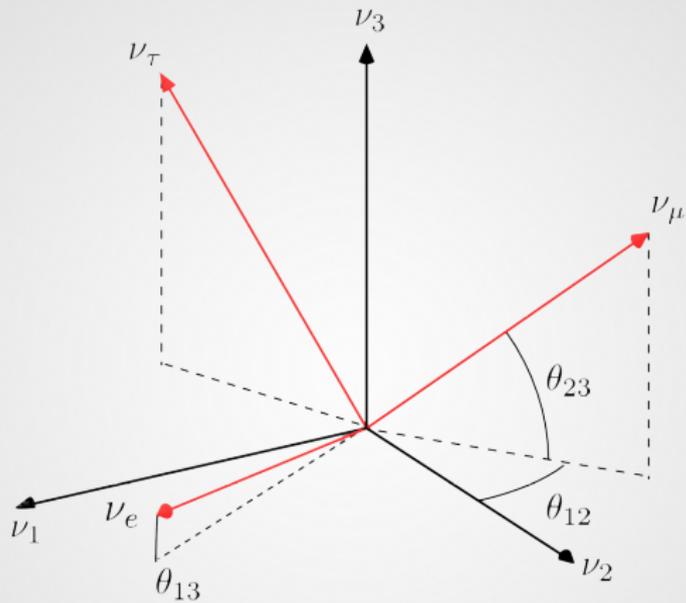
Real scalars:

$$M_a^2 = M_{0a}^2 + \Pi_{aa}(M_{0a}^2)$$

- Finite mass shifts via radiative corrections to couplings

Conclusion/Outlook

- **General idea:** build a tool for studying radiative corrections to masses and mixing angles in neutrino mass models with flavor symmetries
- Contribute to recent discussions on tadpole renormalization
 - ▶ see e. g. [4, 5, 6]
- **Next steps:**
 - ▶ Need gauge interactions to see effect of mixing angle renormalization
 - ▶ Elevate toy model to gauge theory
 - ▶ Relate approach to N-Higgs doublet models (NHDM)
 - ▶ Determine radiative corrections to masses and mixing angles in showcase ν -mass models known from the literature



Thanks!

- [1] S. Stone, *New physics from flavour*, PoS ICHEP **2012** (2013) 033 [arXiv:1212.6374 [hep-ph]].
- [2] M. Fox, W. Grimus and M. Löschner, *Renormalization and radiative corrections to masses in a general Yukawa model*, arXiv:1705.09589 [hep-ph].
- [3] W. Grimus and L. Lavoura, *A Nonstandard CP transformation leading to maximal atmospheric neutrino mixing*, Phys. Lett. B **579** (2004) 113 doi:10.1016/j.physletb.2003.10.075 [hep-ph/0305309].
- [4] A. Denner, L. Jenniches, J. N. Lang and C. Sturm, *Gauge-independent \overline{MS} renormalization in the 2HDM*, JHEP **1609** (2016) 115 [arXiv:1607.07352 [hep-ph]].
- [5] M. Krause, R. Lorenz, M. Mühlleitner, R. Santos and H. Ziesche, *Gauge-independent renormalization of the 2-Higgs-doublet model*, JHEP **1609** (2016) 143 [arXiv:1605.04853 [hep-ph]].
- [6] L. Altenkamp, S. Dittmaier and H. Rzehak, *Renormalization schemes for the two-Higgs-doublet model and applications to $h \rightarrow WW/ZZ \rightarrow 4$ fermions*, arXiv:1704.02645 [hep-ph].