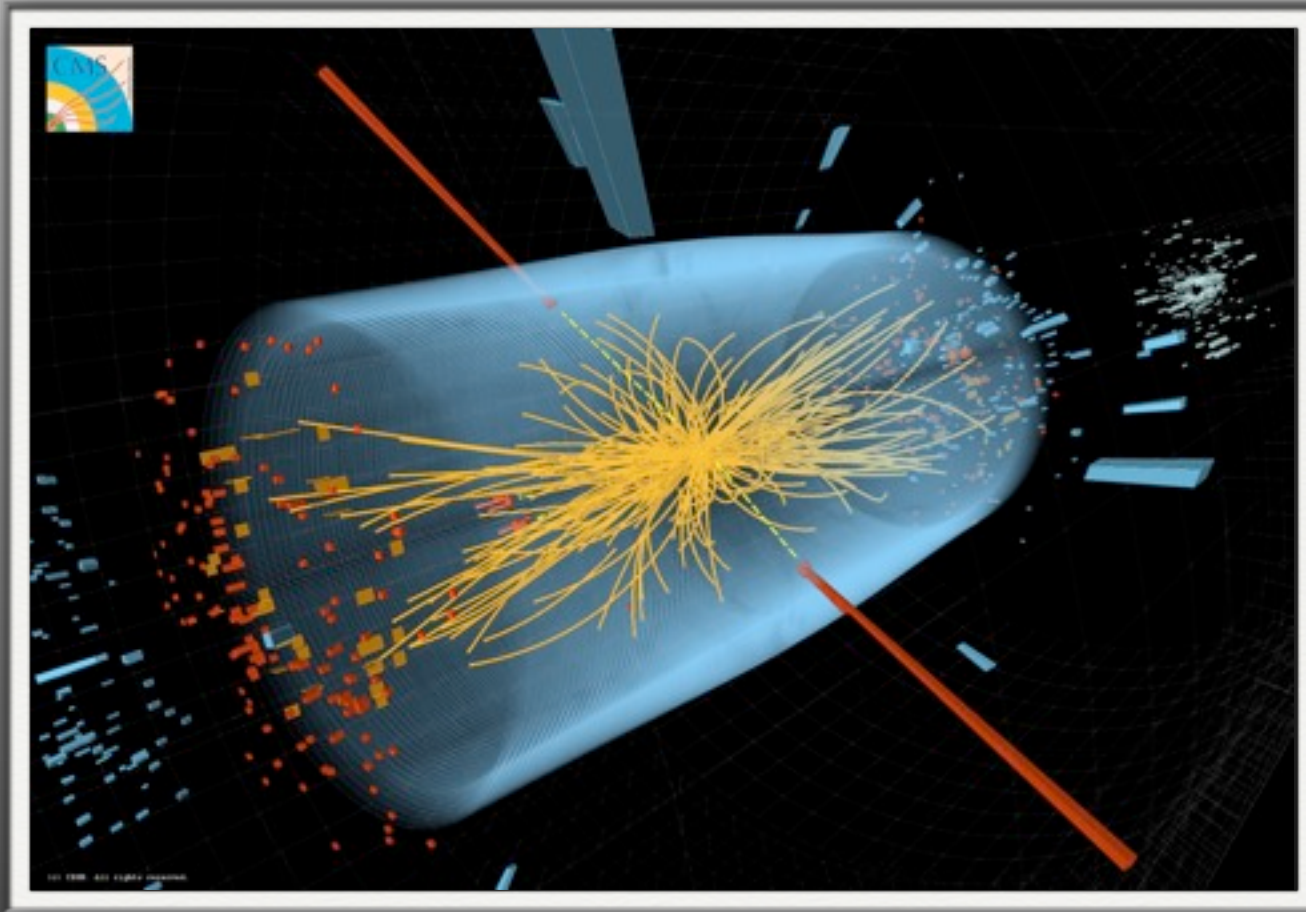


Particle ID: Lecture #2



Richard Cavanaugh, Fermilab & University of Illinois Chicago
LHC Physics Center co-Coordinator

Hadron Collider Physics Summer School
Fermilab, 14 August, 2012



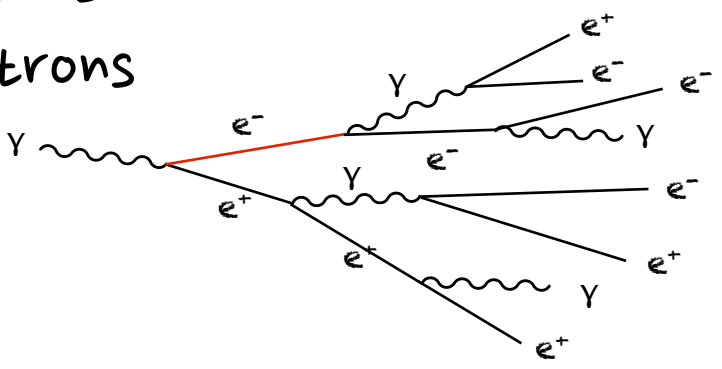
Particle interactions in material



Particle interactions in material

Photons

Electrons

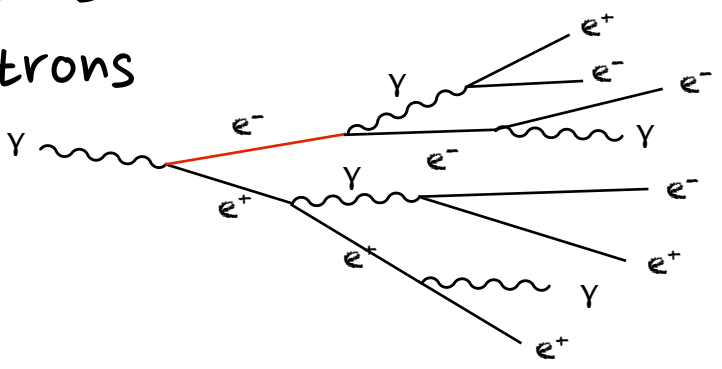




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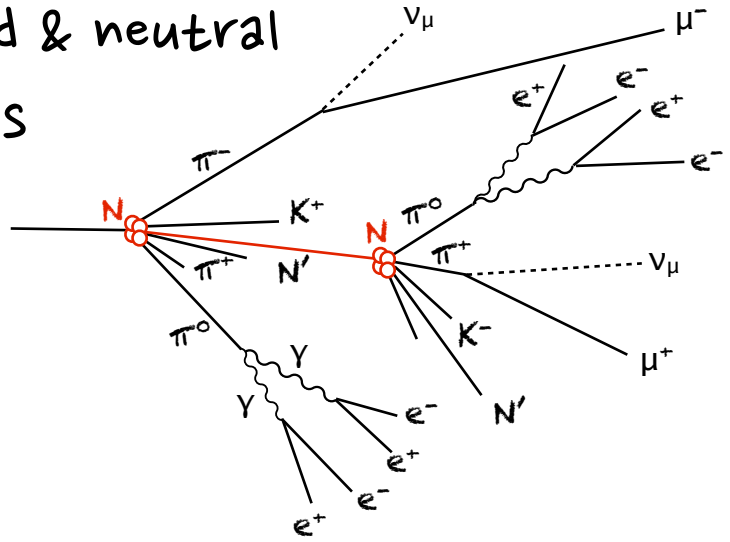
Electrons





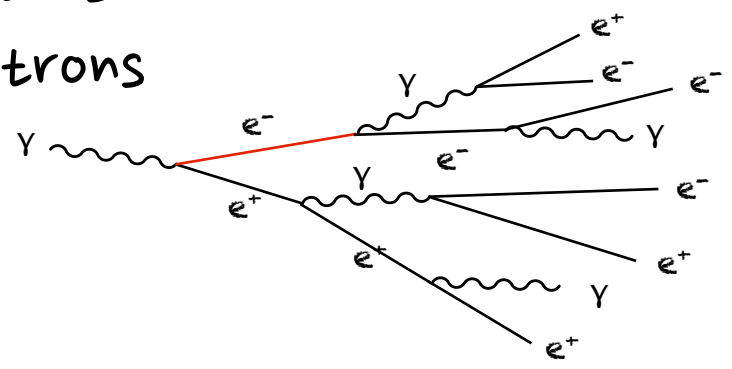
Particle interactions in material

charged & neutral
hadrons



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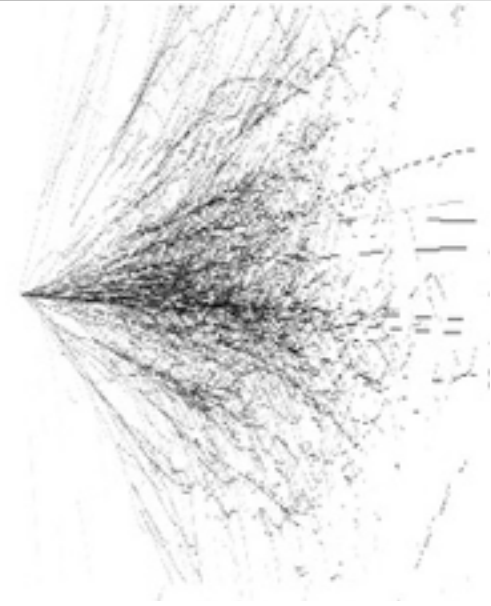
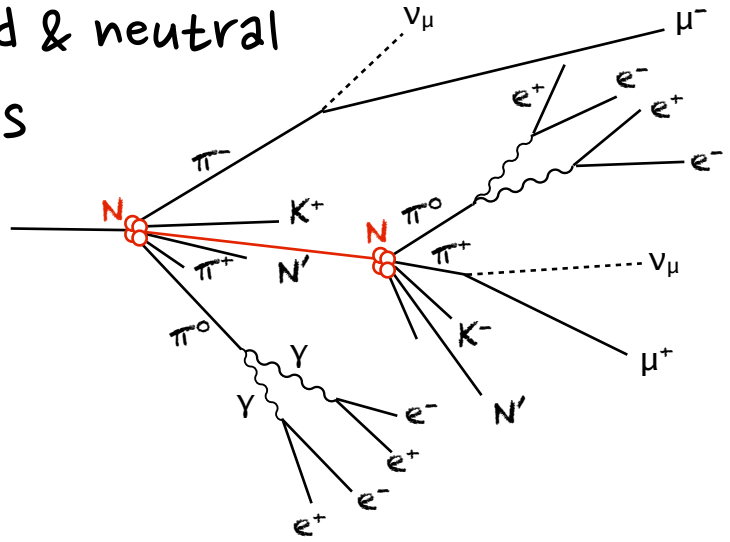
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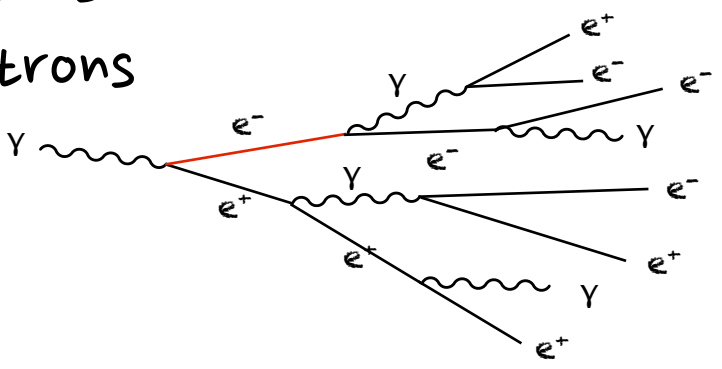


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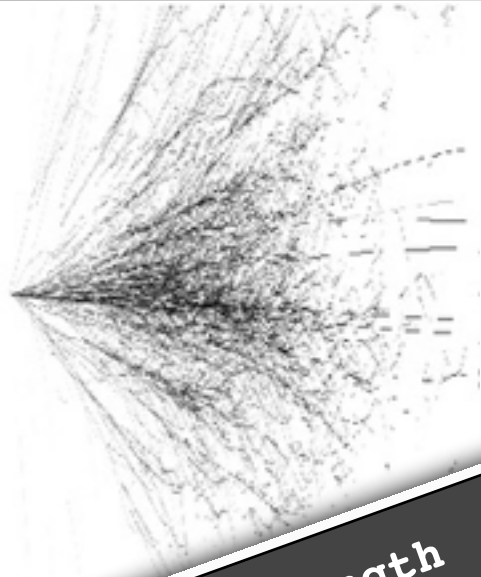
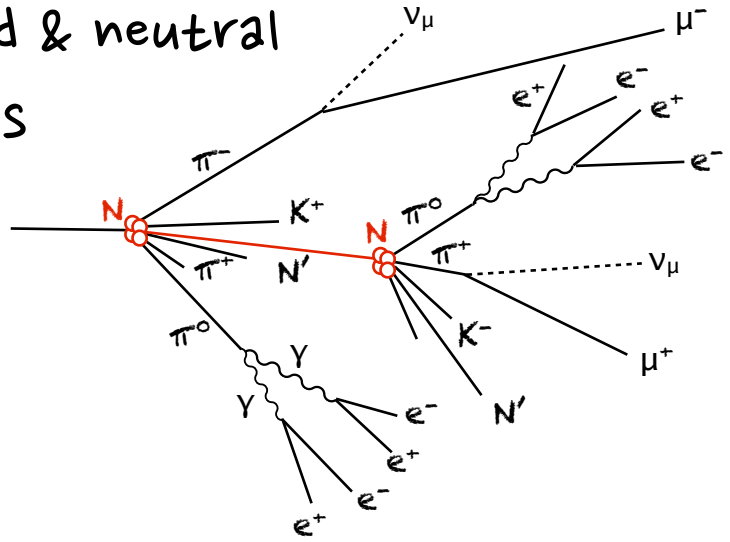
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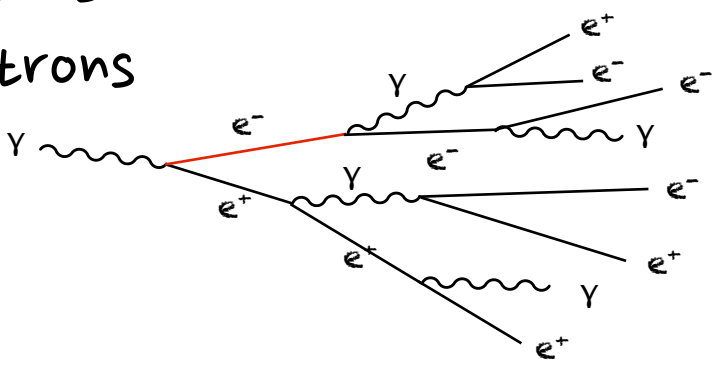


Particle interactions in material

charged & neutral hadrons



Photons
Electrons



Radiation Length

- Characteristic distance over which the electron energy is reduced by a factor of 1/e due to radiation losses only

$$X_0 = \frac{716.4 \text{ g cm}^{-2} A}{Z(Z+1) \ln(287/\sqrt{Z})}$$

- Higher Z materials have short length
- want high-Z material for EM calorimeter

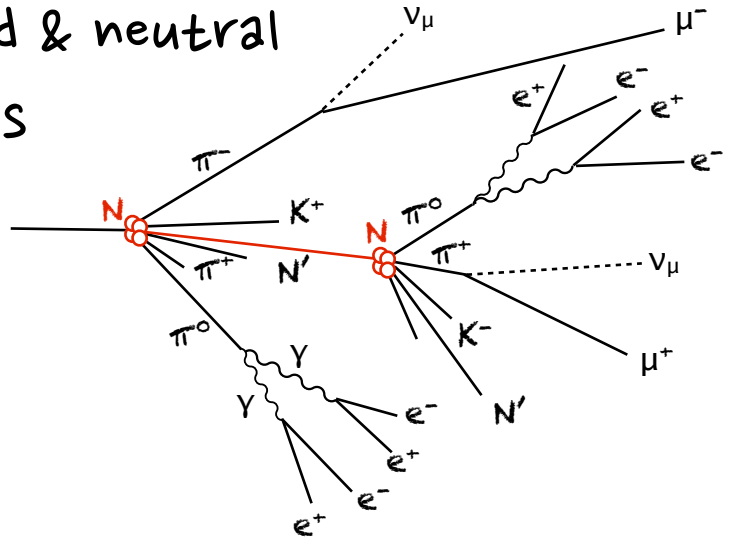
Example: Lead
 $\rho = 11.4 \text{ g/cm}^3$; $X_0 = 5.5 \text{ mm}$



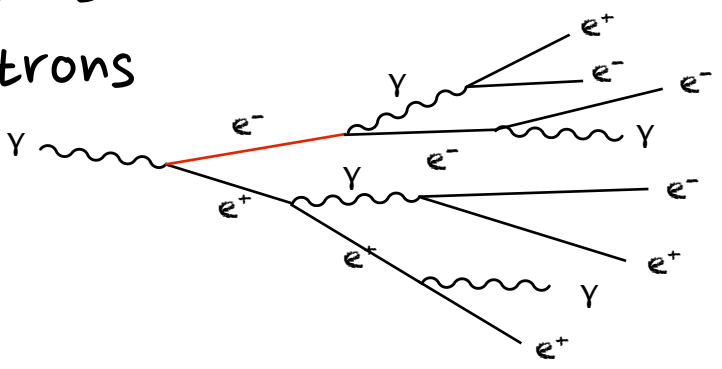
Particle interactions in m

milab

charged & neutral hadrons



Photons
Electrons



Nuclear Interaction Length

- Collisions of hadrons with nuclei produce hadronic showers
- $\lambda \approx 35 \text{ g cm}^{-2} A^{1/3}$
- Nuclear interaction length much longer in high-Z materials than EM radiation length $\lambda > X_0$
- Hadronic showers develop later than EM showers; more diffuse
- Example: Lead $\lambda = 17 \text{ cm}$

Radiation Length

- Characteristic distance over which the electron energy is reduced by a factor of $1/e$ due to radiation losses only
- $$X_0 = \frac{716.4 \text{ g cm}^{-2} A}{Z(Z+1) \ln(287/\sqrt{Z})}$$
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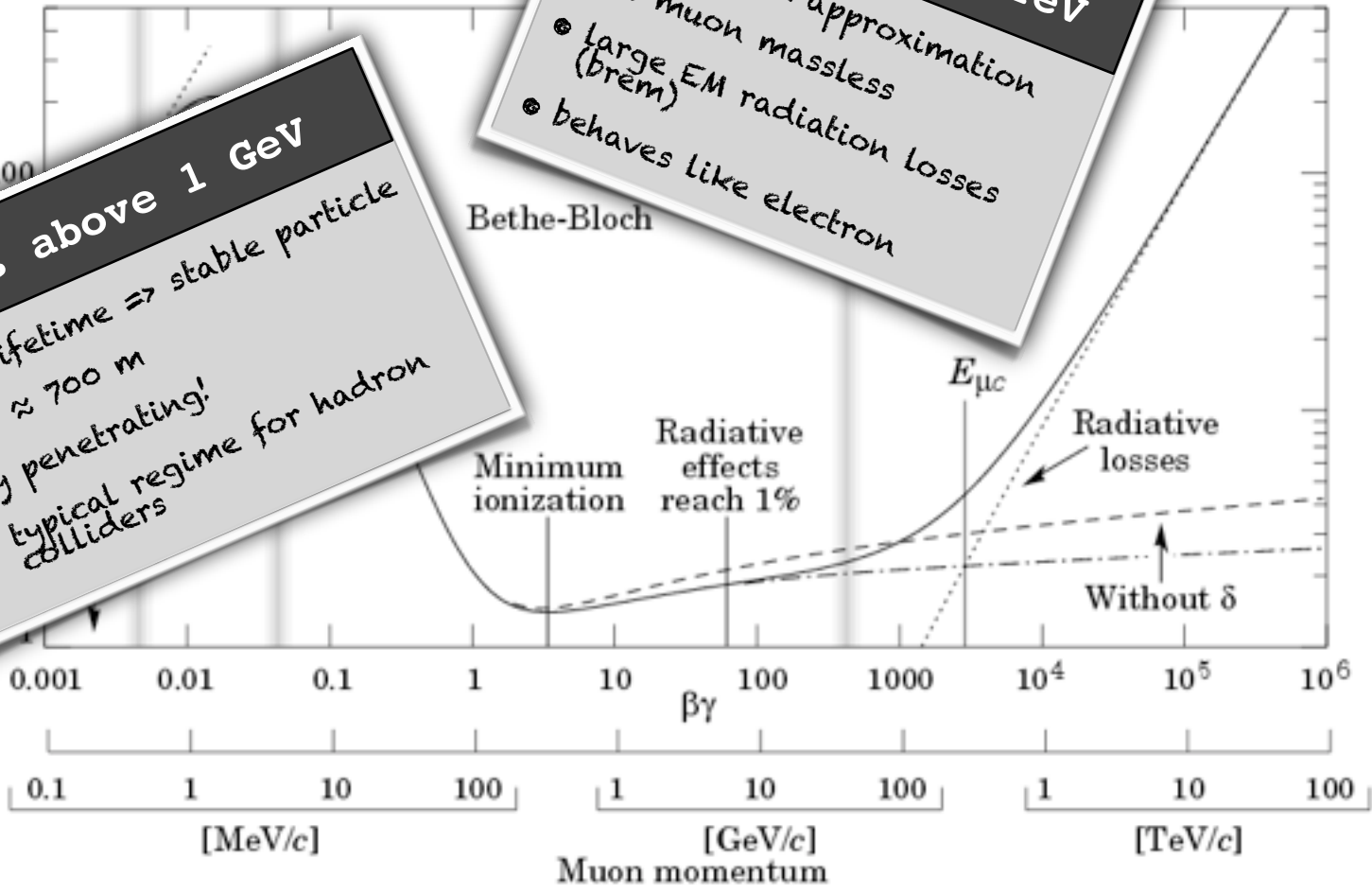
Particle interactions in material

Muons above 1 GeV

- Long lifetime \Rightarrow stable particle
 - $c\tau \approx 700$ m
- Very penetrating!
 - typical regime for hadron colliders

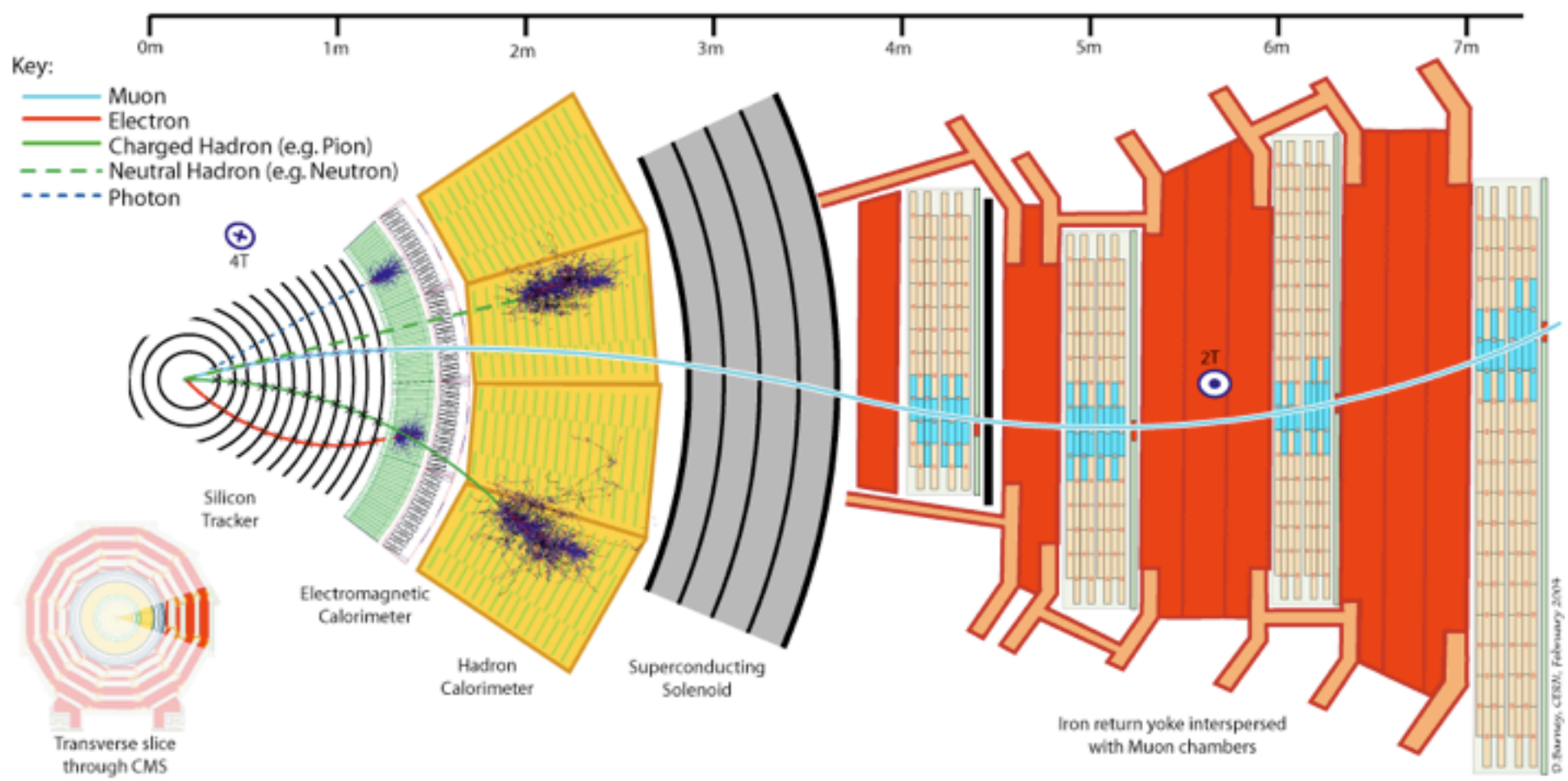
Muons above 1 TeV

- Very good approximation
- muon massless
- large EM radiation losses (brem)
- behaves like electron



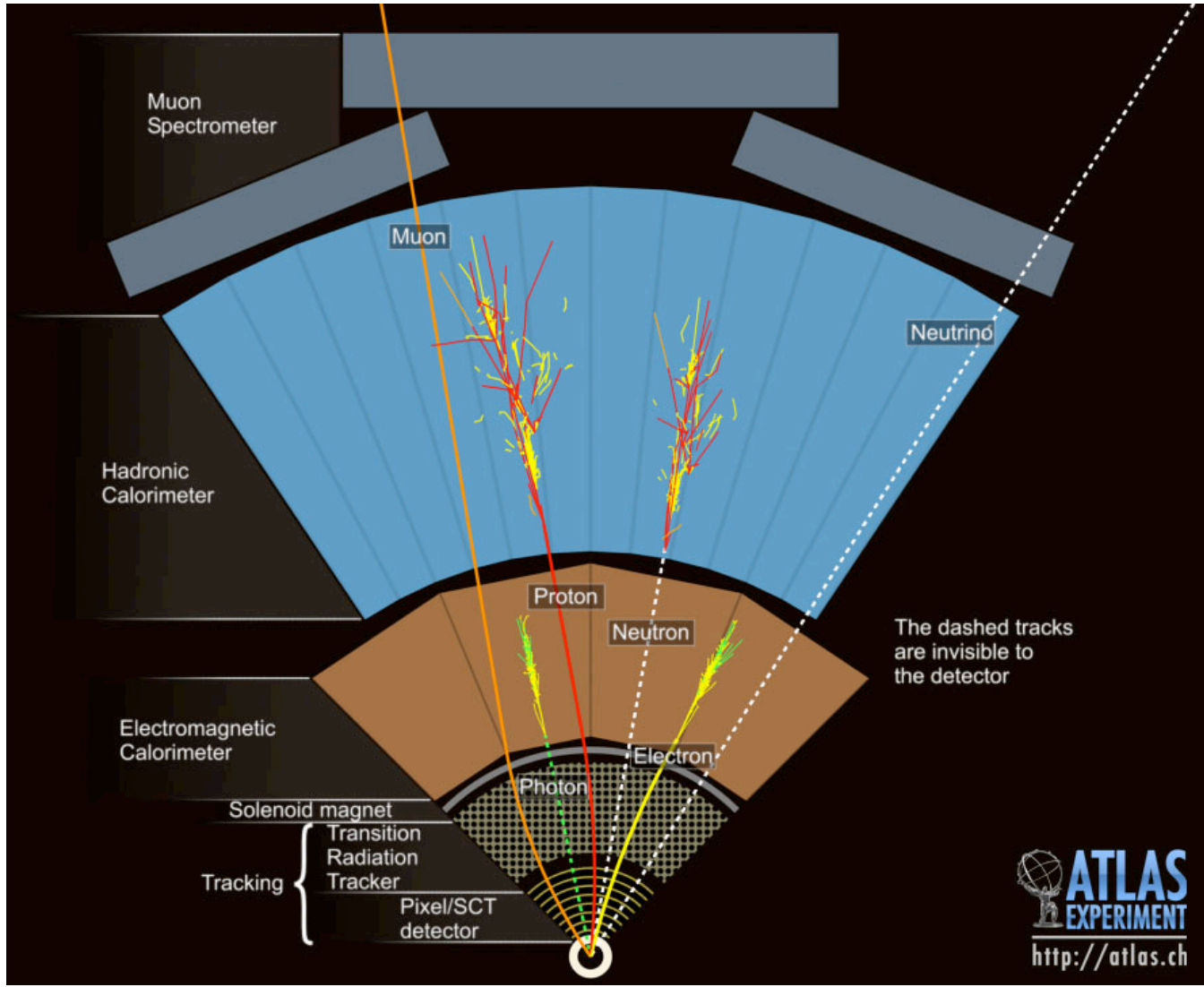


Particle Passage through CMS

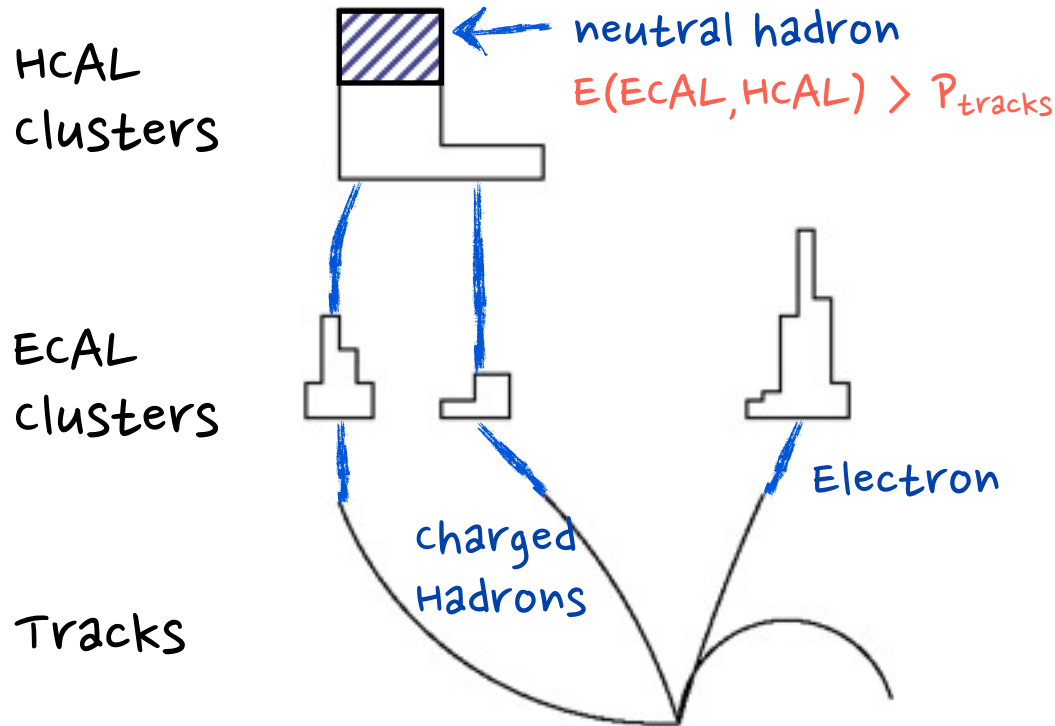




Particle Passage through ATLAS



In a nutshell



- **Aim: Reconstruct and identify all particles**
 - $\gamma, e, \mu, \pi^\pm, K_L^0, \text{pile-up } \pi^\pm, \text{converted } \gamma \& \text{ nuclear interaction } \pi^\pm, \dots$
 - Use best combination of all sub-detectors for E, η , φ , and ID



Last Time

- We arrived at a list of visible...
 - charged particles: could be p^+ , π^\pm , e^\pm , μ^\pm , ...
 - neutral photons: could be prompt γ 's
 - neutral hadrons: could be n^0 , K_L^0 , ...
- Then there was the case of $E \ll p$, when separating merged charged & neutral particles
- **Today:**
 - consider the case when $E \ll p$
 - identify which of the charged particles above are:
 - charged hadrons (1st)
 - electrons (2nd)
 - muons (last)
 - We also need to identify which of the photons are
 - prompt photons (3rd)
 - No need to further identify the neutral hadrons...



Cases with $E \ll p$



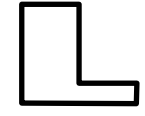
Cases with $E \ll p$

- Let's discuss the "special" case $E \ll p$

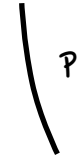


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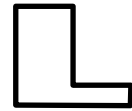
$E = \text{ECAL} + \text{HCAL}$




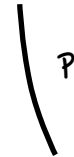


Cases with $E \ll p$

- Let's discuss the "special" case $E \ll p$
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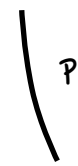
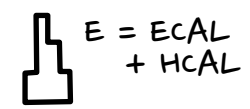
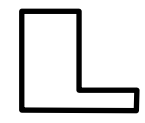
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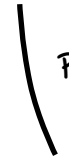
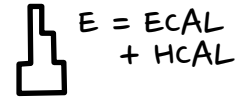
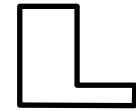
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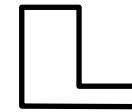
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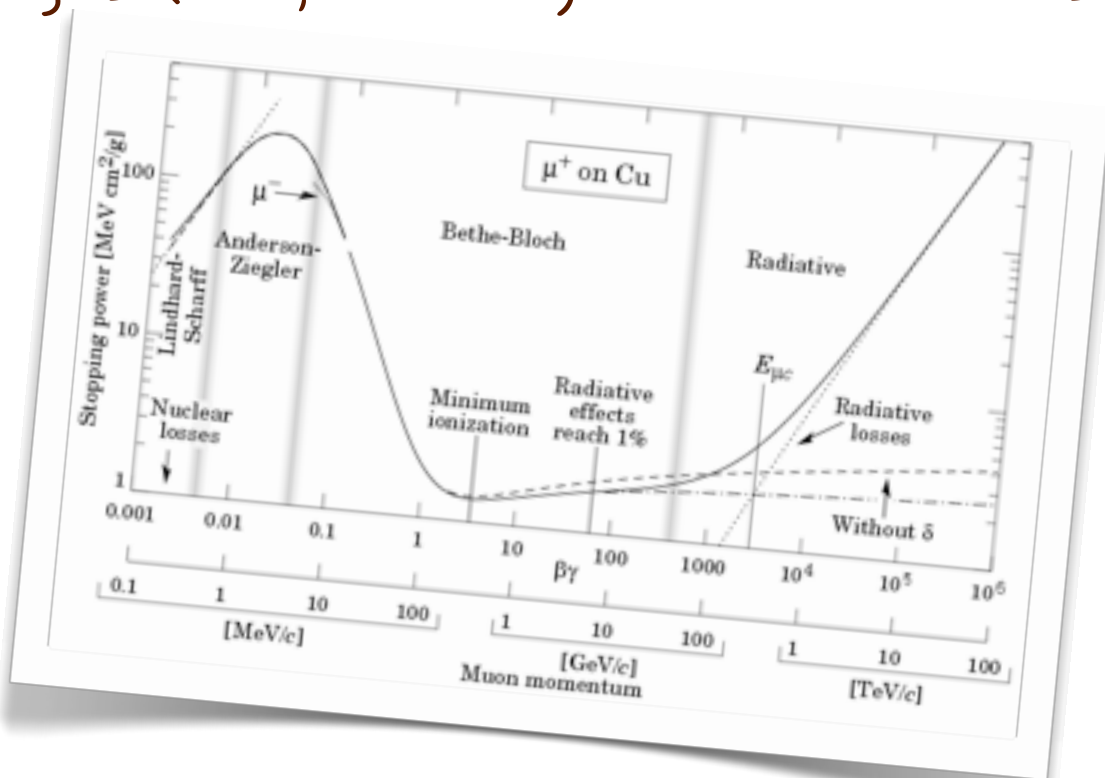
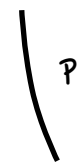


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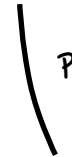
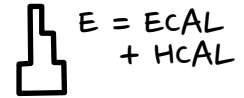
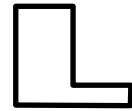
$$E = E_{\text{CAL}} + H_{\text{CAL}}$$





Cases with $E \ll p$

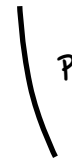
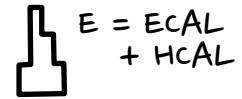
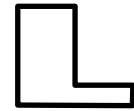
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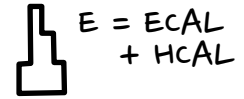
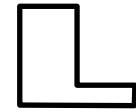
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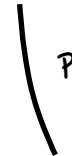
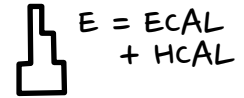
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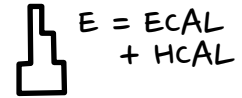
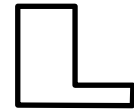
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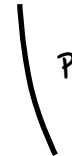
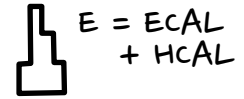
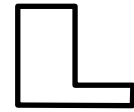
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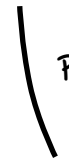
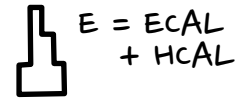
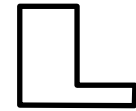
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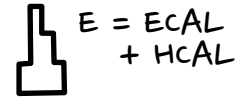
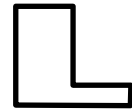
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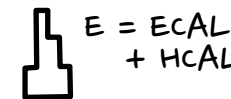
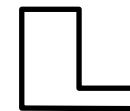
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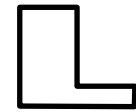
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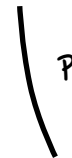


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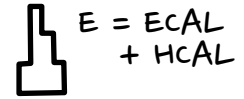
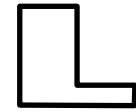
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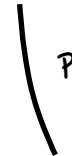
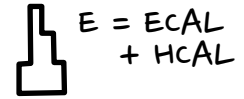
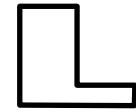
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- Philosophy:





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- Muon ID is very efficient $\geq 98\%$ (more later)
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 - before identifying charged hadrons, photons, neutral hadron
- Remaining $\leq 2\%$ lead to many of the $E \ll p$ cases
 - a much looser muon ID is used to help solve this case
 - isolated tracks (hadrons unlikely to be isolated)
 - tracks in tracker only (plus a few hits in the muon syst.)
 - poor-quality tracks in tracker (but good fit in muon syst.)
- Philosophy:
 - use all detectors to improve particle ID/Reco





Cases with $E \ll p$



Cases with $E \ll p$

- Fake tracks or momentum mis-measurements also lead to $E \ll p$



Cases with $E \ll p$

- Fake tracks or momentum mis-measurements also lead to $E \ll p$
 - despite the tight selection of tracks



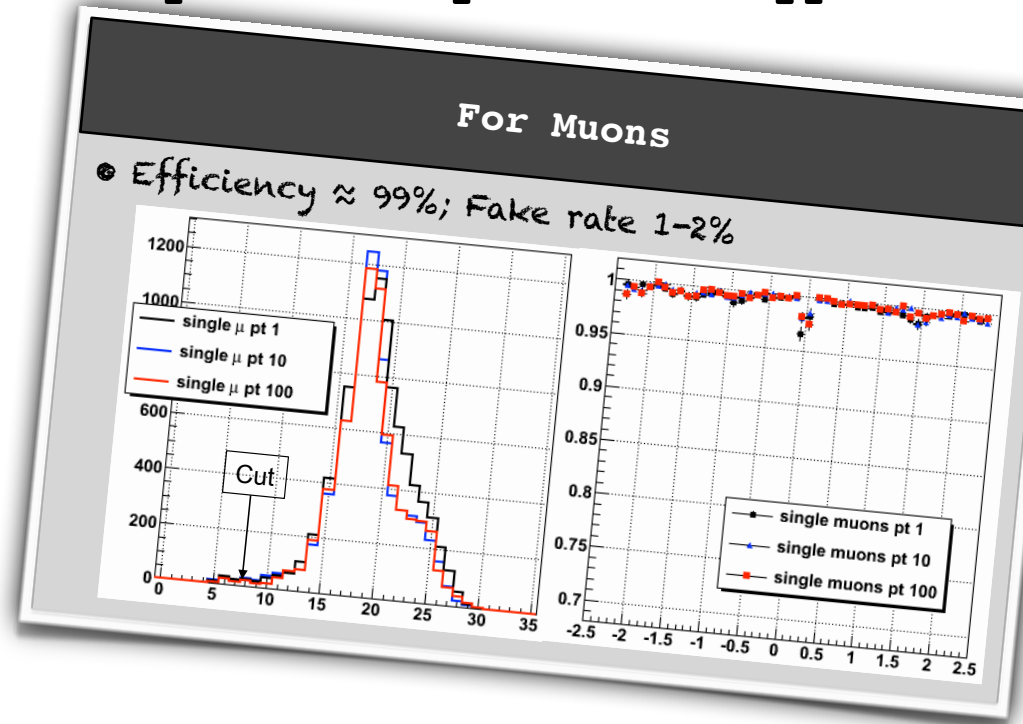
Cases with $E \ll p$

- Fake tracks or momentum mis-measurements also lead to $E \ll p$
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- Tracking eff. & fake rate depends on particle type



Cases with $E \ll p$

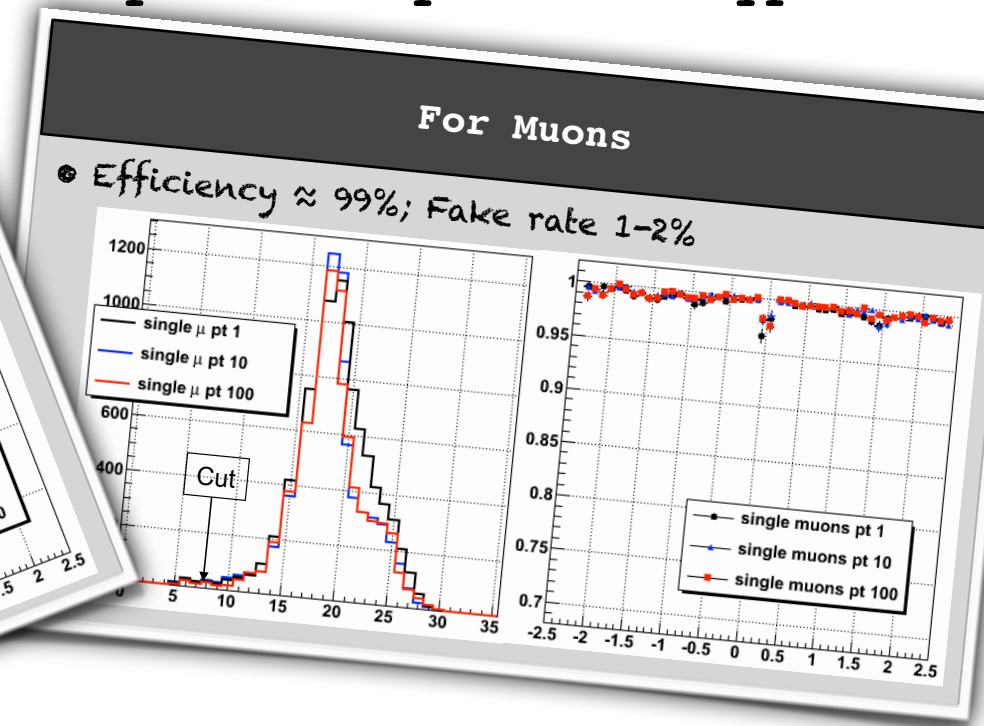
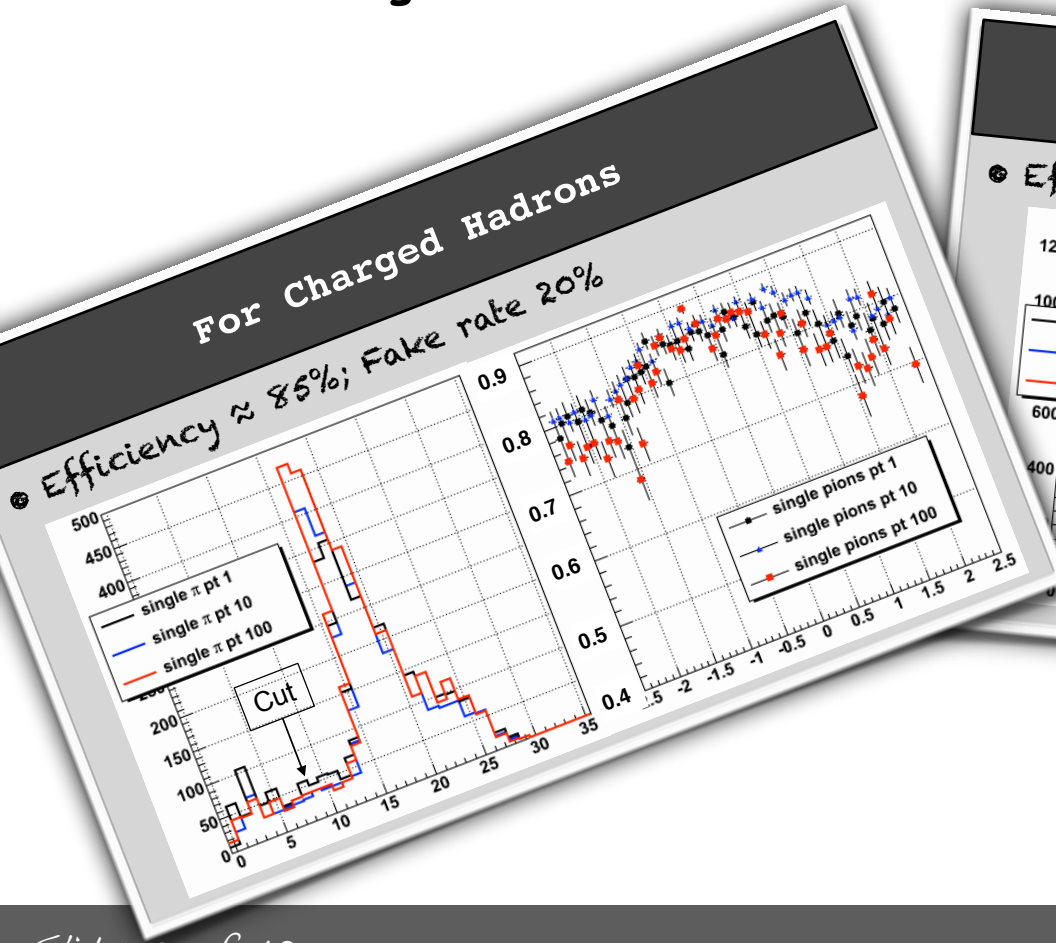
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Cases with $E \ll p$

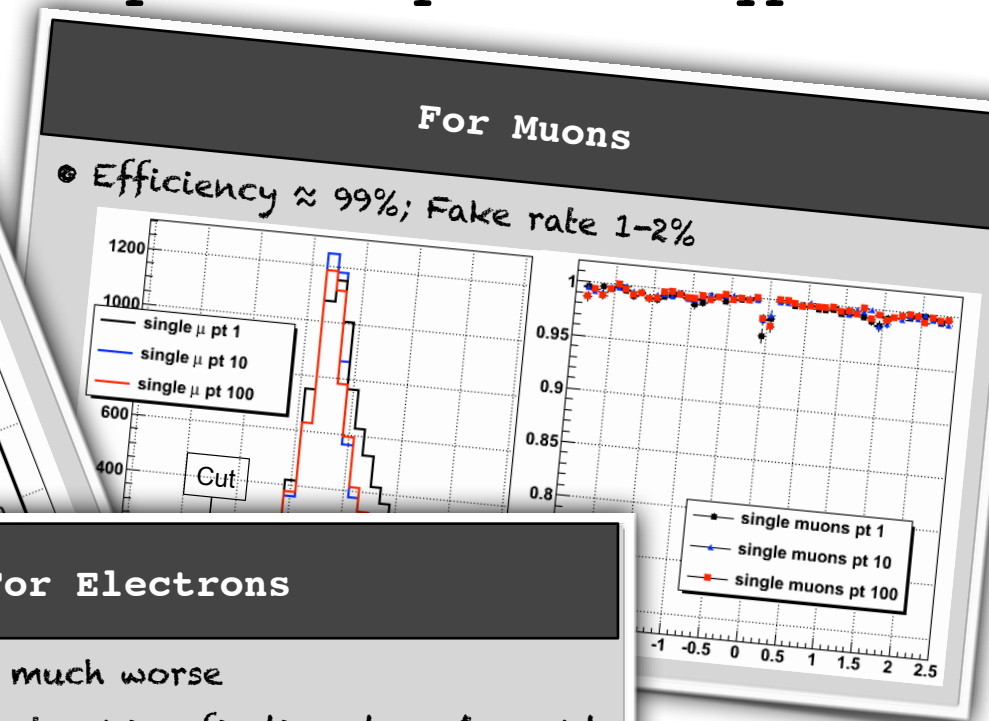
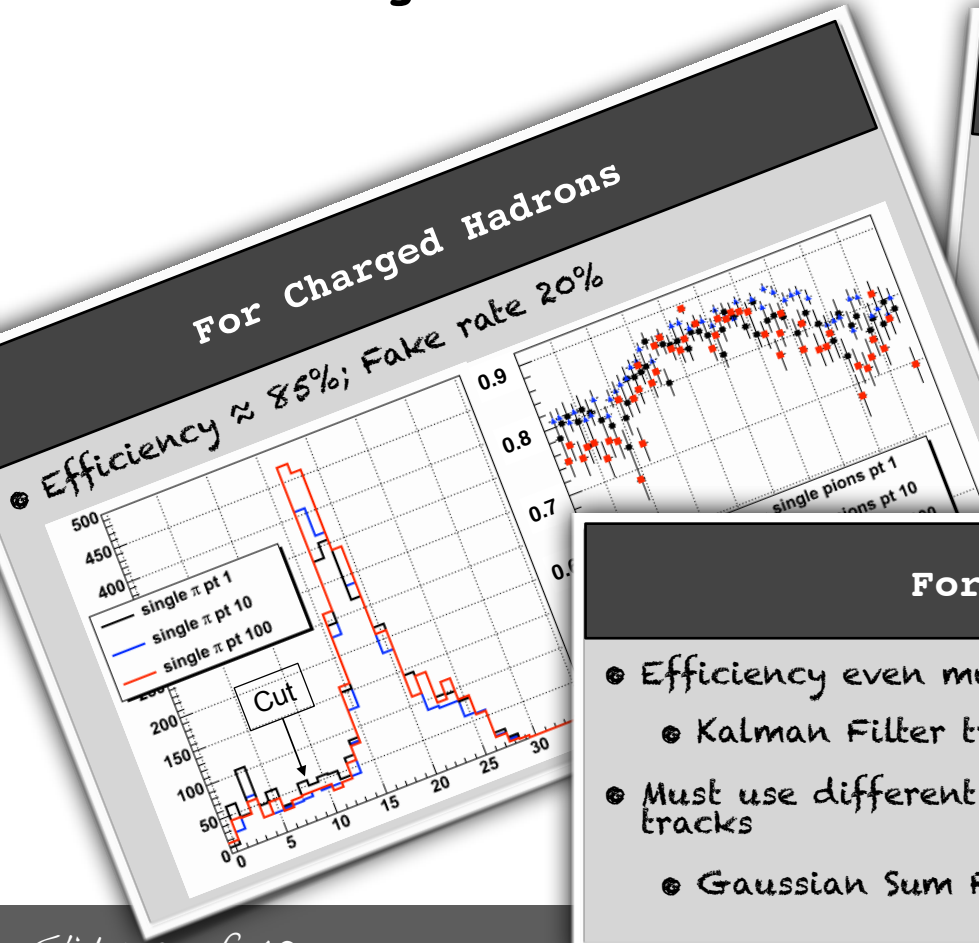
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Cases with $E \ll p$

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For Electrons

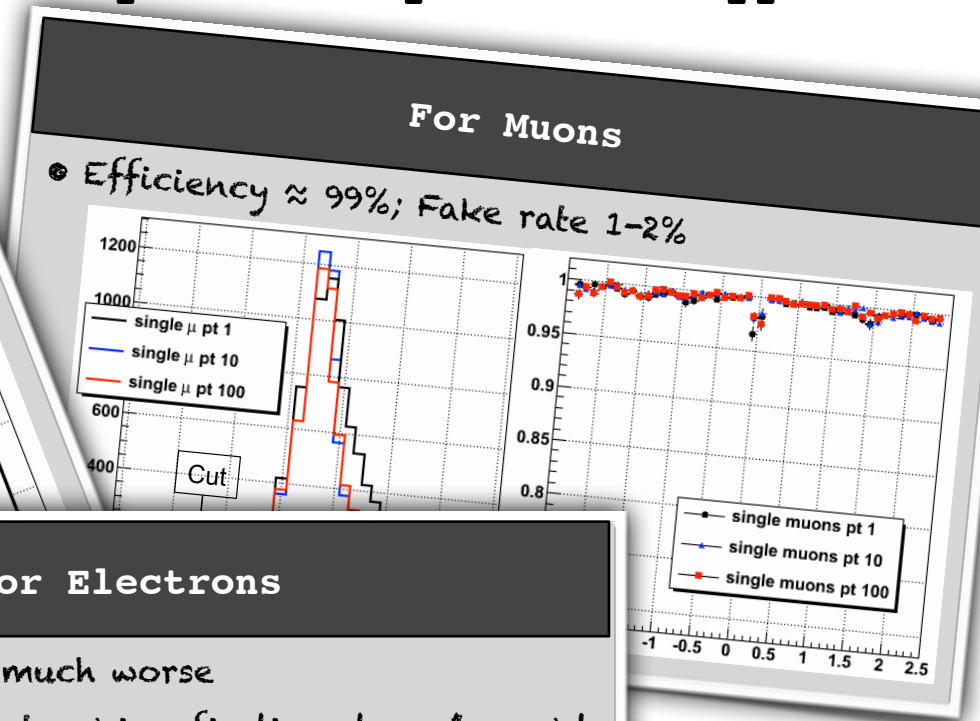
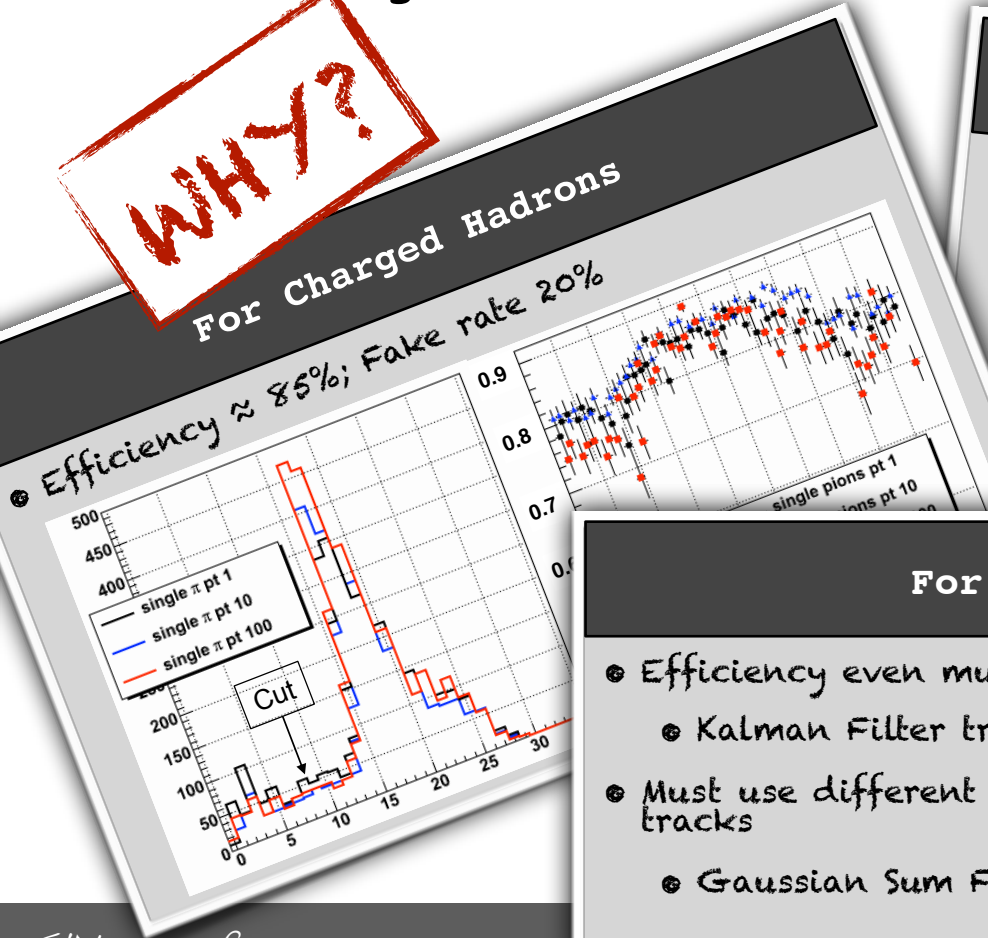
- Efficiency even much worse
 - Kalman Filter tracking finding doesn't work!
- Must use different method to find electron tracks
 - Gaussian Sum Filter (more Later)



Cases with $E \ll p$

- Fake tracks or momentum mis-measurements also lead to $E \ll p$
 - despite the tight selection of tracks
- Tracking eff. & fake rate depends on particle type

WHY?



For Electrons

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 - Kalman Filter tracking finding doesn't work!
- Must use different method to find electron tracks
 - Gaussian Sum Filter (more Later)



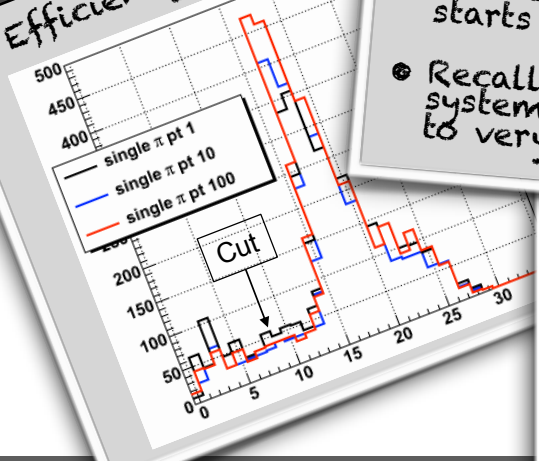
Cases with $E \ll p$

- Fake tracks or momentum mis-measurements also lead to $E \ll p$
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WHY?

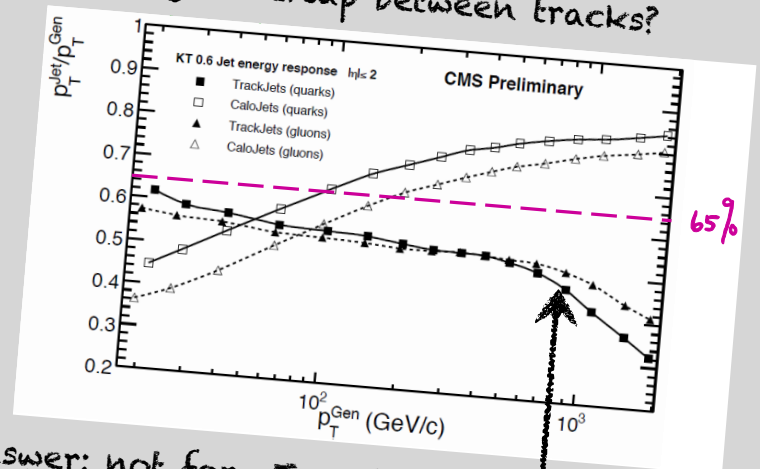
For Charge

Efficiency $\approx 85\%$; Fal



Could it be due to high occupancy inside jets?

Hence large overlap between tracks?



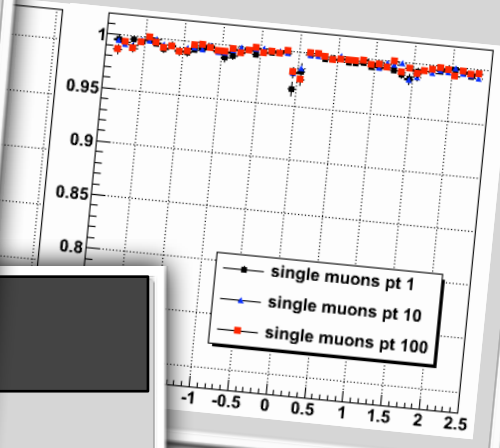
- Answer: not for p_T up to 800 GeV, where it starts to appear
- Recall; 10+60 million channels in tracker system! The tracker has very low occupancy up to very large numbers of particles

- Efficiency even more
- Kalman Filter tracking finding doesn't work!
- Must use different method to find electron tracks
- Gaussian Sum Filter (more later)

particle type

Muons

rate 1-2%





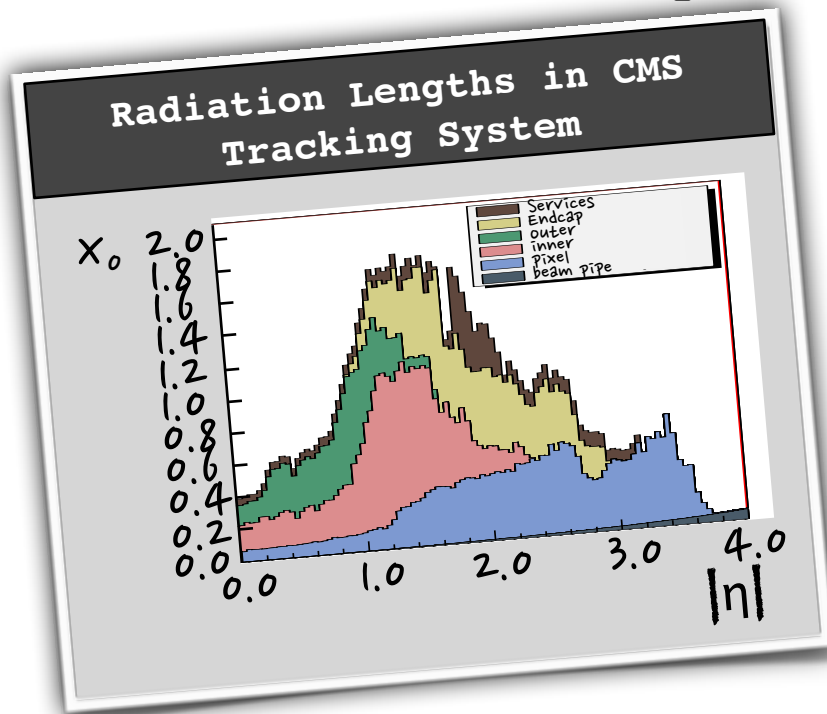
Answer: Tracker Material!

- **LHC Trackers are primarily silicon based => heavy!**



Answer: Tracker Material!

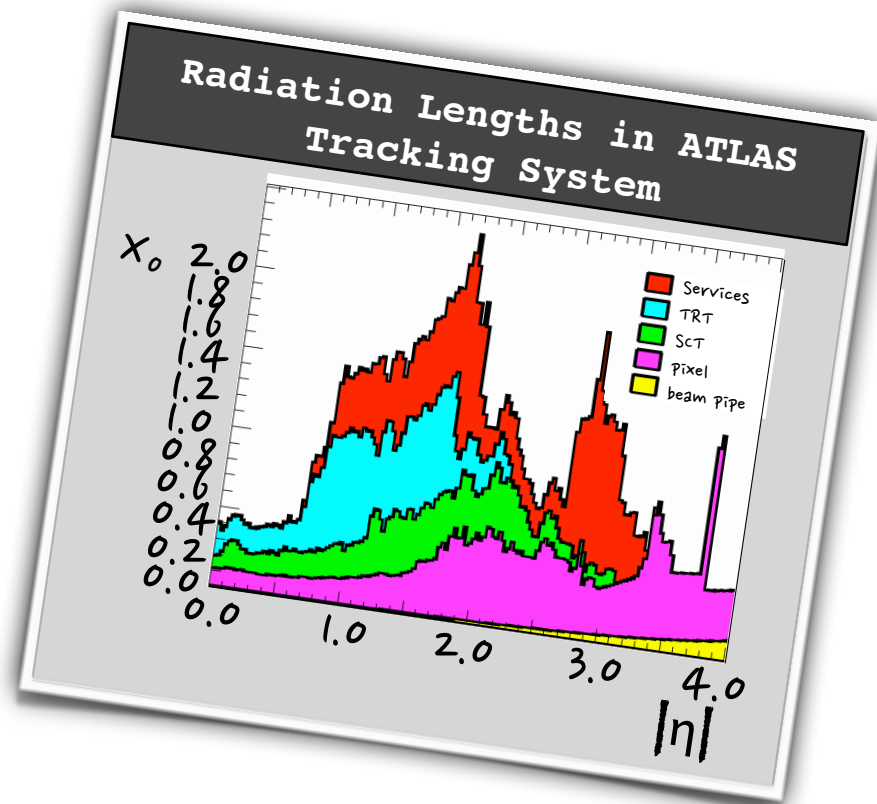
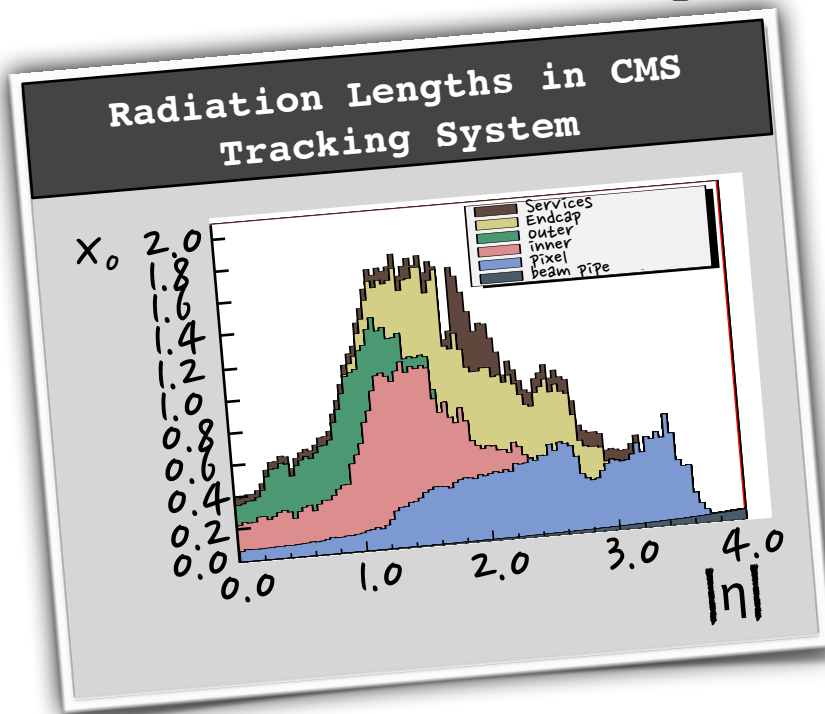
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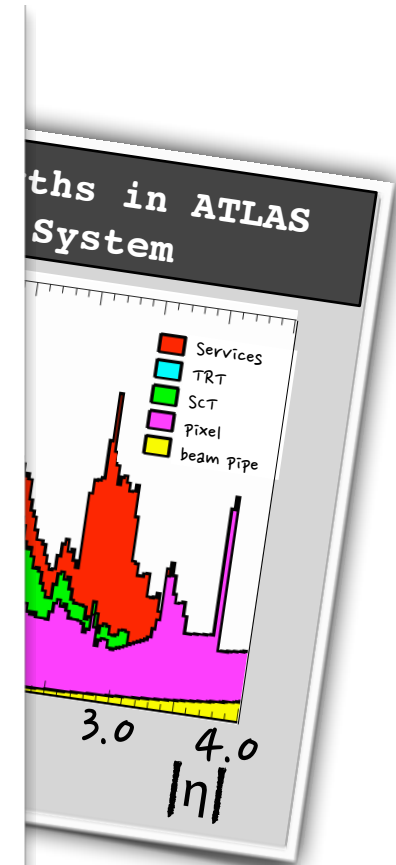
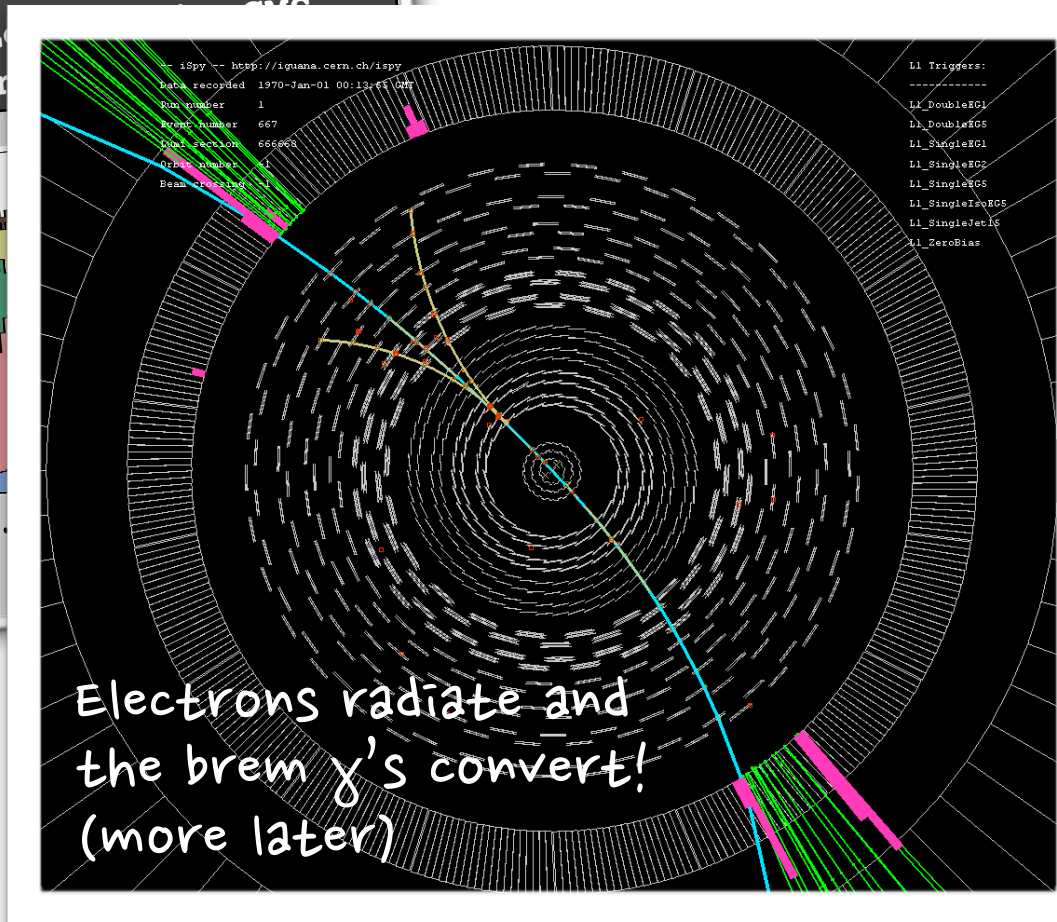
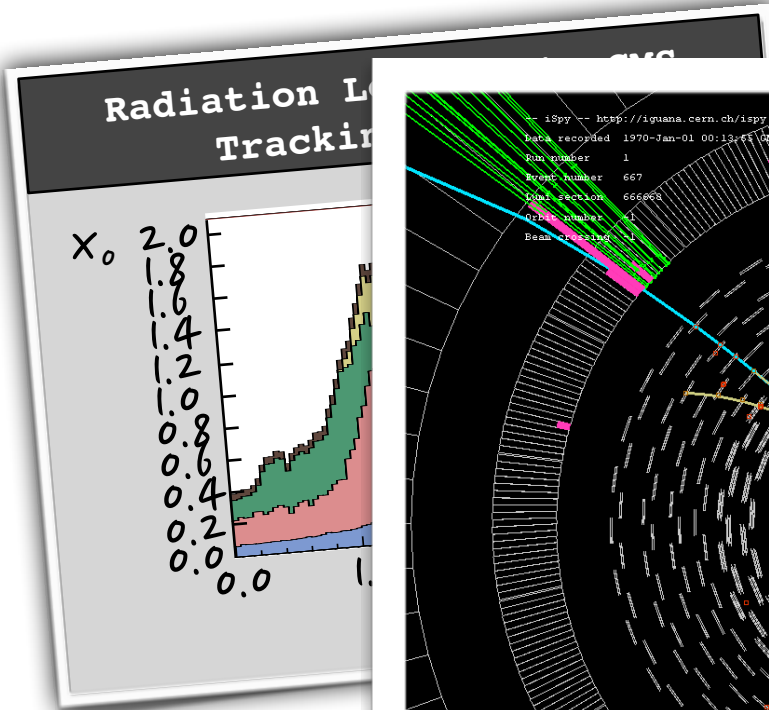
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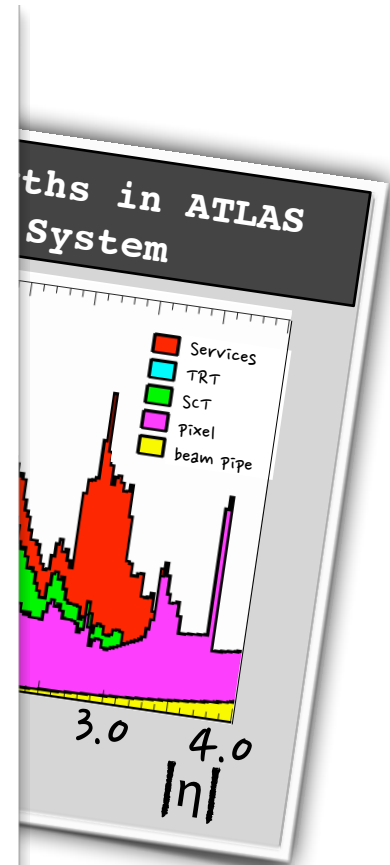
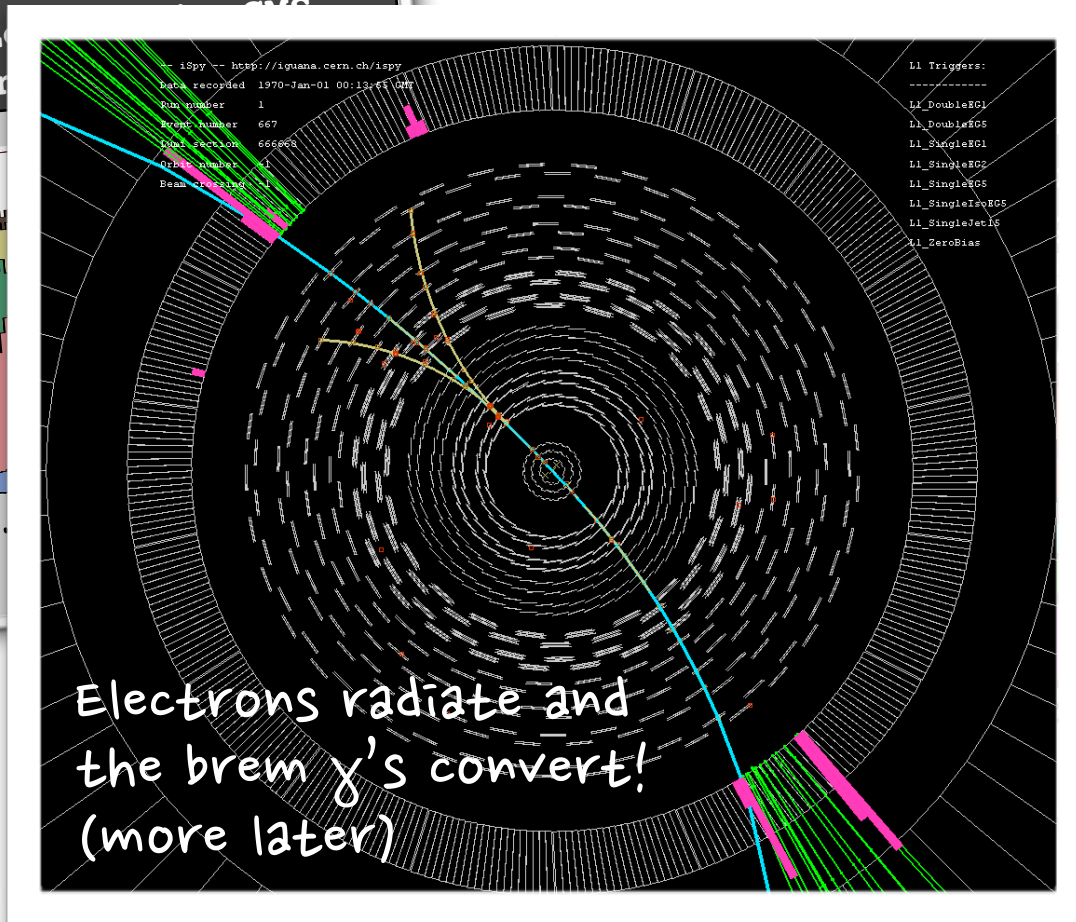
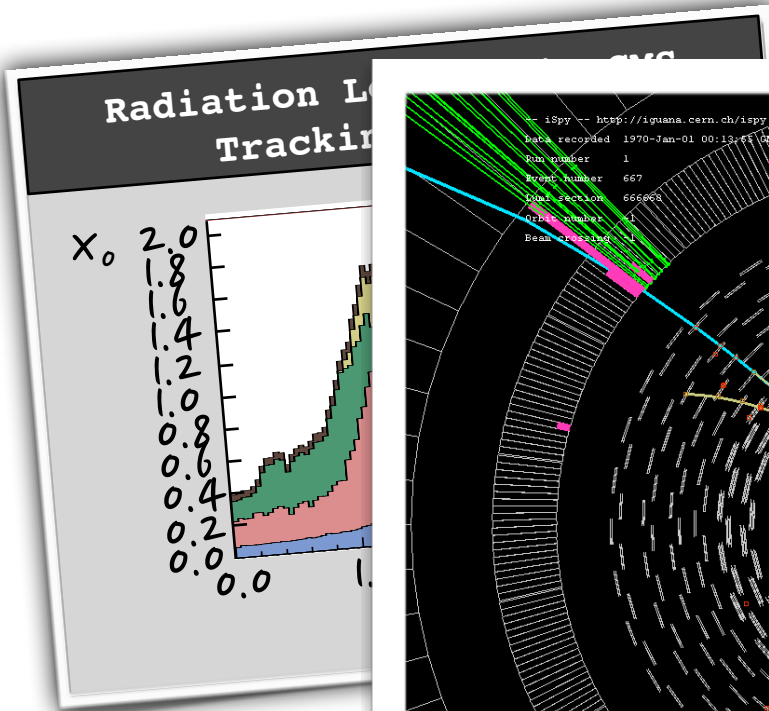
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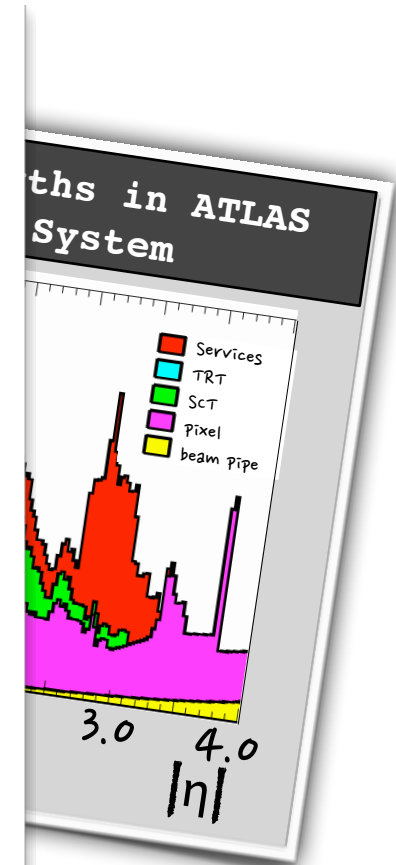
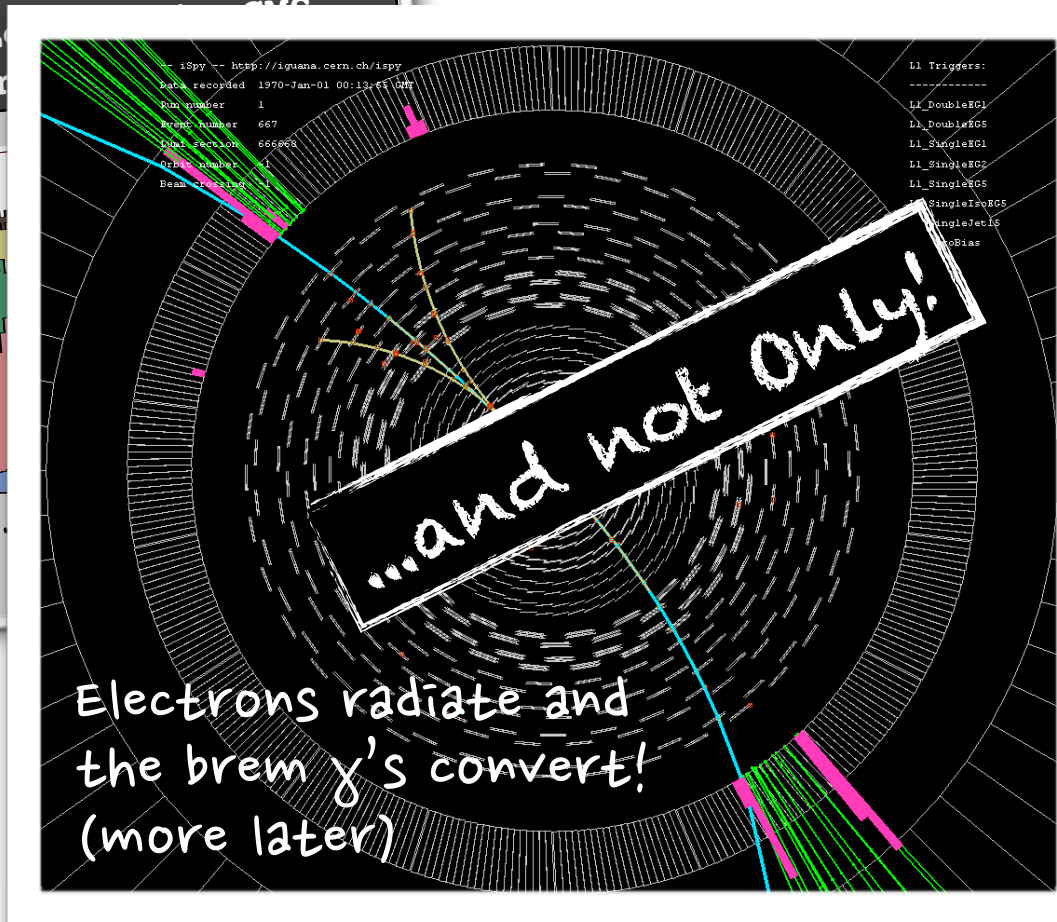
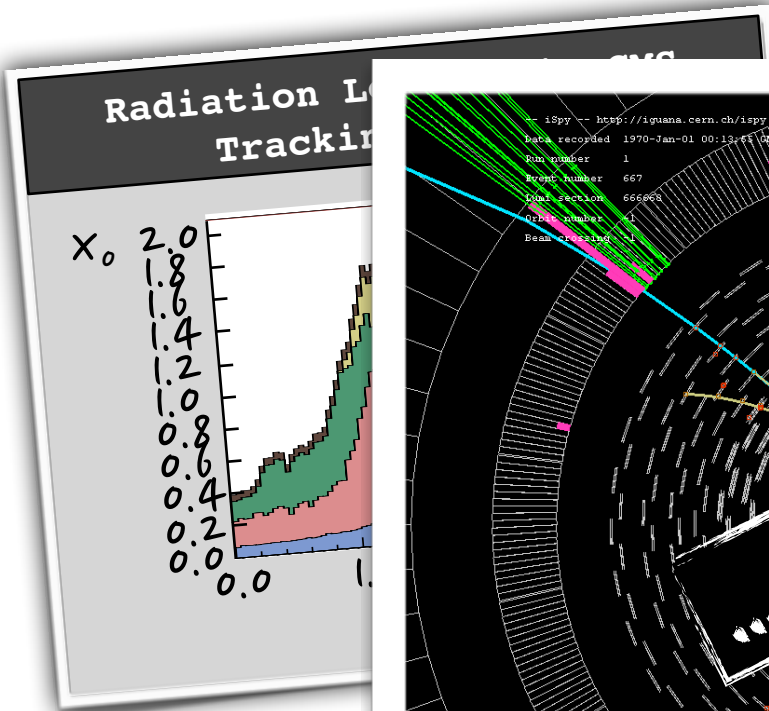


The LHC Trackers act like Electromagnetic pre-showers!



Answer: Tracker Material!

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The LHC Trackers act like Electromagnetic pre-showers!



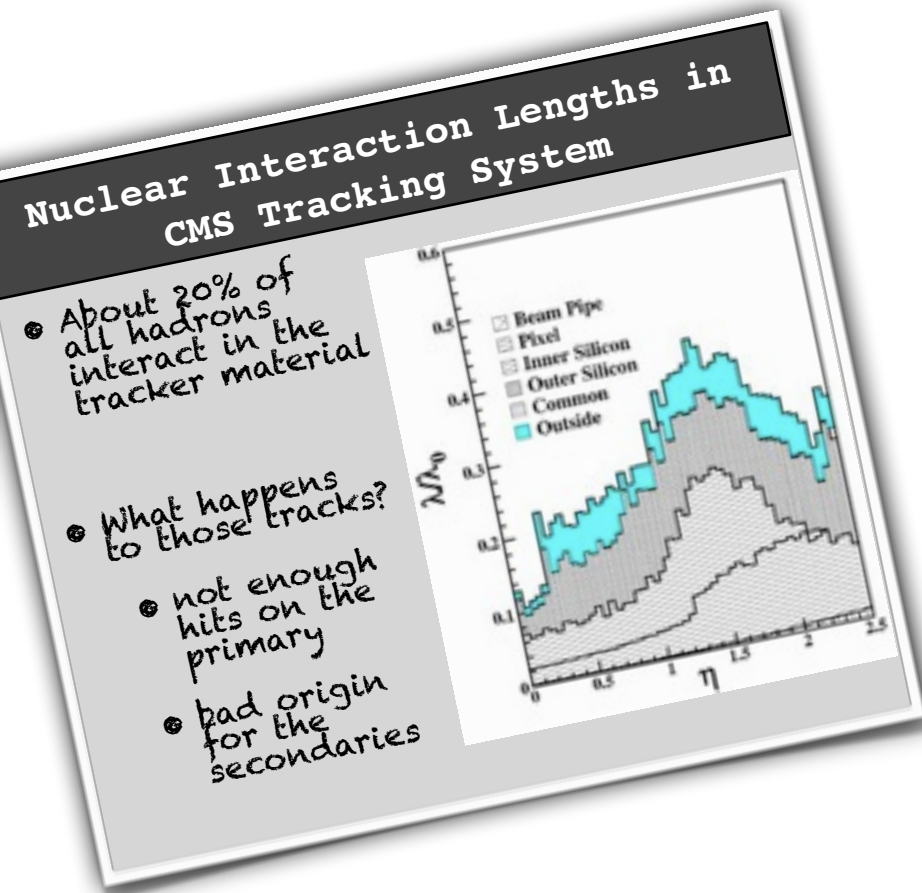
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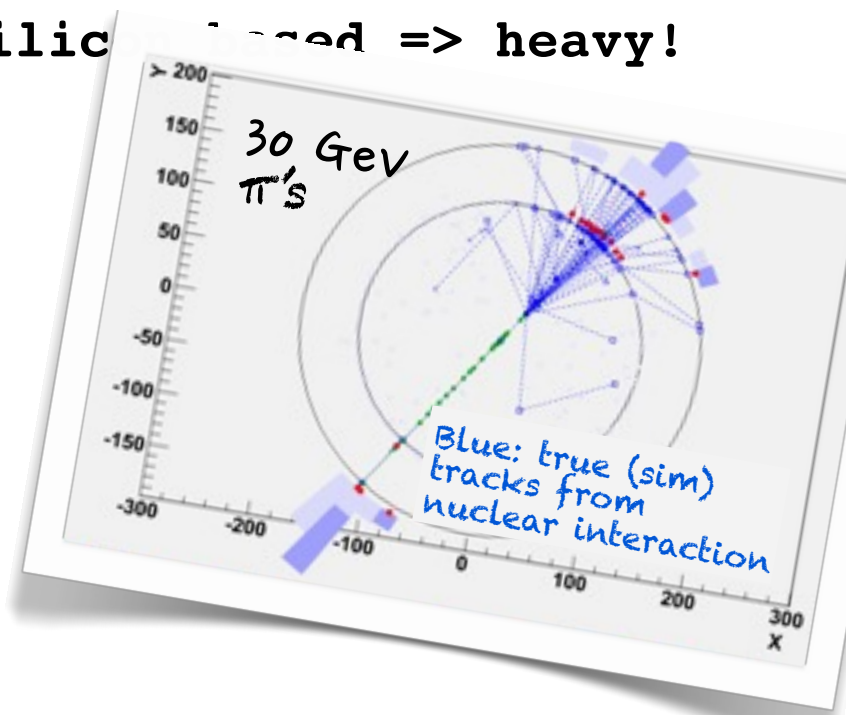
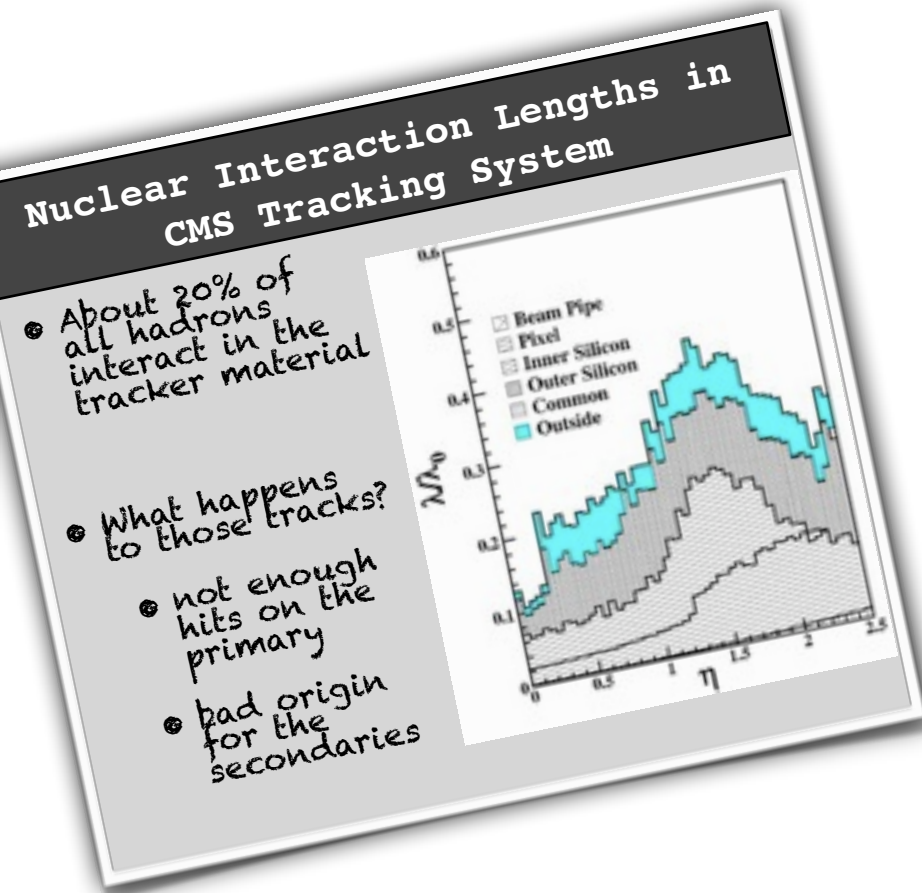
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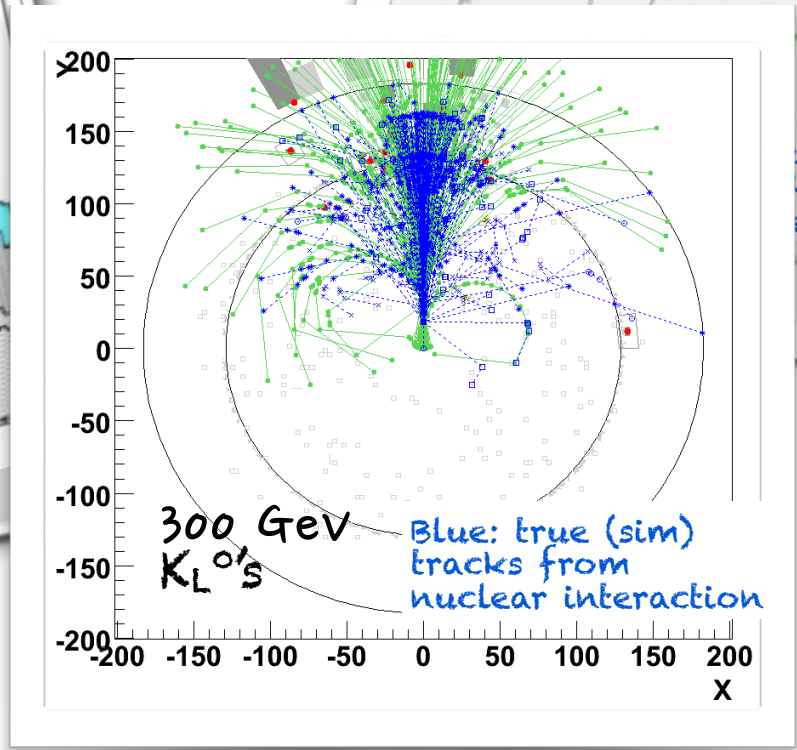
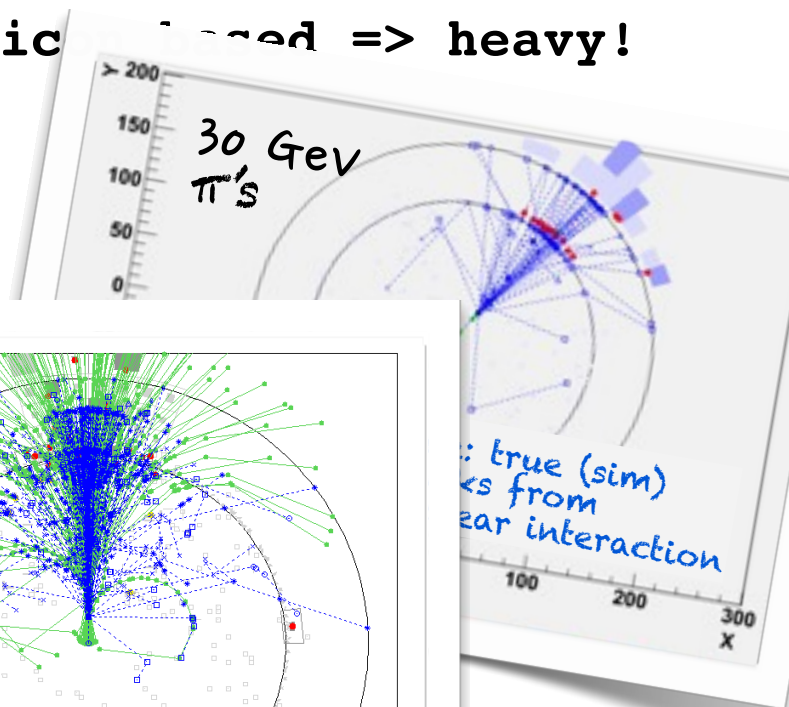
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Nuclear Interaction Lengths in CMS Tracking System

- About 20% of all hadrons interact in the tracker material
- What happens to those tracks?
 - not enough hits on the primary
 - bad origin for the secondaries

Legend for interaction lengths:

- Beam Pipe
- Pixel
- Inner Silicon
- Outer Silicon
- Common
- Outside



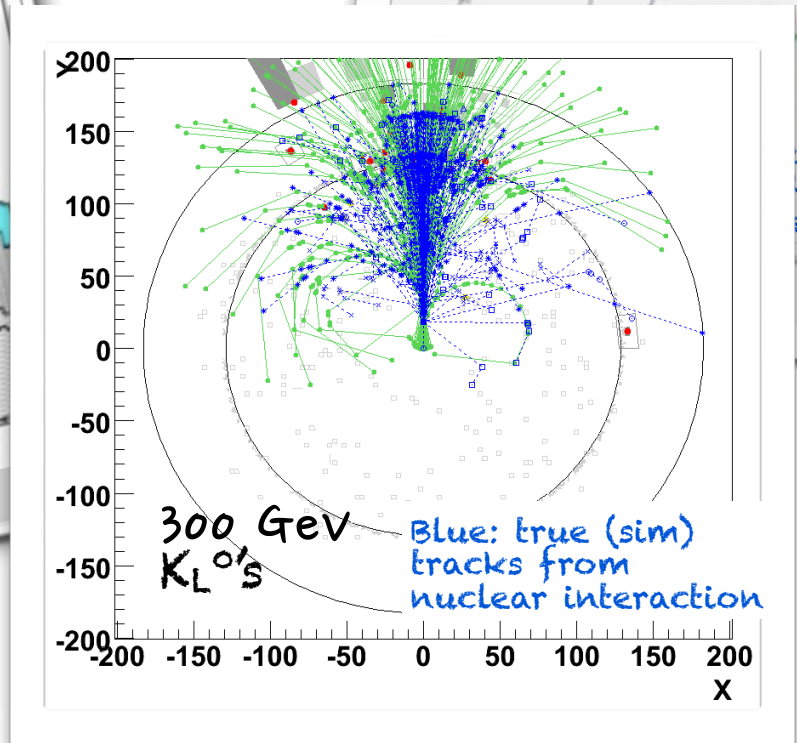
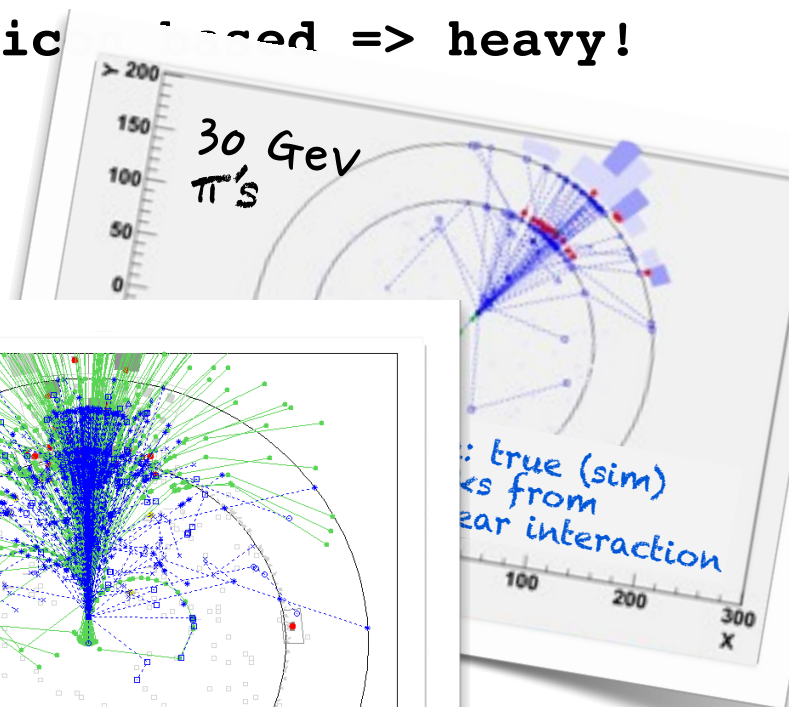


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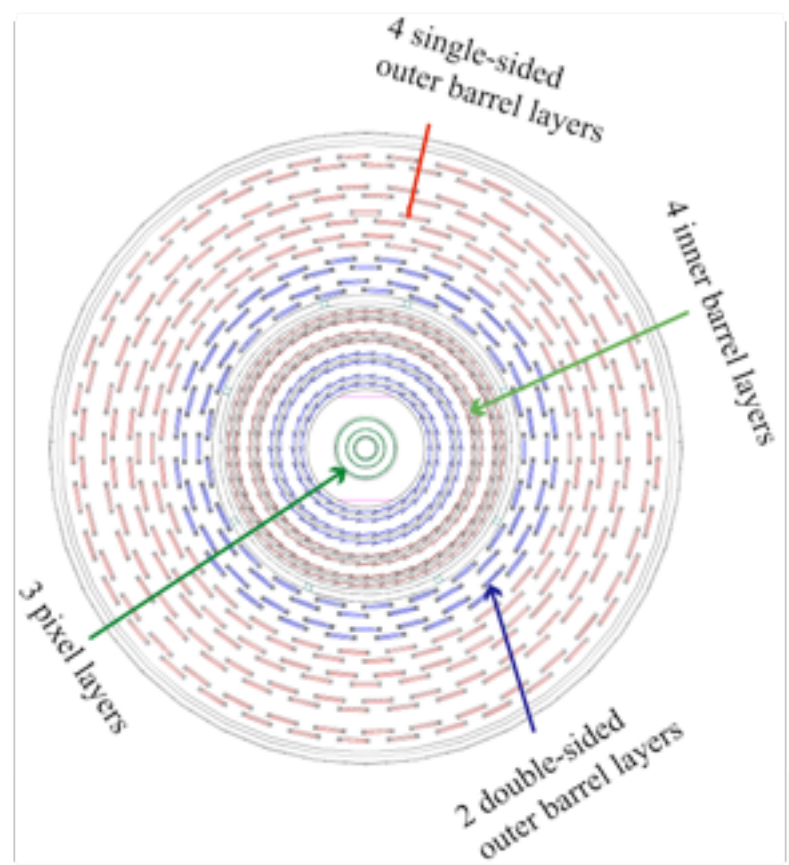


- **Fake tracks come from wrong combinations of hit associations**
 - So, reduce the number of hits fed to the track finder



Nuclear Interactions in Tracker

- Fake tracks come from wrong combinations of hit associations
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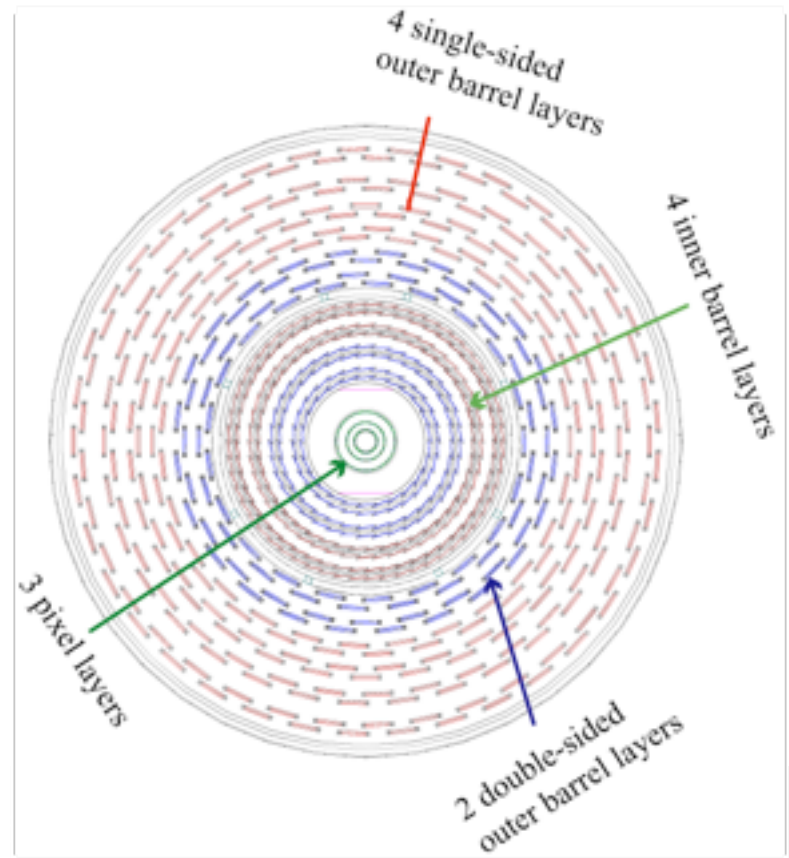


Nuclear Interactions in Tracker

- Fake tracks come from wrong combinations of hit associations
- So, reduce the number of hits fed to the track finder

Start from a very pure track seeding

- Example: 3 pixel hits, very tight origin constraint, $p_T > 0.9 \text{ GeV}$
- 75% efficiency, less than 1% fakes
- Reconstruct corresponding tracks (≥ 3 hits) & "remove" the used hits
- 40% of the hits are removed





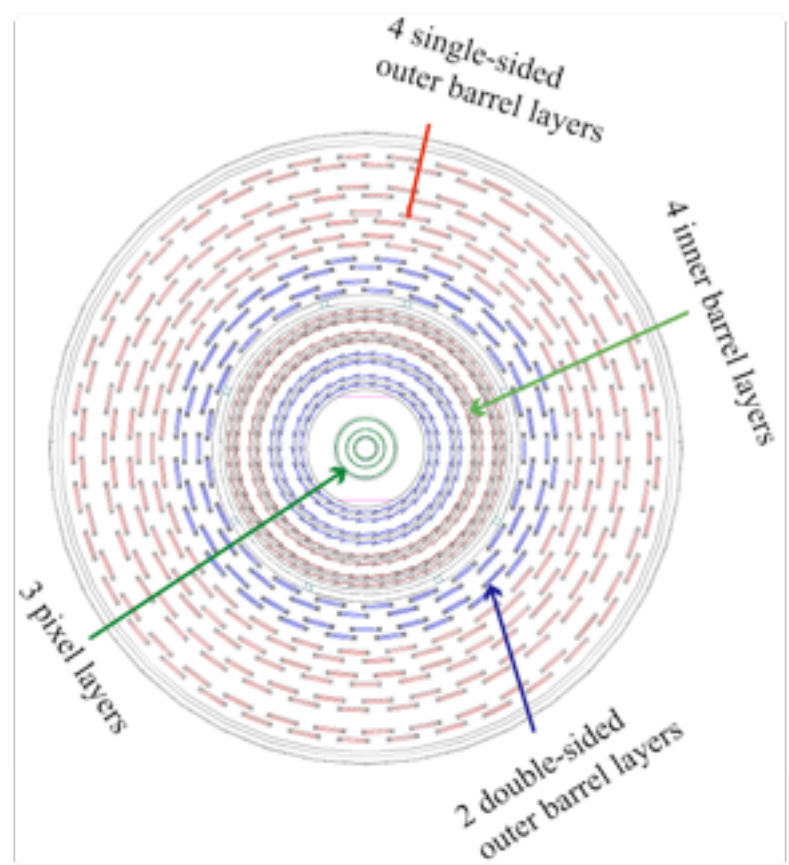
Nuclear Interactions in Tracker

- Fake tracks come from wrong combinations of hit associations
- So, reduce the number of hits fed to the track finder

Start from a very pure track seeding
 2 pixel hits, very tight origin

Next, try a looser seeding on
 60% remaining hits

- Example: 2 pixel hits, very tight origin constraint, $p_T > 0.9$ GeV
- adds 15% efficiency, still less than 1% fakes
- Combinatorial possibilities much less
- Reconstruct the corresponding tracks (≥ 3 hits) & remove the used hits
- 10% of the hits are removed in this 2nd iteration





Nuclear Interactions in Tracker

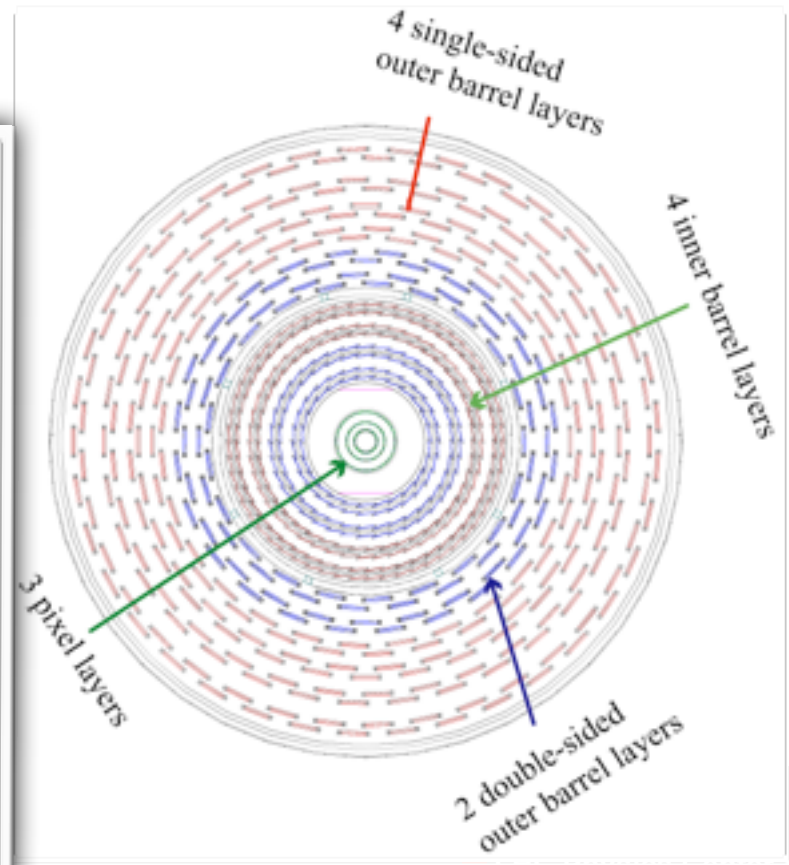
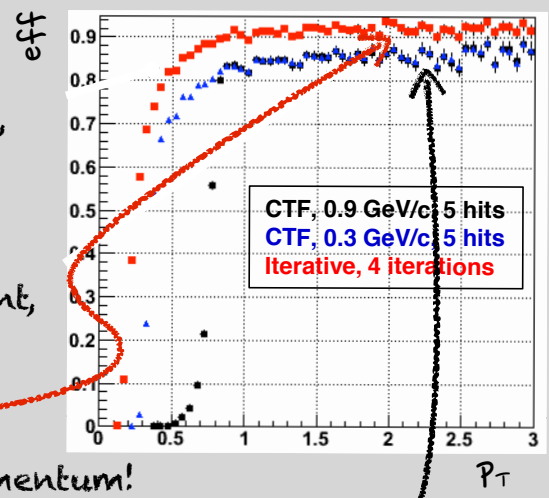
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60% remaining hits
3 pixel hits, very tight origin

And so on...with more iterations

- 3rd iteration:
 - 3 pixel hits, tight origin constraint, $p_T > 0.2 \text{ GeV}$
- 4th iteration:
 - 2 pixel hits, looser origin constraint, $p_T > 0.3 \text{ GeV}$
 - 93% efficiency, 1-2% fake rate
 - down to very low momentum!
- Recall original situation
 - 85% efficiency, 20% fake rate





Nuclear Interactions

- Fake tracks come from wrong combinations of hits

• So, reduce number of hits for seeding

Start from a very pure track seeding

3 pixel hits, very tight origin constraint

Next, try a looser seeding on 60% remaining hits

2 pixel hits, very tight origin

And so on...with more iterations

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- 3 pixel hits, tight origin constraint, $p_T > 0.2 \text{ GeV}$

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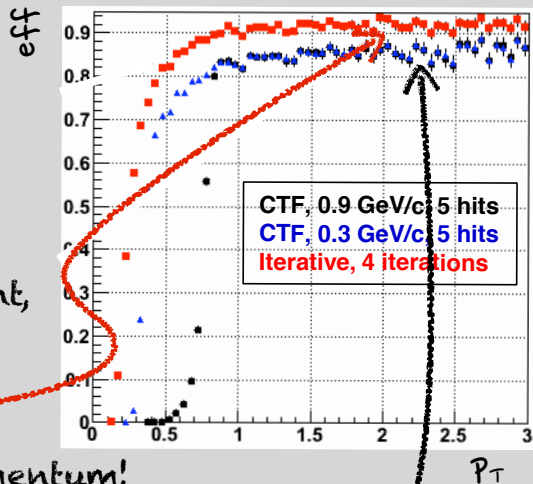
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And even try to catch secondary tracks

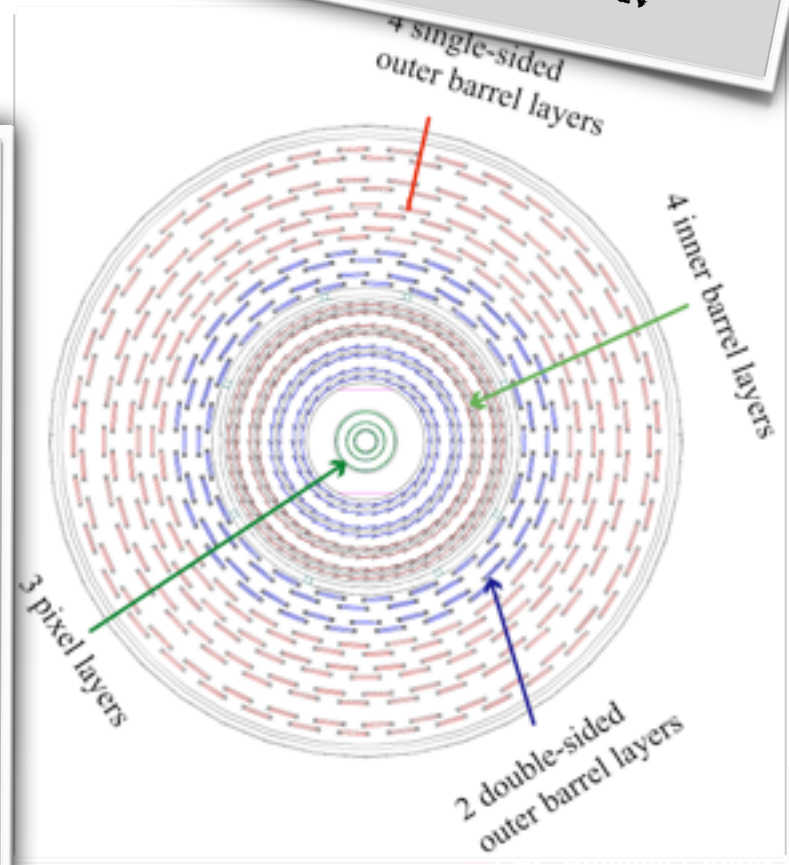
- Interactions, conversions, decays,...

• 5th iteration:

- Inner tracker seeding, loose origin constraint, $p_T > 0.5 \text{ GeV}$

• 6th iteration:

- Outer tracking seeding, very loose origin constraint, $p_T > 0.8 \text{ GeV}$





Nuclear Interactions

- Fake tracks come from wrong correlations
- So, reduce the number of hits for each track

And even try to catch secondary tracks

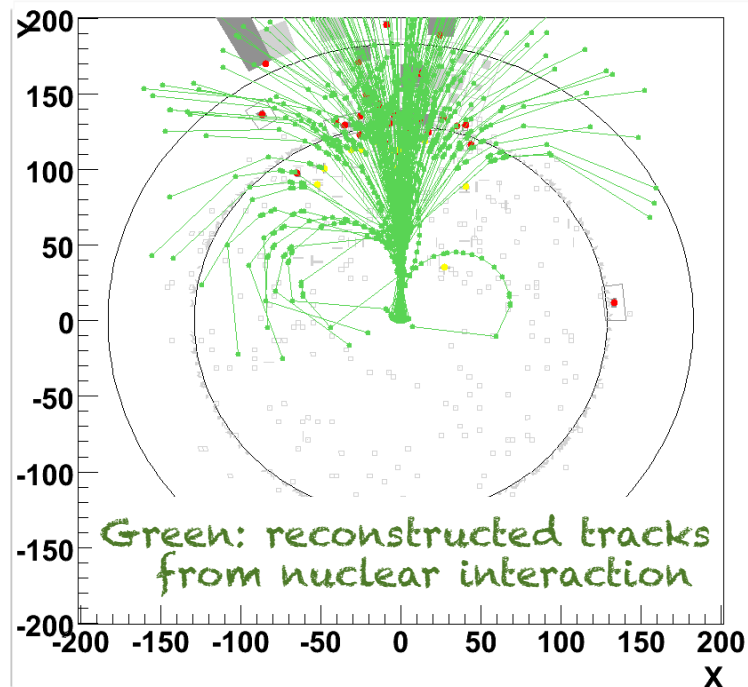
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 - Outer tracking, seeding, very loose origin constraint, $p_T > 0.8 \text{ GeV}$

That's a lot of work!
 what's the reward?



Nuclear Interactions

- Fake tracks come from wrong combinations
- So, reduce the number of hits for



That's a lot of work!
 What's the reward?

And even try to catch secondary tracks

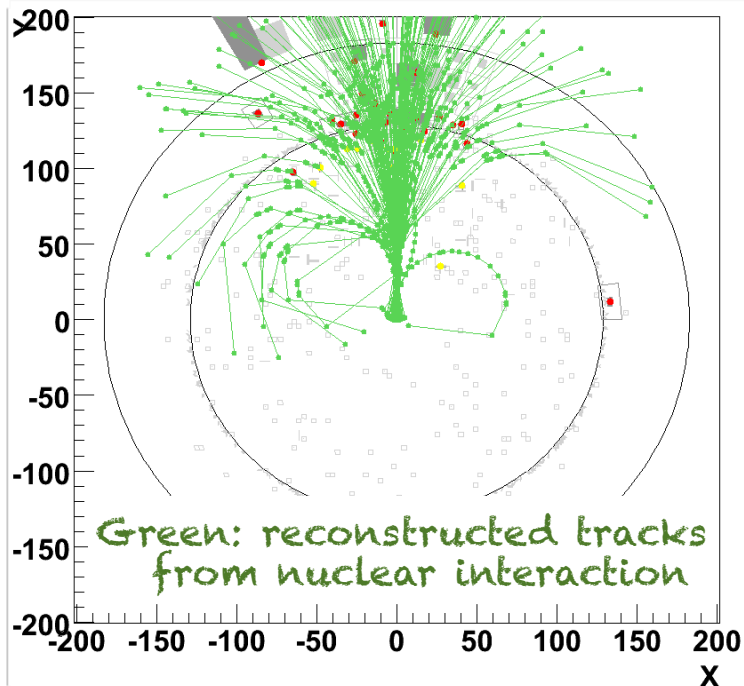
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Nuclear Interactions

- Fake tracks come from wrong combinations
- So, reduce the number of hits for



That's a lot of work!
 What's the reward?

Fermilab
 And even try to catch secondary tracks

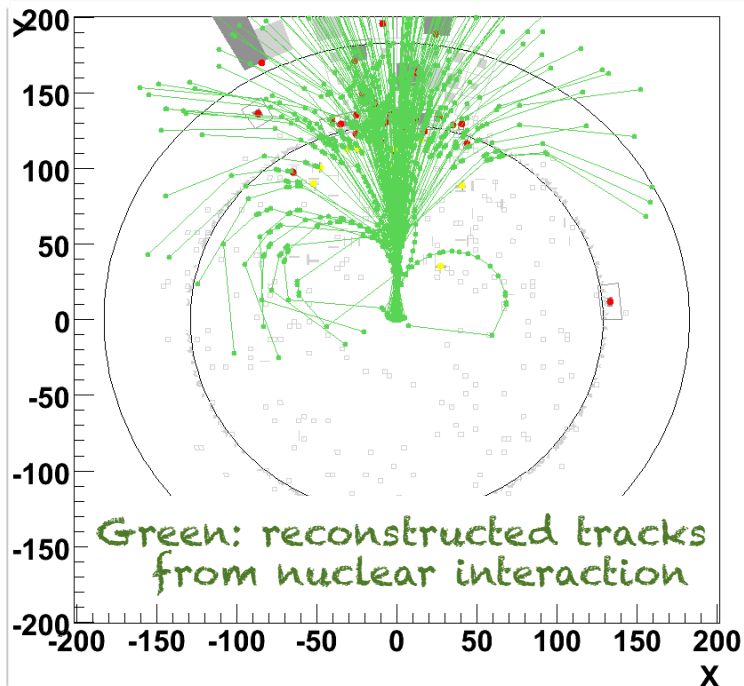
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Many real tracks reconstructed with inner vs outer seeding



Nuclear Interactions

- Fake tracks come from wrong combinations
- So, reduce the number of hits for



That's a lot of work!
What's the reward?



And even try to catch secondary tracks

- Interactions, conversions, decays,...
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Next problem: avoid double counting primary vs secondaries

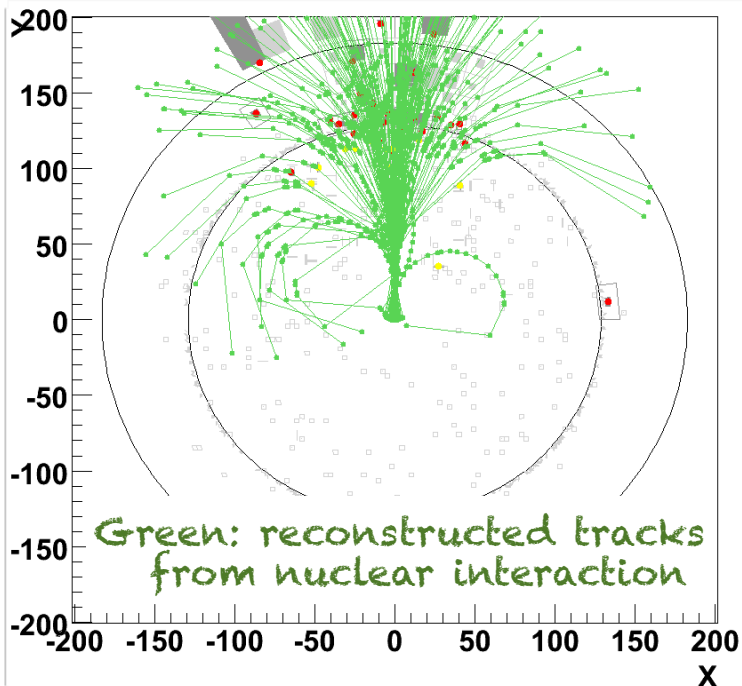
- Create a "link by vertex" between primaries and secondaries
 - choose the best energy determination
 - primary track usually has more hits (5-6)
 - secondaries usually have less hits

Min



Nuclear Interactions

- Fake tracks come from wrong combinations
- So, reduce the number of hits for



That's a lot of work!
What's the reward?

Fermilab
And even try to catch secondary tracks

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M:

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Philosophy: reconstruct & identify as many particles as possible!



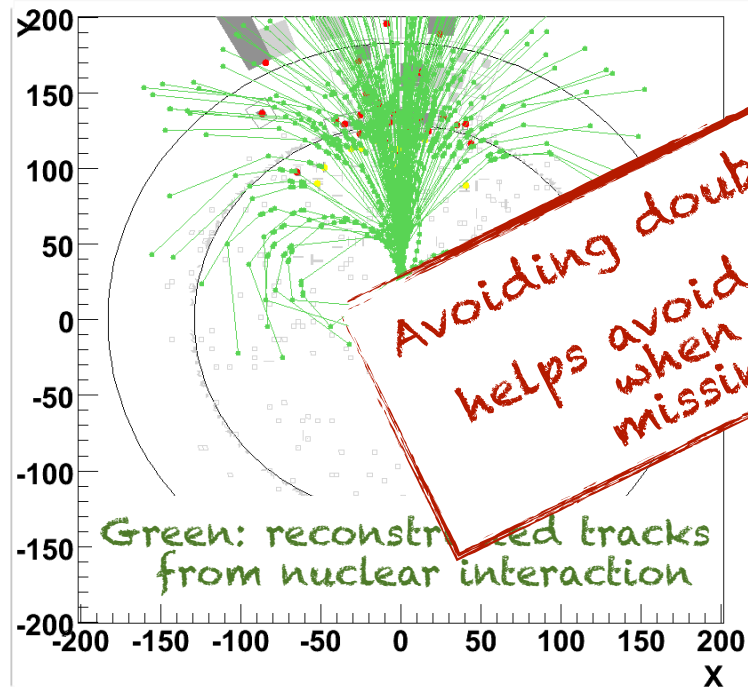
Nuclear Interactions



- Fake tracks come from wrong correlations
- So, reduce the number of hits from wrong associations

And even try to catch secondary tracks

- Interactions, conversions, decays, ...
- 5th iteration:
 - Inner tracker seeding, loose origin
 - 6th iteration:
 - PT > 0.5 GeV



Avoiding double counting energy => helps avoid non-gaussian tails when reconstructing missing energy (later)

Problem: avoid double counting primary vs secondaries

- Create a "link by vertex" between primaries and secondaries
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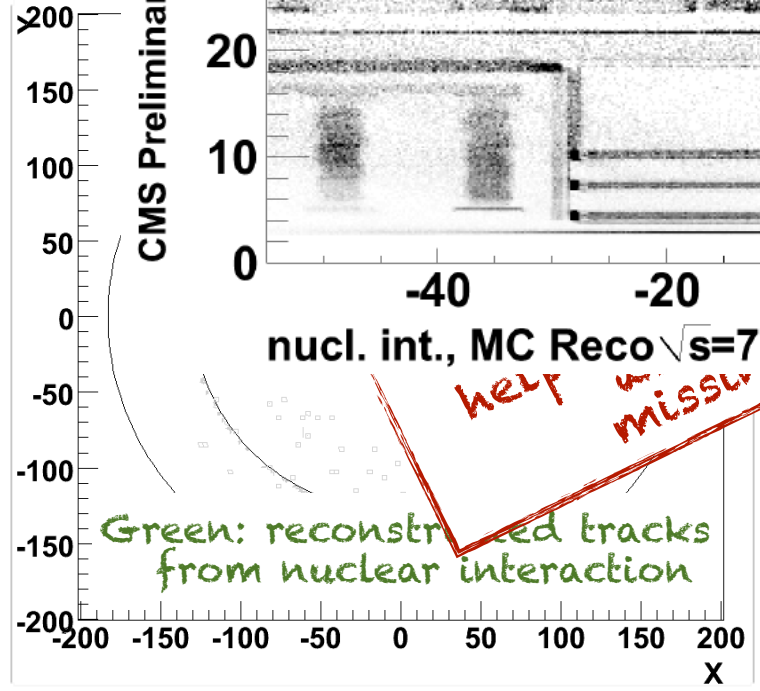
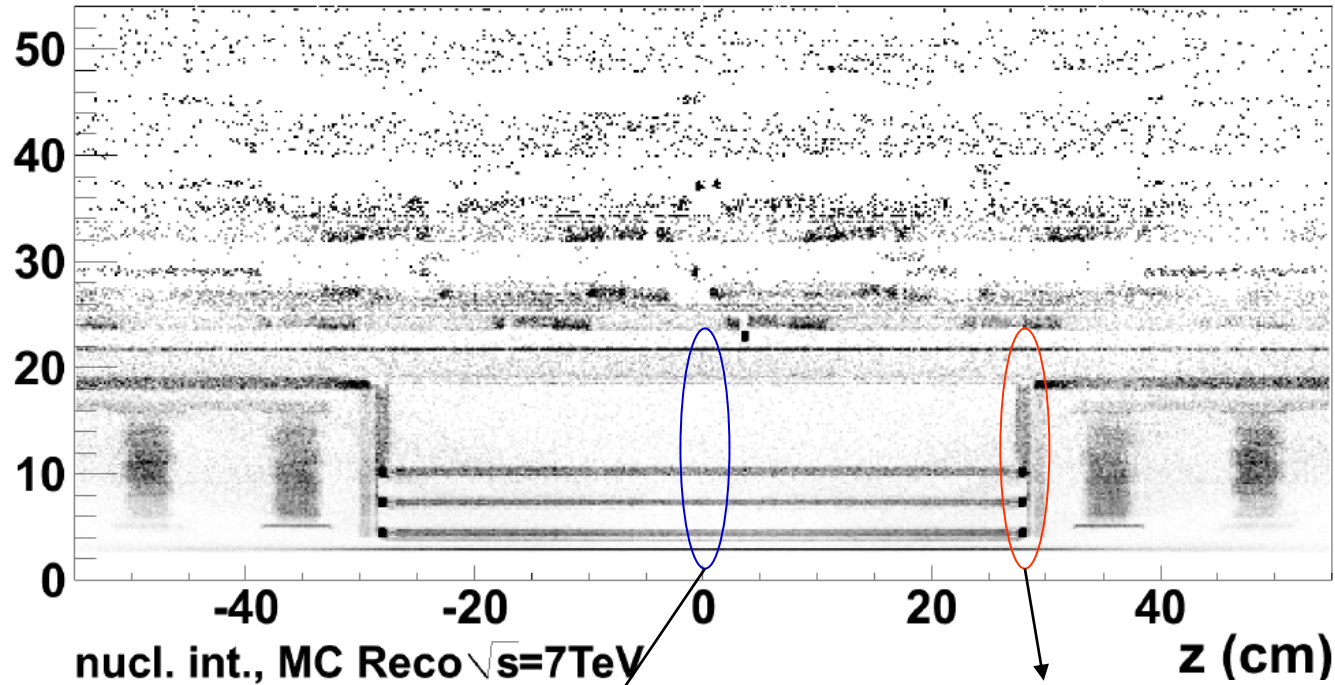
Philosophy: reconstruct & identify as many particles as possible!



Nu

- Fake associations
- Secondary decays, ...
- Trace back to primary loose GeV

CMS Preliminary 2010 R (cm)



- M:
- primaries and secondaries
 - choose the best energy determination
 - primary track usually has more hits (5-6)
 - secondaries usually have less hits

And even ...

Fermilab

secondary

decays, ...

Trace origin

primary loose GeV

double boundaries

That's a lot of work!
 What's the reward?

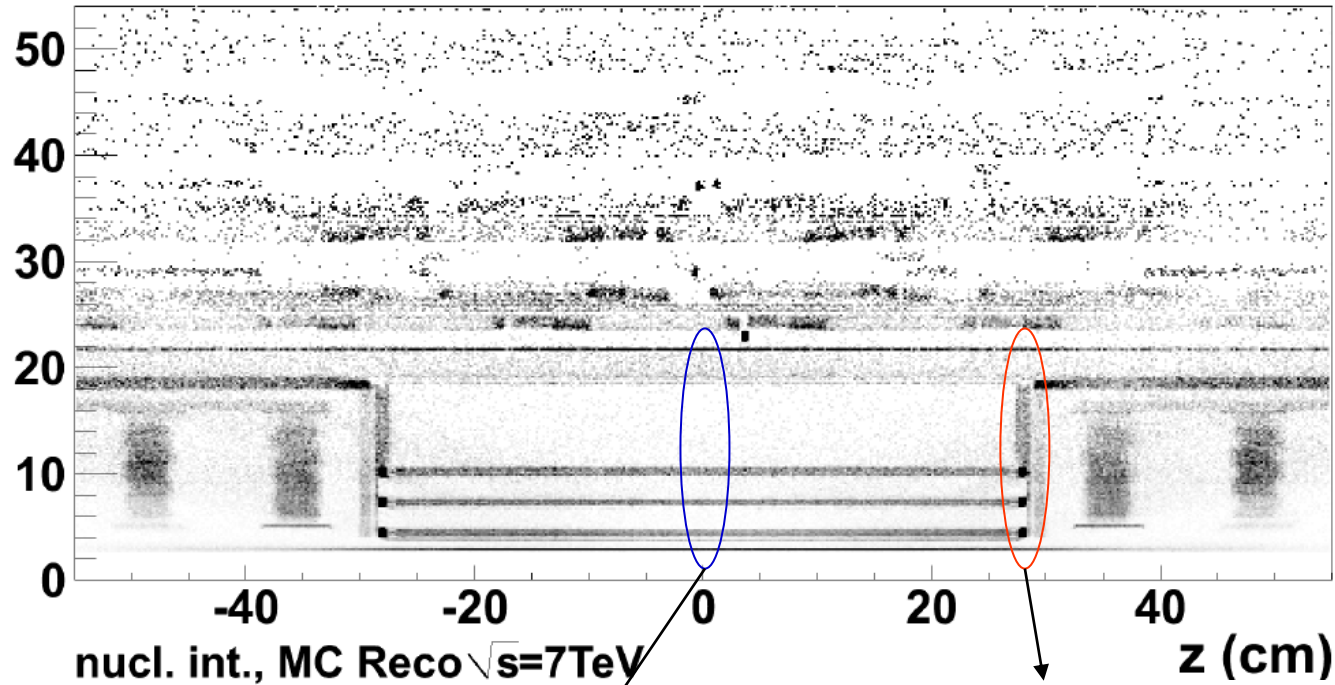
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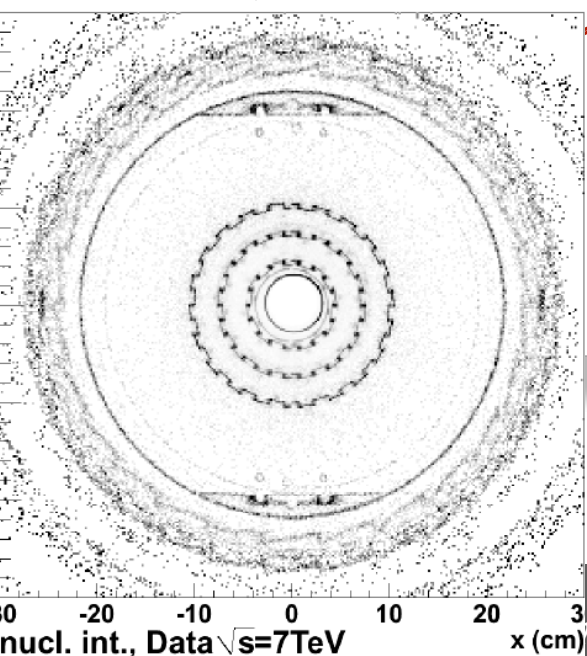
CMS Preliminary 2010 R (cm)



nucl. int., MC Reco $\sqrt{s}=7\text{TeV}$

Y 200
150
100
50
0
-50
-100
-150
-200

CMS Preliminary 2010



nucl. int., Data $\sqrt{s}=7\text{TeV}$

x (cm)

Mini

- primaries and secondaries
- choose the best energy determination
 - primary track usually has more hits (5-6)
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Philosophy: reconstruct & identify as many particles as possible!

secondary

decays,...

lose origin

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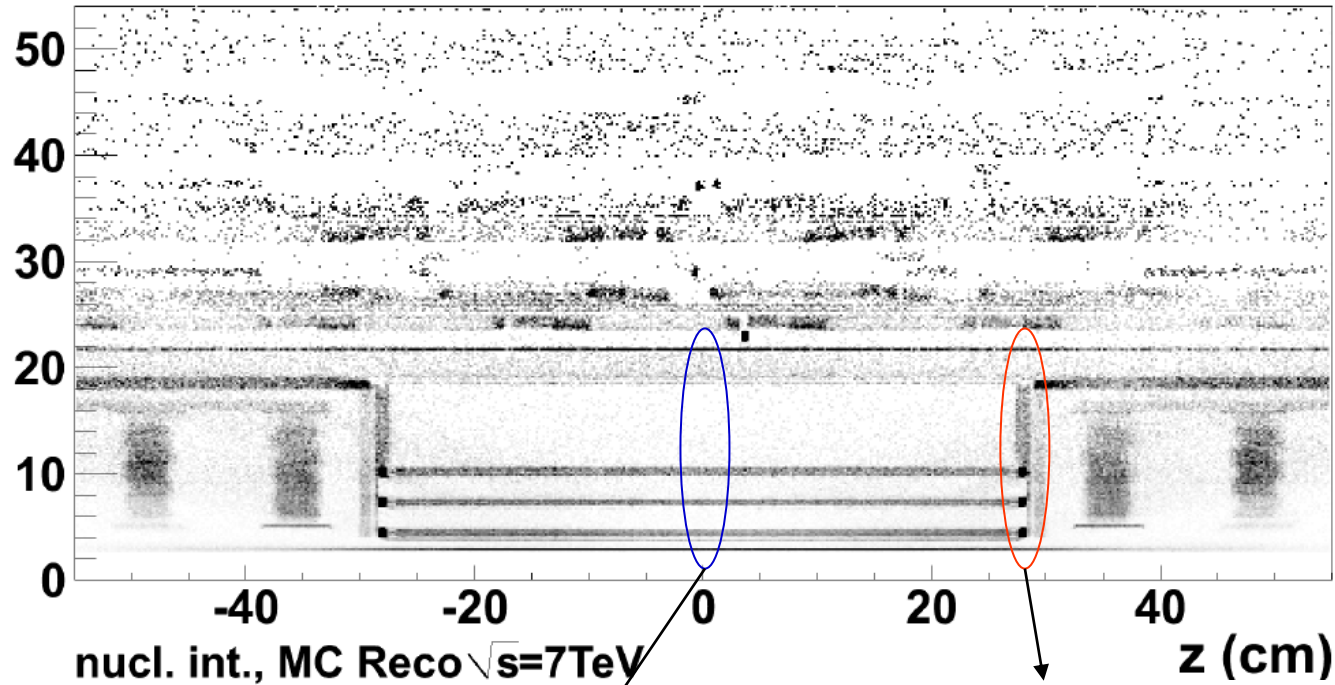
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Nu

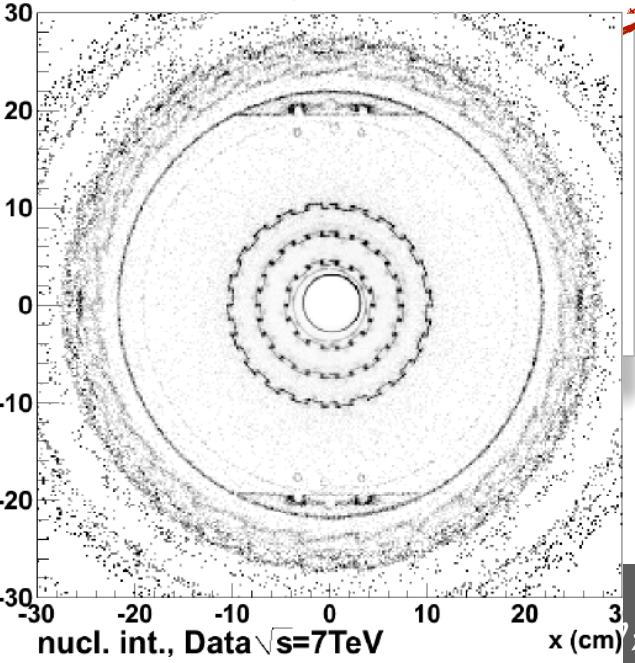
- Fake
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- Sc

CMS Preliminary 2010 R (cm)



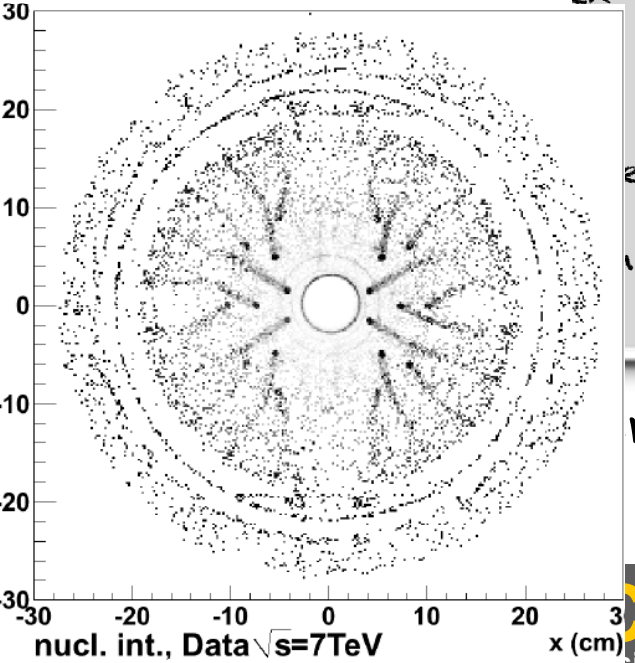
Y 200
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CMS Preliminary 2010



Phil
as v

CMS Preliminary 2010



And even

Fermilab

secondary

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lose origin

ry loose GeV

uble boundaries



ID Individual Hadrons



ID Individual Hadrons

Particle	m [MeV]	Quarks	Main decay	Lifetime	$c\tau$ [cm]
π^\pm	140	$u\bar{d}$	$\mu\nu_\mu$	2.6×10^{-8} s	780
K^\pm	494	$u\bar{s}$	$\mu\nu_\mu, \pi^\pm\pi^0$	1.2×10^{-8} s	370
K_S^0	498	$d\bar{s}$	$\pi\pi$	0.9×10^{-10} s	2.7
K_L^0	498	$d\bar{s}$	$\pi\pi\pi, \pi l\nu$	5×10^{-8} s	1550
p	938	uud	stable	$> 10^{25}$ years	∞
n	940	udd	$p e \nu_e$	890 s	2.7×10^{13}
Λ	1116	uds	$p\pi$	2.6×10^{-10} s	7.9



ID Individual Hadrons

Why?

- Instead of making do with jet reconstruction
 - some physics requires the ID of individual hadrons
- Most are unstable, and decay into a few long-lived particles

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K_L^0	498	$d\bar{s}$	$\pi\pi\pi, \pi l\nu$	5×10^{-8} s	1550
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π^\pm	140	$u\bar{d}$	$\mu\nu_\mu$	2.6×10^{-8} s	780
K^\pm	494	$u\bar{s}$	$\mu\nu_\mu, \pi^\pm\pi^0$	1.2×10^{-8} s	370
K_S^0	498	$d\bar{s}$	$\pi\pi$	0.9×10^{-10} s	2.7
K_L^0	498	$d\bar{s}$	$\pi\pi\pi, \pi l\nu$	5×10^{-8} s	1550
p	938	uud	stable	$> 10^{25}$ years	∞
n	940	udd	$p e \nu_e$	890 s	2.7×10^{13}
Λ	1116	uds	$p\pi$	2.6×10^{-10} s	7.9

How?

- Since the interactions of charged hadrons are similar, the most direct way to distinguish them is to determine their (rest) mass
- Their momentum is measured by the tracking system, so this is equivalent to determining their velocity, since $p = \gamma m v$, so $m = p/\gamma v = p/\gamma\beta c$



ID Individual Hadrons

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Why?

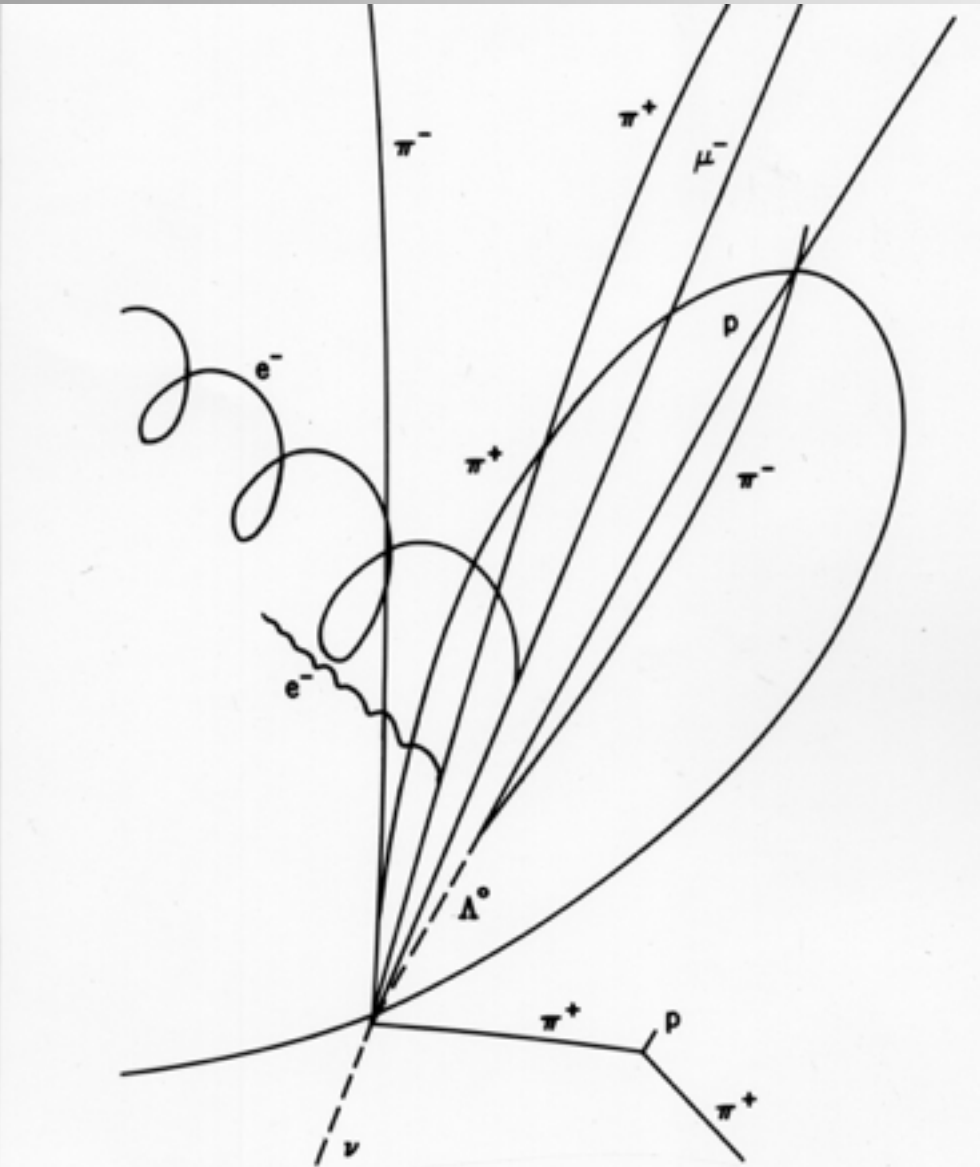
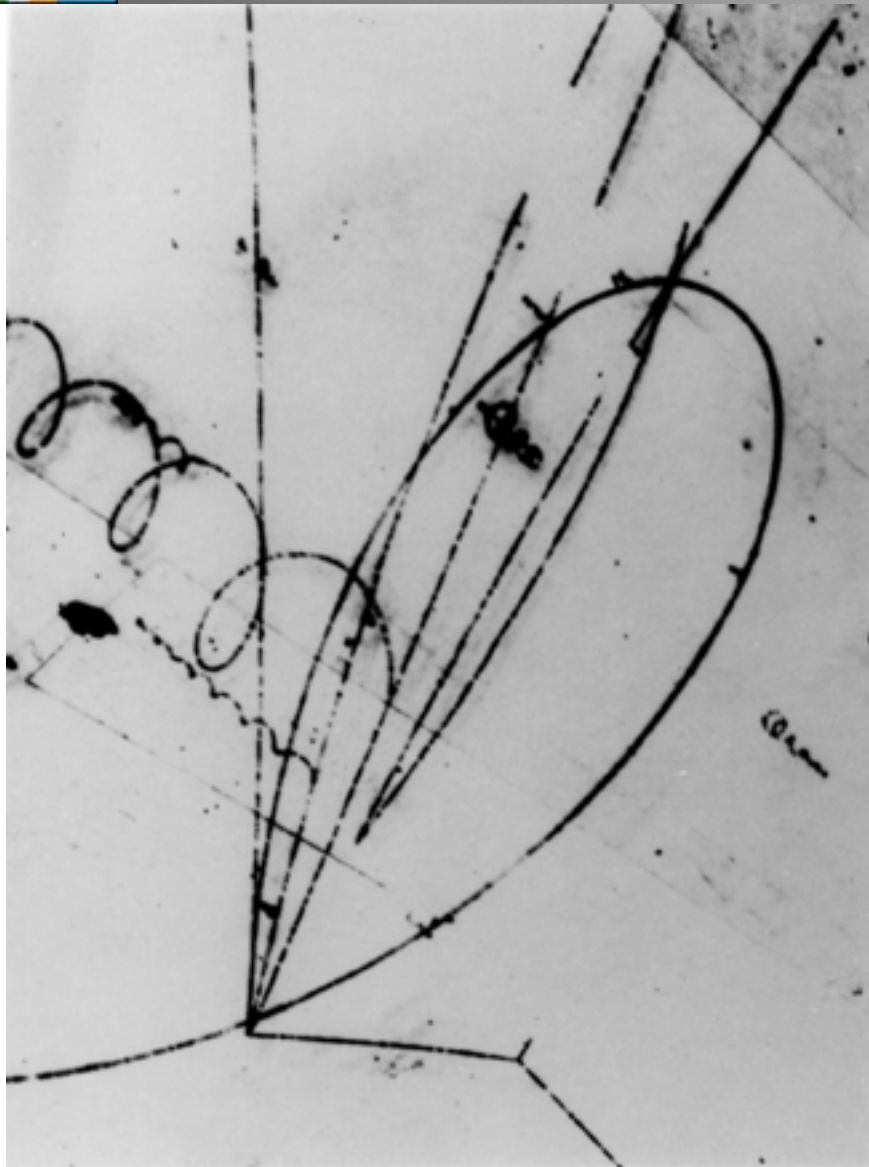
- Instead of making do with jet reconstruction
 - some physics requires the ID of individual hadrons
- Most are unstable, and decay into a few long-lived particles

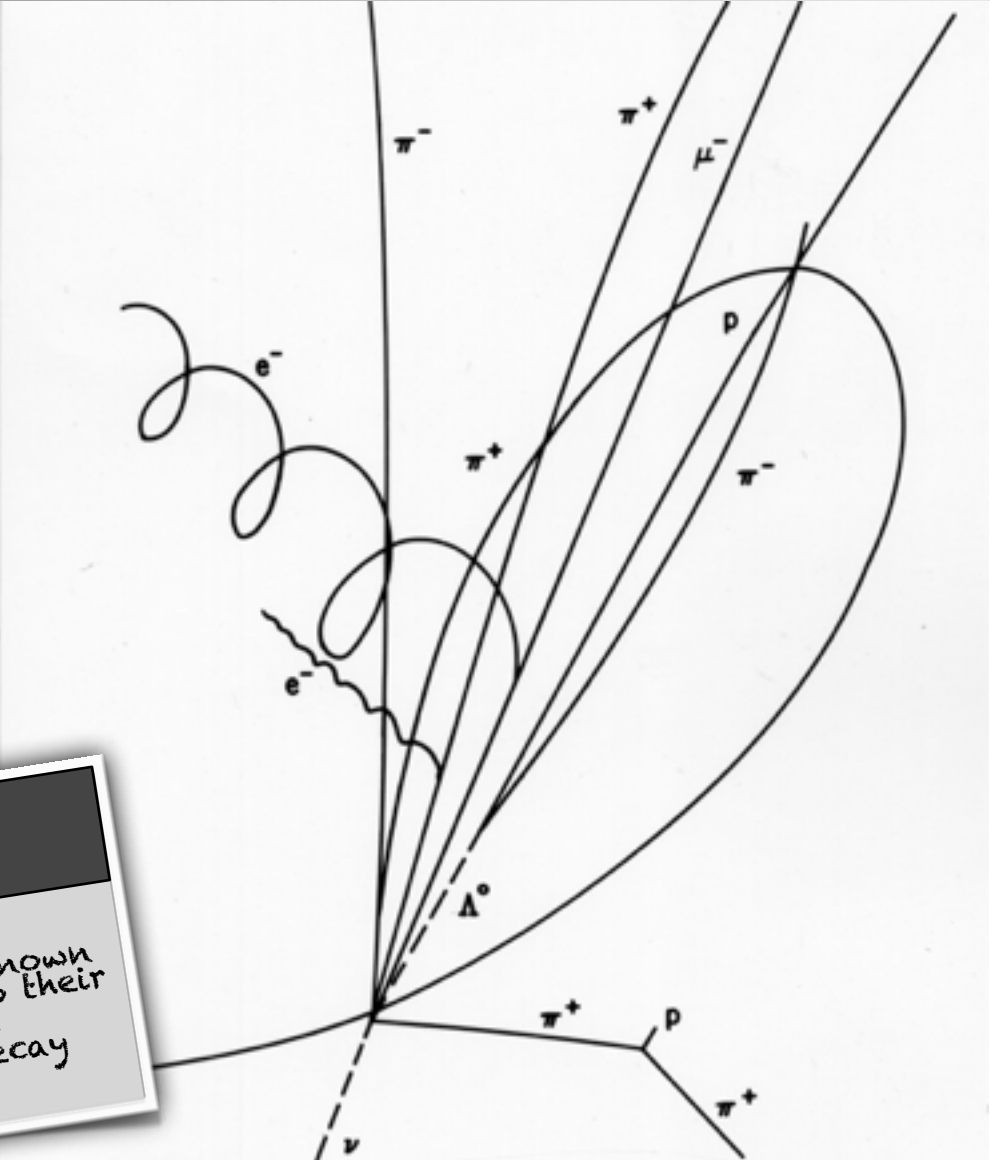
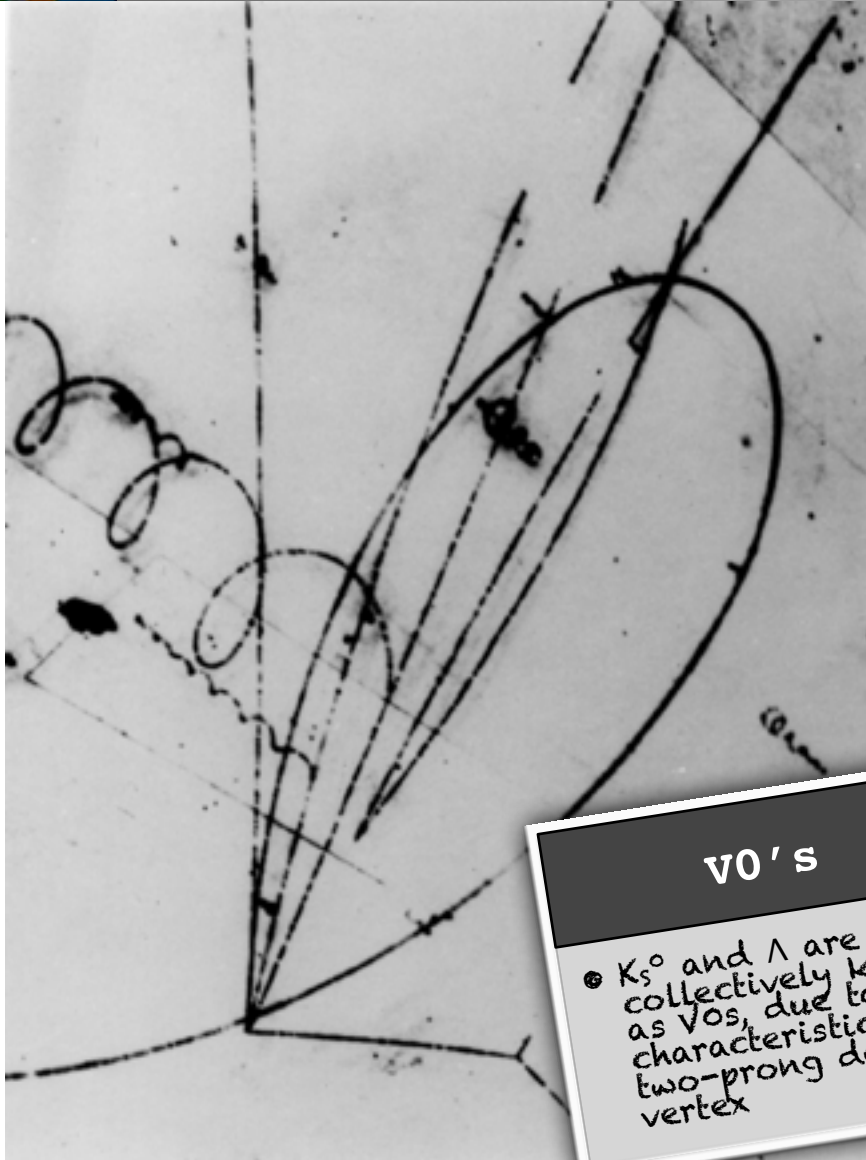
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Four main processes

- Most direct is to measure the Time Of Flight (TOF) of the particles over a fixed distance
- Alternatively one can look at the detail of their interaction with matter. The main source of energy loss is via Ionization (dE/dx)
- If the velocity of the particle changes compared to the local speed of light it will radiate photons, detected as Transition radiation
- If a particle travels at greater than the local speed of light, it will radiate Cherenkov radiation

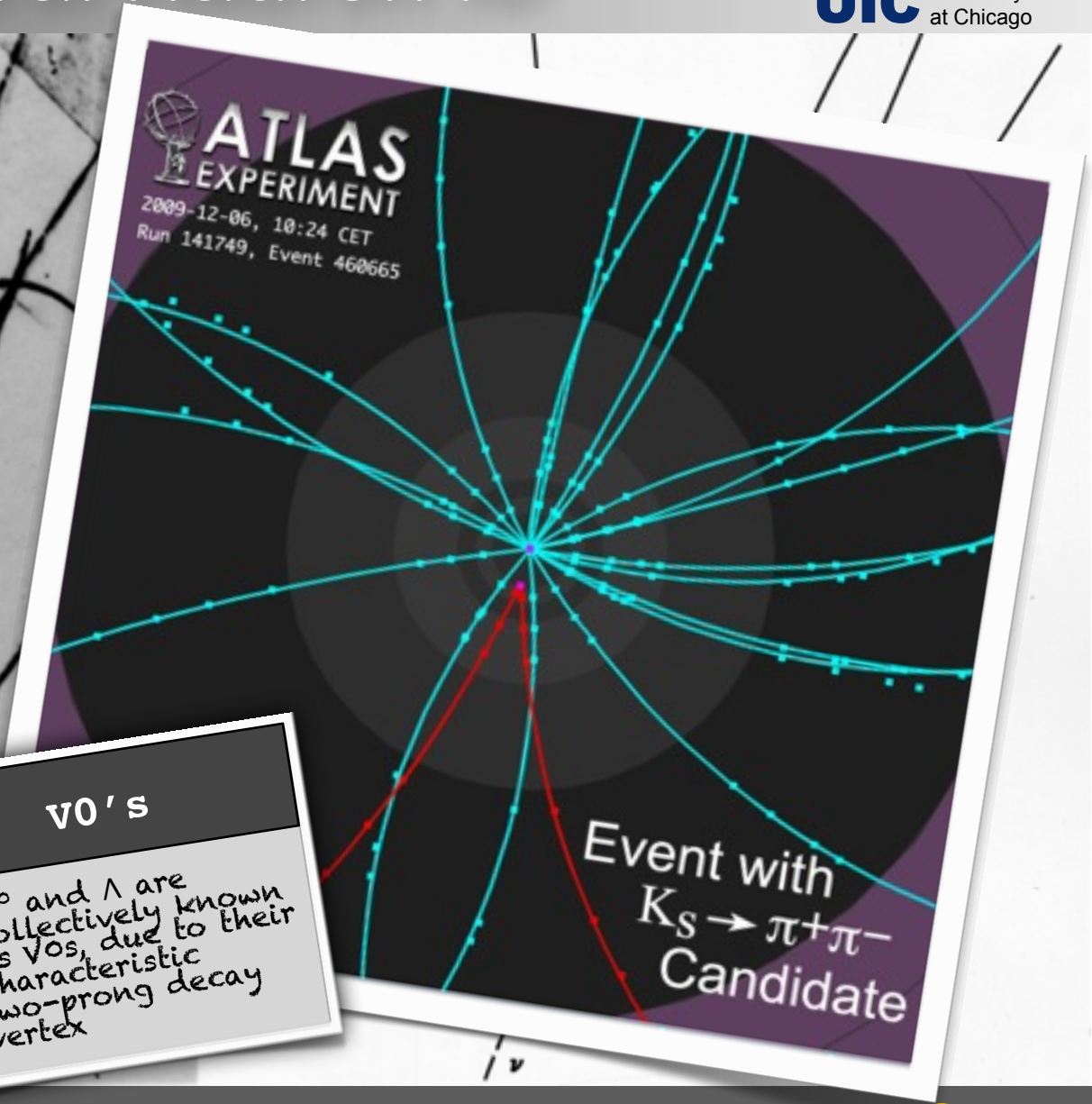
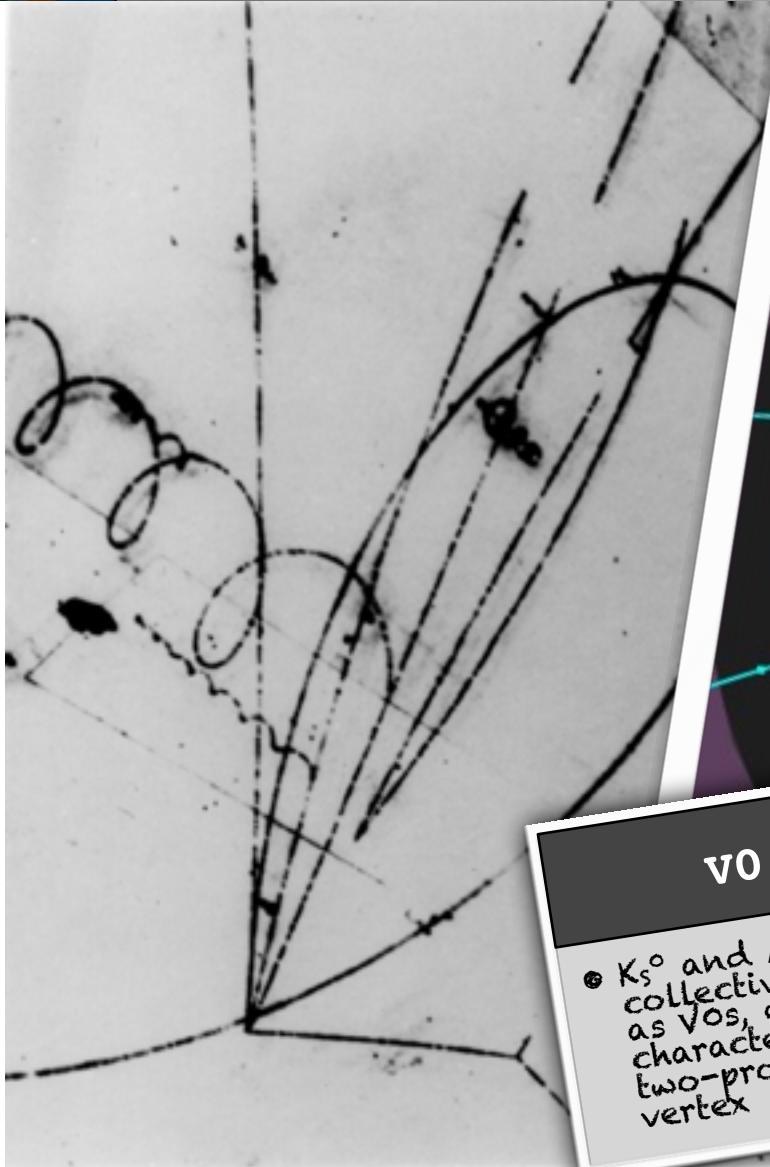




V0'S

- K_S^0 and Λ are collectively known as V0's, due to their characteristic two-prong decay vertex

Neutral Hadron ID



V0'S

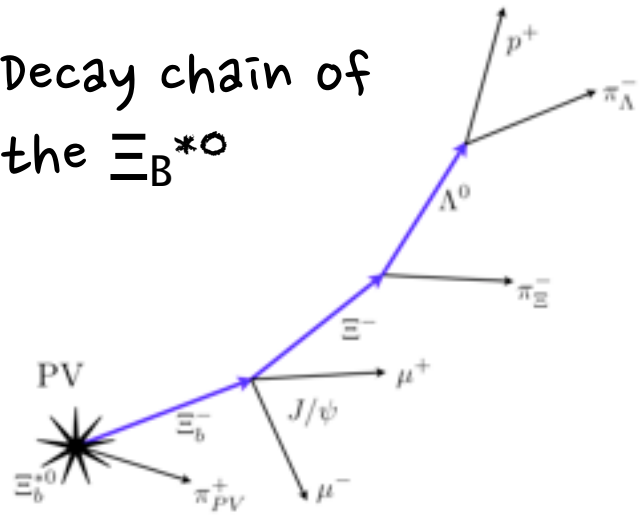
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Event with
 $K_S \rightarrow \pi^+ \pi^-$
Candidate



Extreme Hadron ID!

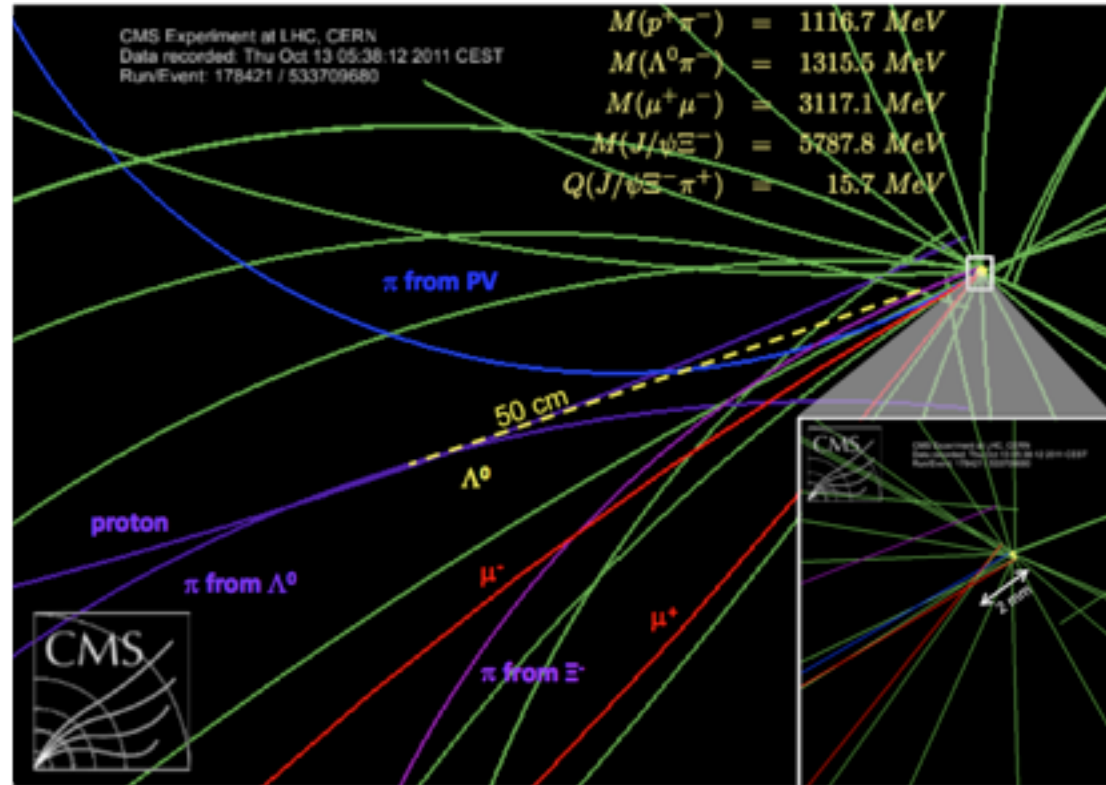
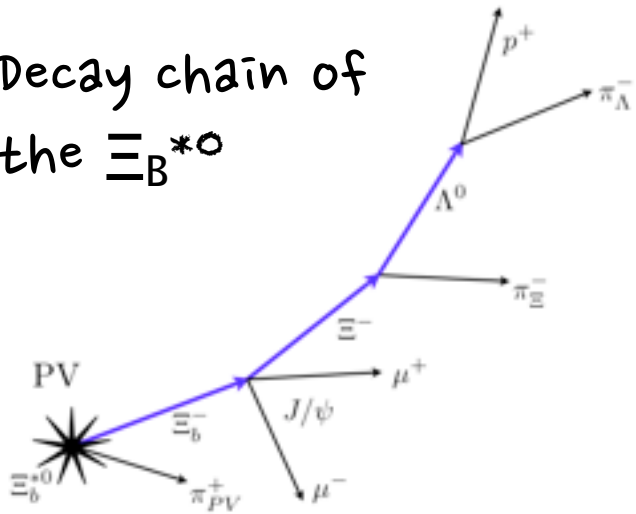
Decay chain of
the Ξ_B^{*0}





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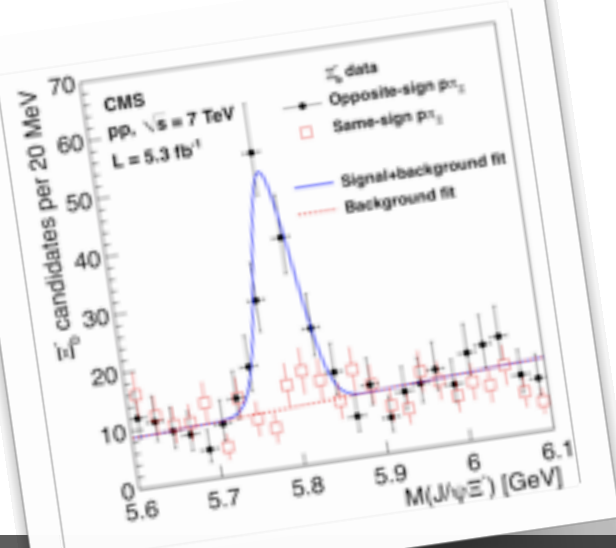
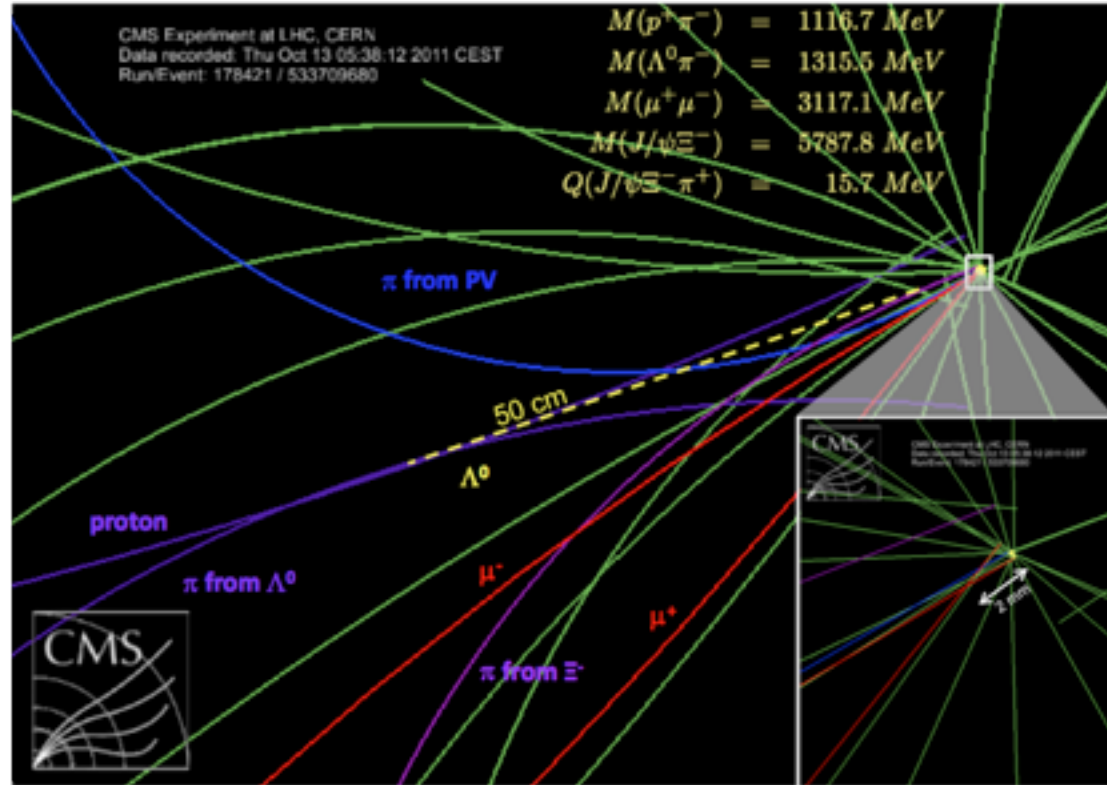
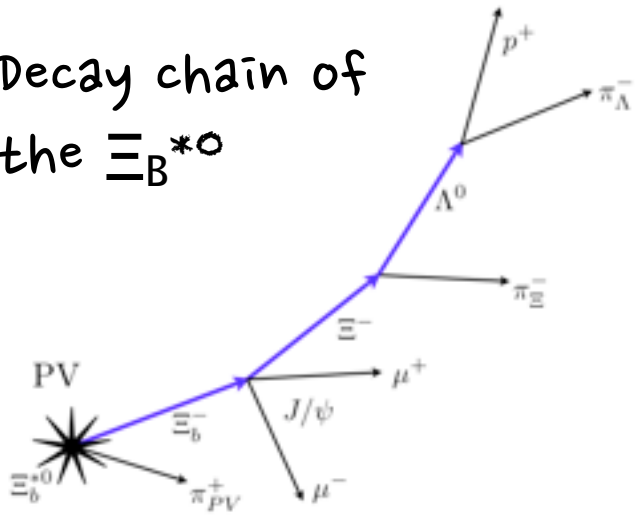
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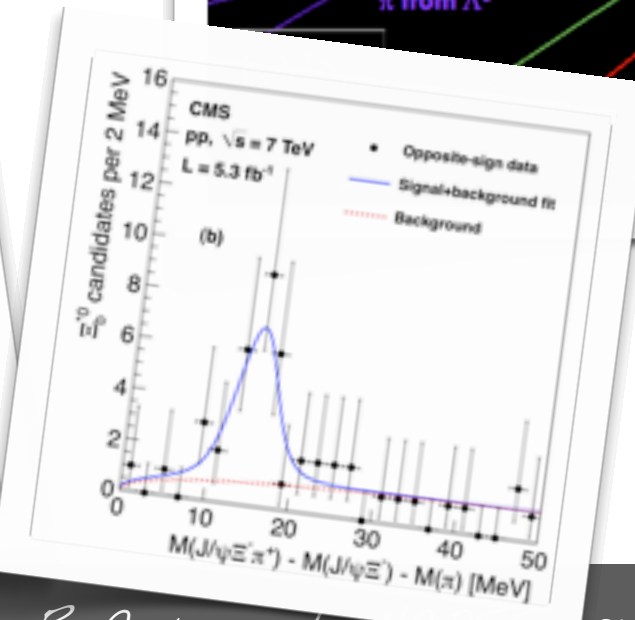
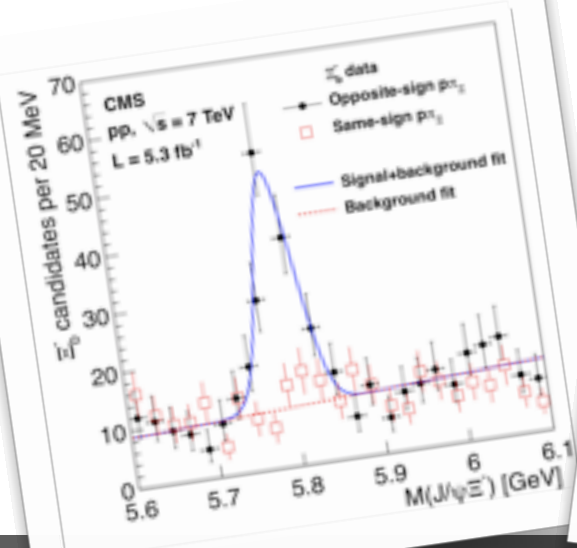
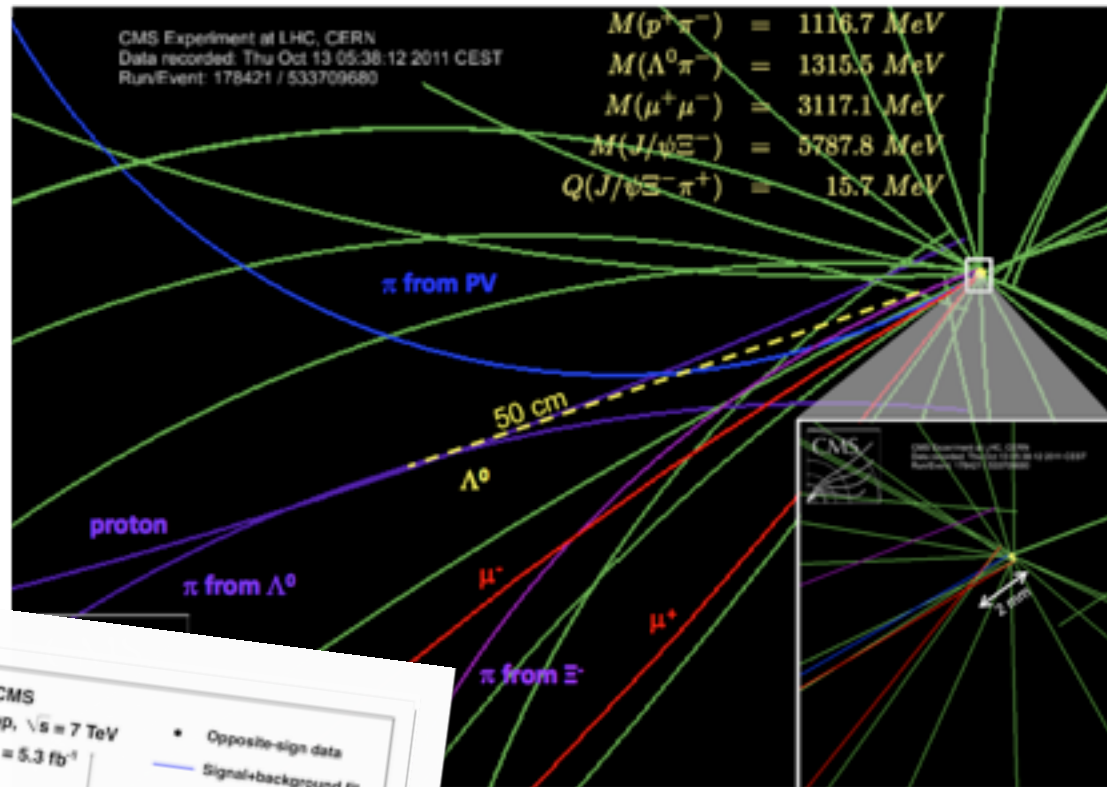
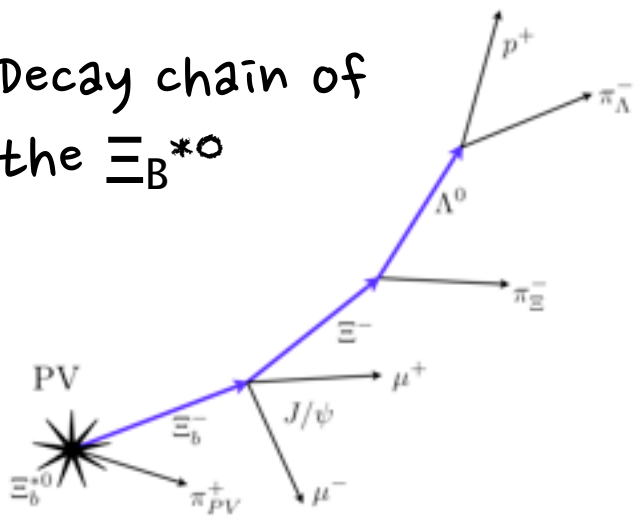
S... 50
Date: 14.08.2012

R. Cavanaugh, HCPSS 2012



Extreme Hadron ID!

Decay chain of the Ξ_B^{*0}



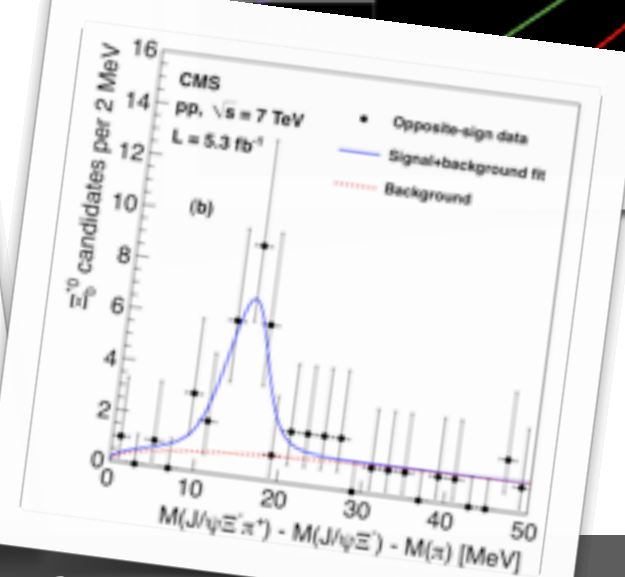
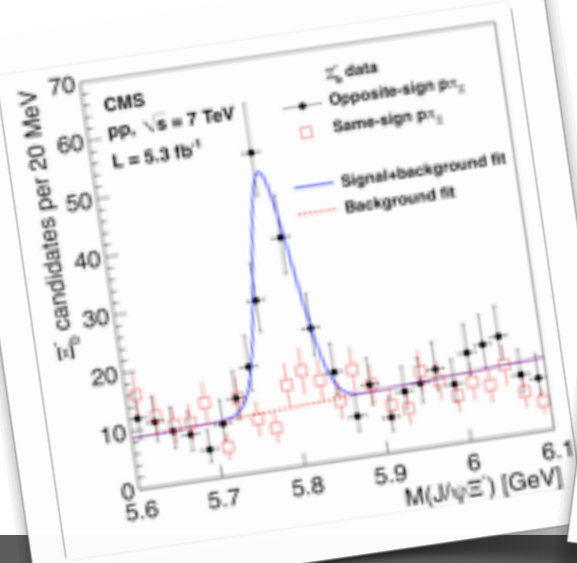
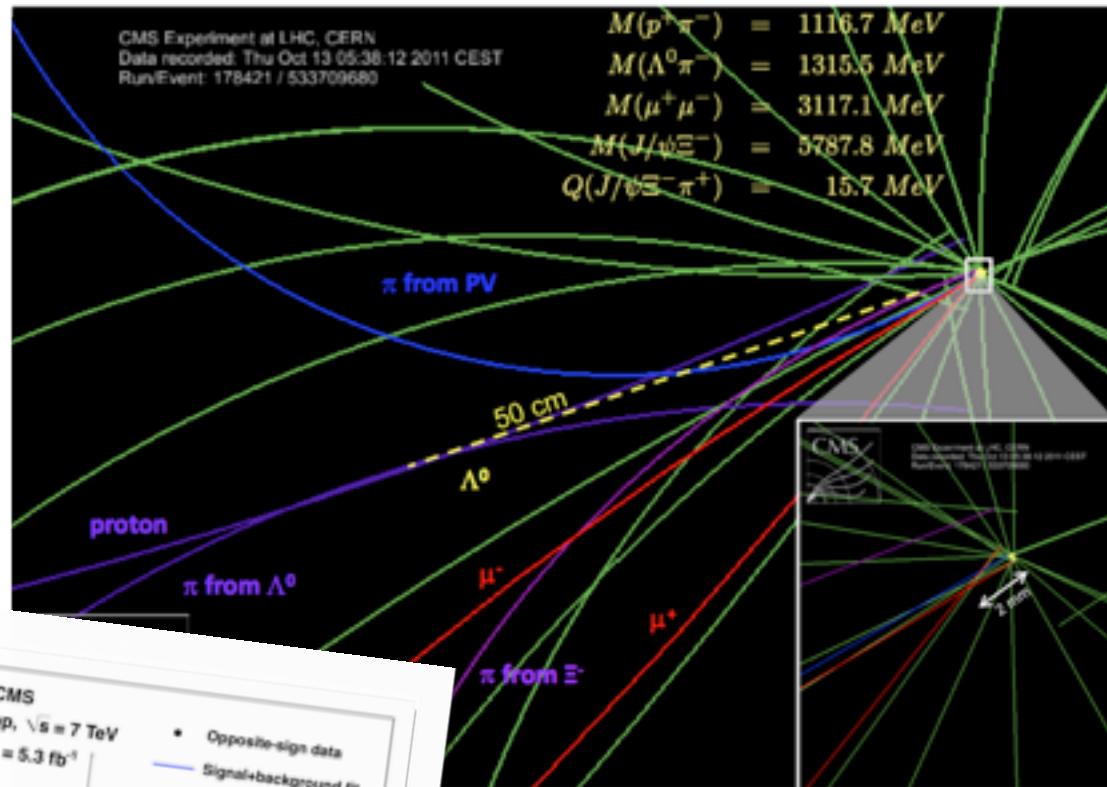
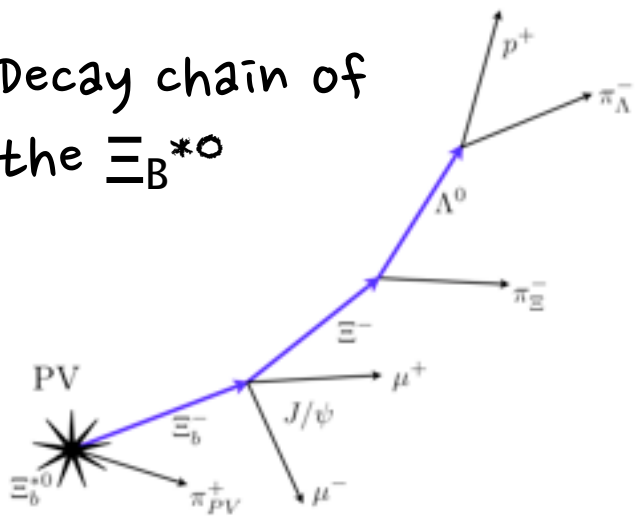
Date: 14.08.2012

R. Cavanaugh, HEP-2012-012



Extreme Hadron ID!

Decay chain of the Ξ_B^{*0}



New baryon discovered!

Date: 14.08.2012

R. Cavanaugh, HEP 2012



Charged Hadron ID



Charged Hadron ID

Motivation

- Charged hadrons (π , K , p) are all effectively stable, and have similar interactions
 - track + hadronic shower
- However, identifying them can be crucial, in particular for the study of hadronic decays



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Example: $\varphi \rightarrow K^+ K^-$

- Make all two-track combinations in an event and calculate their invariant mass
 - large combinatoric background (most tracks are pions, from other sources)
- By identifying the two tracks as kaons, signal to background ratio is much improved



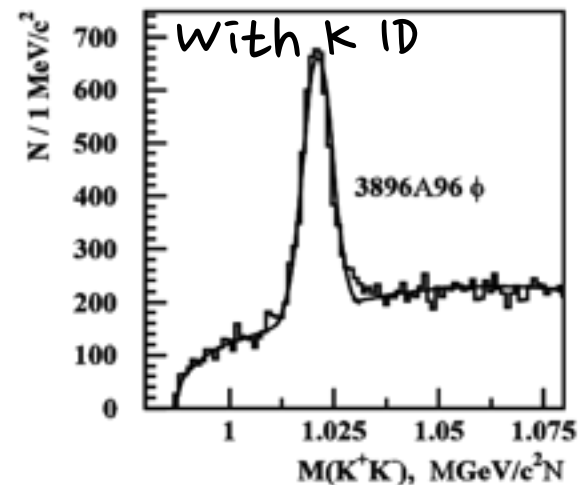
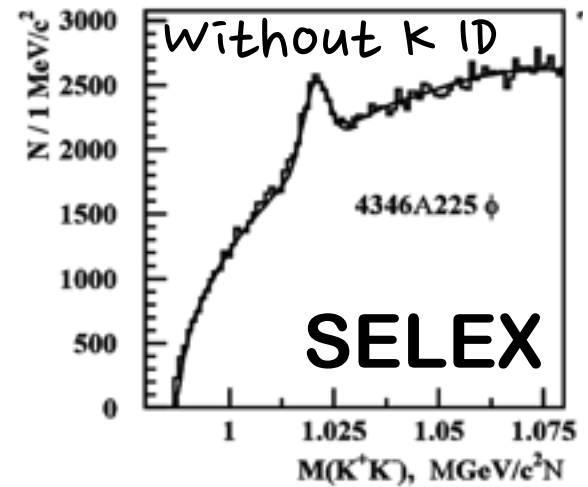
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Time Of Flight



Time Of Flight

Simple Concept

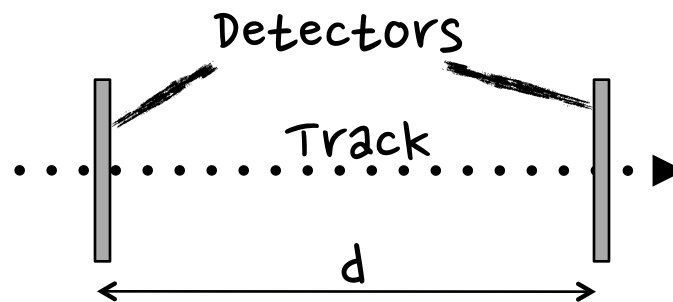
- measure the time difference between two detector planes
 $\beta = d / c \Delta t$
- At high energy, particle speeds are relativistic, closely approaching to c
- Example:
 - For a 10 GeV K,
 - time to travel 12 m is 40.05 ns,
 - whereas for a π
 - it would be 40.00 ns,
 - so difference is only 50 ps



Time Of Flight

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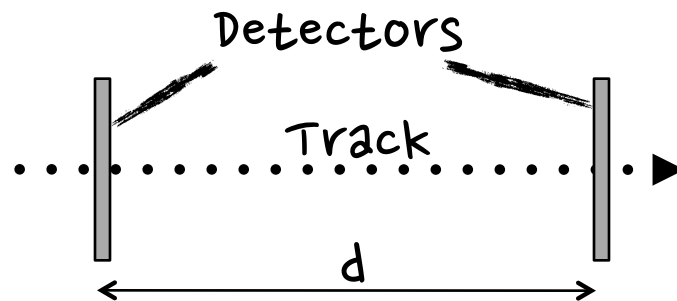
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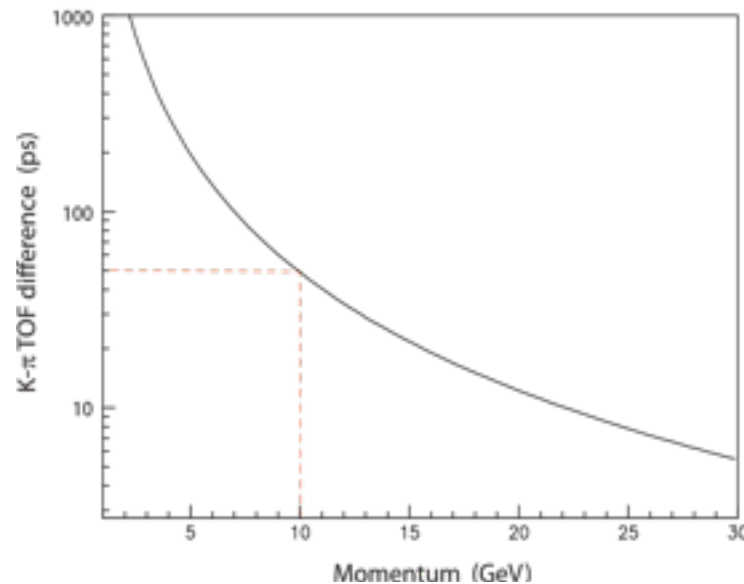
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TOF difference for $d = 12$ m

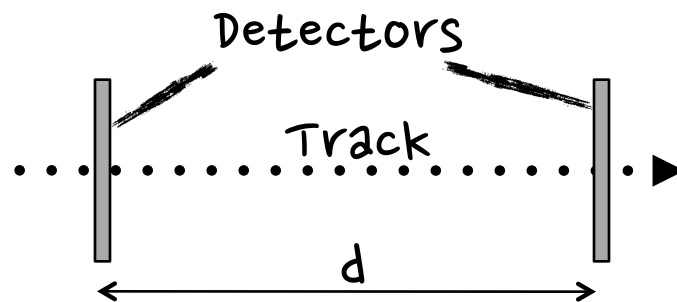




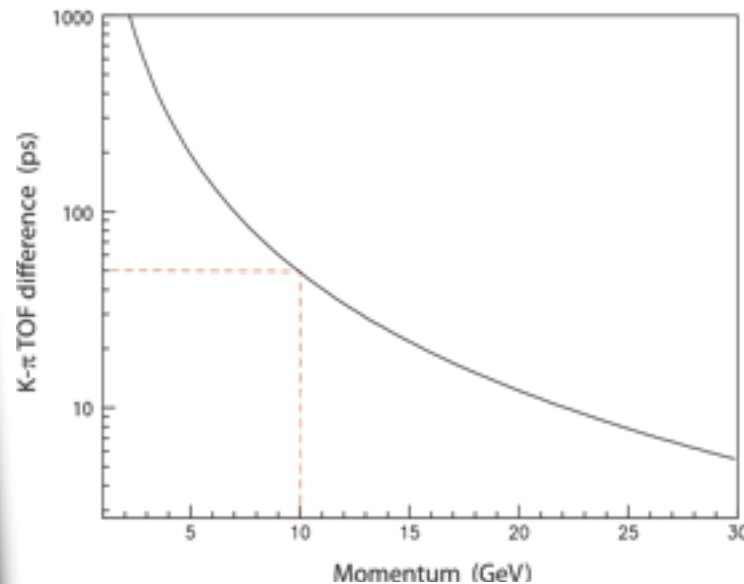
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TOF difference for $d = 12$ m



Modern Detectors

- (+ readout electronics) have resolution $\sigma(t) \sim 10$ ns, fast enough for the LHC (bunch crossings 25 ns apart) but need $\sigma(t) < 1$ ns to do useful TOF
- TOF gives good ID at low momentum
Very precise timing required for $p > 5$ GeV



Ionization



Ionization

Concept

- Charged particles passing through matter can knock out electrons from atoms of the medium: ionization

- Energy loss described by the Bethe-Bloch formula, which gives the universal velocity dependence:

$$dE/dx \propto \log(\beta^2 \gamma^2) / \beta^2$$

- This can be used to identify particles, particularly at low momentum where dE/dx varies rapidly



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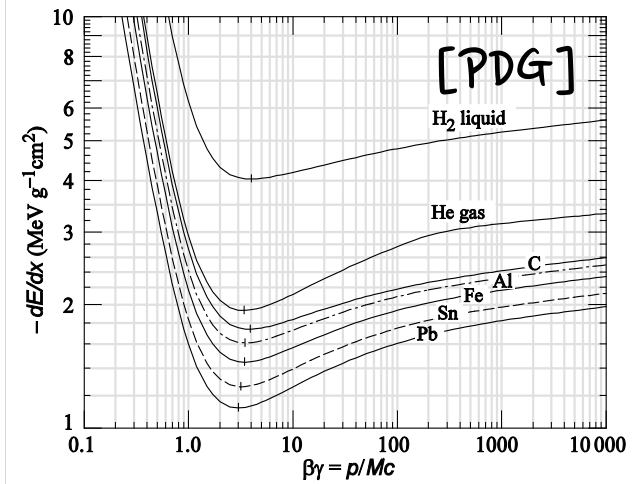
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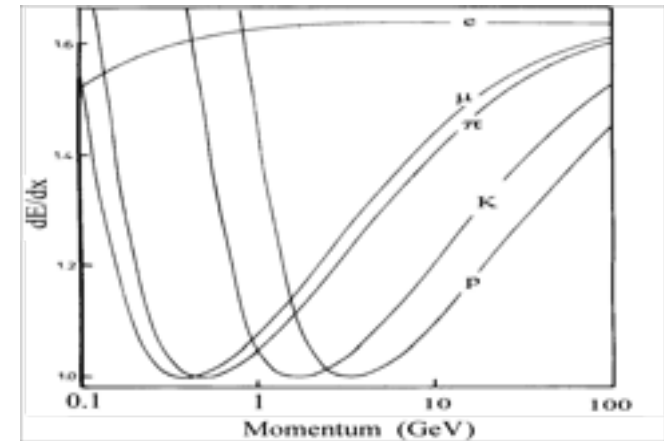
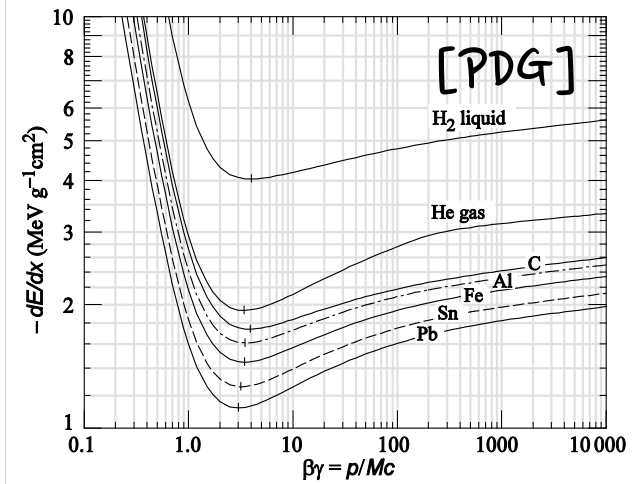
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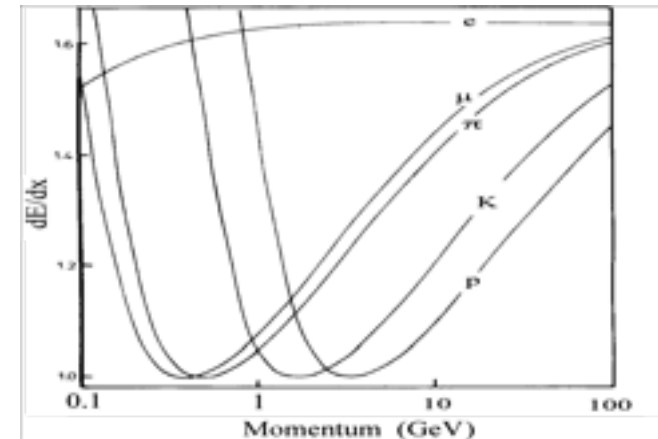
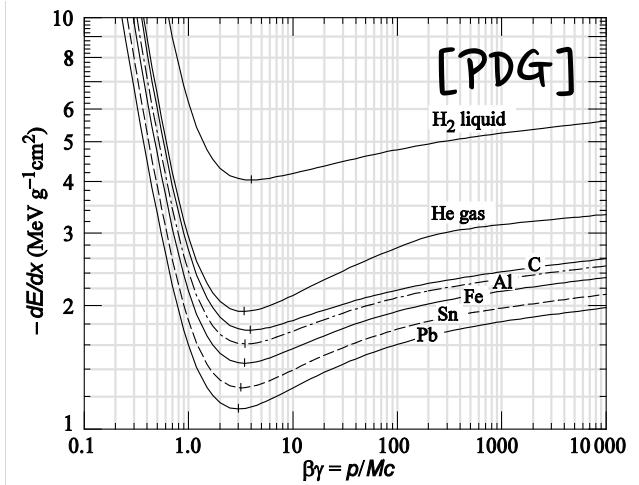
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- This can be used to identify particles, particularly at low momentum where dE/dx varies rapidly

Considerations

- *Advantage:* uses existing detectors needed for tracking (but requires the accurate measurement of the charge)

- *Note:* these techniques all provide signals for charged leptons e, μ as well as π, K, p . But $m(\mu) \approx m(\pi)$, so they are not well separated (dedicated detectors do a better job)





Transition Radiation



Transition Radiation

Transition Radiation

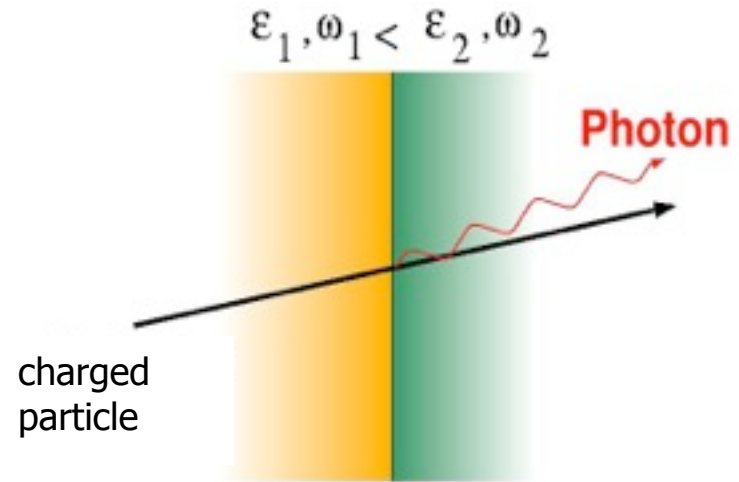
- Local speed of light in medium with refractive index n is $c' = c/n$
- If particle's relative velocity v/c' changes, particle will radiate photons
 - change of direction \rightarrow Synchrotron rad
 - change of speed \rightarrow Bremsstrahlung rad
 - change of index $n \rightarrow$ Transition rad
- Transition rad emitted whenever relativistic charged particle traverses border between two media having different dielectric constants $n = \sqrt{\epsilon}$
- Emitted energy proportional to boost γ of particle
 - can discriminate electrons from pions!
 - (also discriminate hadrons at high energy)



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Черенков Radiation



Considerations

- Named after the Russian scientist P. Cherenkov who was the first to study the effect in depth (he won the Nobel Prize for it in 1958)
- From Relativity, nothing can go faster than the speed of light c (in vacuum)
- However, due to the refractive index n of a material, a particle can go faster than the local speed of light in the medium $c' = c/n$
- This is analogous to the bow wave of a boat travelling over water or the sonic boom of an aeroplane travelling faster than the speed of sound



Черенков Radiation

Considerations

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Concept

- Cherenkov Light is produced when charged particle goes faster than the speed of light
 - $\cos \theta_c = 1 / \beta n$
- Produced in three dimensions, so the wavefront forms a cone of light around the particle direction
- Measuring the opening angle of cone
 - particle velocity can be determined



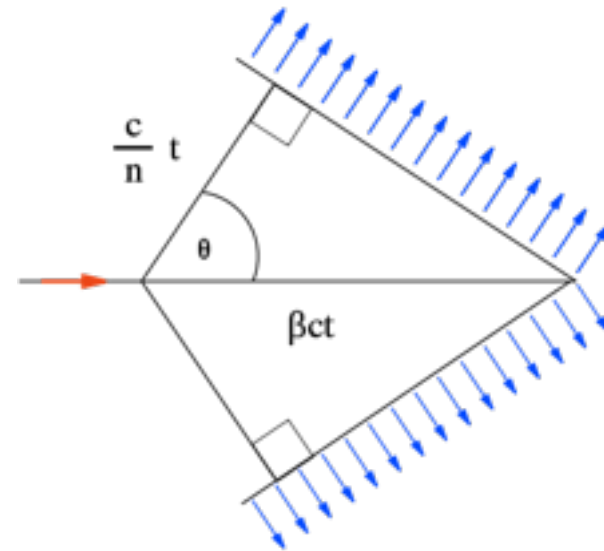
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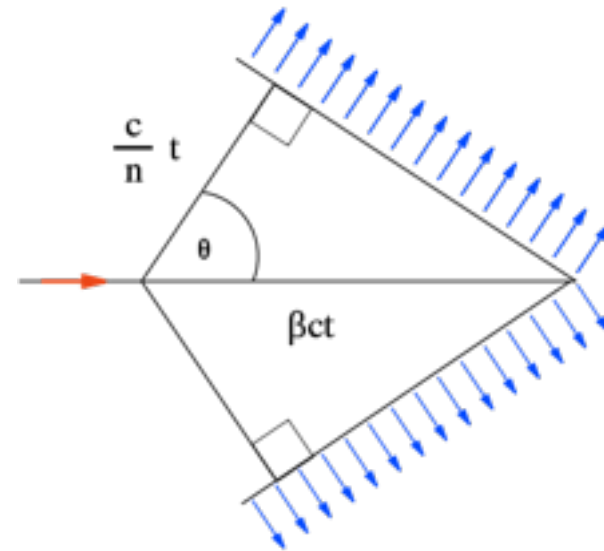
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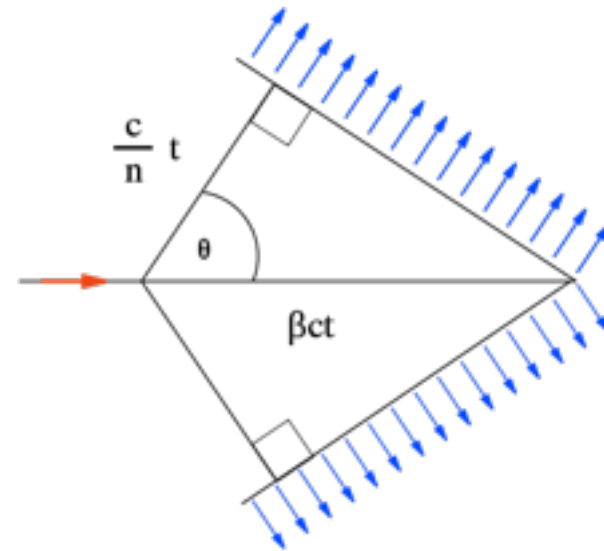
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Example:

- Using this construction, we can determine (roughly) the boat speed:
 - $\theta \approx 70^\circ$, $v_{\text{wave}} = 2$ knots on water
 - $v_{\text{boat}} = v_{\text{wave}} / \cos \theta \approx 6$ knots



K. Cavanaugh, HCPSS 2012



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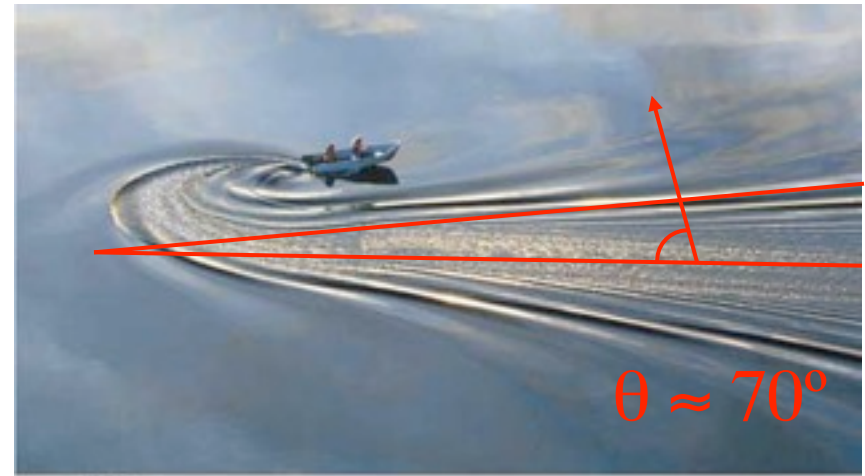
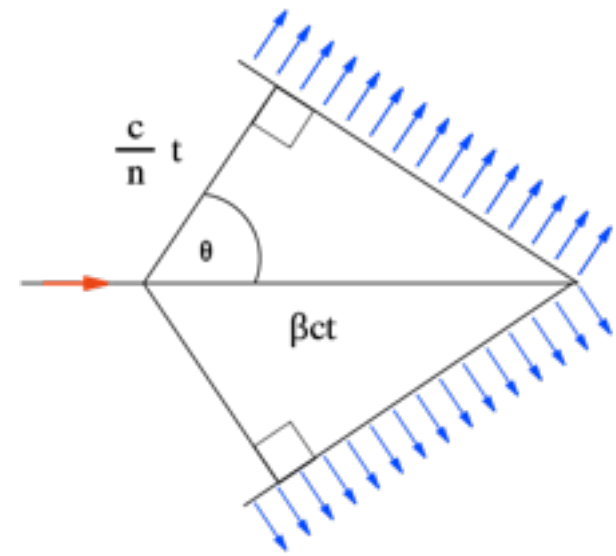
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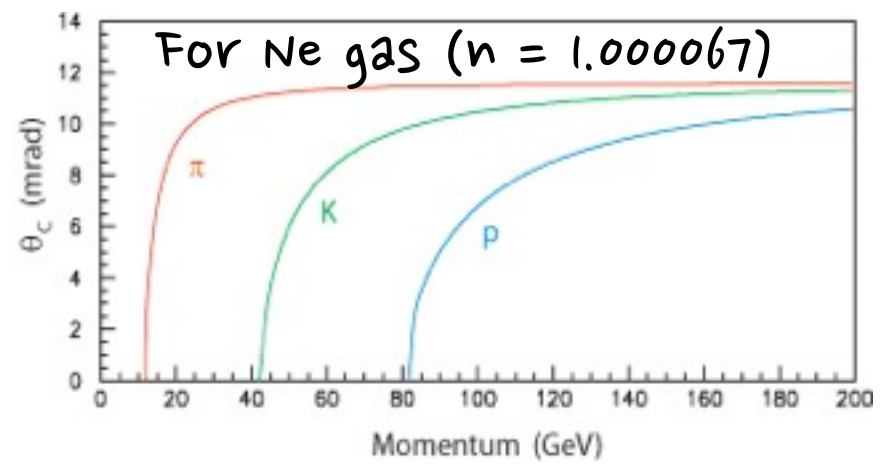
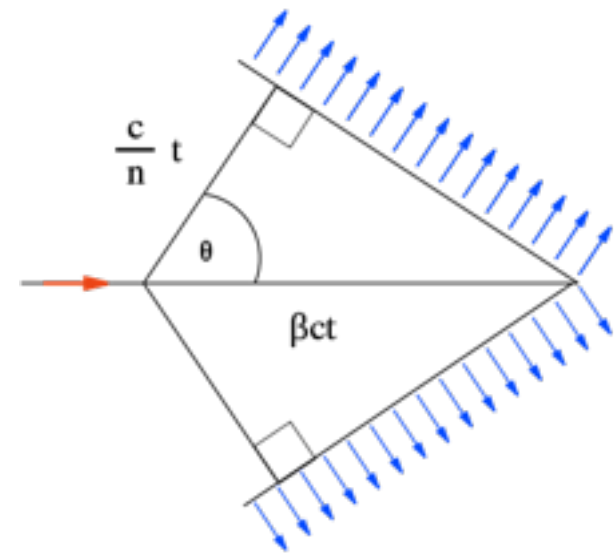
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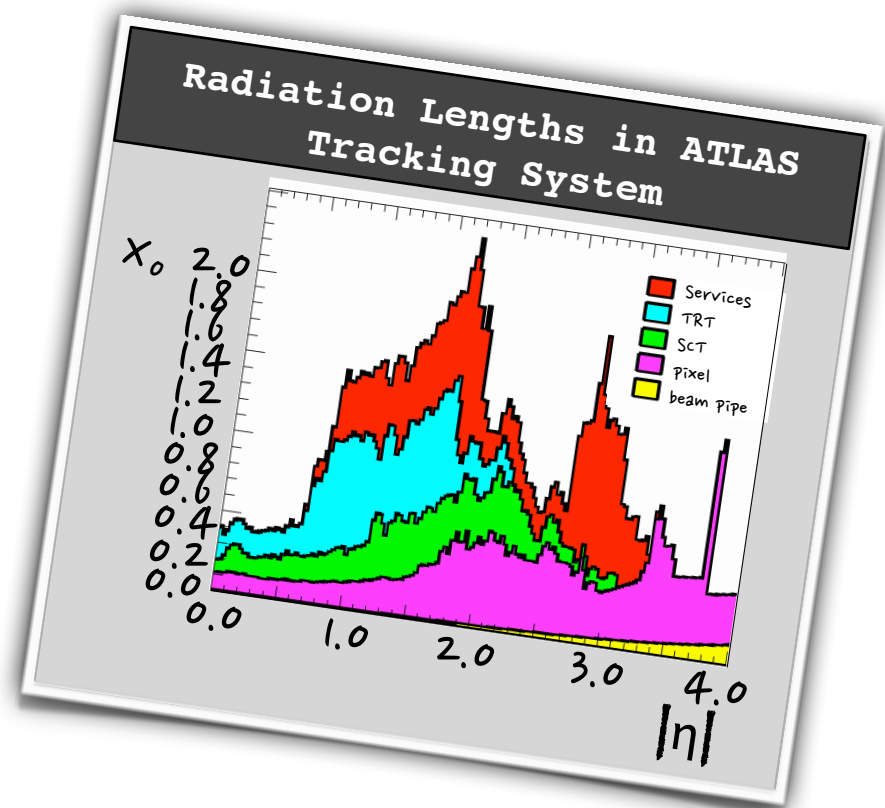
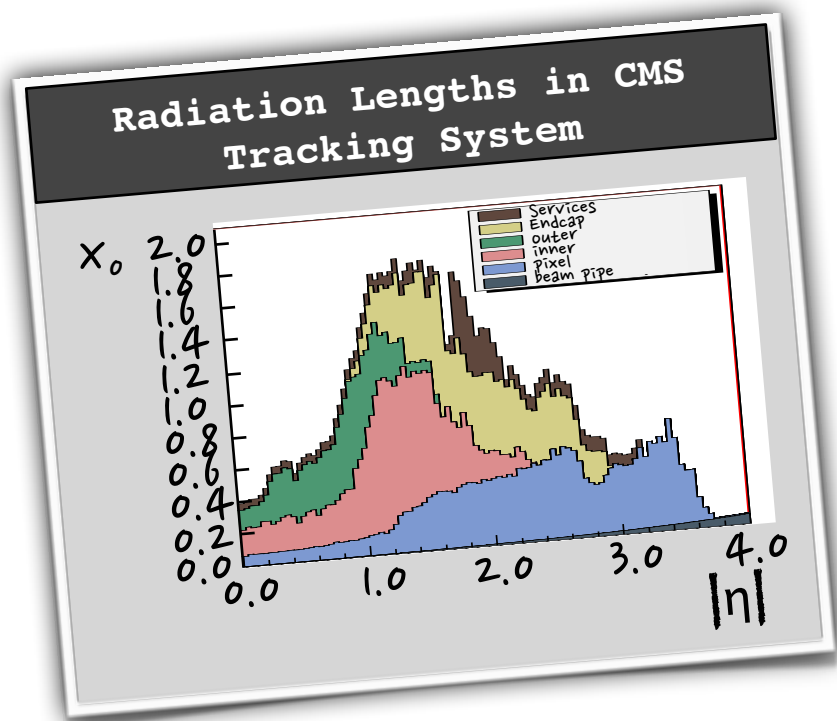
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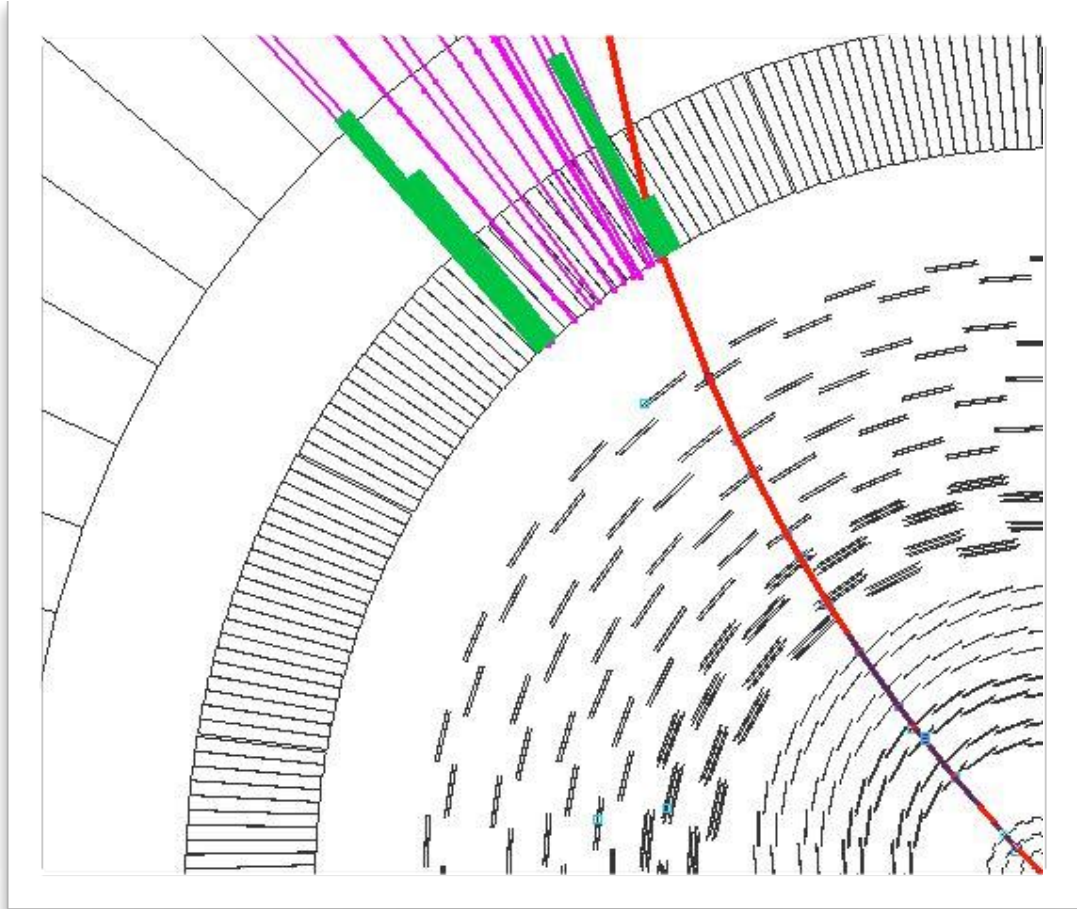


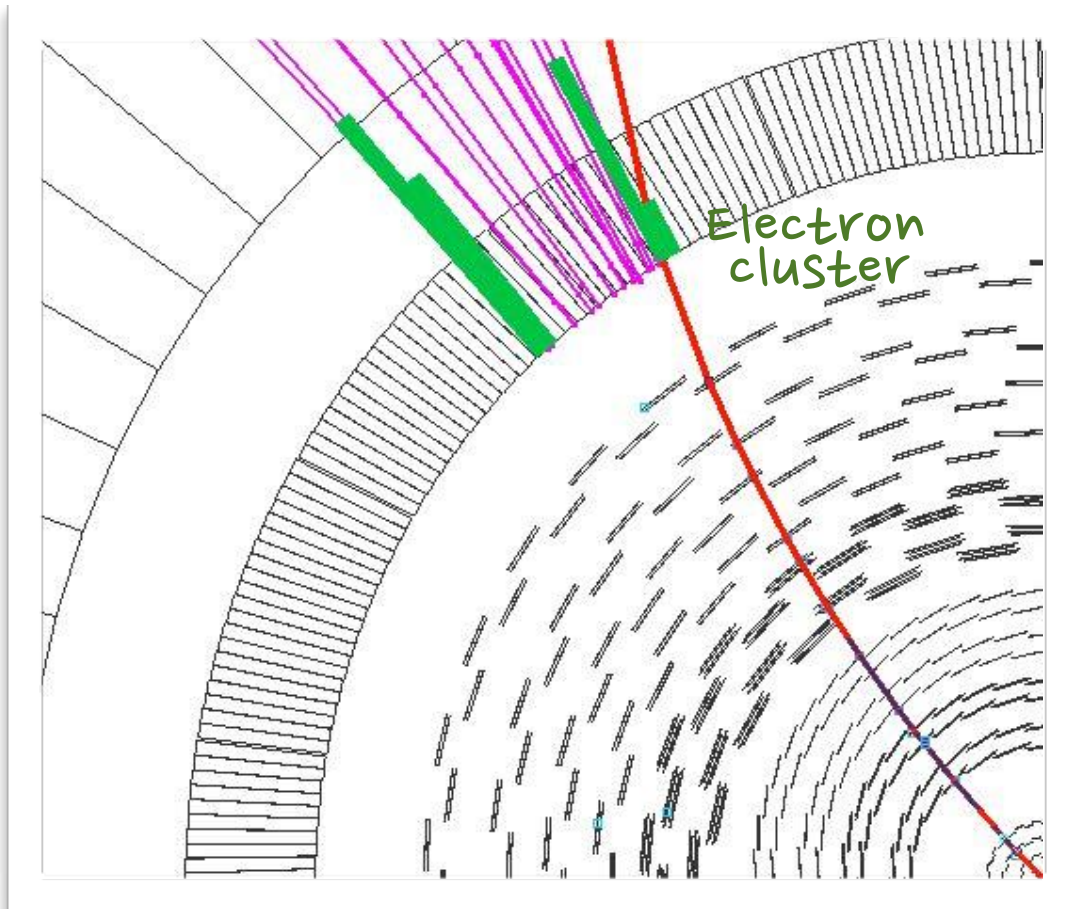
Come back to Electrons & Photons

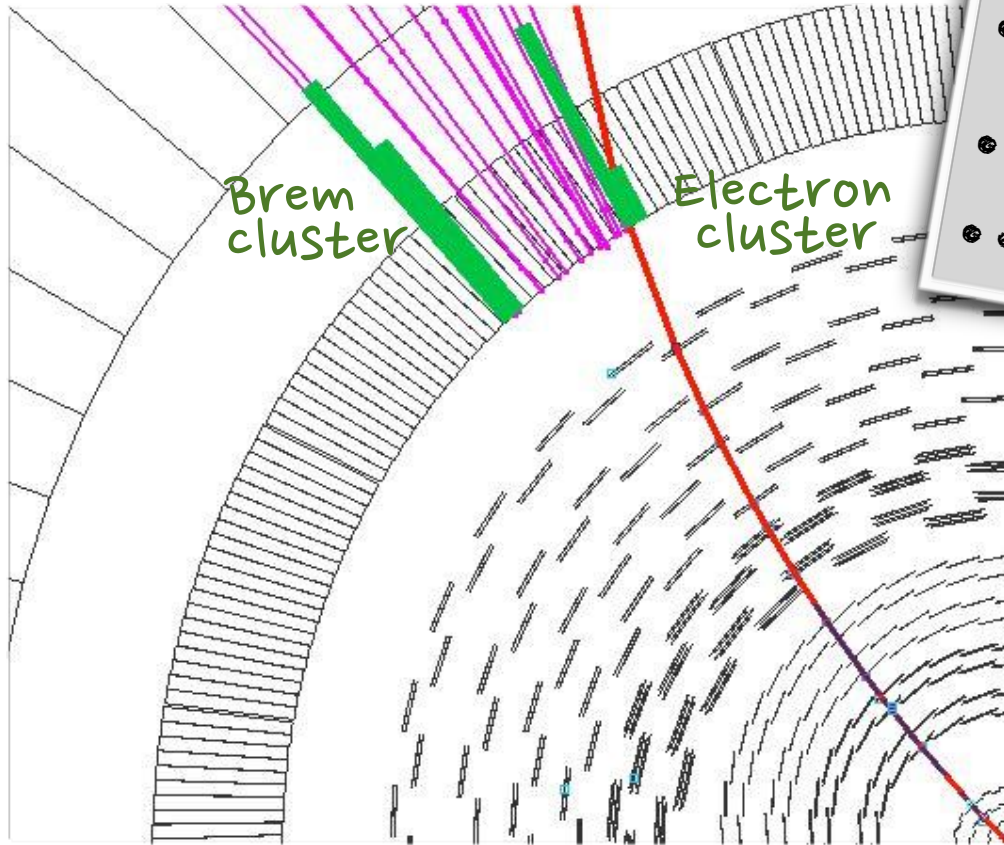
- Indeed, the LHC Trackers are heavy!



The LHC Trackers act like Electromagnetic pre-showers!

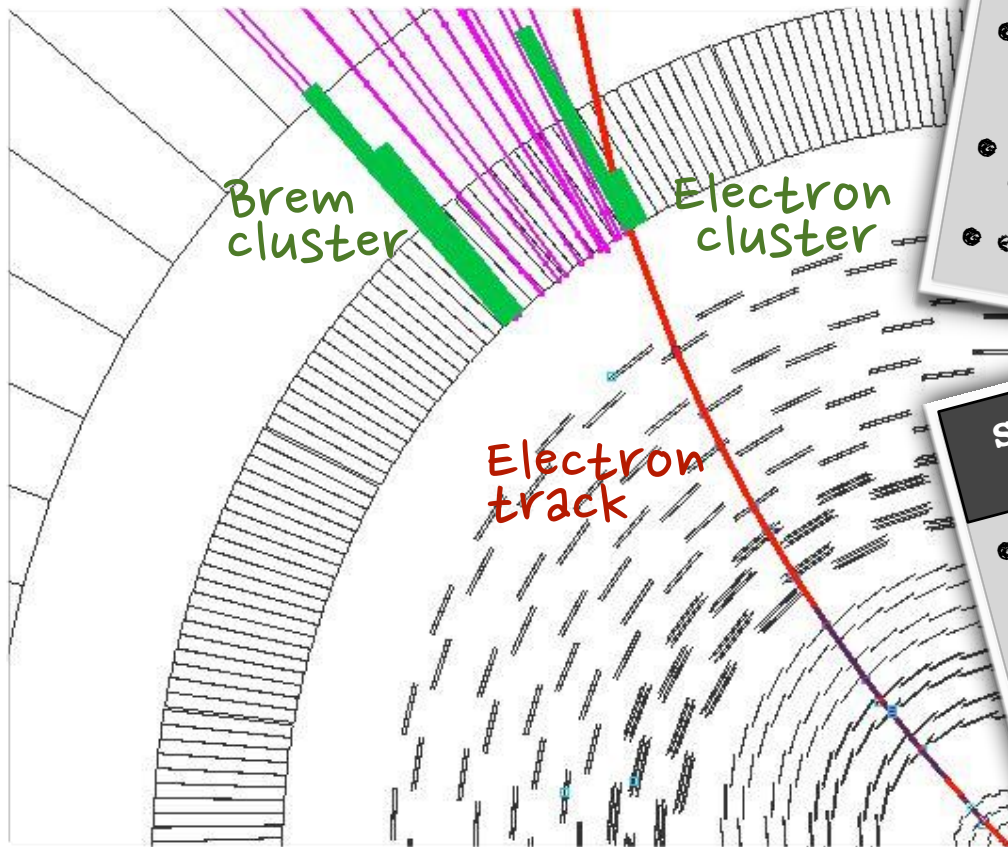






Huge Effect!

- Electrons radiate on average $\approx 70\%$ of their energy in the track by bremsstrahlung
- photons have $> 50\%$ probability to convert to e^+e^- pair
- energy spreads in φ due to B-field

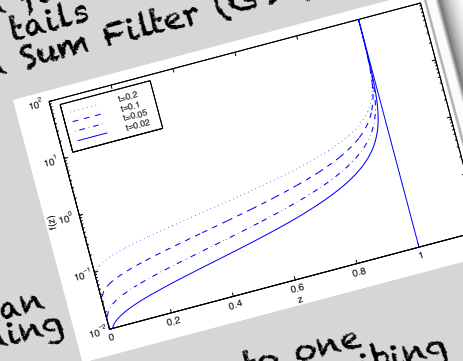


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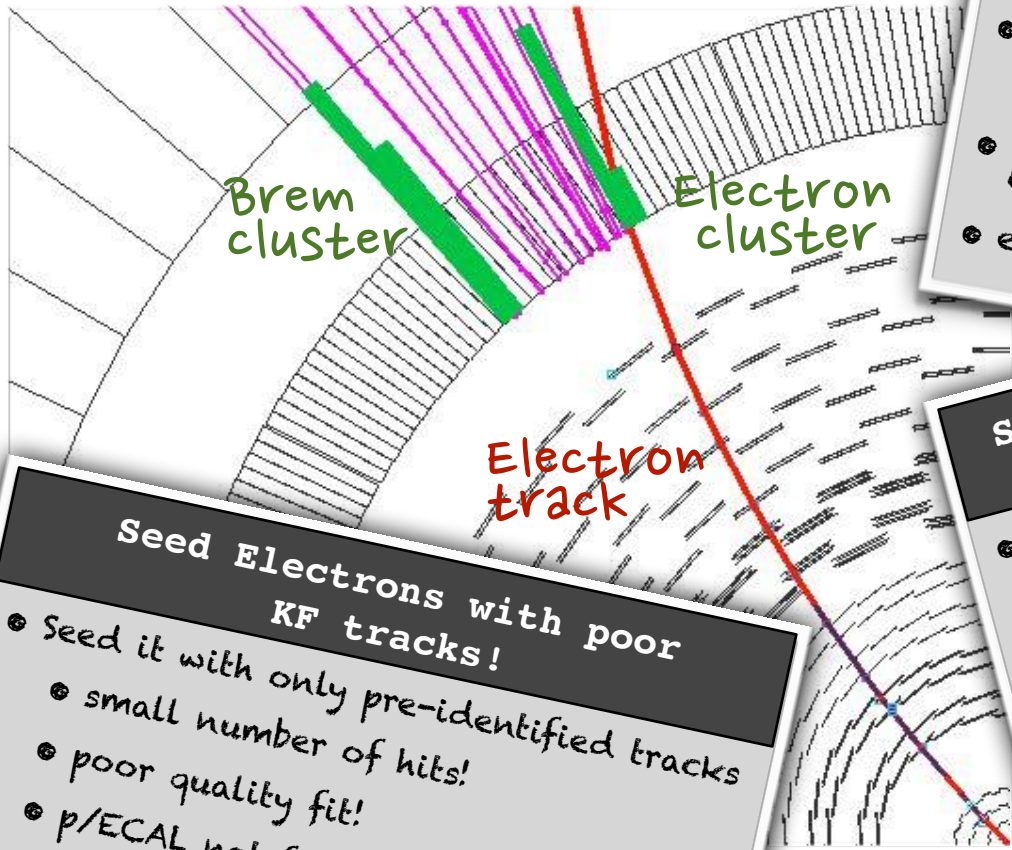
Standard KF tracking fails for electrons!

- Because they radiate! Kalman filter pattern recognition quickly gives up
- Use sum of Kalman filters to account for non-Gaussian tails \Rightarrow Gaussian Sum Filter (GSF)
- attempts to approximate Bethe-Heitler energy loss
- resembles a set of Kalman filters running in parallel
- each corresponding to one gaussian component describing BH energy loss
- Computationally expensive!





Electron Identification



Huge Effect!

- Electrons radiate on average $\approx 70\%$ of their energy in the track by bremsstrahlung
- photons have $> 50\%$ probability to convert to e^+e^- pair
- energy spreads in

Standard KF tracking fails for electrons!

- Because they radiate! Kalman filter pattern recognition quickly gives up
- Use sum of Kalman filters to account for non-Gaussian tails \Rightarrow Gaussian Sum Filter (GSF)
- attempts to approximate Bethe-Heitler energy loss
- resembles a set of Kalman filters running in parallel
- each corresponding to one gaussian component describing BH energy loss
- Computationally expensive!

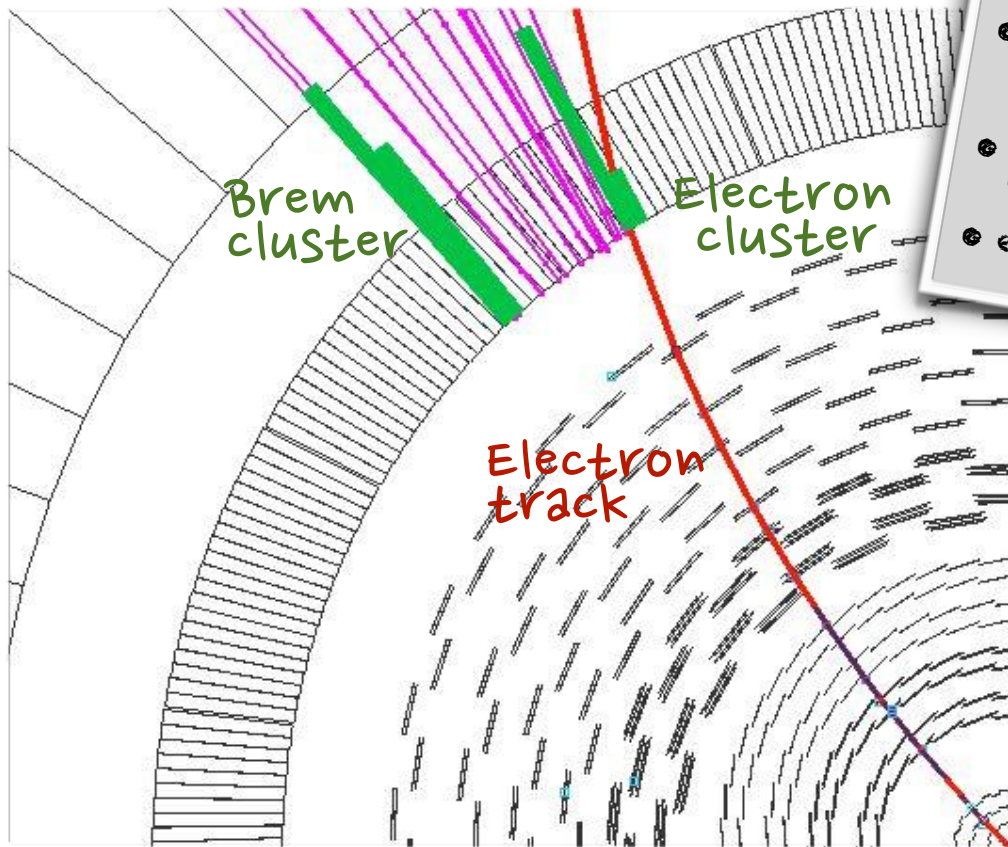
The graph plots energy loss on a logarithmic scale (y-axis, 10^0 to 10^2) against a parameter (x-axis, 0 to 1). It shows several curves for different values of α : $\alpha=0.2$ (dotted), $\alpha=0.1$ (dashed), $\alpha=0.05$ (solid), and $\alpha=0.02$ (dash-dot). The curves show an exponential-like increase as the parameter approaches 1.

Seed Electrons with poor KF tracks!

- Seed it with only pre-identified tracks
- small number of hits!
- poor quality fit!
- p/ECAL not far from unity
- With the above seeds
- only 5% of all tracks are seeds
- sufficient to run GSF tracking
- 95% efficiency for electrons

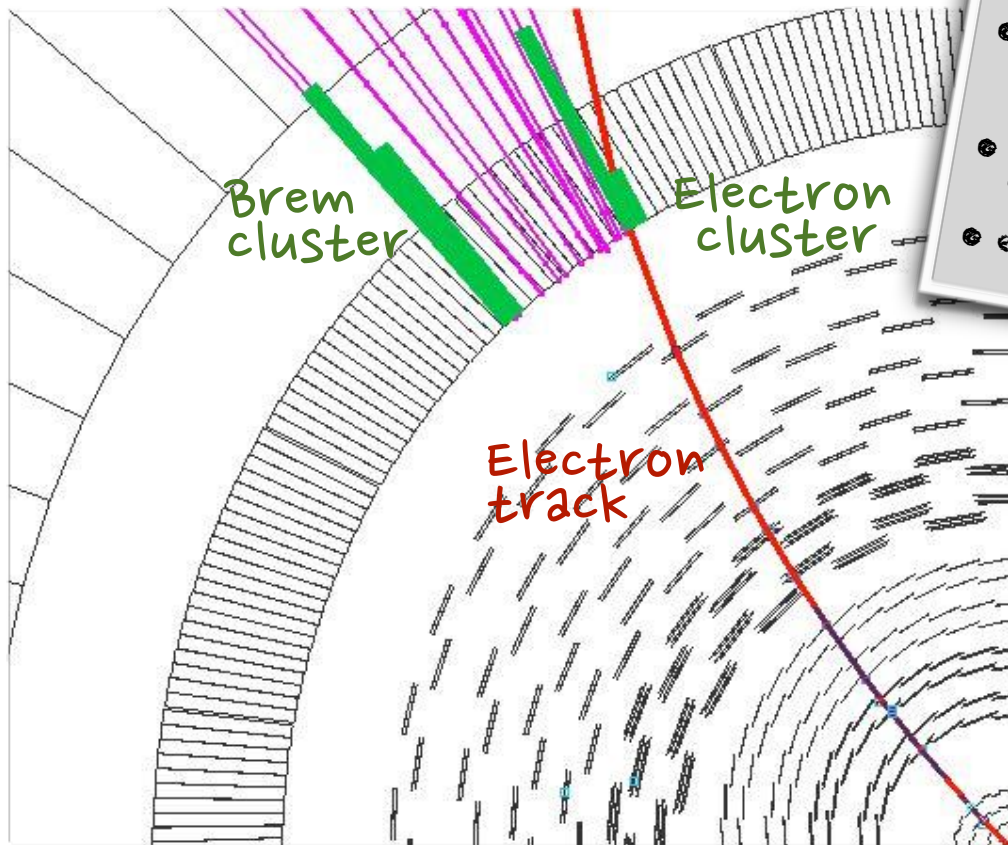
Slide: 20
Date: 14.08.2012

vanough, HCPSS 2



Huge Effect!

- Electrons radiate on average $\approx 70\%$ of their energy in the track by bremsstrahlung
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- energy spreads in φ due to B-field

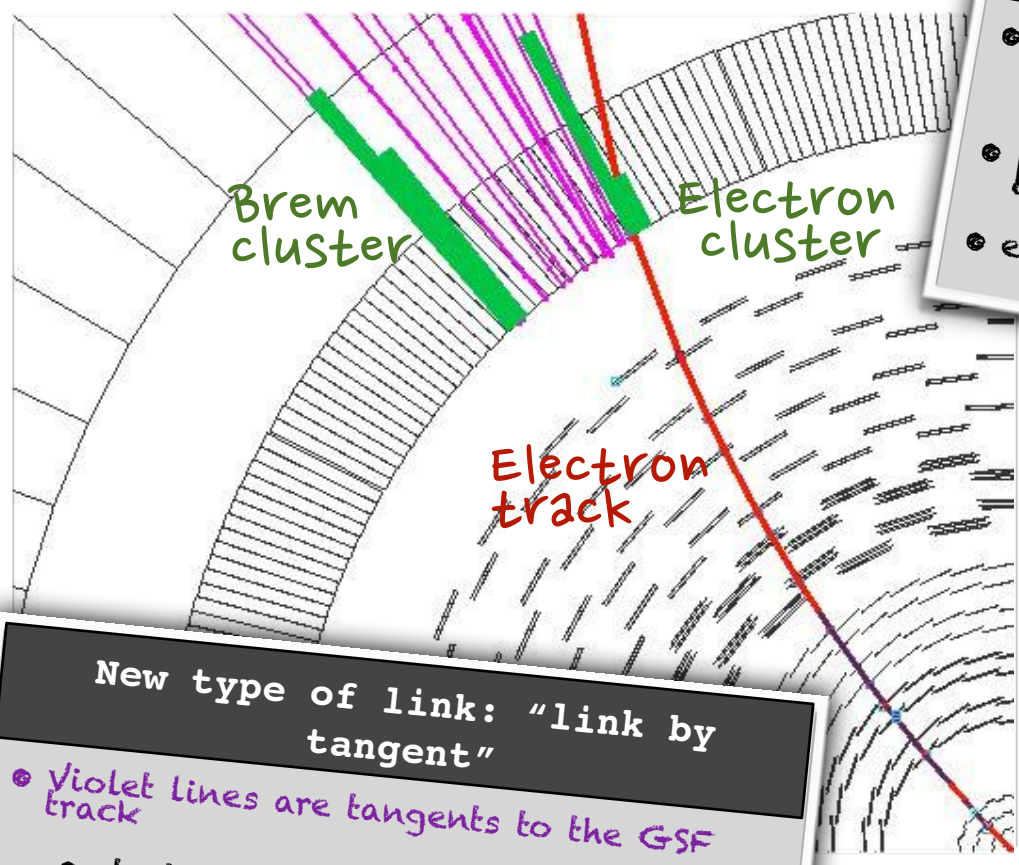


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Recovering the Bremsstrahlung photons

- If nothing is done, radiated photon energy counted twice
- once from electron initial track momentum p_{in}
- once from the energy corresponding to the ECAL cluster(s), E_{brem}



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New type of link: "link by tangent"

- Violet lines are tangents to the GSF track
- starting from each tracker layer

If tangent points to a ECAL cluster

- link cluster to track

Another possible handle

- test compatibility between ECAL cluster and $\Delta\phi$ along GSF track



Huge Effect!

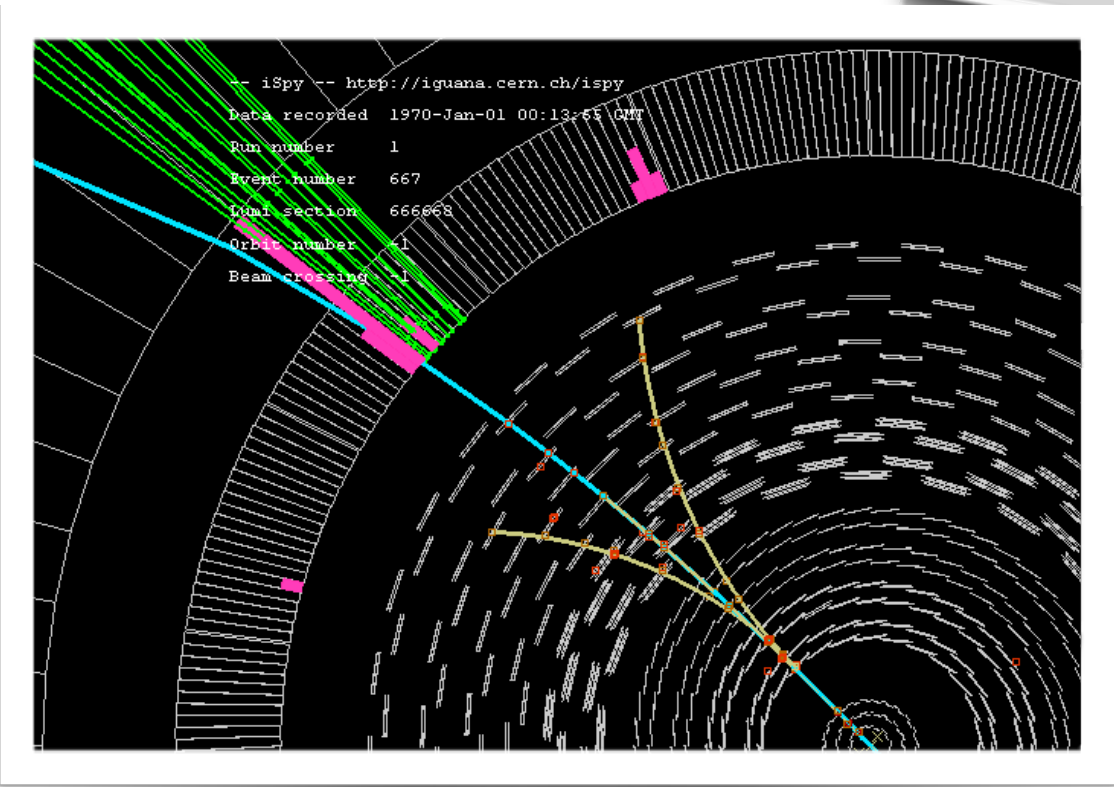
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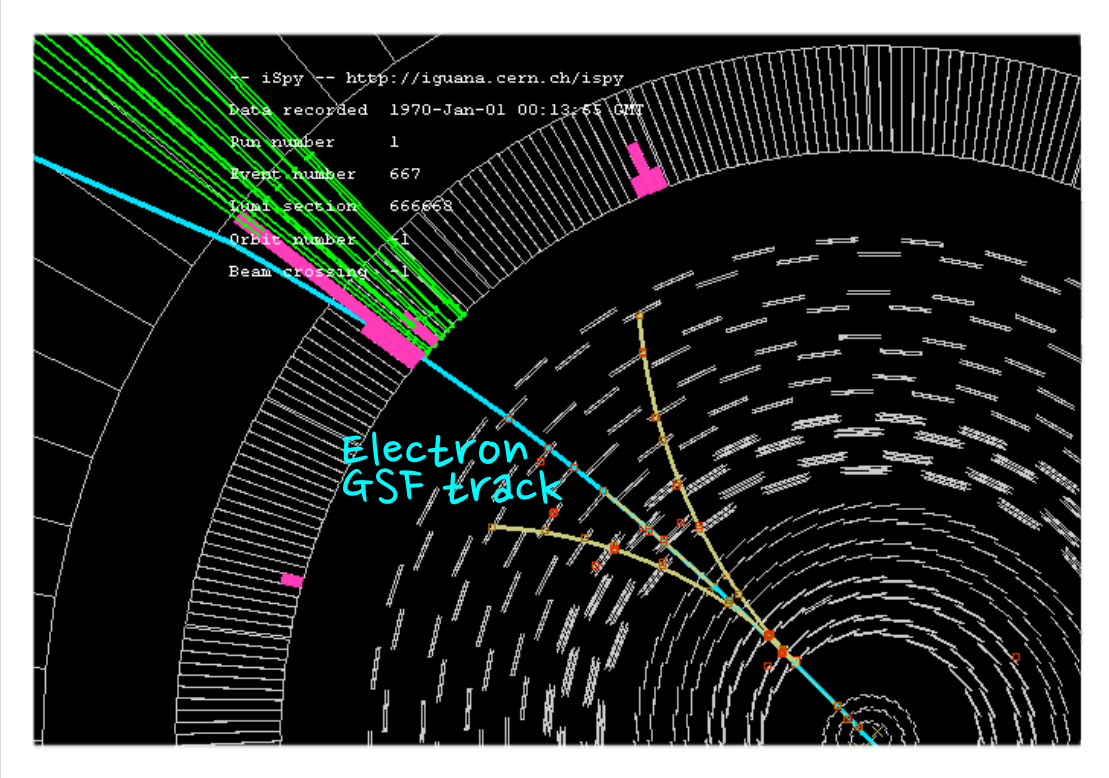
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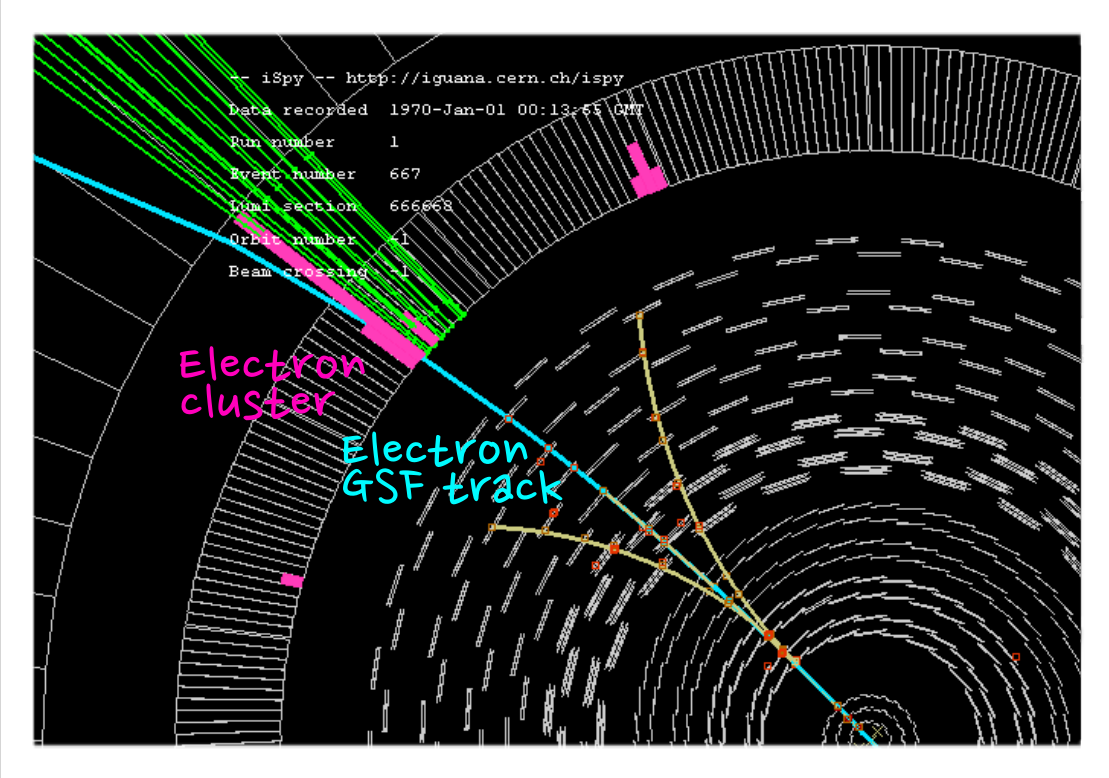
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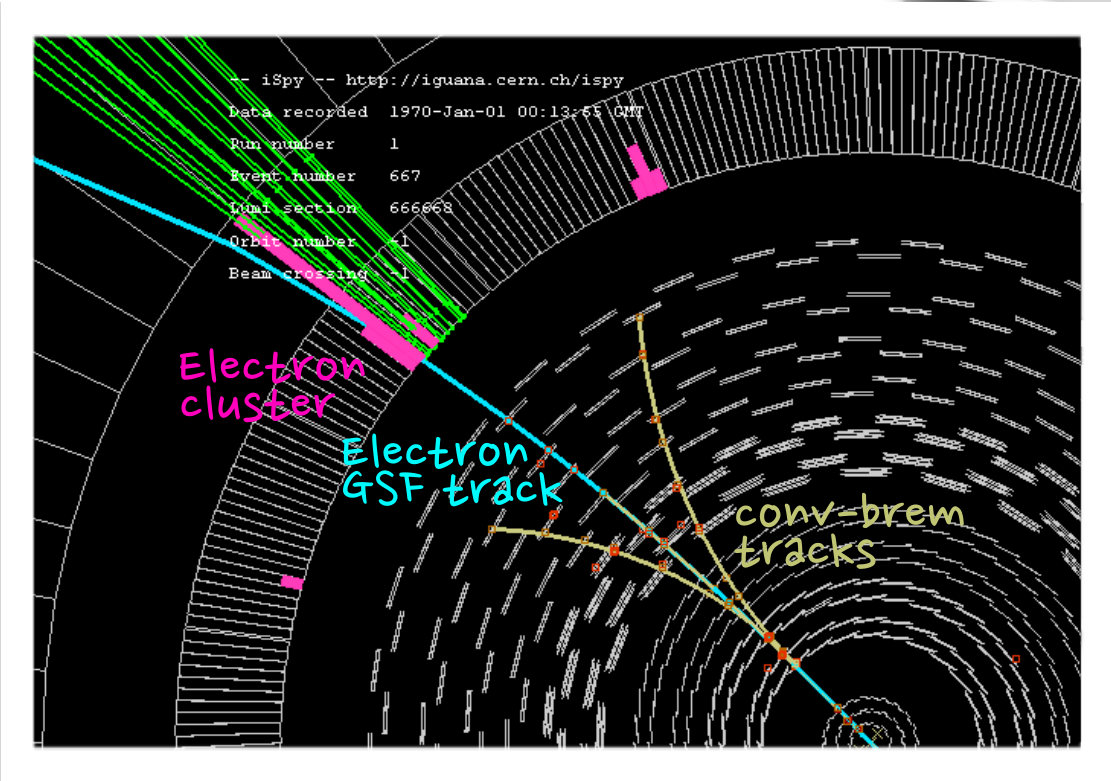
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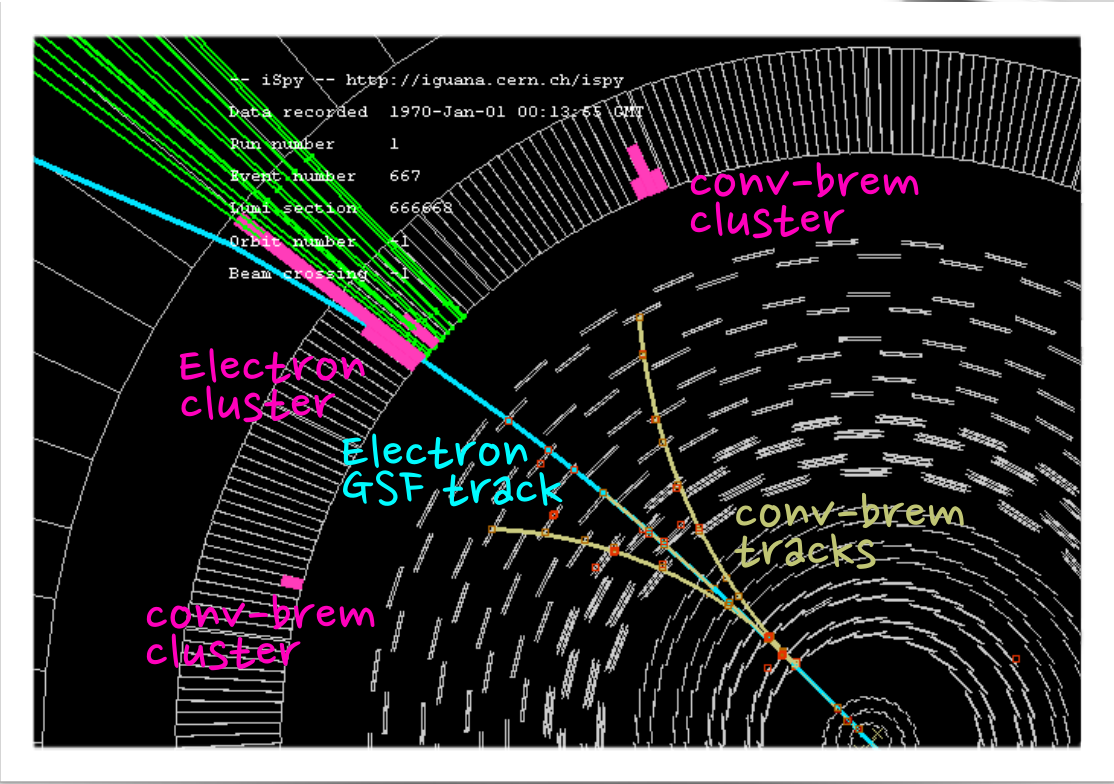
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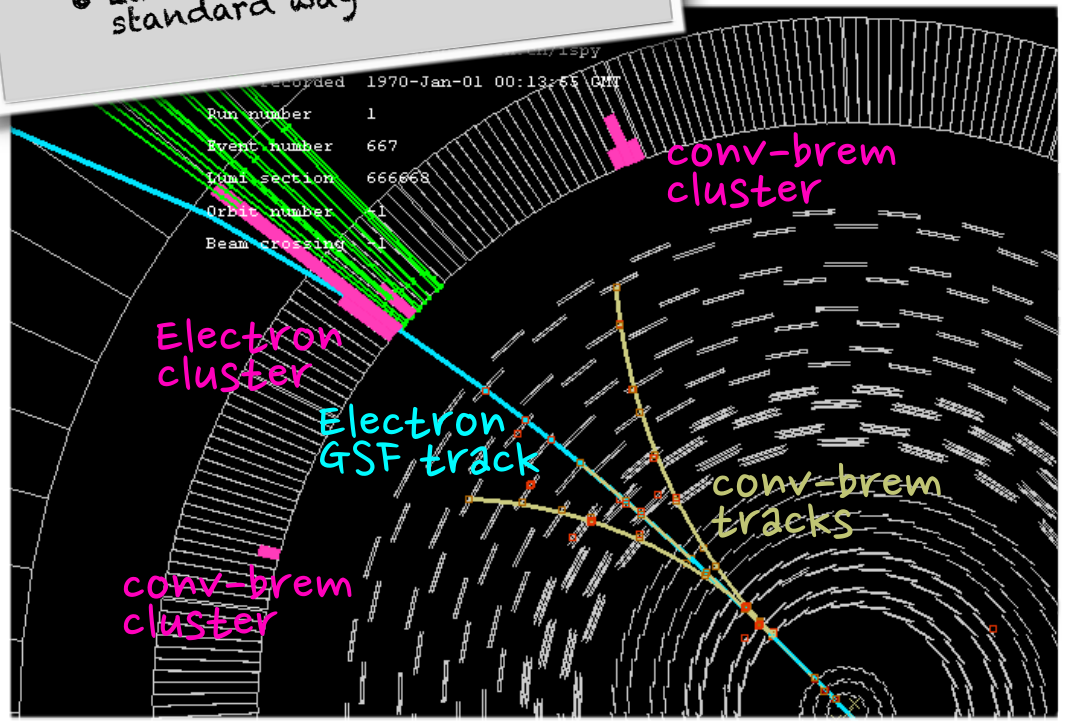
Electron Identification

Recovering the converted Bremsstrahlung photons

- Tracks recovered by the 5th and 6th step of iterative tracking
- Linked by vertex to the original electron track (similar to π nuclear interactions)
- Linked to ECAL clusters in standard way

Huge Effect!

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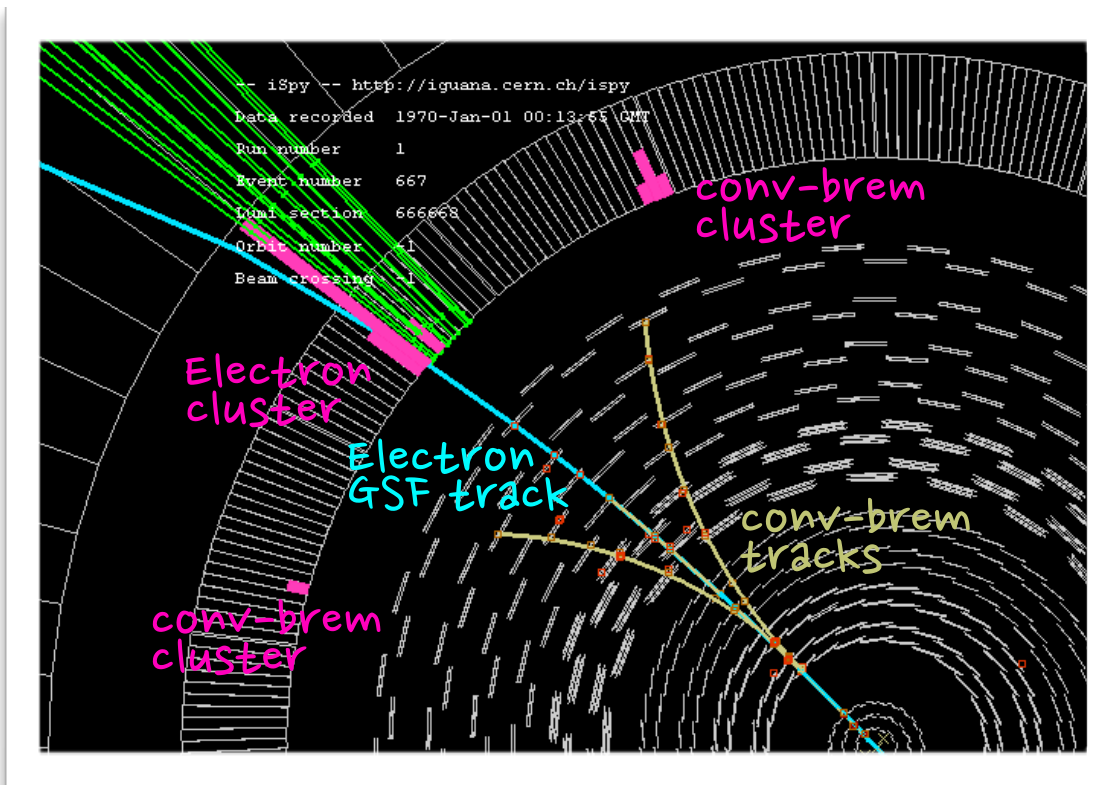
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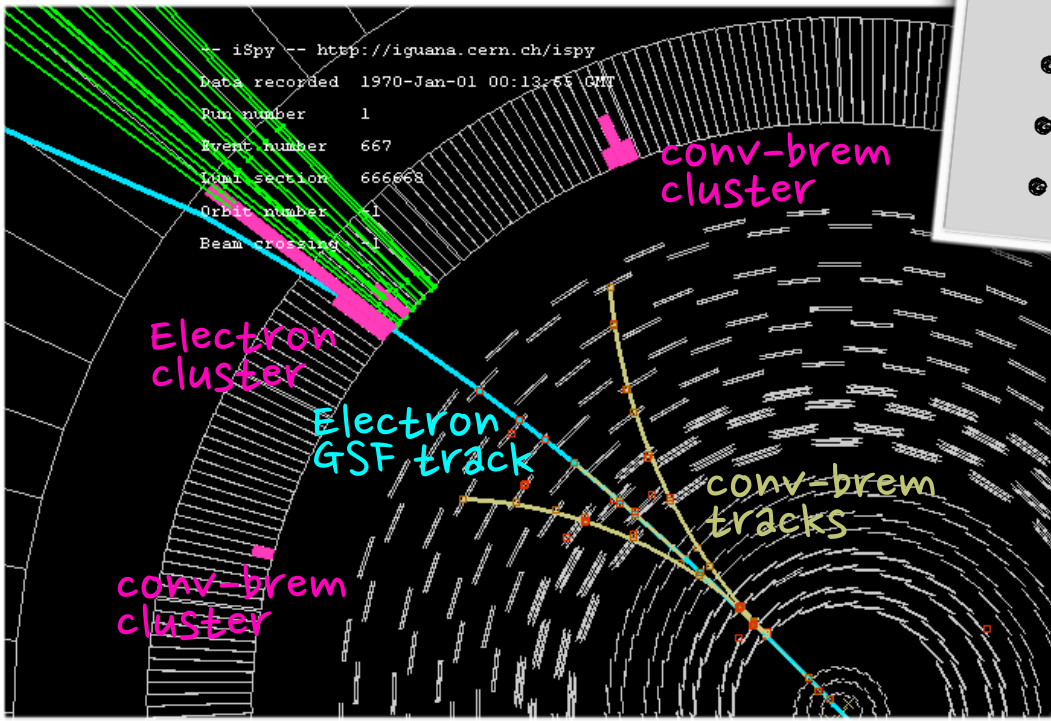
Electron Identification



Use Tracker as Preshower!

- electrons radiate; pions don't!
- Use the following to discriminate
 - Number of hits of the KF tracks
 - Energy loss along GSF track:

$$\Delta p = p_{in} - p_{out}$$
 - Number of Brem γ 's associated with track
- Compare E_{brem} and $\Delta p = p_{in} - p_{out}$
- Compare $E_{electron} + E_{brem}$ and p_{in}
- etc



```

-- iSpy -- http://iguana.cern.ch/ispy
Data recorded 1970-Jan-01 00:13:55 GMT
Run number 1
Event number 667
Lumi section 666668
Orbit number 1
Beam crossing 1
    
```



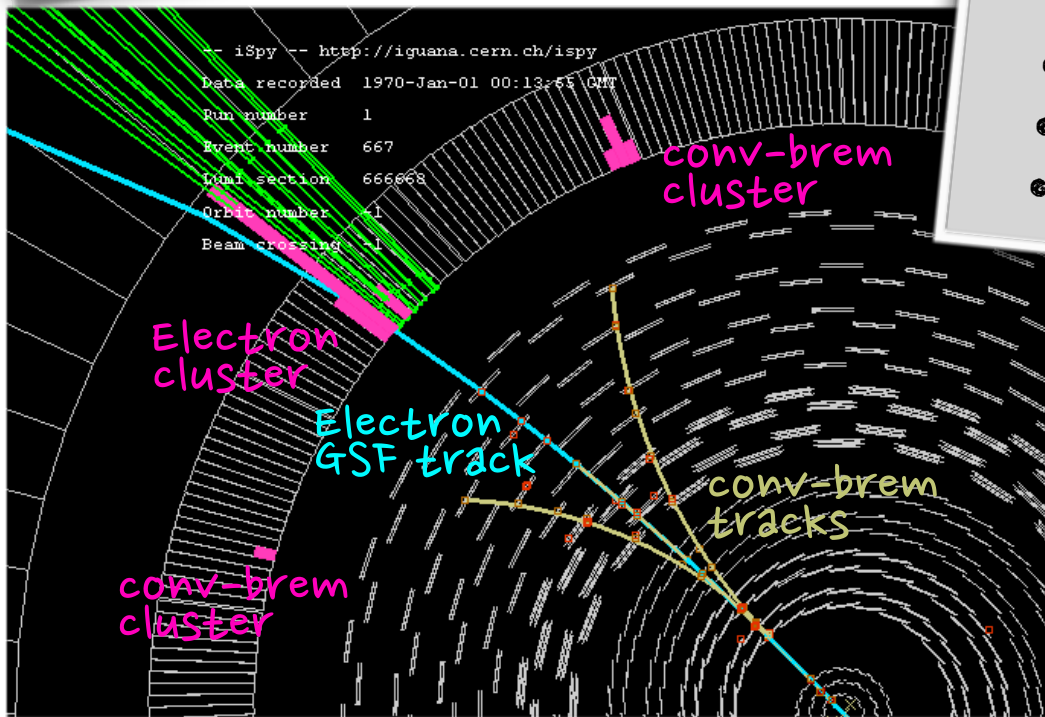
Electron Identification

Use cluster information!

- Shower shape, such as width along η
- any possible linked HCAL energy
- etc

Use Tracker as Preshower!

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 - Number of hits of the KF tracks
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Electron Identification

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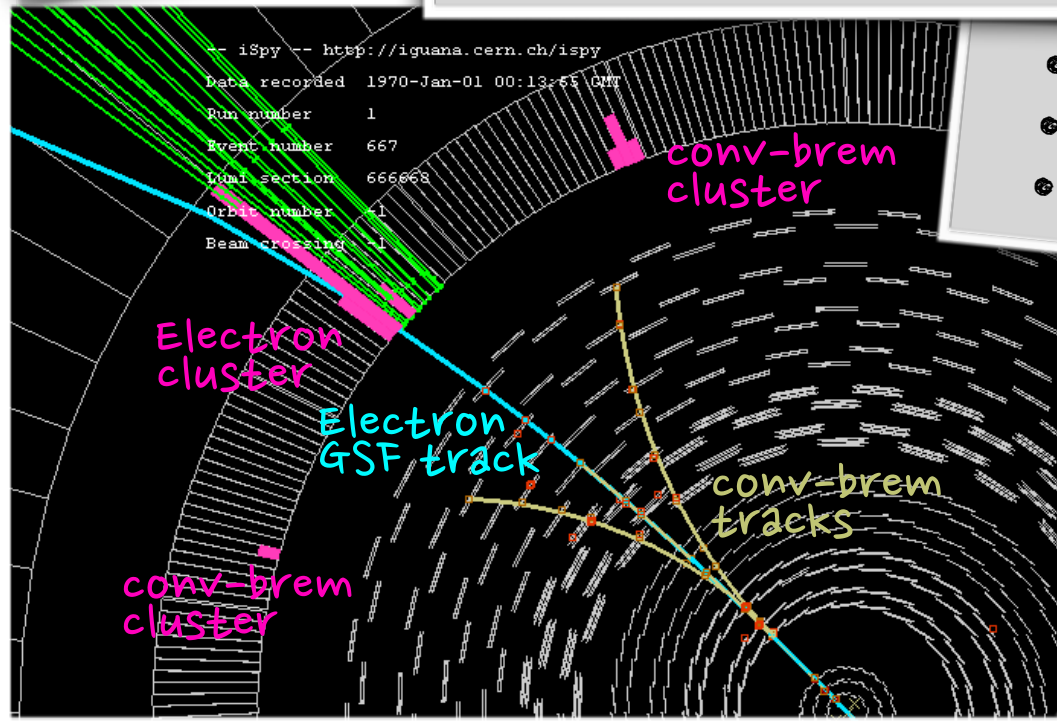
Combine into Multivariate Discriminator

- 95% efficiency for isolated electrons
- 70-80% efficiency in jets

Use Tracker as Preshower!

- electrons radiate; pions don't!
- allowing to discriminate
- of hits of the KF tracks
- loss along GSF track:
 $\Delta p = p_{in} - p_{out}$
- of Brem γ 's associated with

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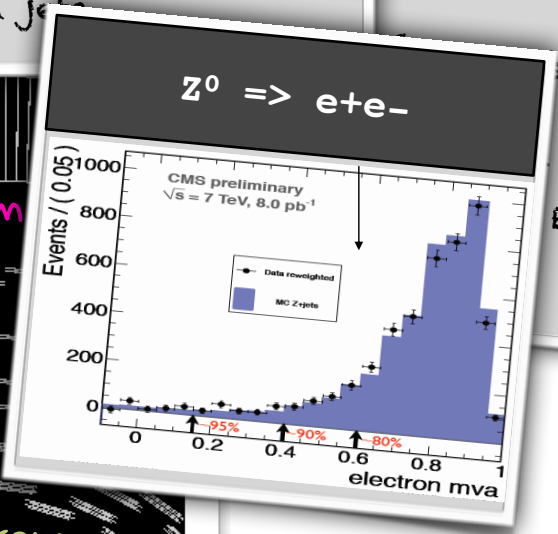
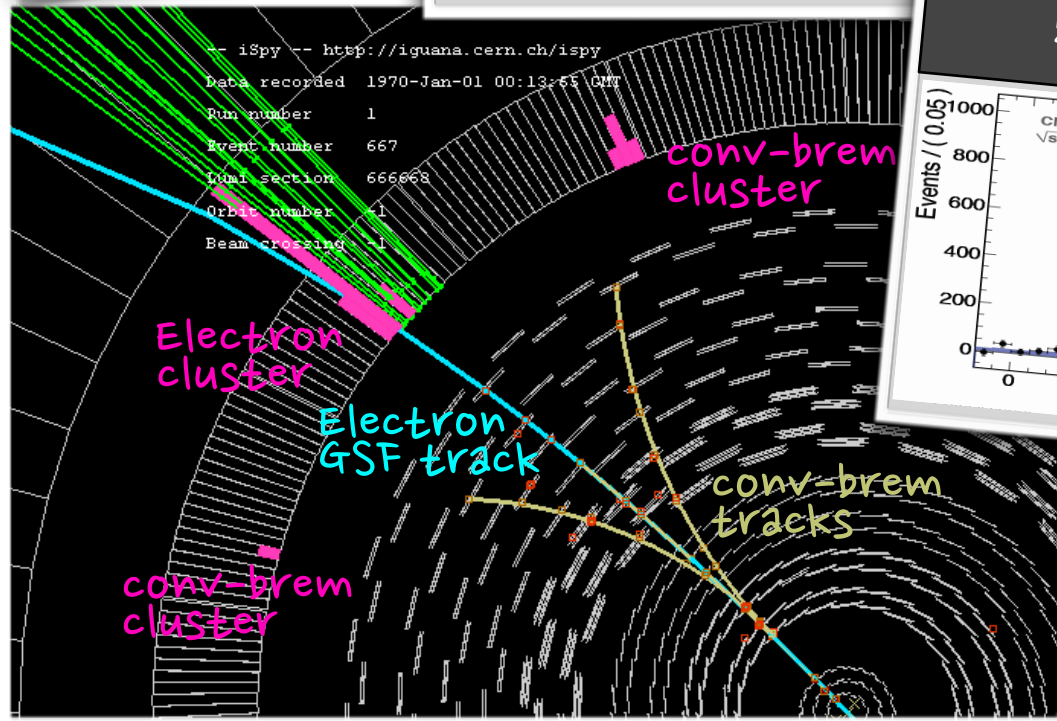
- Shower shape, such as
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- of hits of the KF tracks
- loss along GSF track:
- $\Delta p = p_{in} - p_{out}$
- associated with
- $\Delta p = p_{in} - p_{out}$
- Ebrem and p_{in}





Electron Identification

Use cluster information!

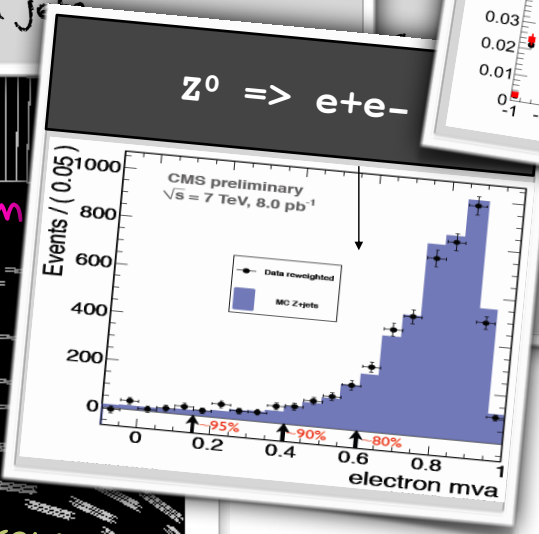
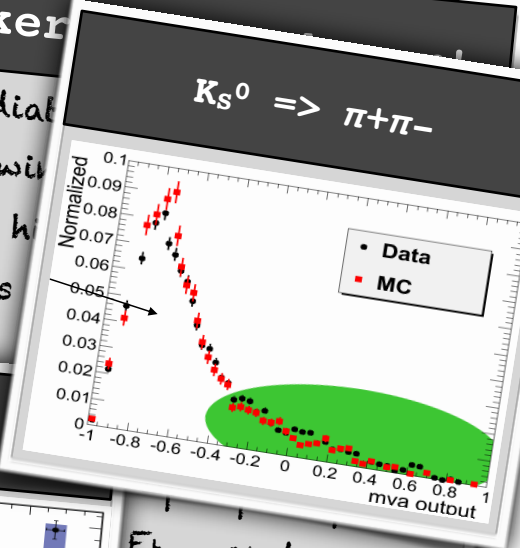
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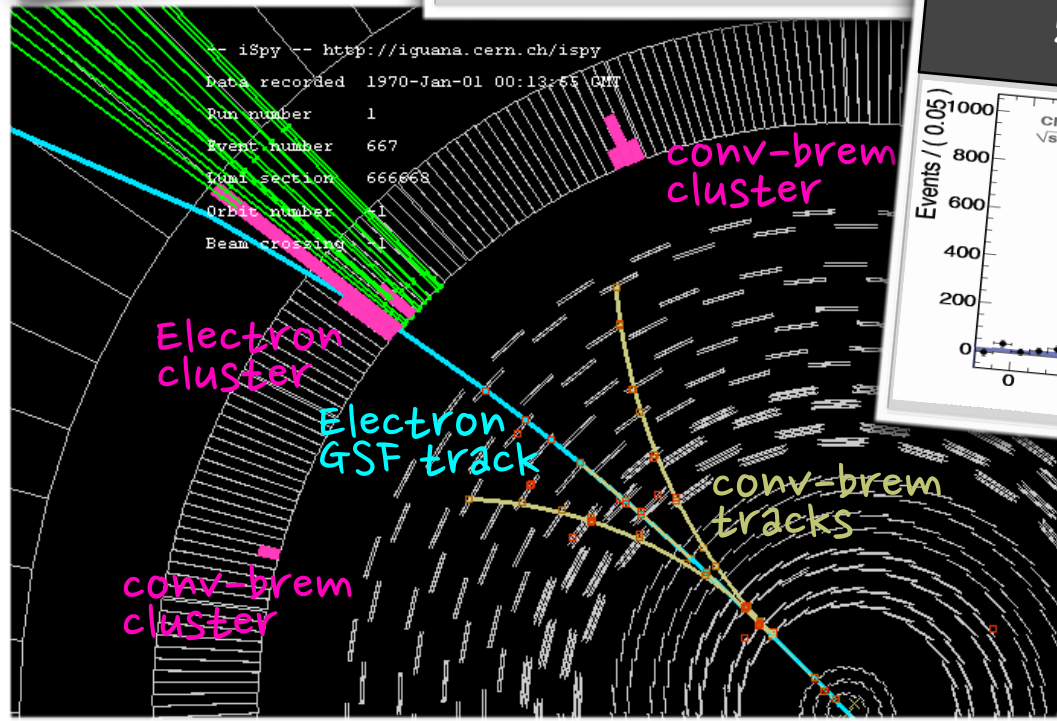
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Use Tracker

- electrons radiate
- low loss



Ebrem and pin





Electron Identification

Use cluster information!

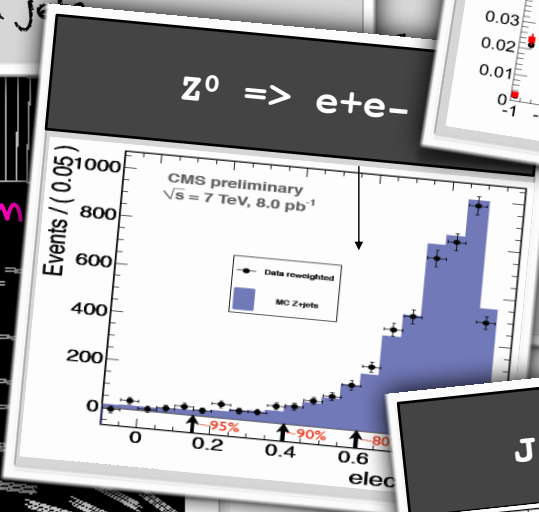
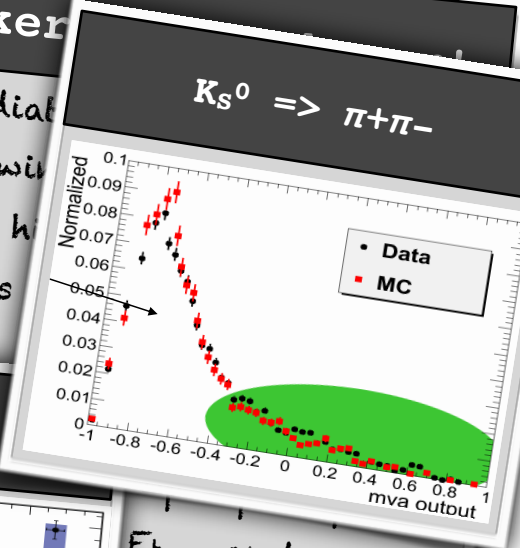
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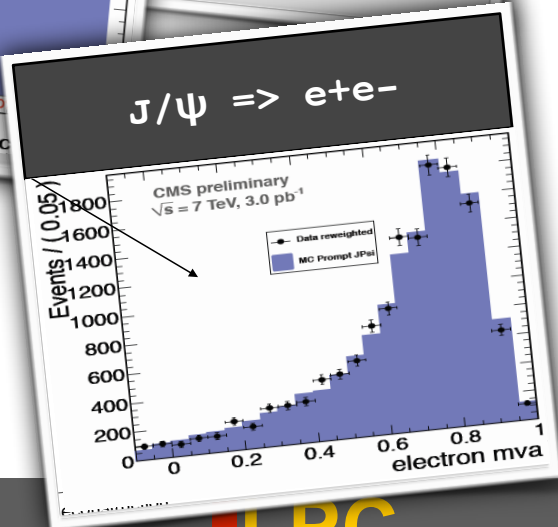
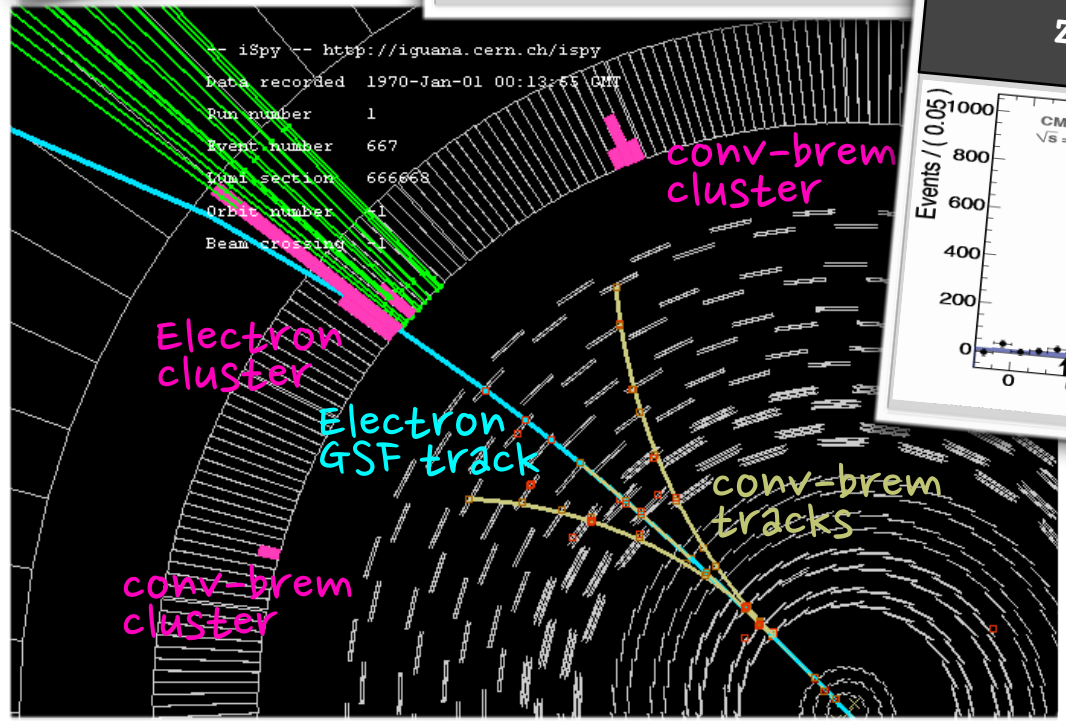
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Use Tracker

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Ebrem and pin





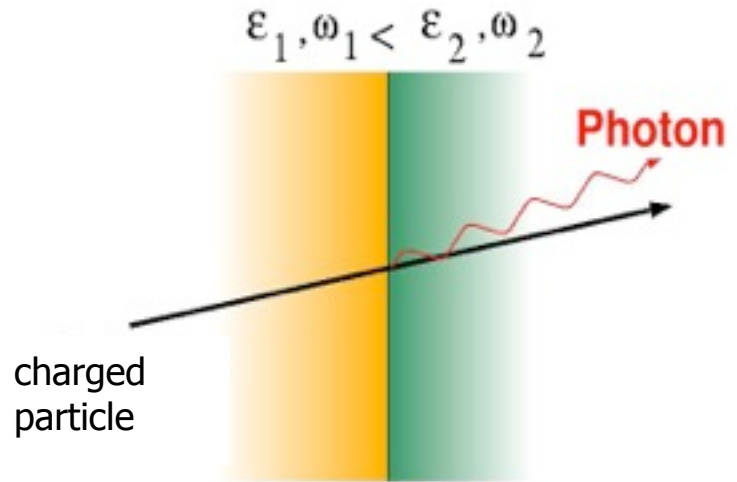
Electron Identification

Transition Radiation

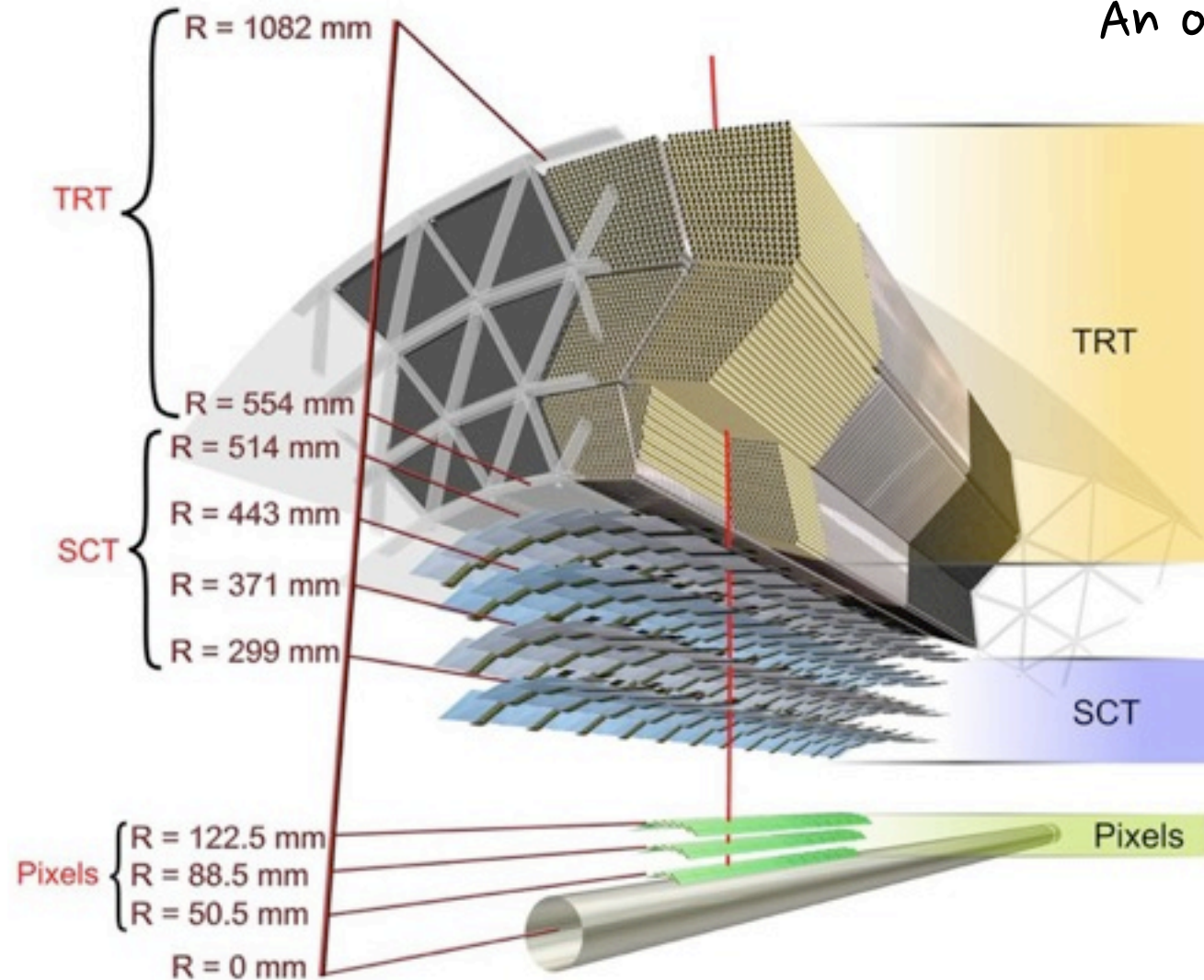
- Local speed of light in medium with refractive index n is $c' = c/n$
- If particle's relative velocity v/c' changes, particle will radiate photons
 - change of direction \rightarrow Synchrotron rad
 - change of speed \rightarrow Bremsstrahlung rad
 - change of index $n \rightarrow$ Transition rad
- Transition rad emitted whenever relativistic charged particle traverses border between two media having different dielectric constants $n = \sqrt{\epsilon}$
- Emitted energy proportional to boost γ of particle
 - can discriminate electrons from pions!
 - (also discriminate hadrons at high energy)

An other way...

Recall:



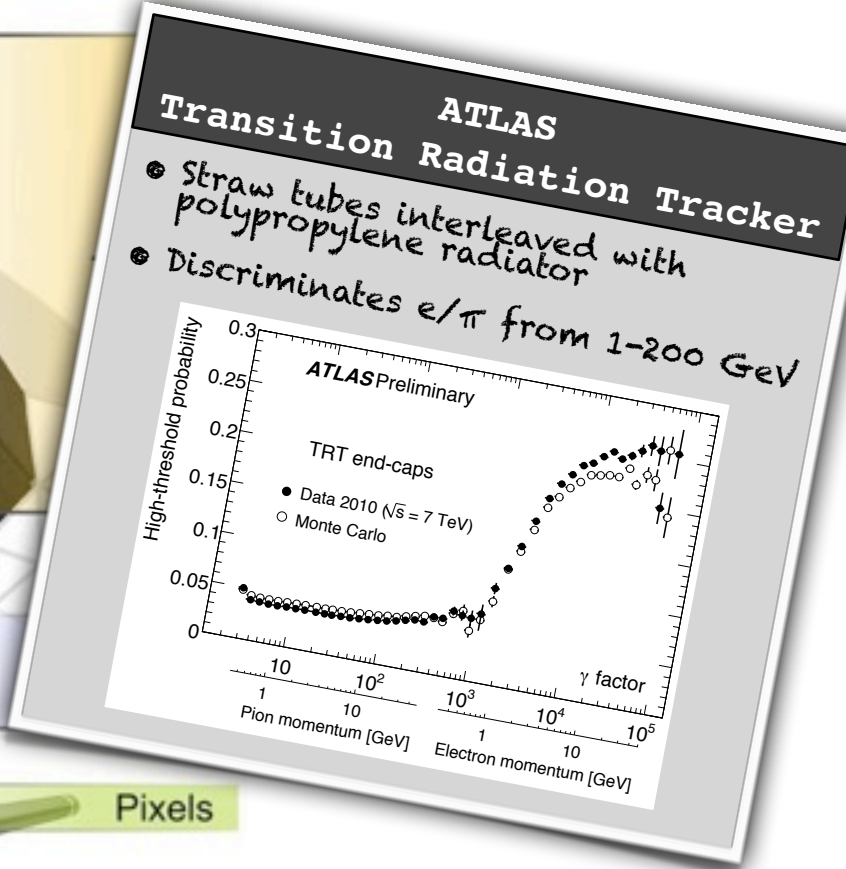
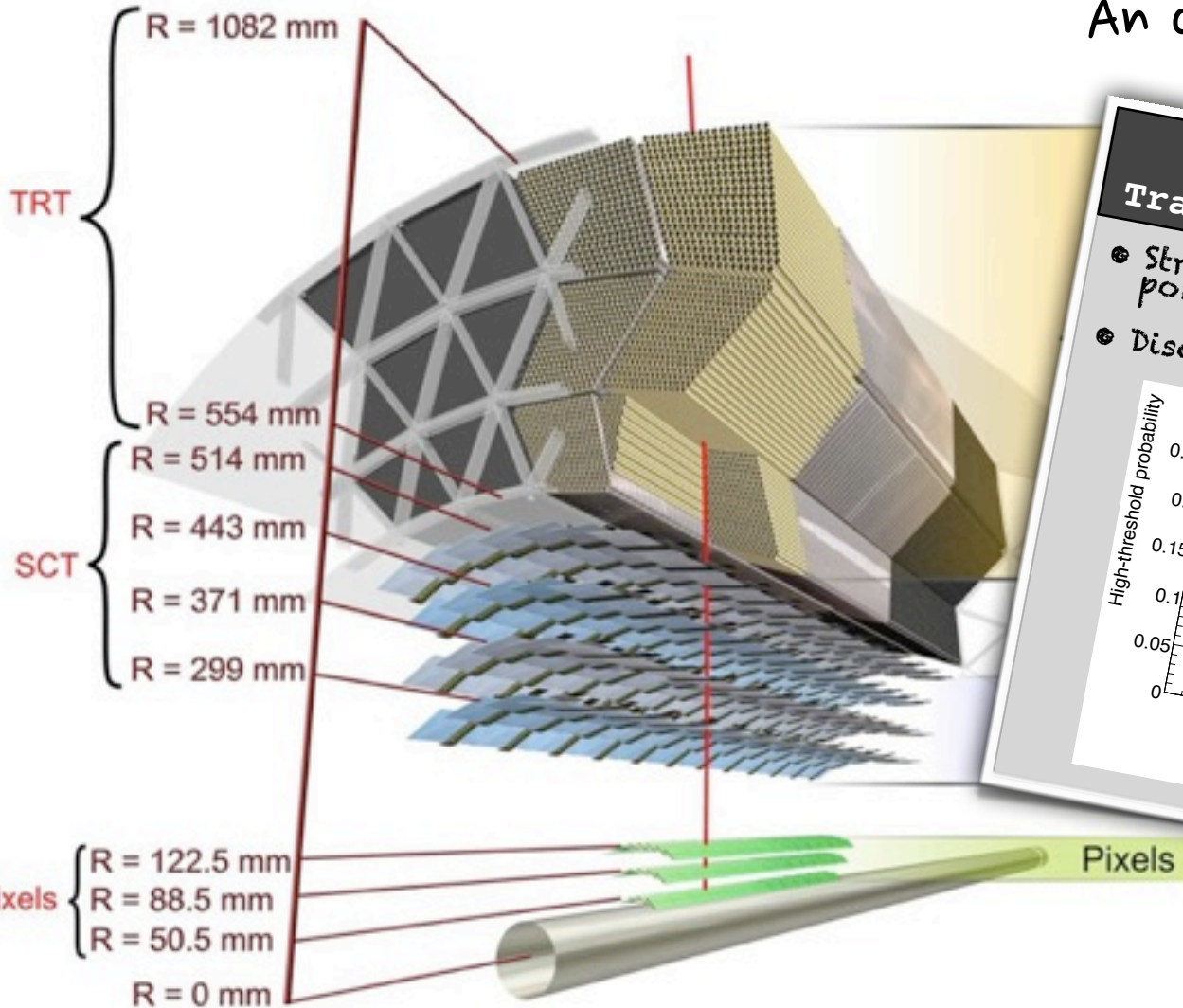
An other way...





Dedicated Electron ID

An other way...





Electron Identification

An other way...

EM calorimeter
clusters



Form "Super" clusters

- Different strategy to recover energy of brem photons and conversion pairs
- Seed clusters can be
 - nearest neighbor algorithm
 - rectangular window (brem recovery)
- SuperCluster
 - collect clusters in long ϕ -road
- Classify cluster as electron, photon, or converted photon
 - match cluster to track(s)
 - Use MVA (cluster shape, etc)
- Calculate energy & direction
 - Energy weighted sum
 - corrected for det. effects

An other way...

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EM calorimeter
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Seeding electron reconstruction

this way works well for isolated electrons; tends to miss non-isolated ones

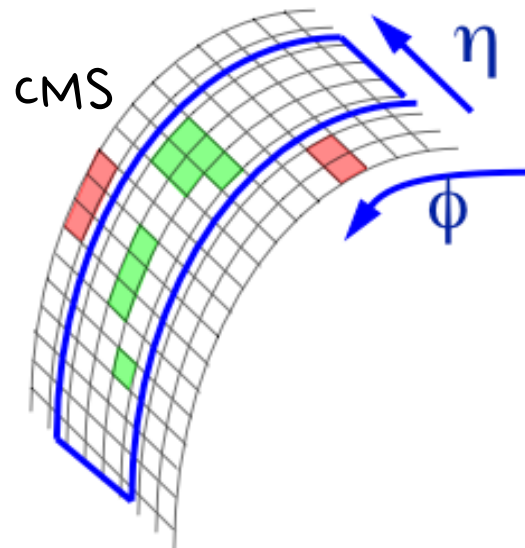


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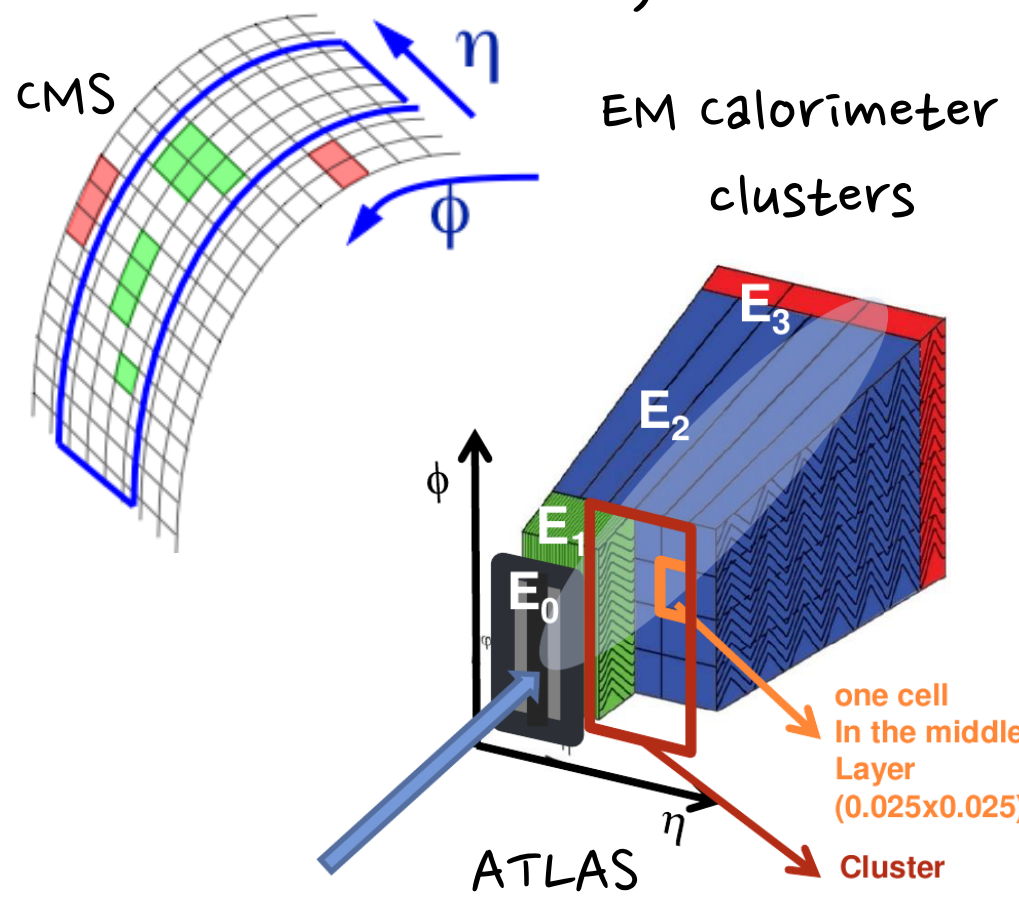


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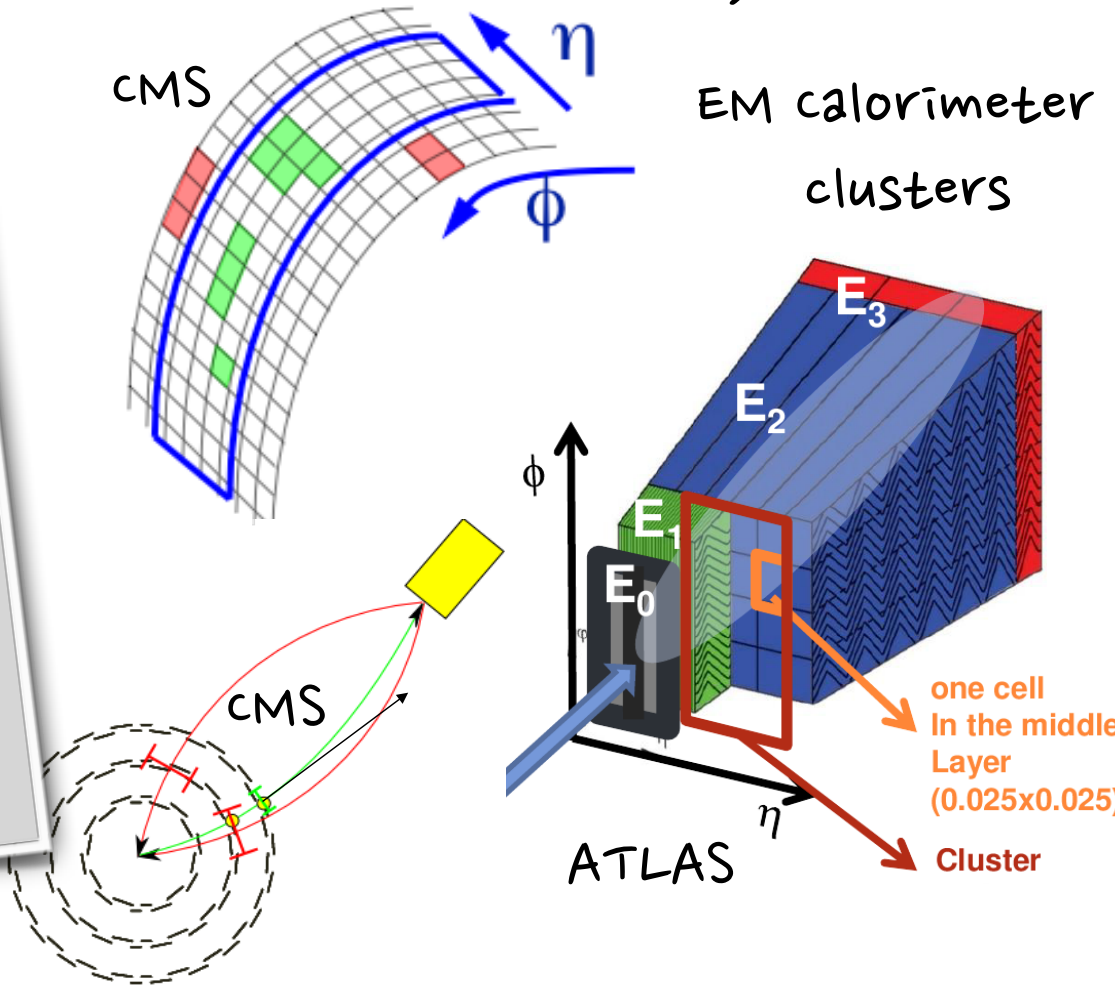


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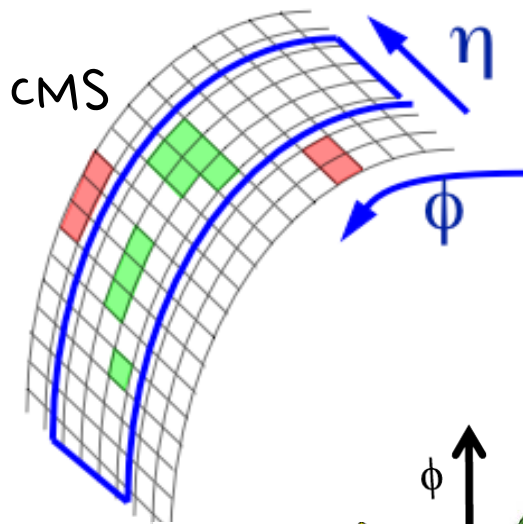
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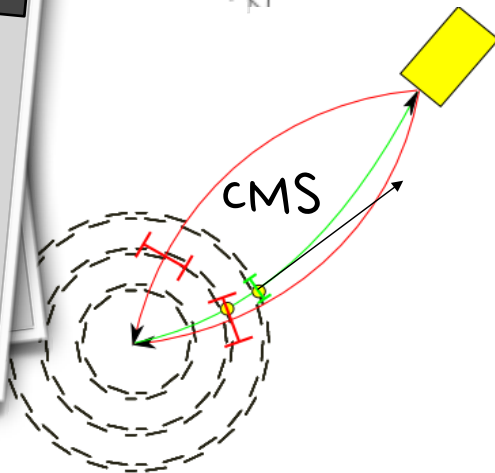
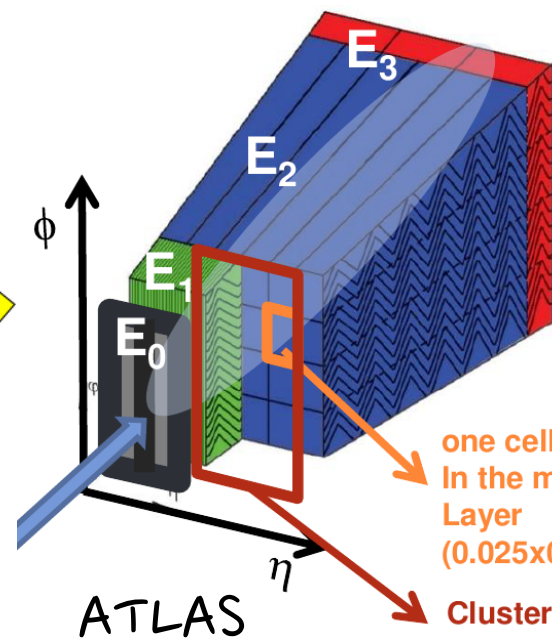
Discriminating electrons

- electrons vs hadrons
 - e.g. fraction of TRT hits, fraction of energy in layer 1, etc
- electrons vs conversions
 - e.g. number of hits in innermost

An other way...



EM calorimeter clusters

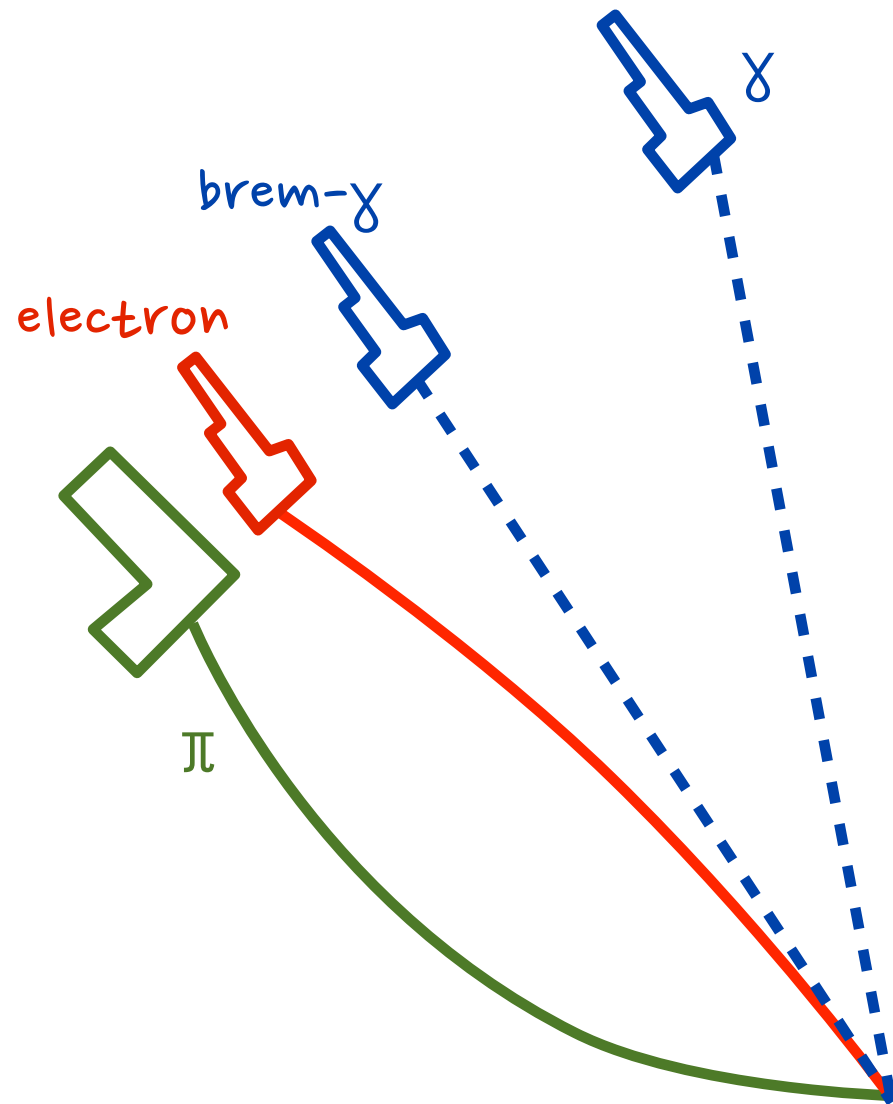


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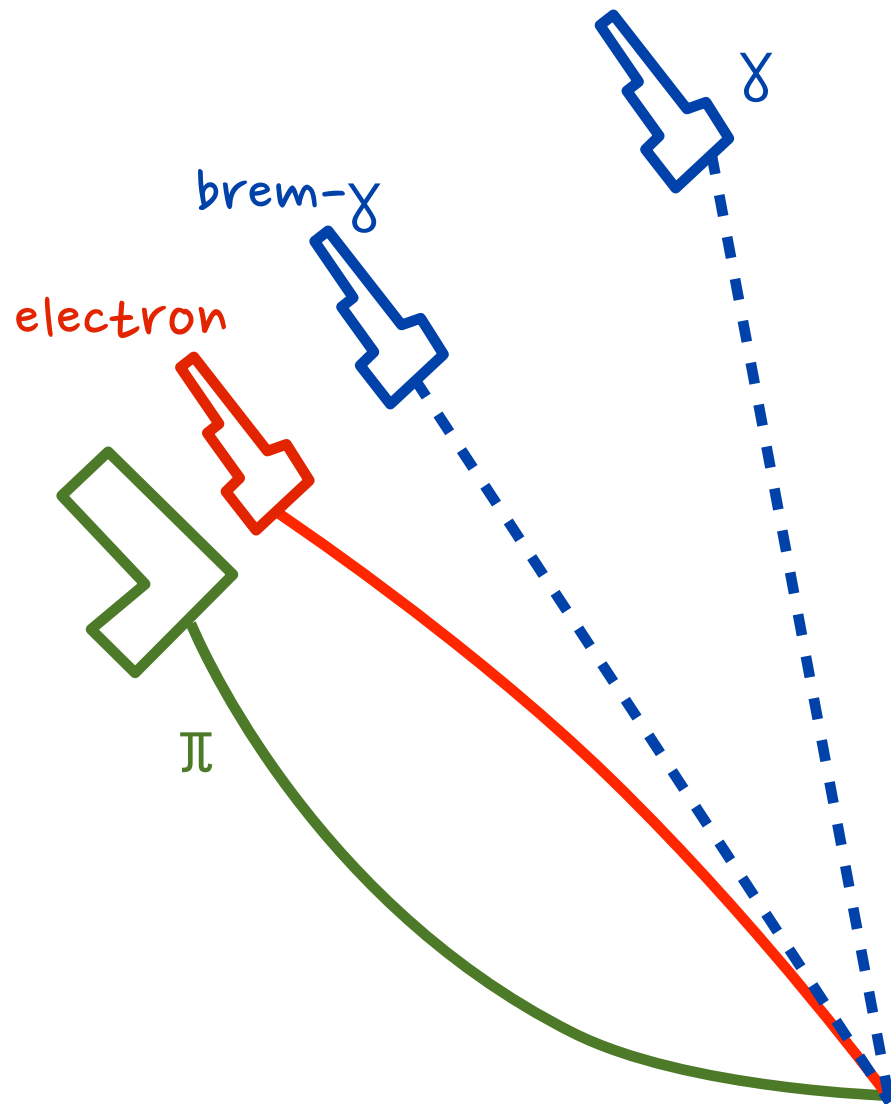




Electron Identification

Isolation from other particles

- Use simple $\Delta R \equiv \sqrt{(\Delta\phi^2 + \Delta\eta^2)} \approx 0.3$ angular distance
- between electron and surrounding particles





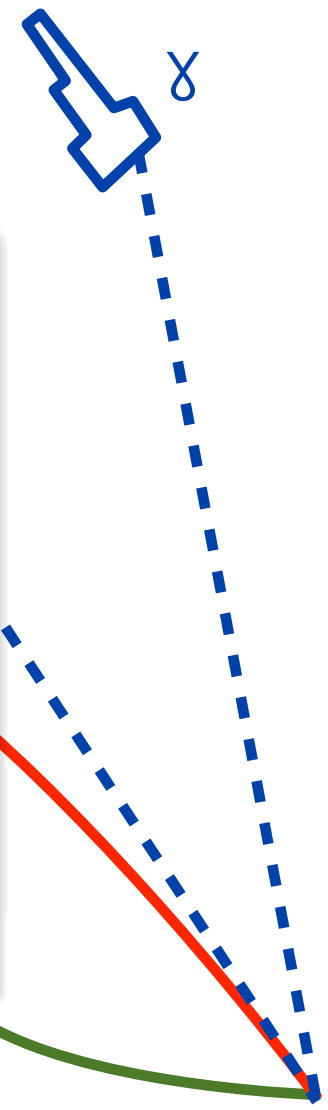
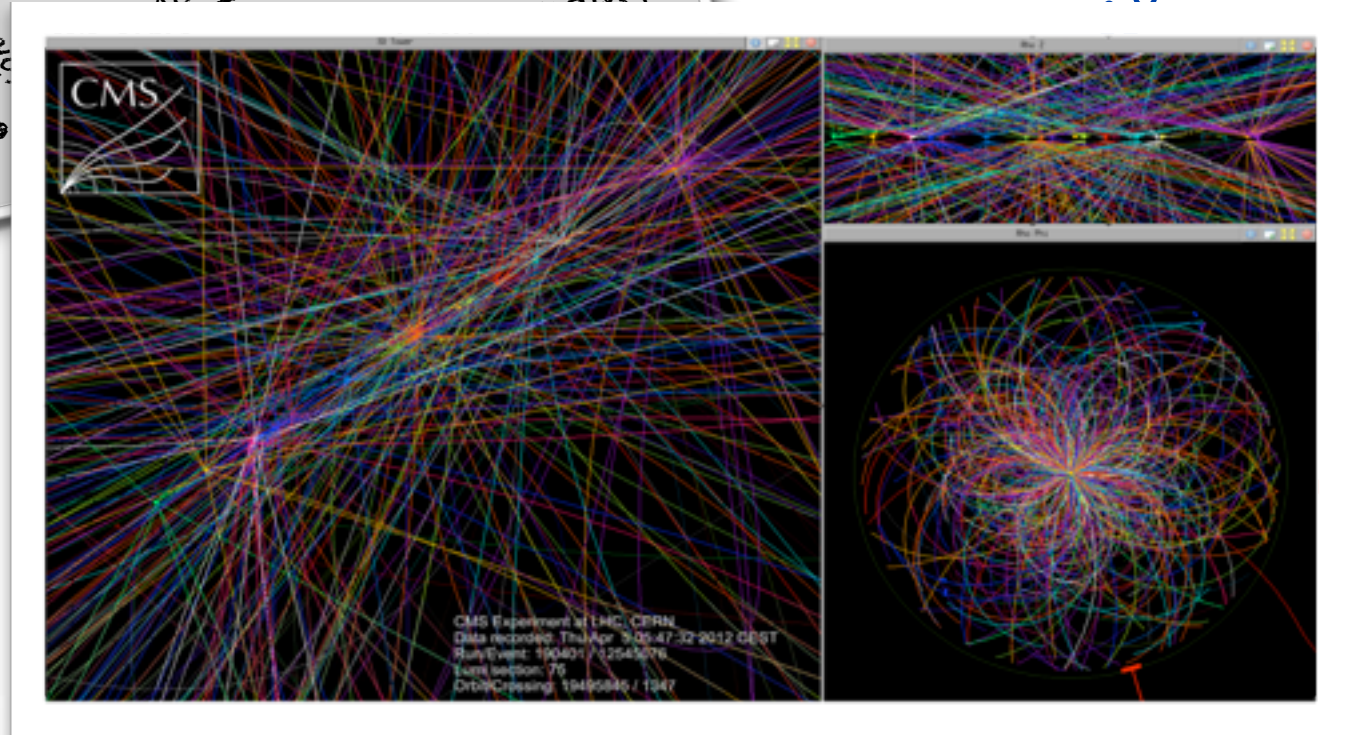
Electron Identification

Isolation from other particles

$$I = \sqrt{(\Delta\phi^2 + \Delta\eta^2)} \approx 0.3$$

- Use and

brem- γ

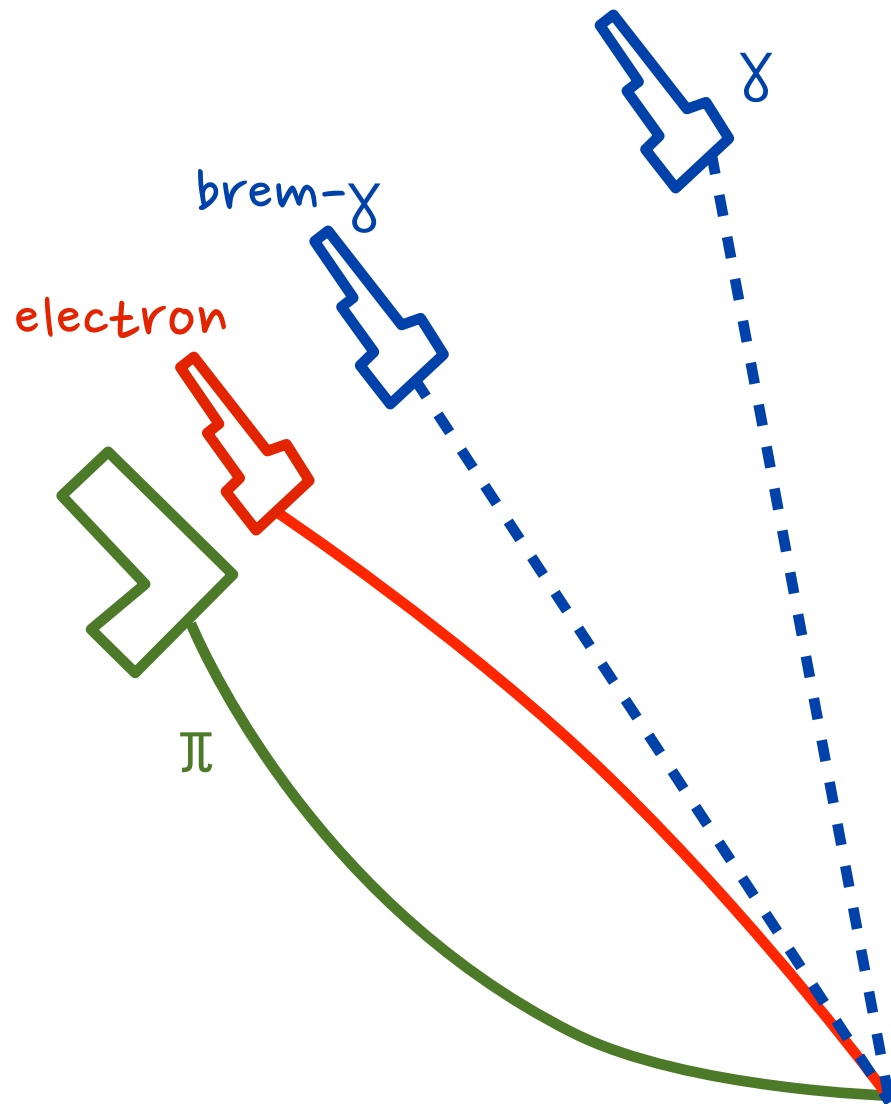




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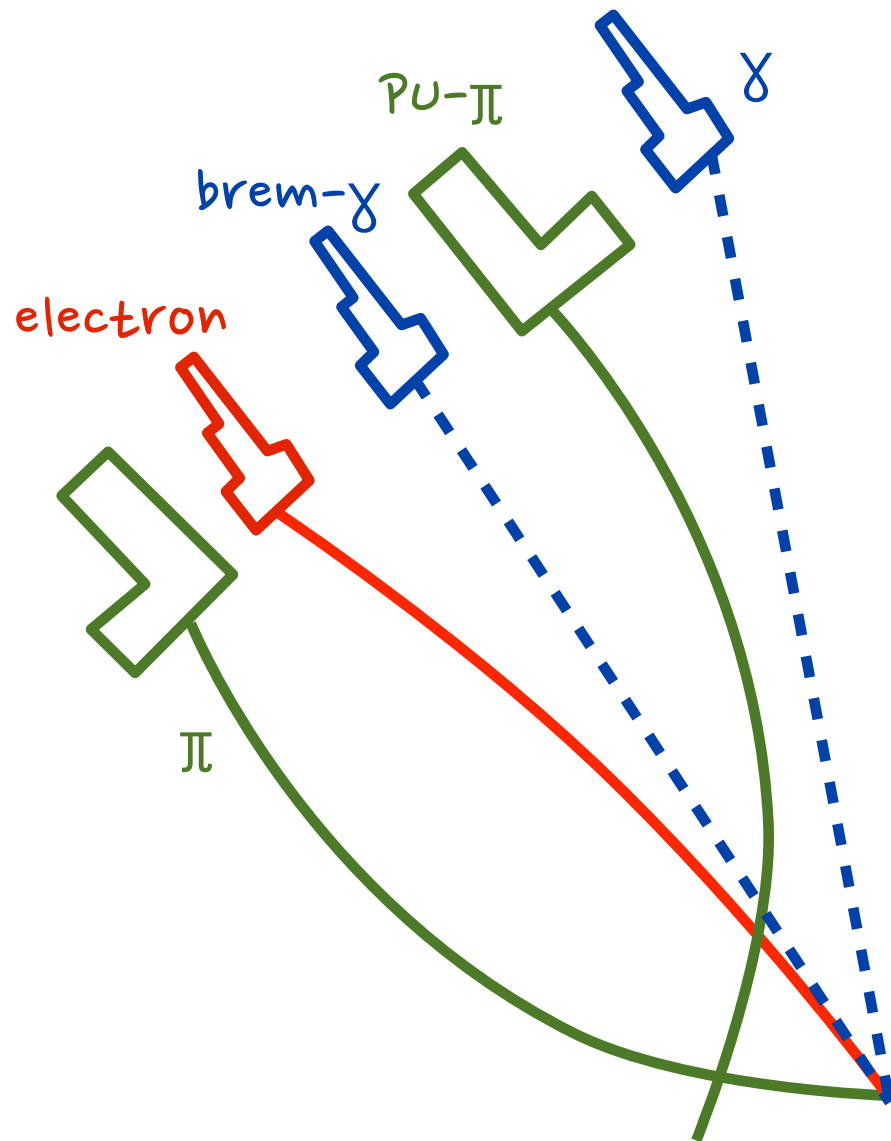




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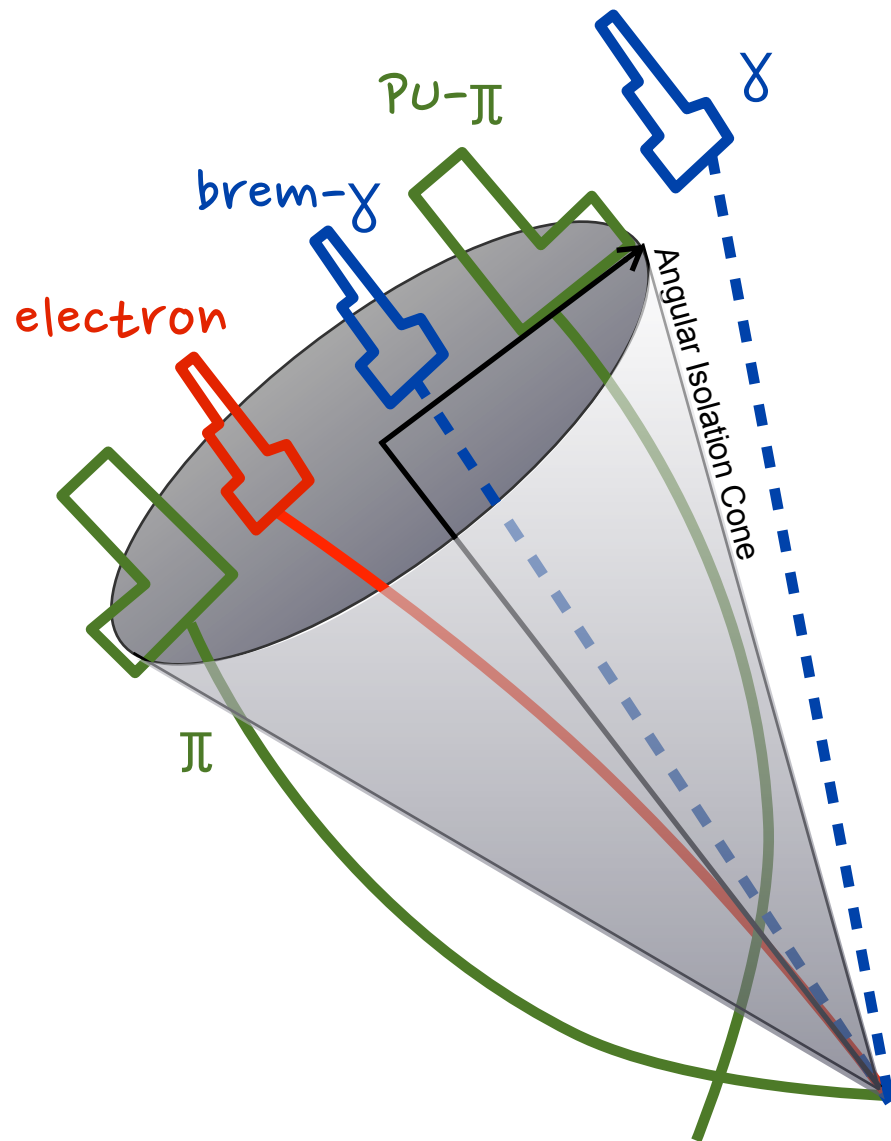




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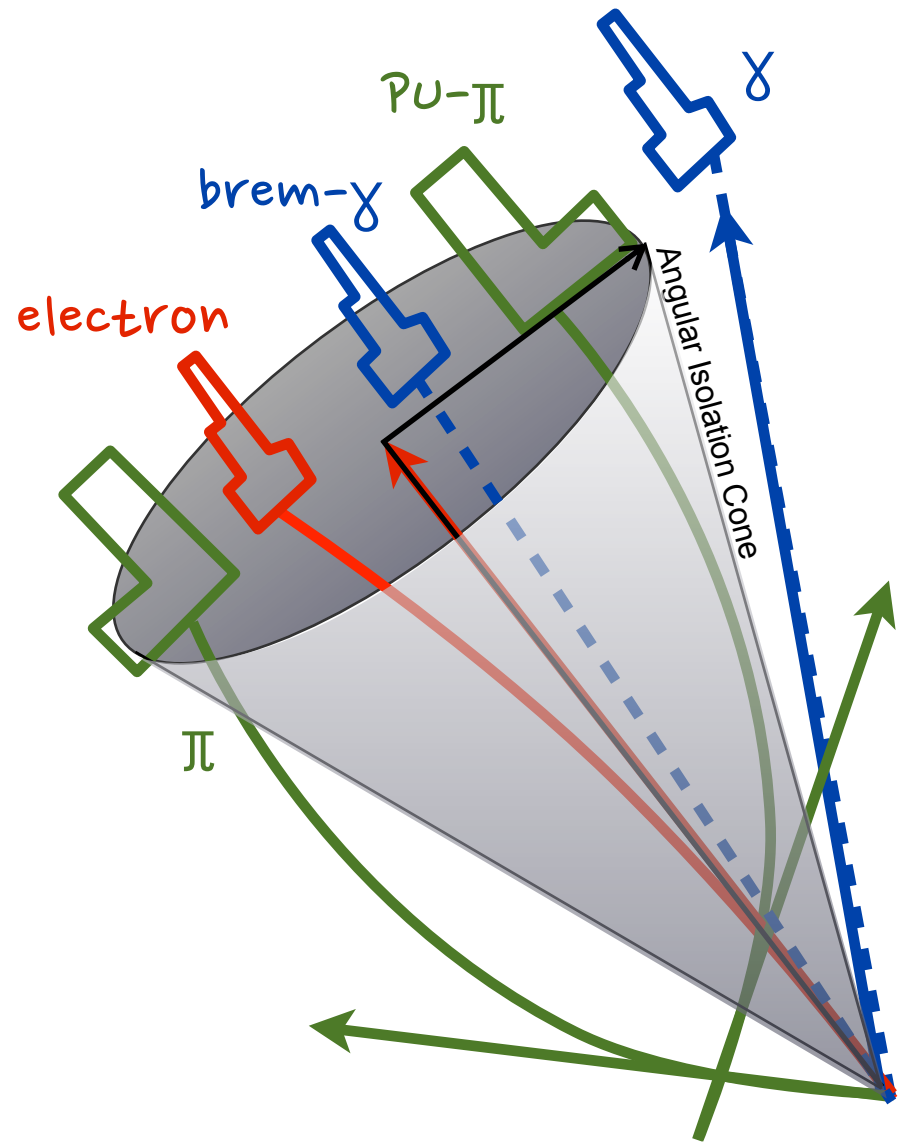




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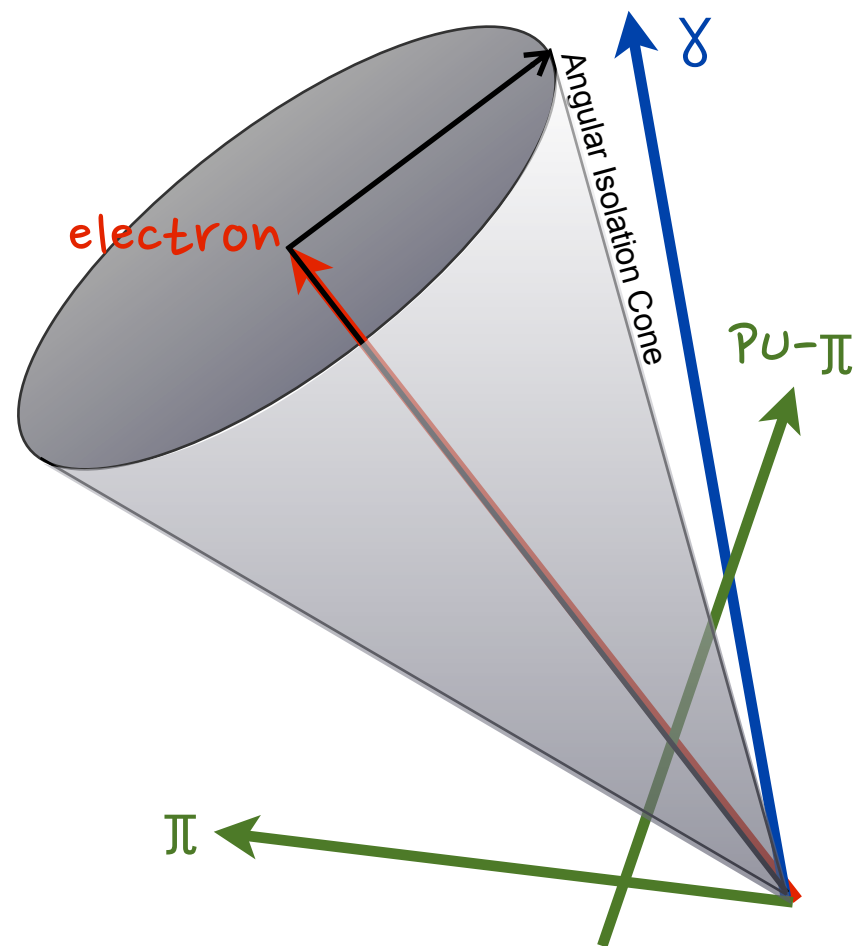




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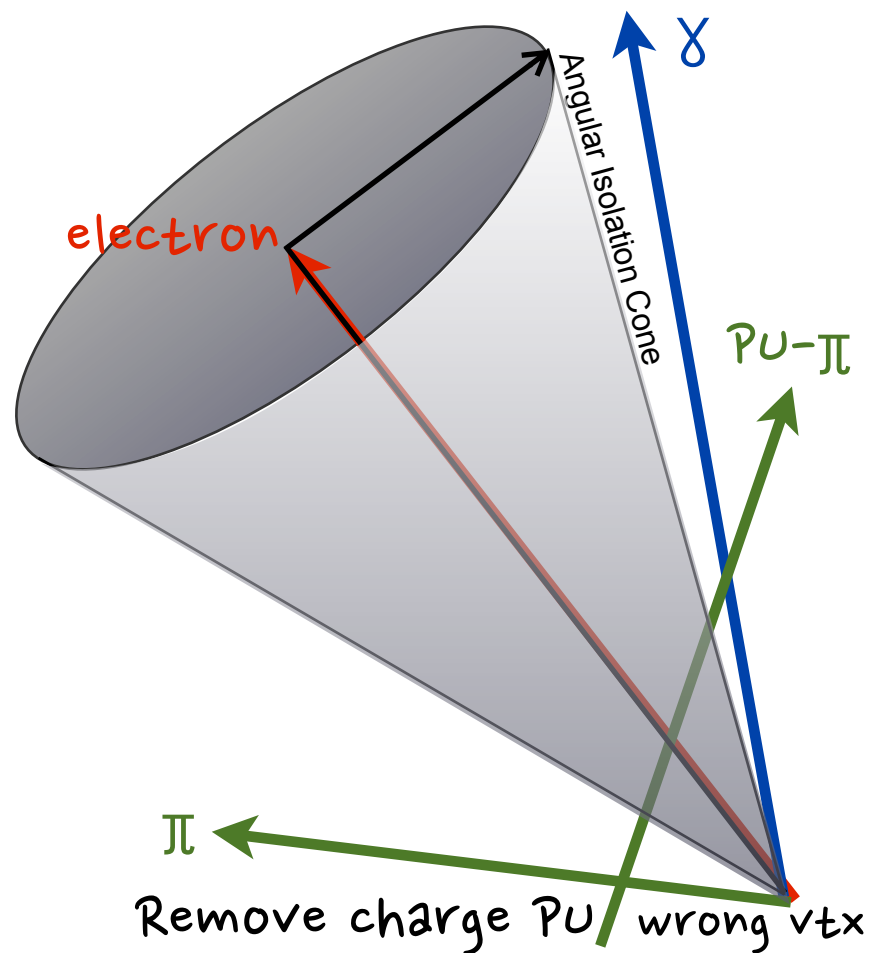




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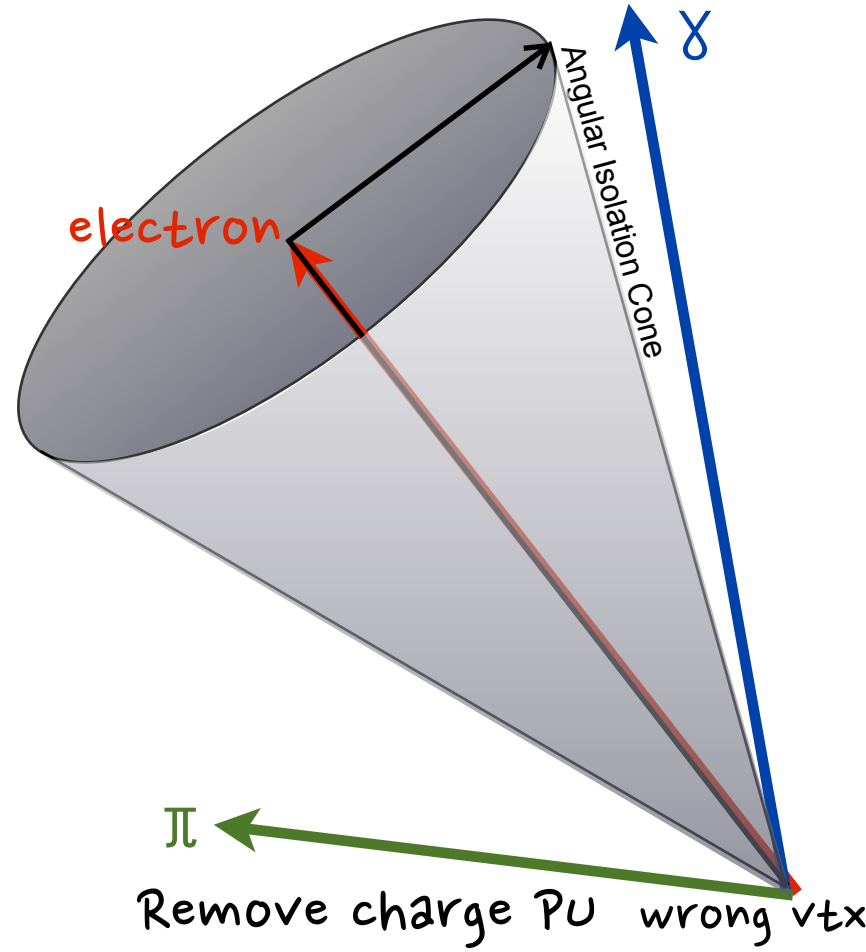




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Isolation from other particles

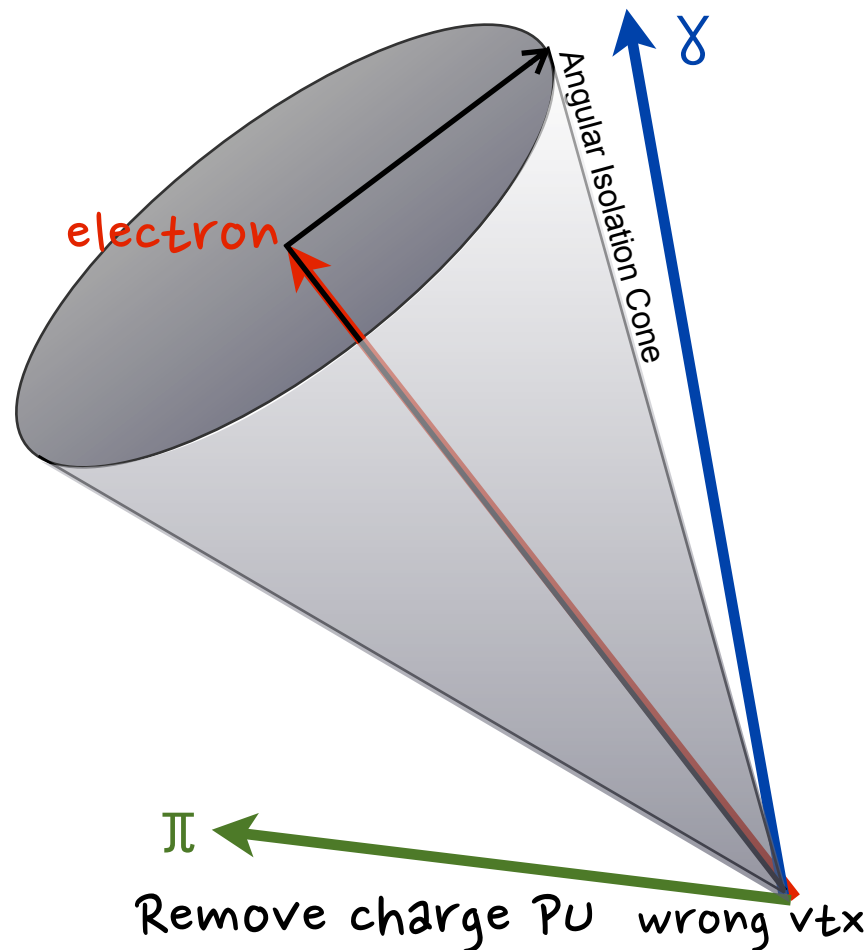
- Absolute energy in cone:

$$I = \sum p_T(h^\pm) + \sum E_T(\gamma) + \sum E_T(h^0)$$

- Relative energy in cone:

$$I = \frac{\sum p_T(h^\pm) + \sum E_T(\gamma) + \sum E_T(h^0)}{p_T(e^\pm)}$$

- Apply average correction for neutral PU
 - determine from event energy density
- Typically require $I < 15\%$
 - efficiency \approx independent of PU!





Electron Identification

Isolation from other particles

- Use simple $\Delta R = \sqrt{(\Delta\phi^2 + \Delta\eta^2)} \approx 0.3$ angular distance between electron and surrounding particles

Isolation from other particles

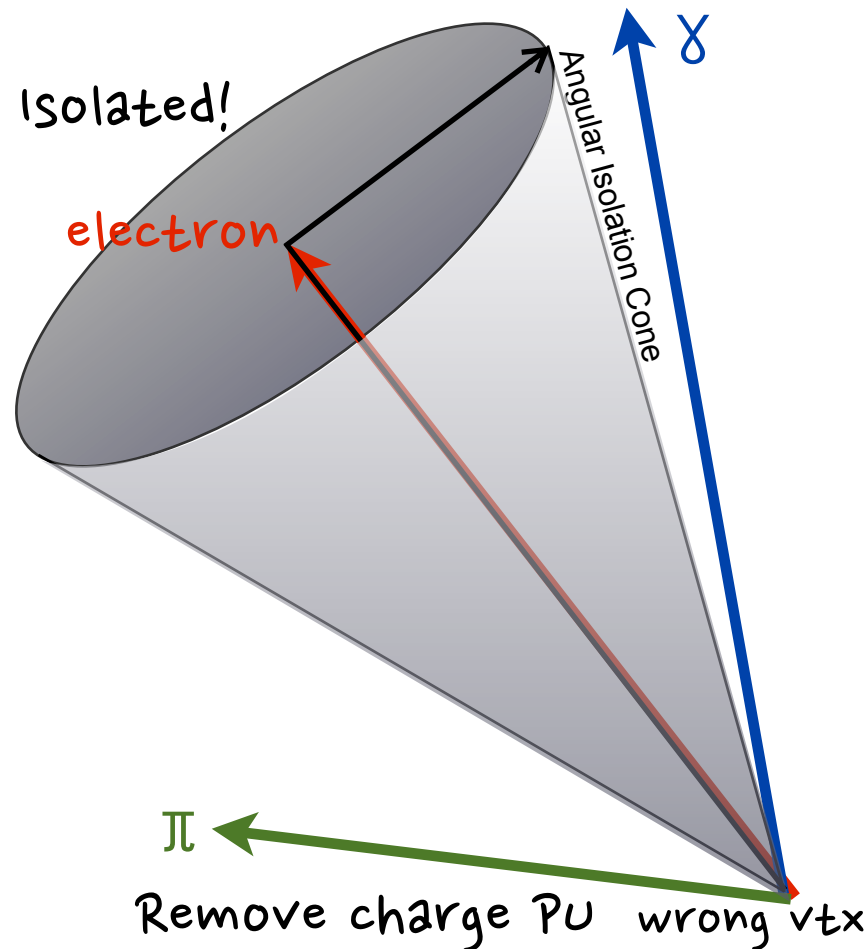
- Absolute energy in cone:

$$I = \sum p_T(h^\pm) + \sum E_T(\gamma) + \sum E_T(h^0)$$

- Relative energy in cone:

$$I = \frac{\sum p_T(h^\pm) + \sum E_T(\gamma) + \sum E_T(h^0)}{p_T(e^\pm)}$$

- Apply average correction for neutral PU
 - determine from event energy density
- Typically require $I < 15\%$
 - efficiency \approx independent of PU!





Electron Identification

Isolation from other particles

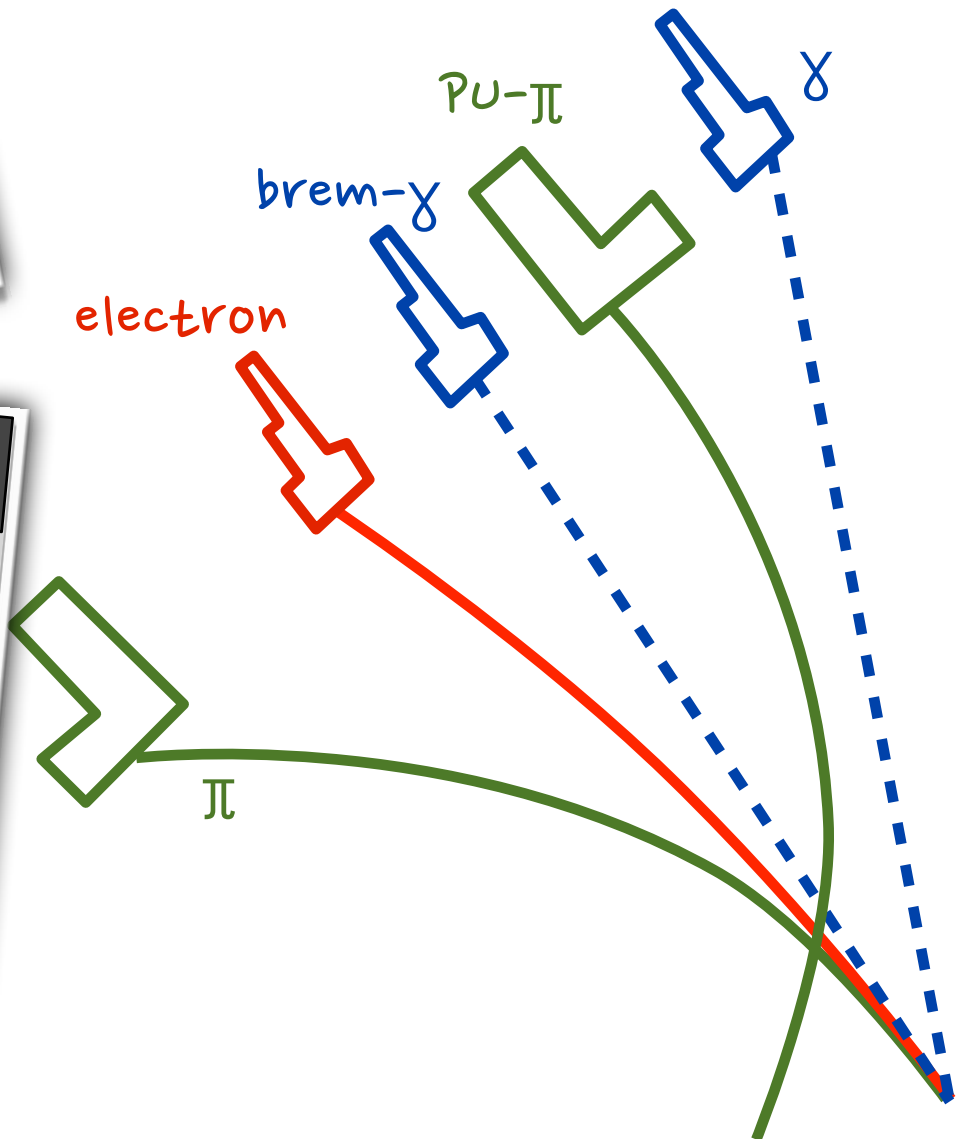
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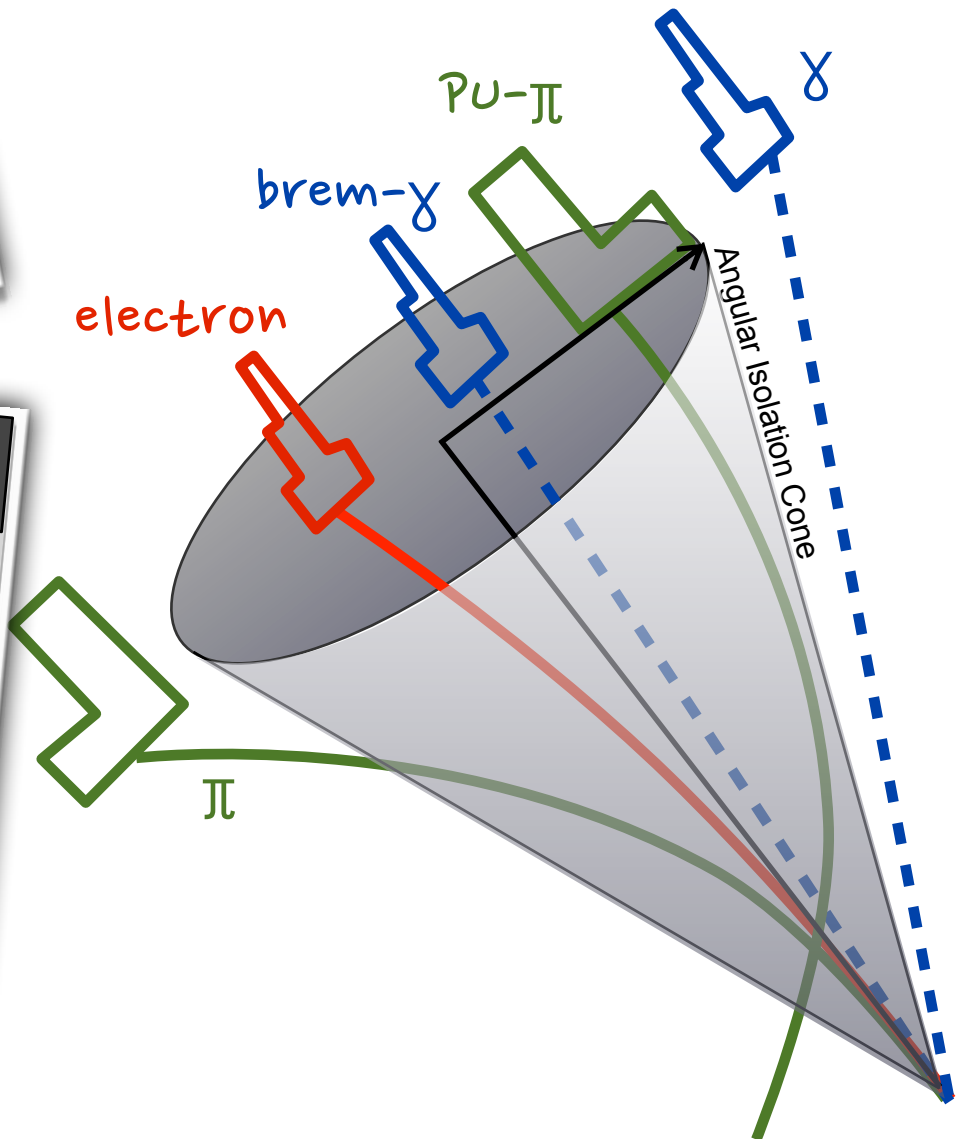
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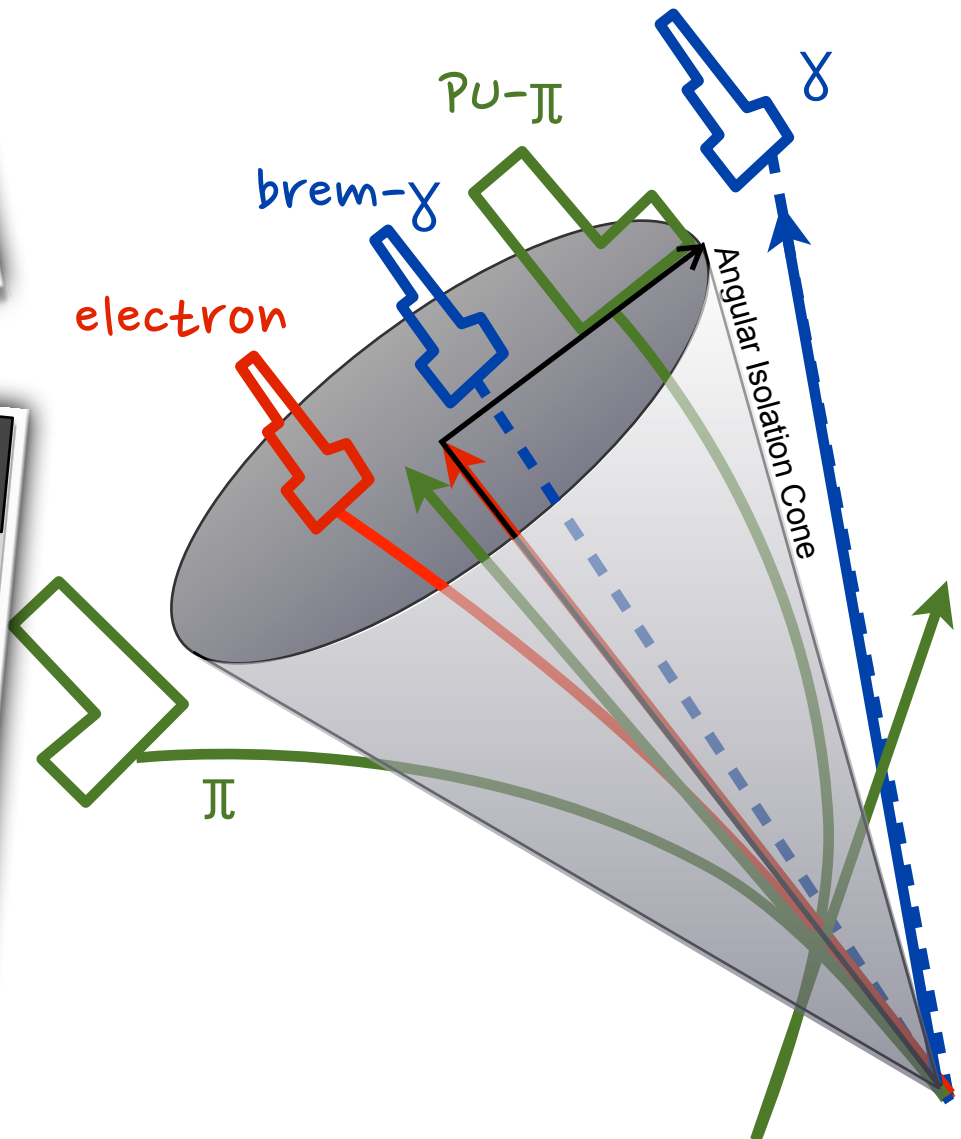
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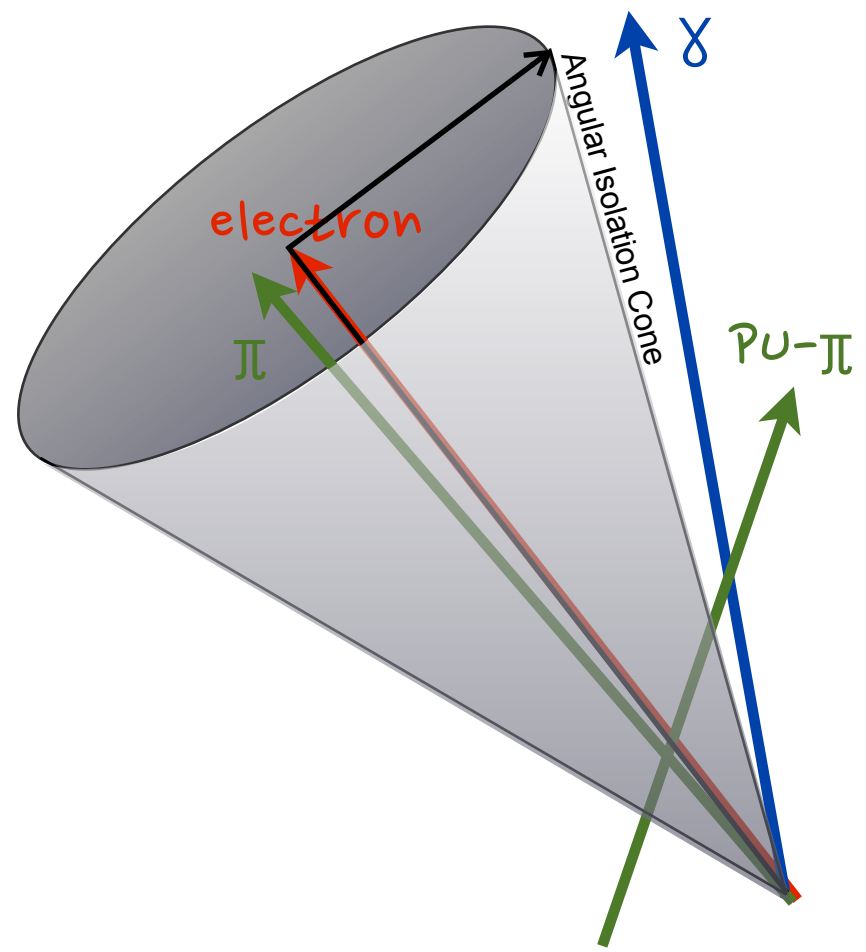
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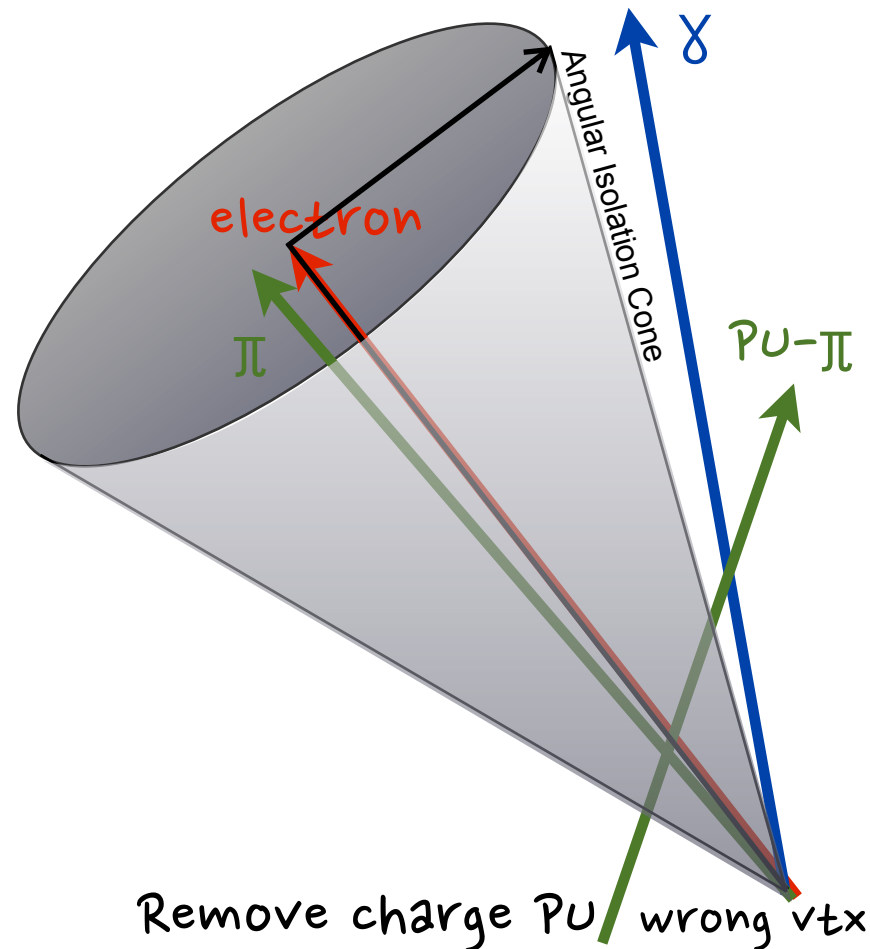
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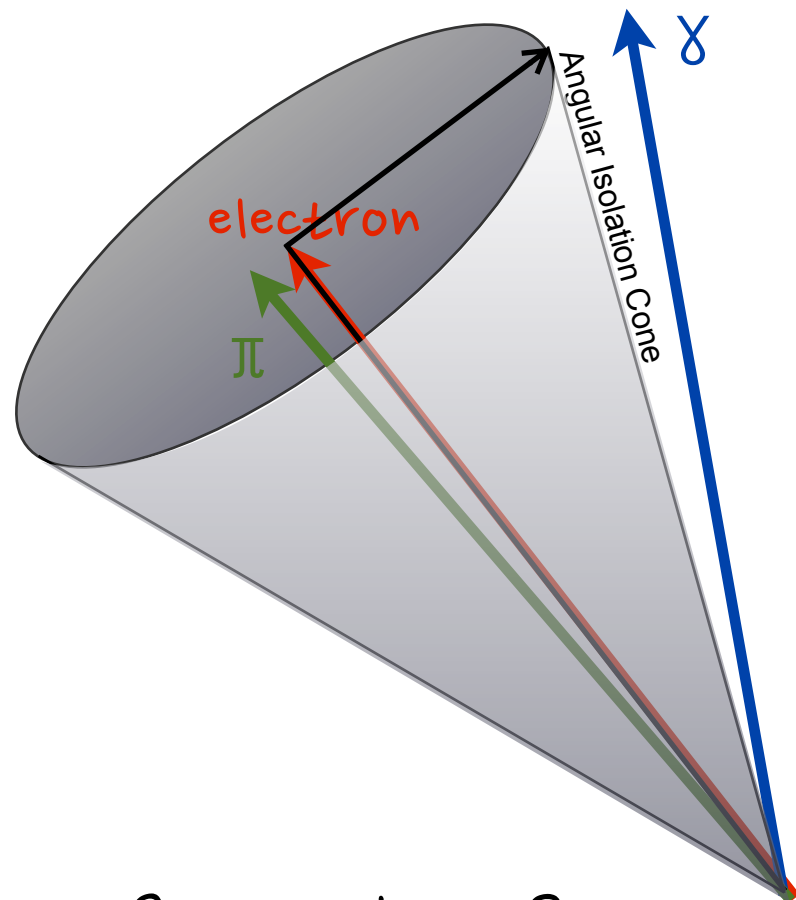
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Remove charge PU wrong vtx



Electron Identification

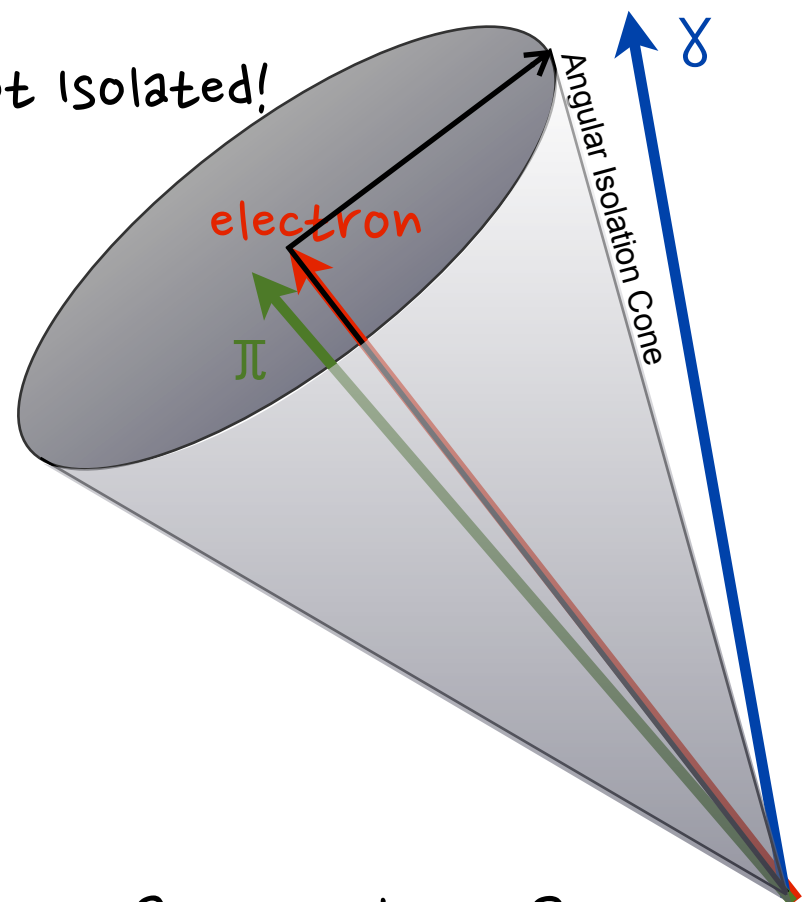
Isolation from other particles

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Not Isolated!



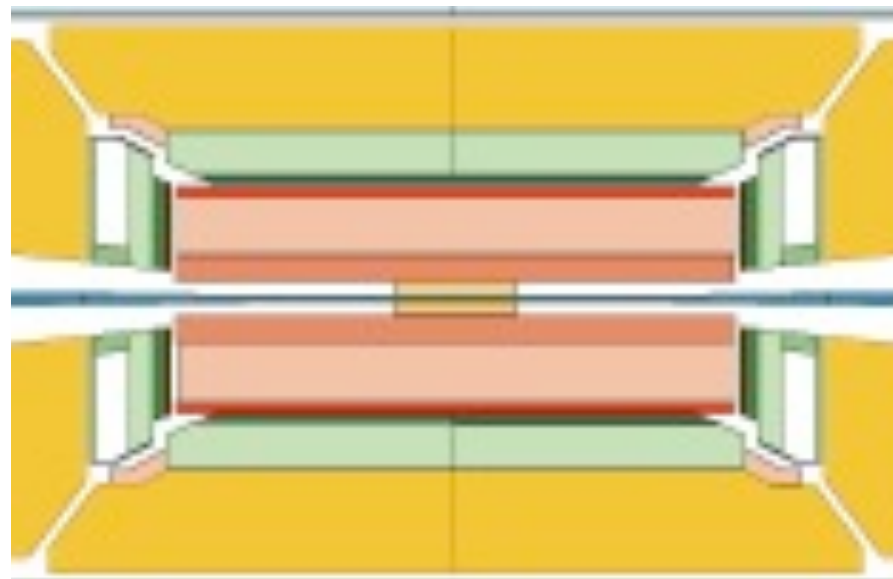
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Tag and Probe Method

- Tag: Select a well reconstructed, tight selected electron from a sample of single electron triggered events
- Probe: Select an appropriate object (e.g. ECAL SuperCluster)
- Select criteria to define passing probes
 - depends on what you want to study!
- Compute invariant mass of Tag and Probe, separately
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- perform fit around Z^0 or J/ψ mass peak
- Determine yields as a function of electron p_T

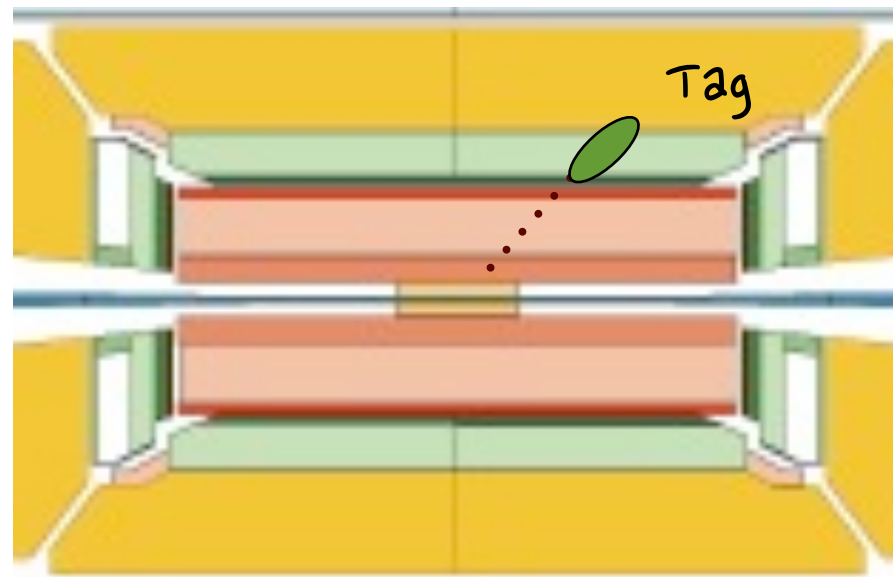
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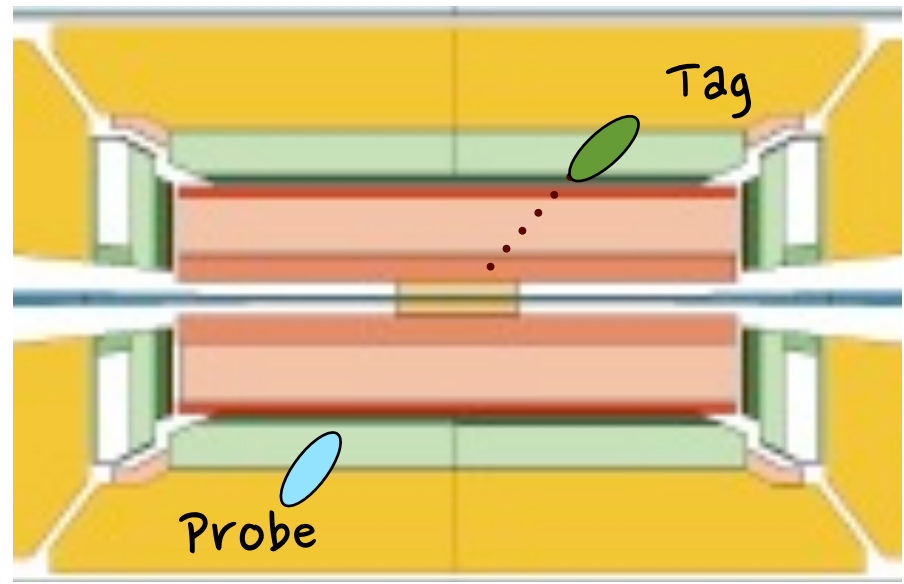


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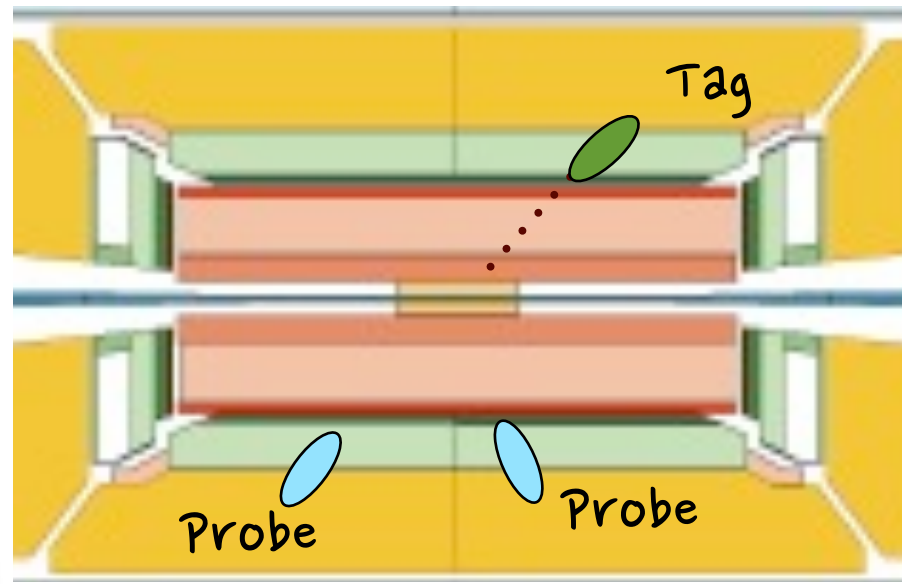


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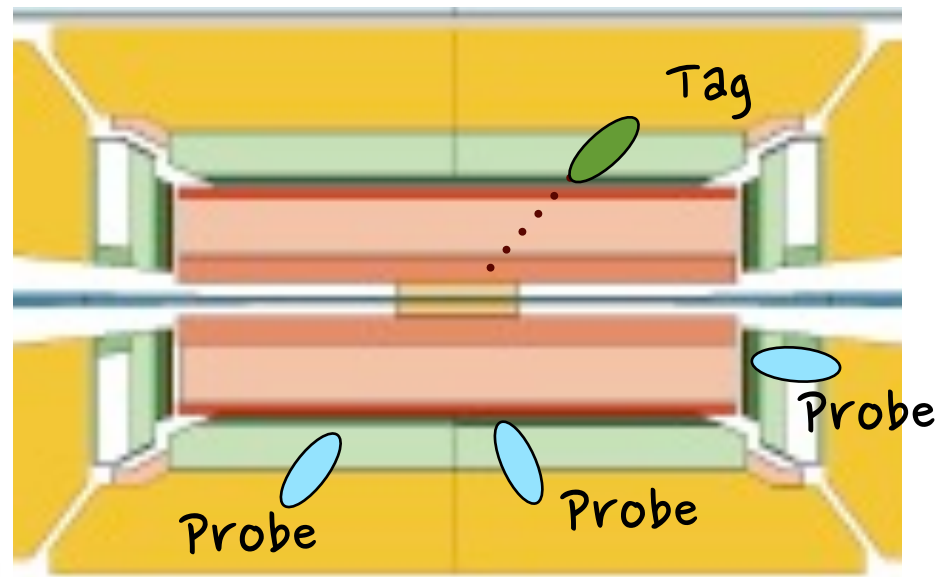
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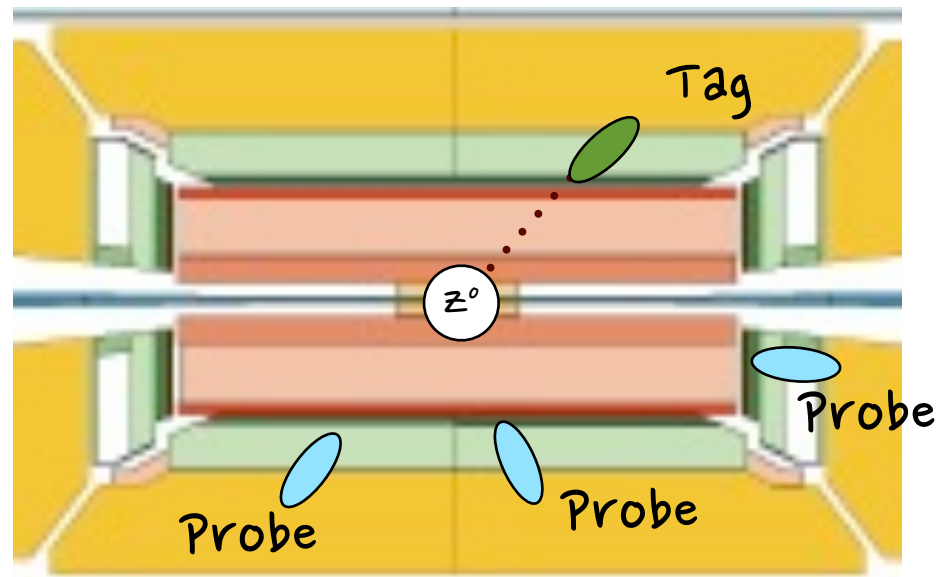
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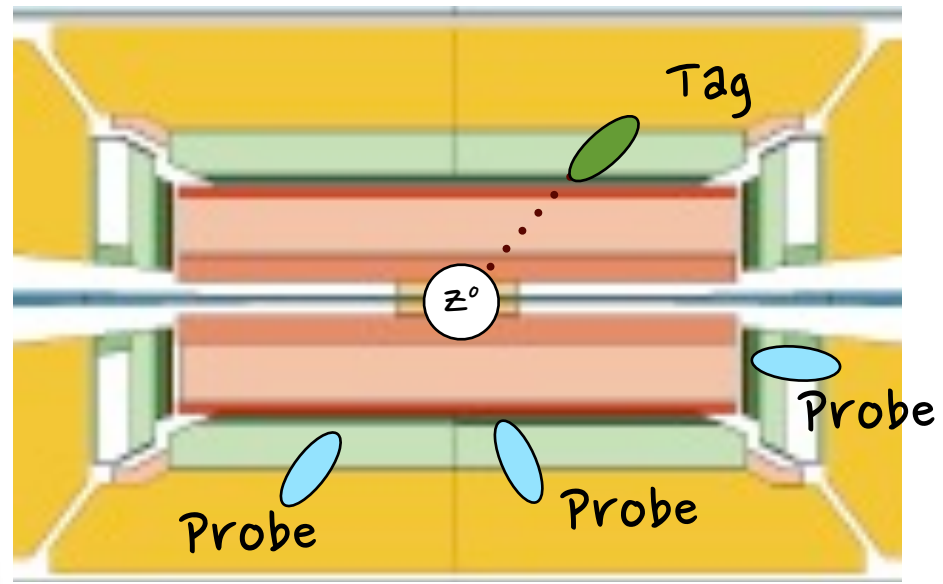
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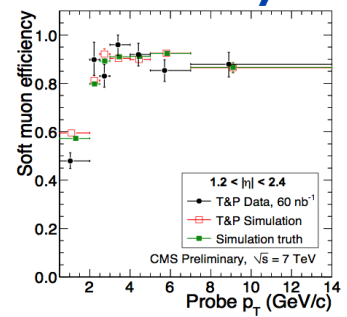


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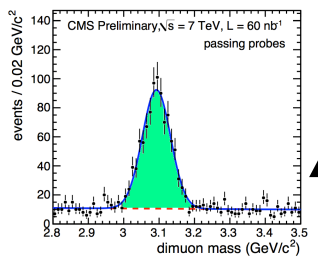
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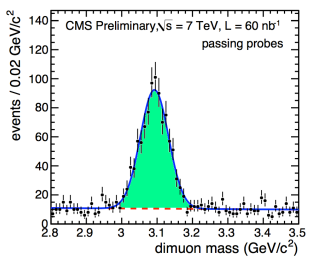
Efficiency



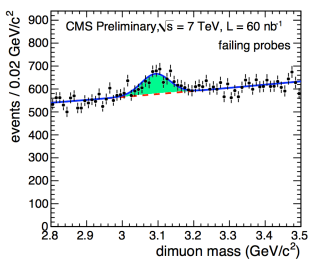
Passing probes



Passing probes

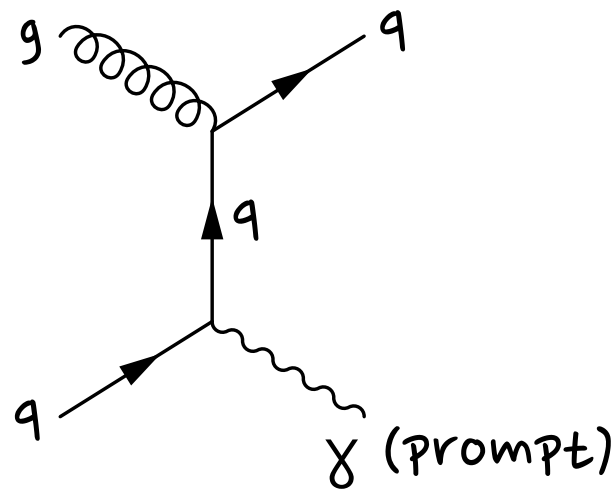


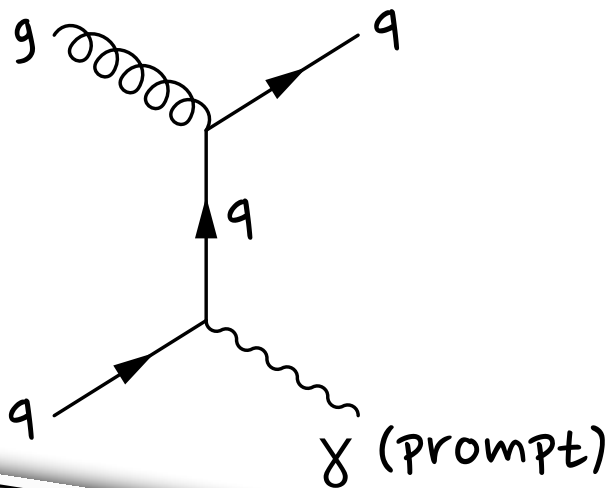
Failing probes



$$= \frac{\text{Passing probes}}{\text{Passing probes} + \text{Failing probes}}$$

Prompt Photon ID





Challenges

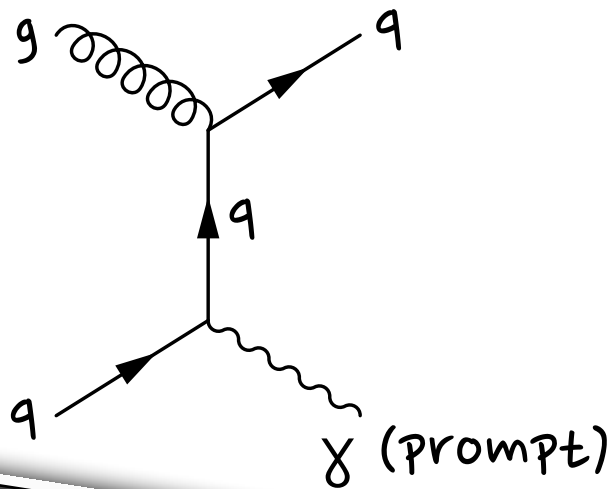
- Unconverted photons have no redundancy: calorimeter only!
- Not many handles
 - Shower shape consistent with EM
 - use fine segmentation
- Isolation!
 - can be particle-based or detector based
 - can be absolute or relative
- plus special tricks for cosmics/beam halo

Date: 14.08.2012

Cavanaugh, HCPSS 2012



Prompt Photon ID



Keys to prompt γ ID

- 3 important backgrounds
 - jets, electrons, cosmics / beam halo (muon brems)
- Discriminate between prompt γ and
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- Note: none of the γ 's above are fake!

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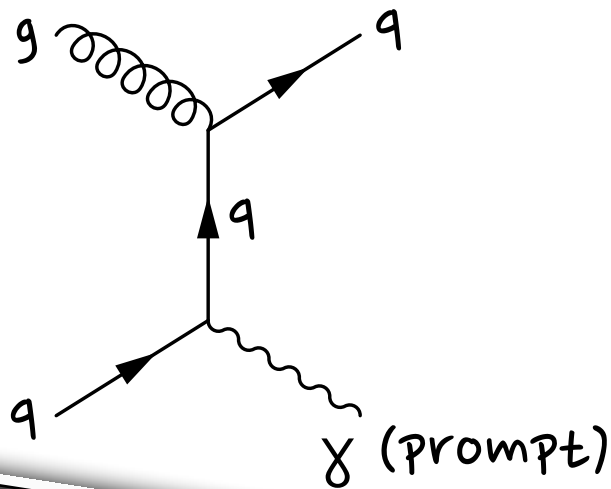
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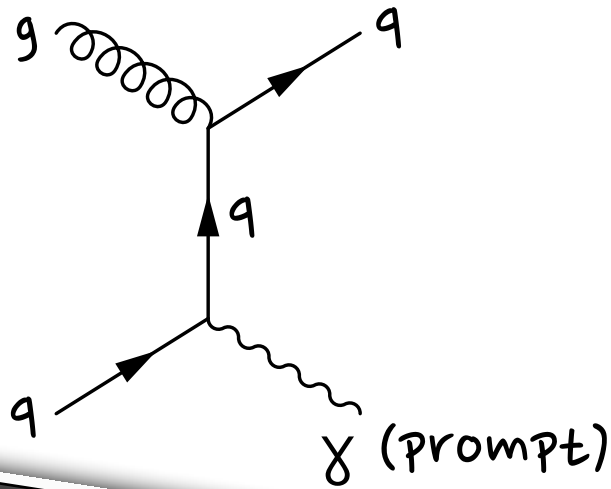
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Different jet backgrounds

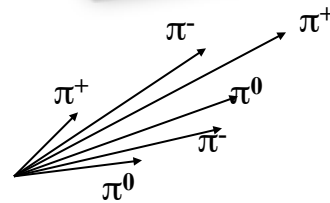


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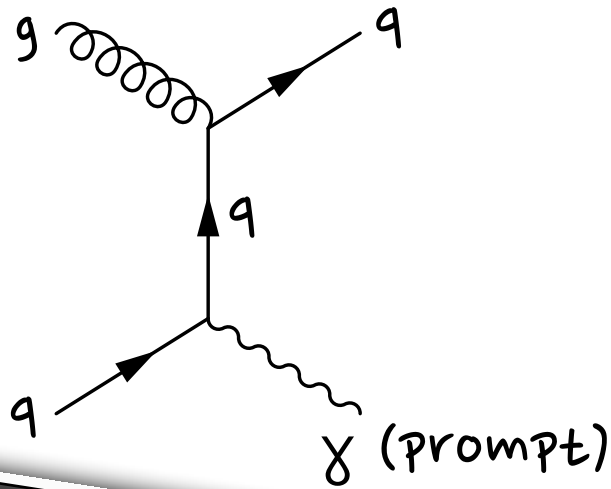
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normal jet

Different jet backgrounds

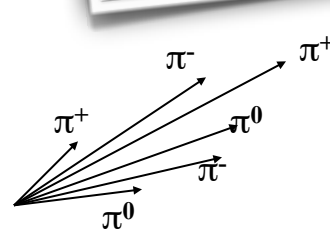


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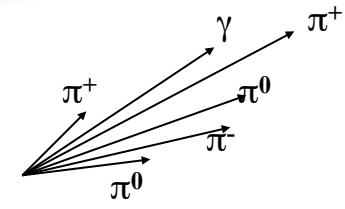
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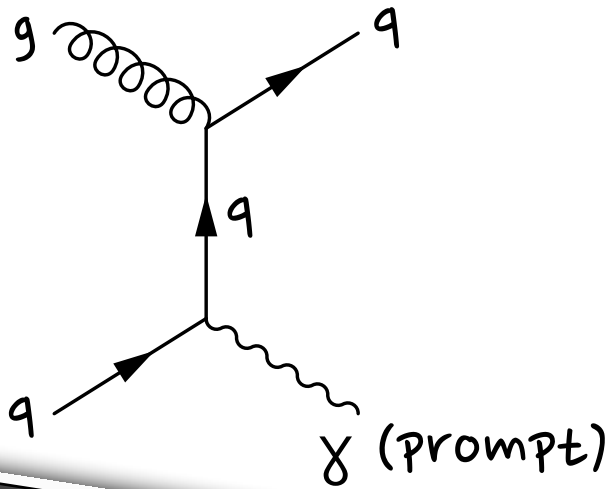


normal jet



fragmentation γ

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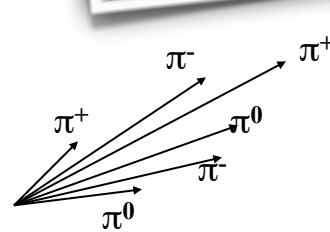


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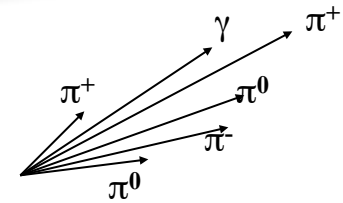
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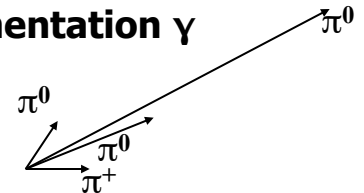


normal jet



fragmentation γ

Different jet backgrounds



fluctuated jet: most energy is carried by π^0



Prompt Photon ID



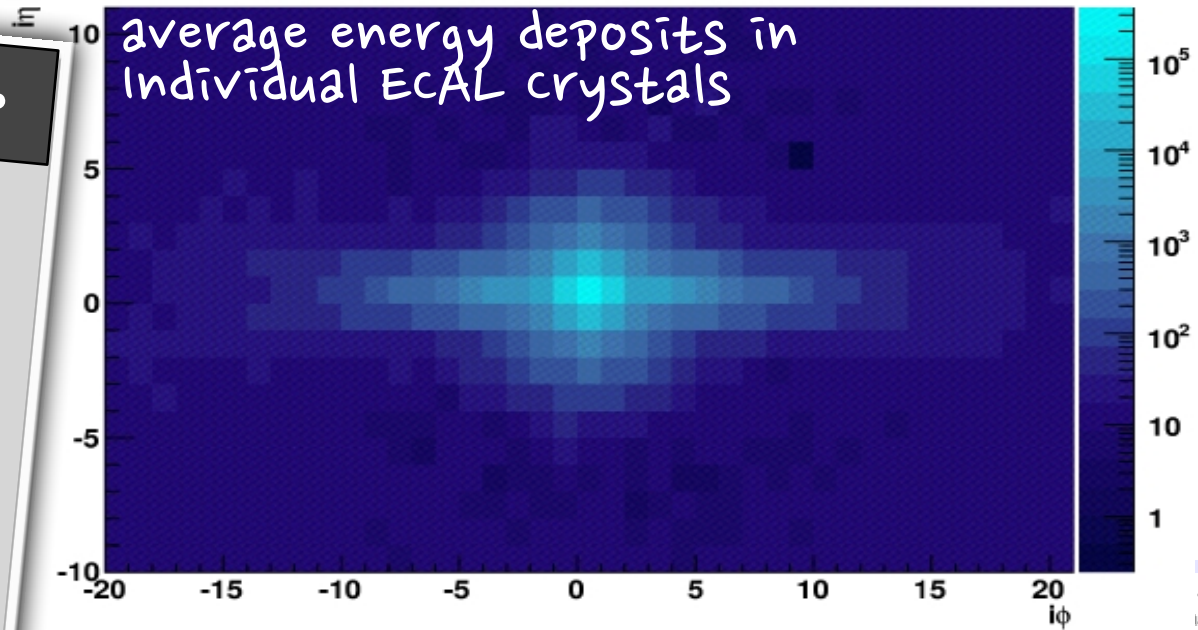
Prompt Photon ID

What do γ 's look like?

- Depends on the experiment details!
- CMS ECAL lives inside the B-field
- Electrons in ECAL shower
 - spread radially out from impact point
 - also bend in ϕ
- Photons often pair convert
 - bend in ϕ
- Effect: "Jurasic symbol"

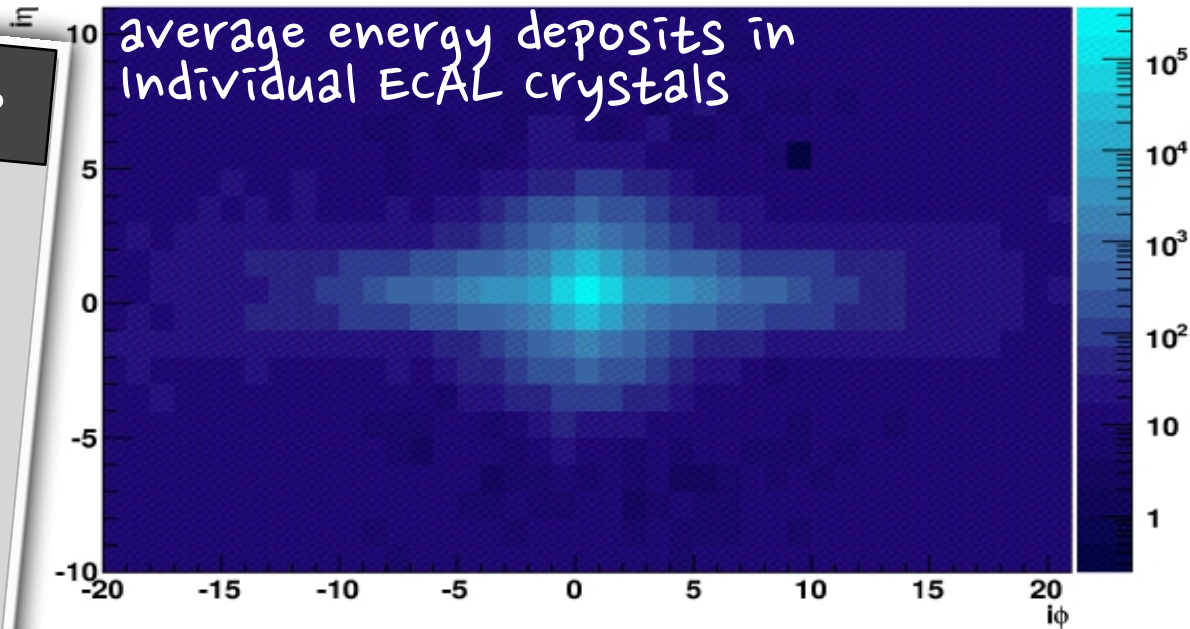
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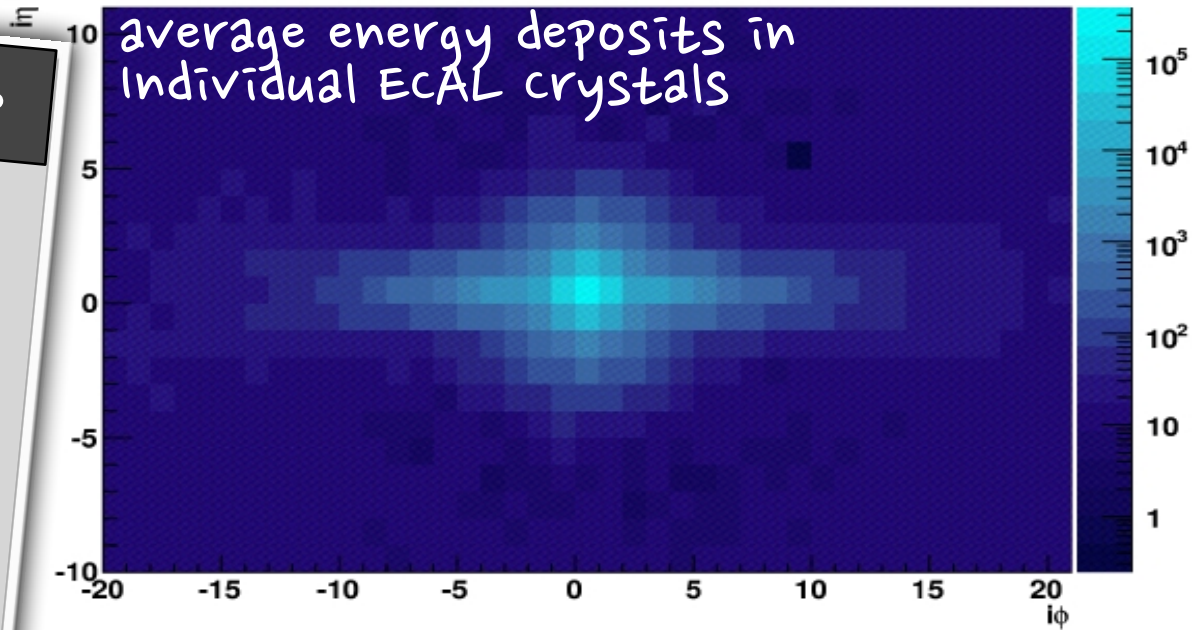


Clusters & Superclusters

- Cluster nearby cells together
- Ensure all energy is gathered
 - due to converting photons (or electron brem)
 - make clusters of clusters "superclusters" along ϕ
- Not particle-flow based algorithm!

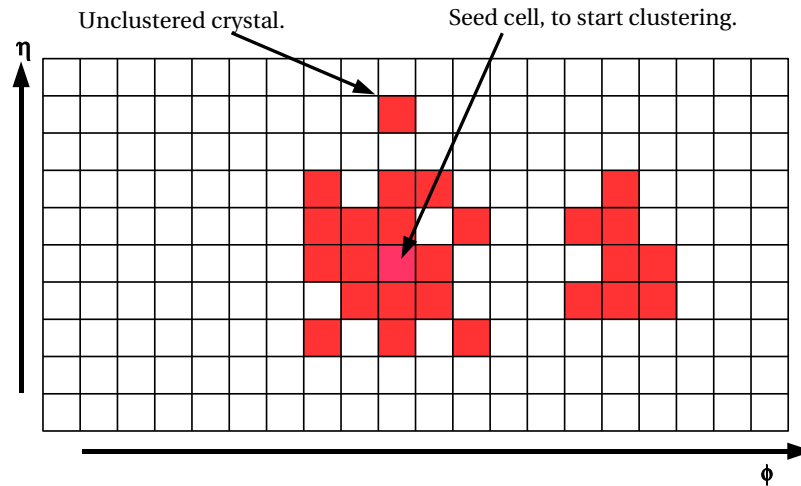
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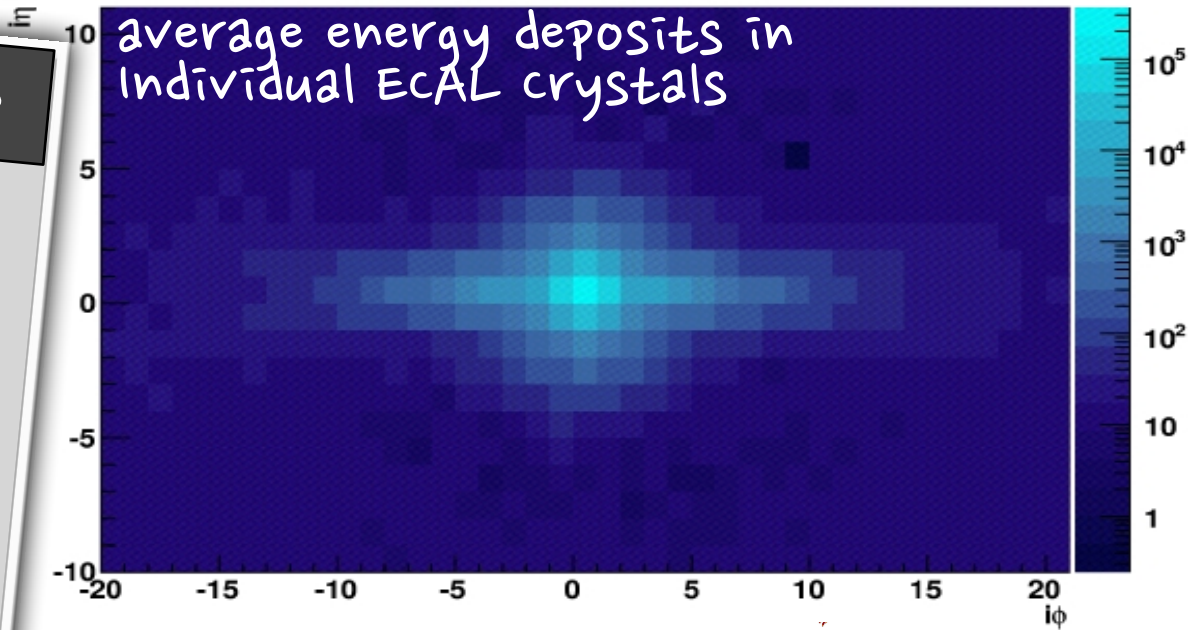
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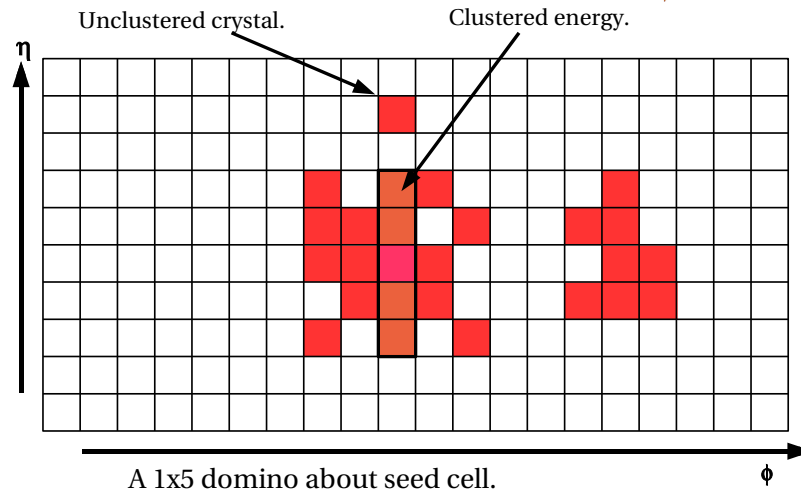
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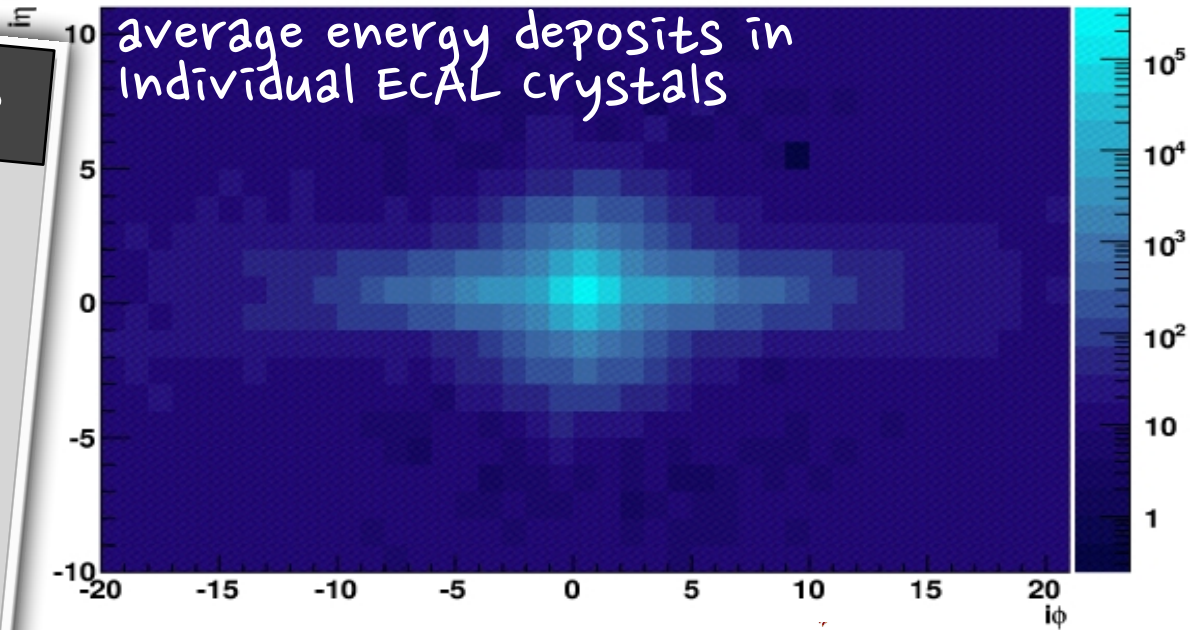
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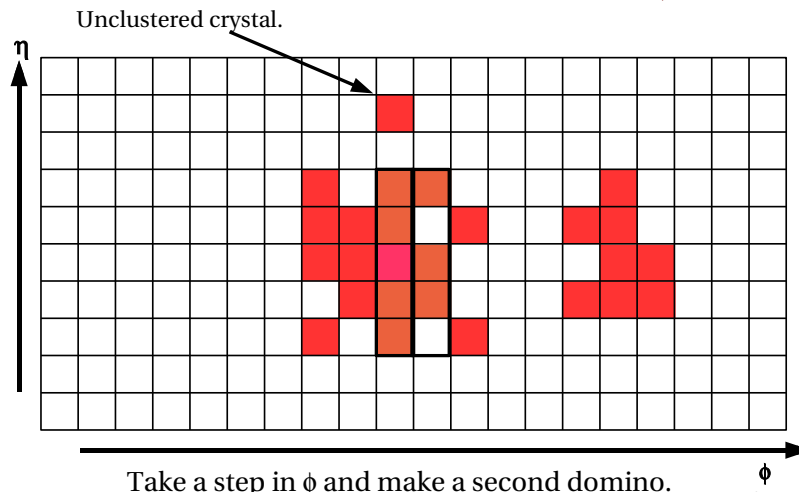
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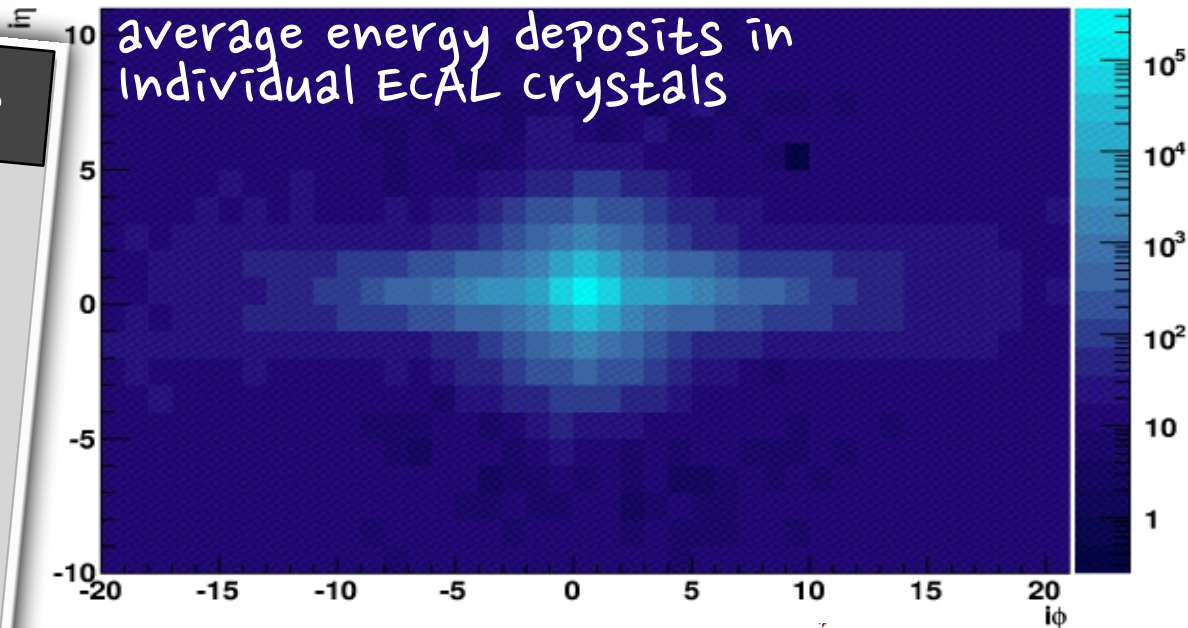
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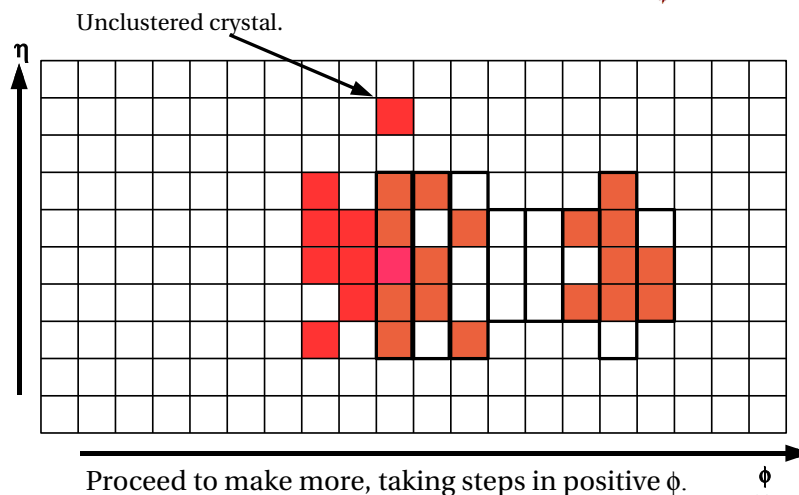
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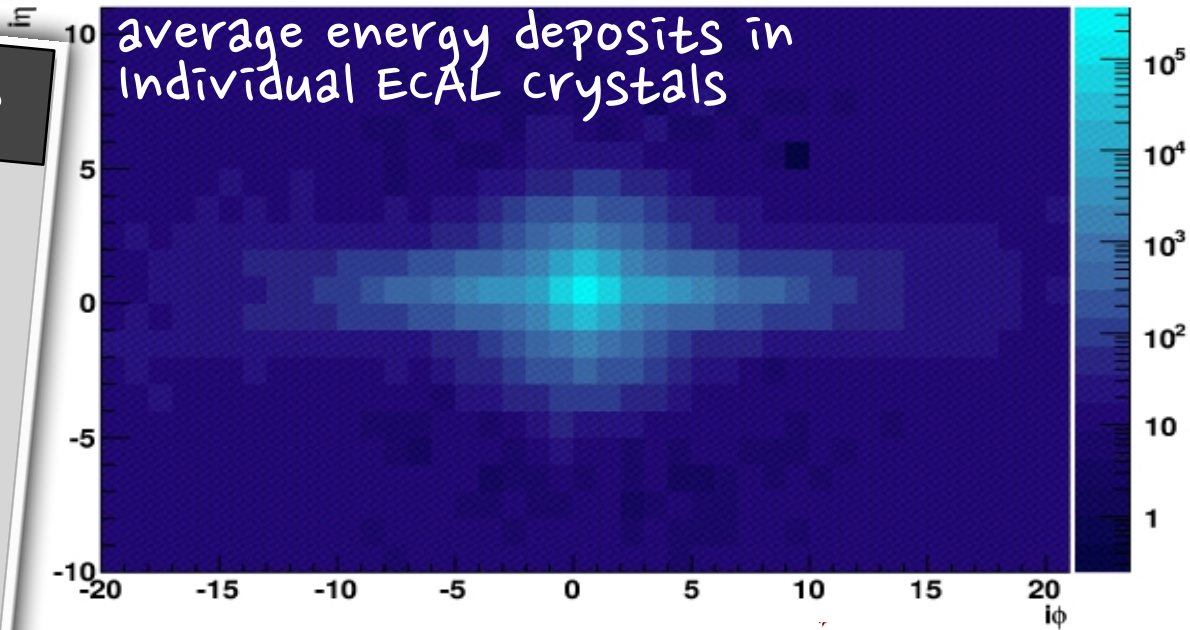
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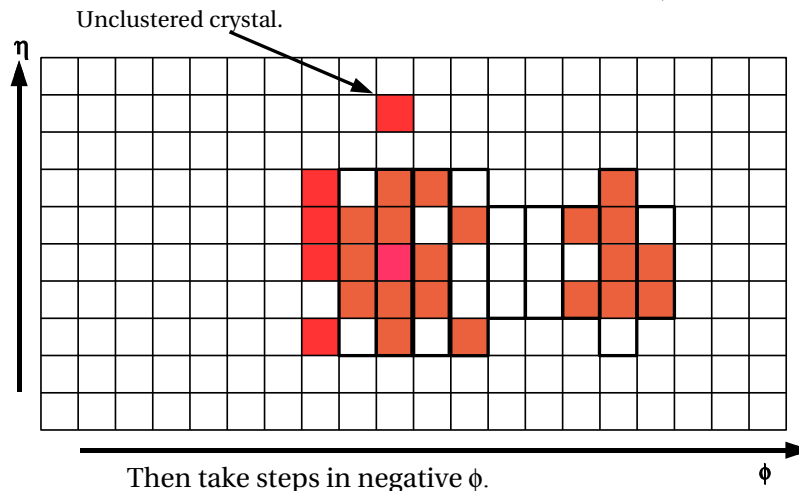
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- Depends on the experiment details!
- CMS ECAL lives inside the B-field
- Electrons in ECAL shower
 - spread radially out from impact point
 - also bend in ϕ
- Photons often pair convert
 - bend in ϕ
- Effect: "Jurasic symbol"



Clusters & Superclusters

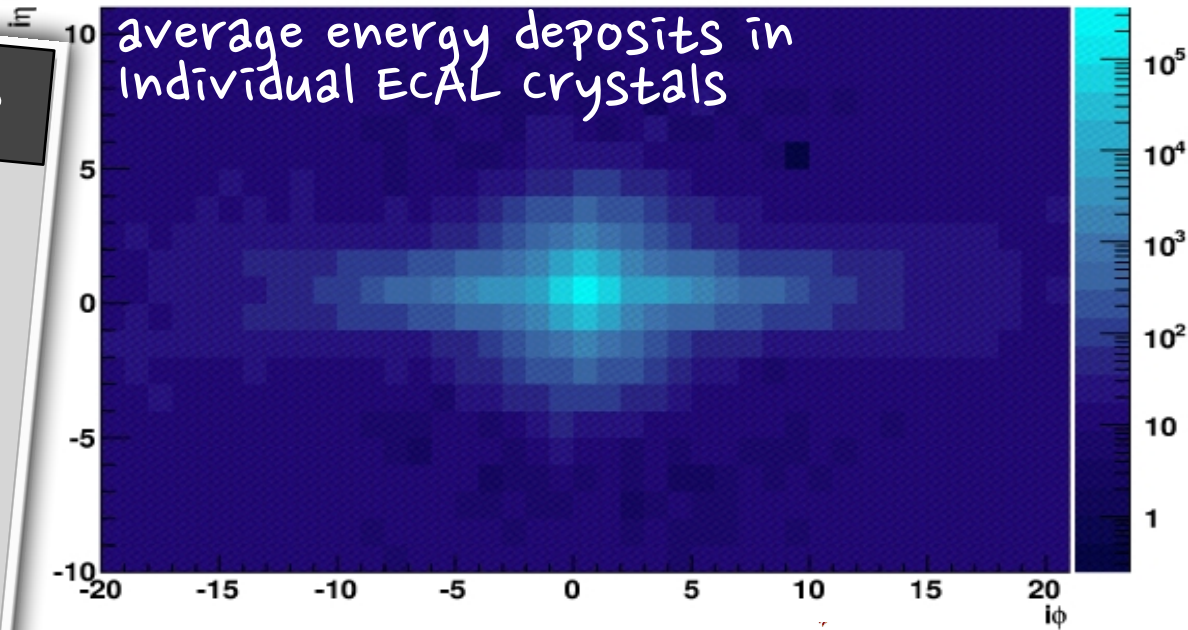
- Cluster nearby cells together
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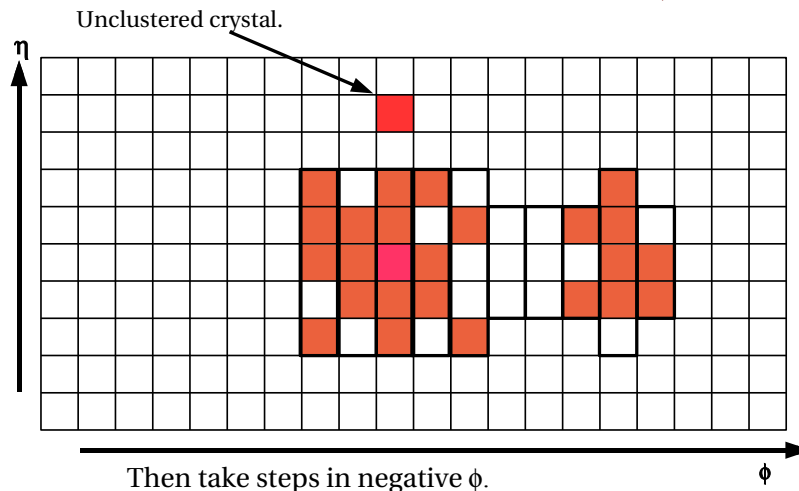
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average energy deposits in individual ECAL crystals



Clusters & Superclusters

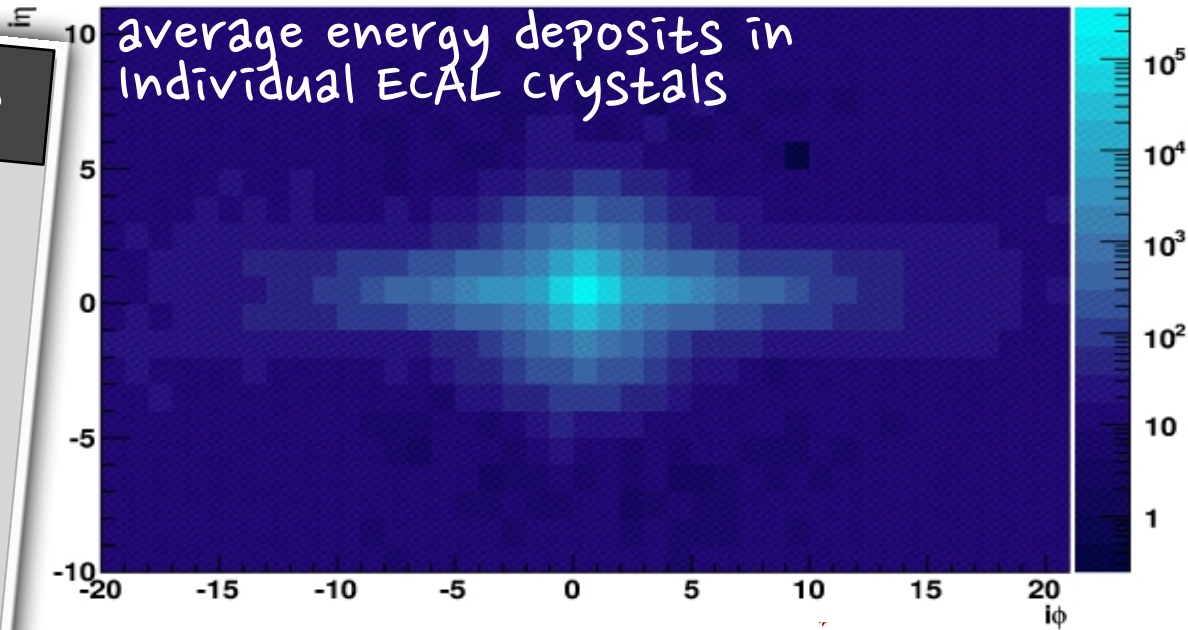
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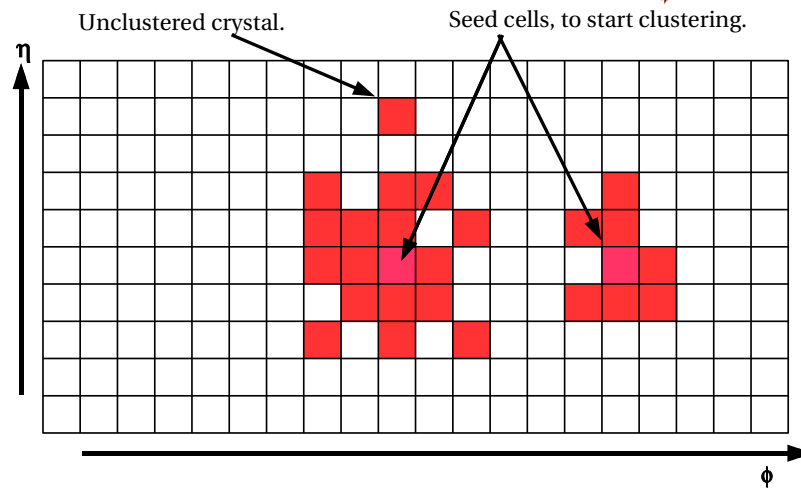
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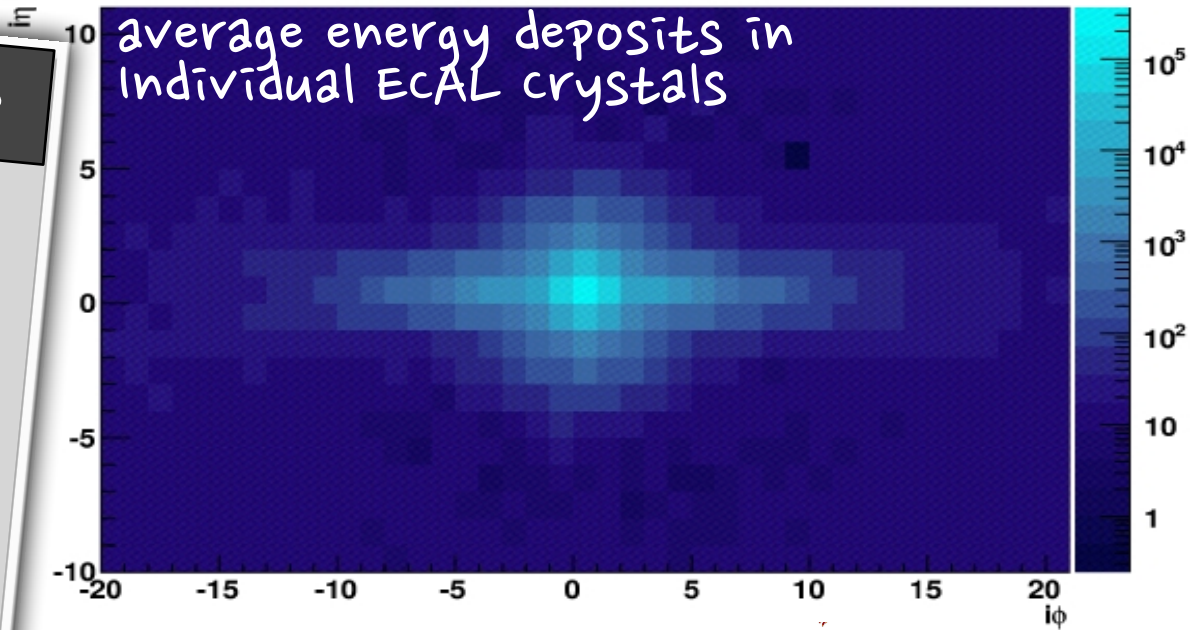
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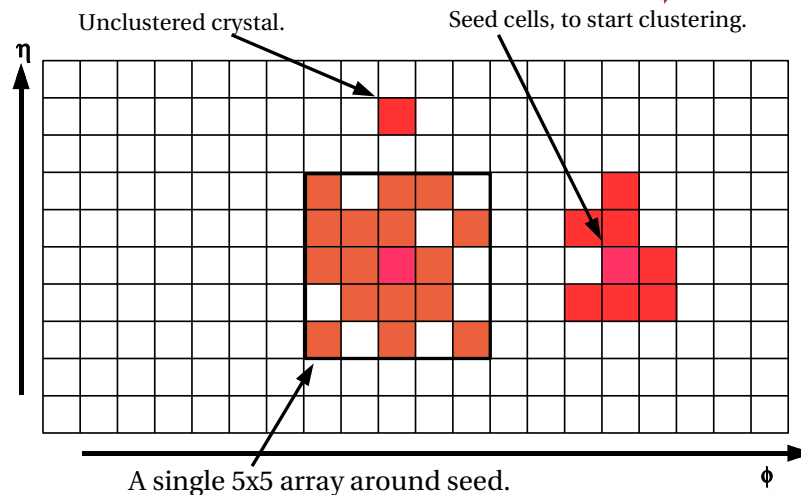
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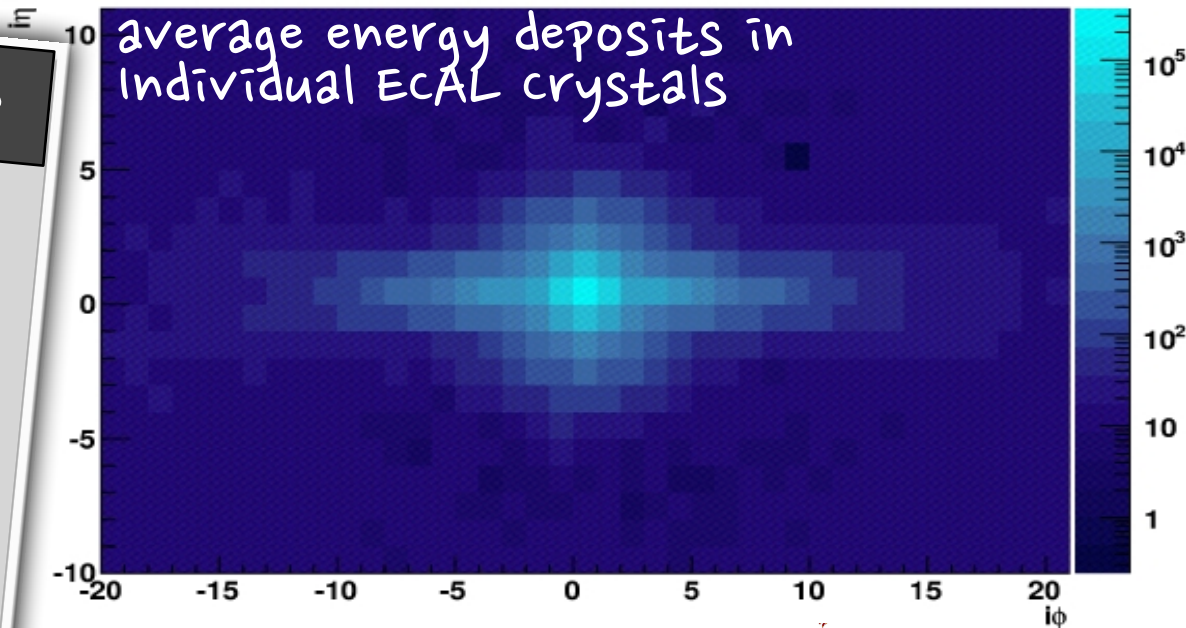
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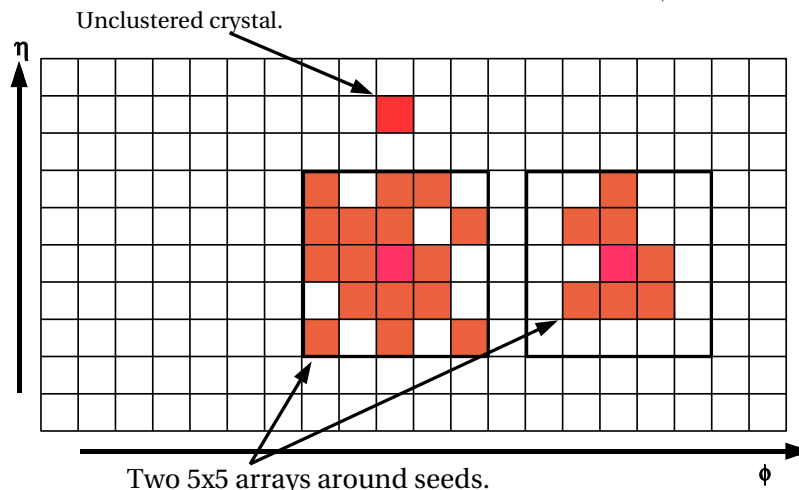
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Clusters & Superclusters

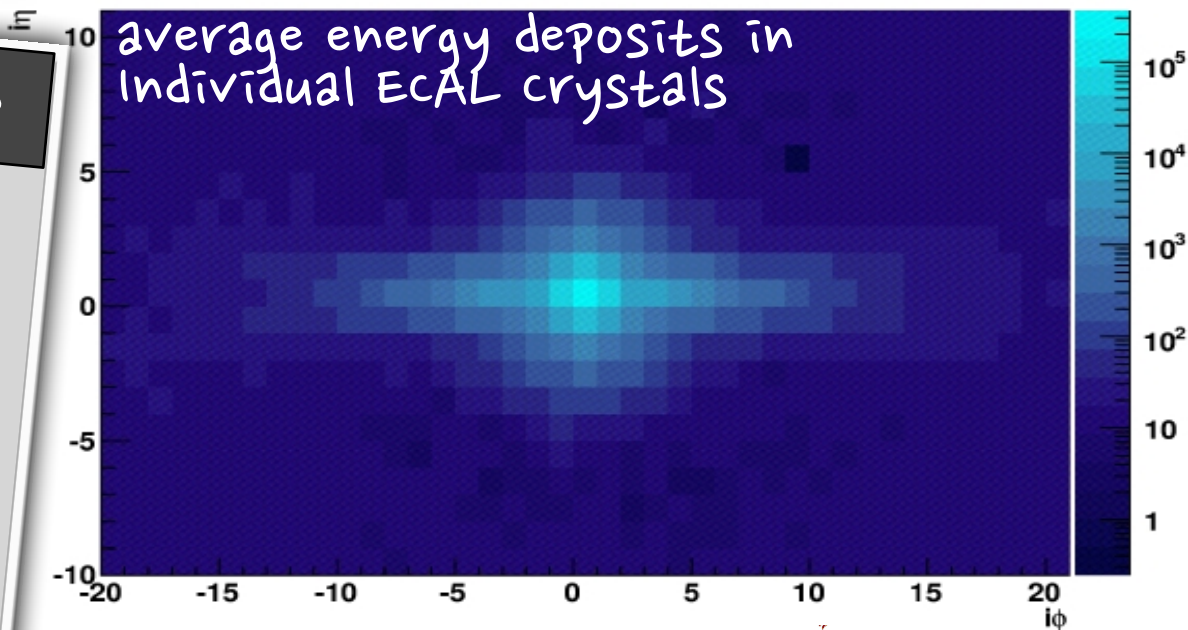
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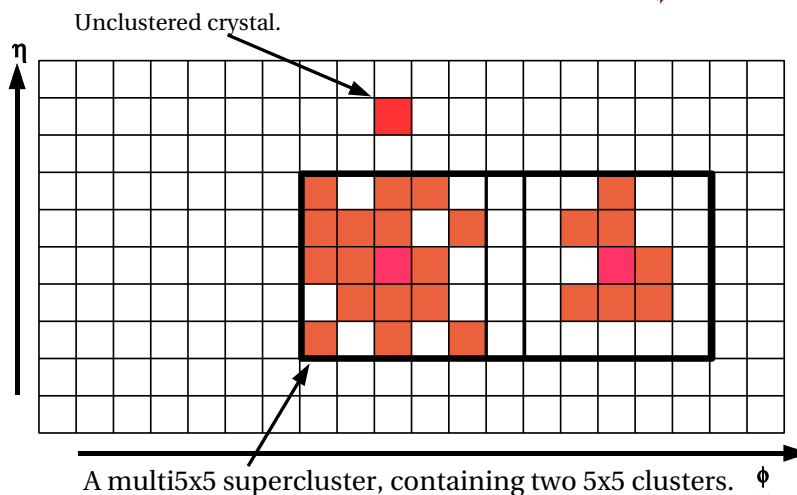
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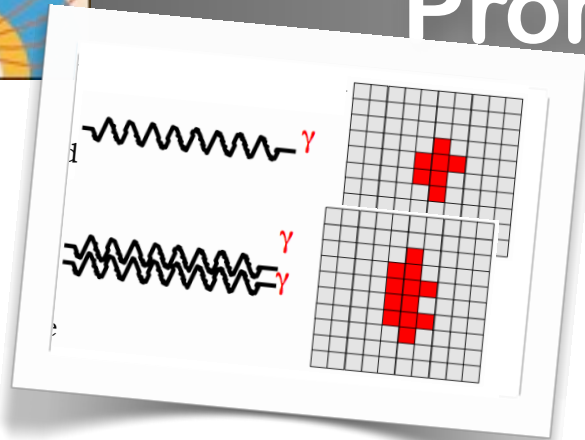




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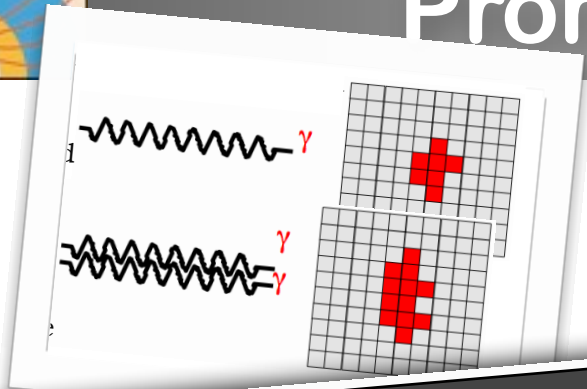


Prompt Photon ID





Prompt Photon ID



Shower shape variables

Pre-select γ 's

- $R_\eta = 3 \times 7 / 7 \times 7$ ratio of energies in η
- $R_\phi = 3 \times 7 / 7 \times 7$ ratio of energies in ϕ
- $w_\eta =$ RMS width along η of cluster
- $E_T(\text{had})/E_T =$ hadronic E fraction

ID π^0 giving two sep. showers

- $E_{\text{ratio}} =$ Ratio of 1st & 2nd maxima along η
- $\Delta E =$ diff between 2nd max and min

ID π^0 giving two merged showers

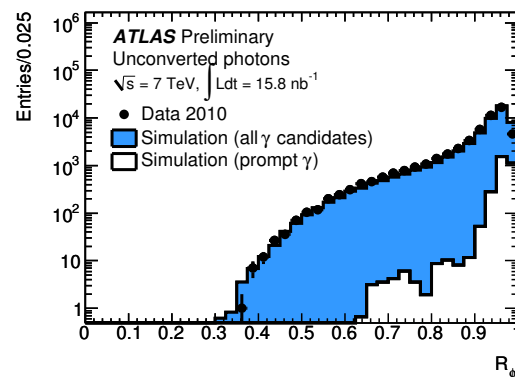
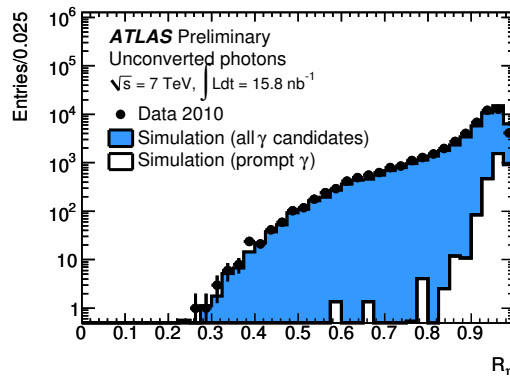
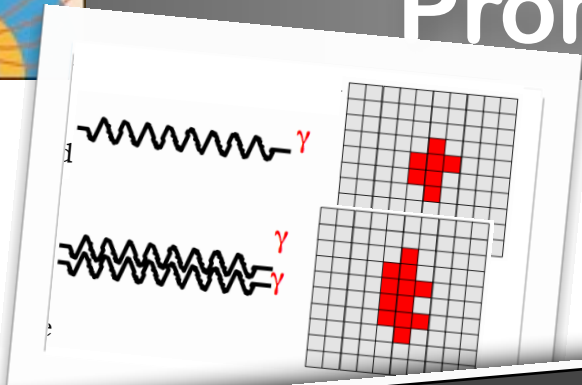
- $F_{\text{side}} =$ energy in 7 cells centered around 1st max (excluding the core cells)
- $w_{s3} =$ RMS width of core cells

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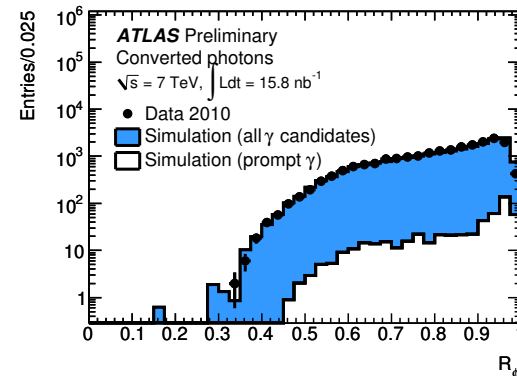
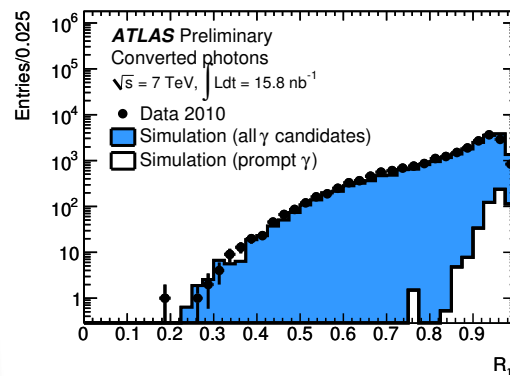
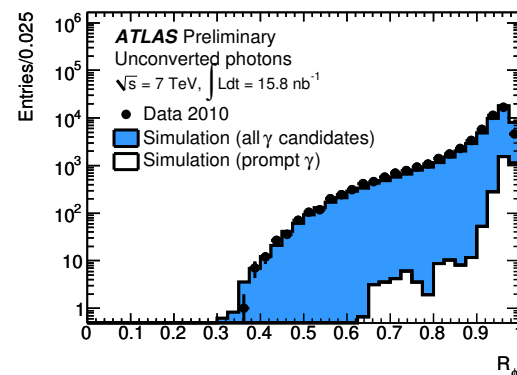
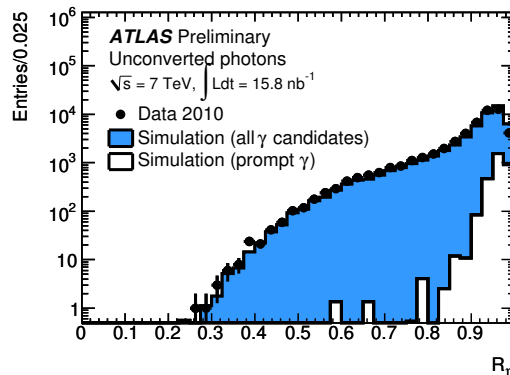
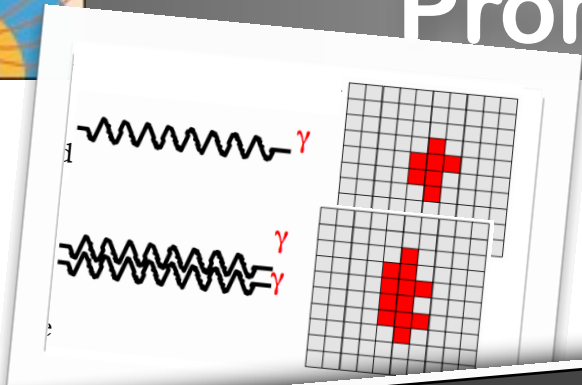
ID π^0 giving two merged showers

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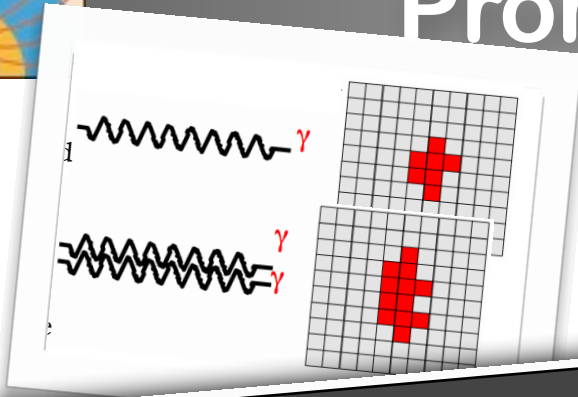
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Shower shape variables

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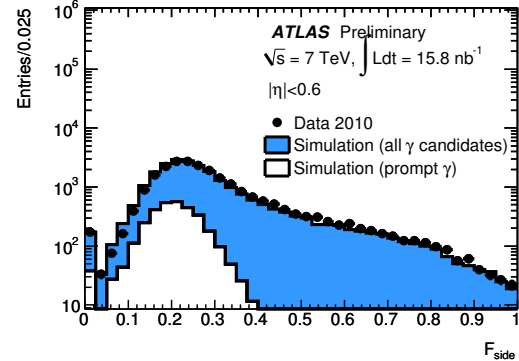
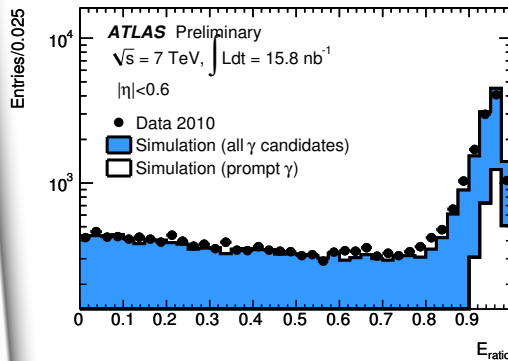
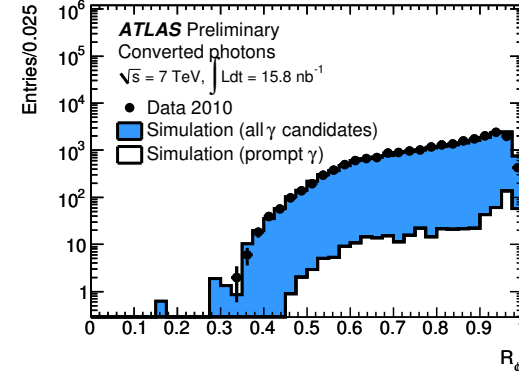
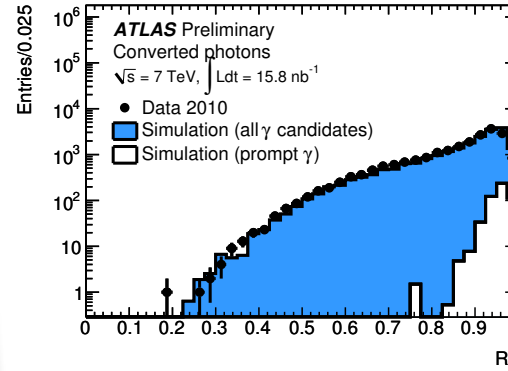
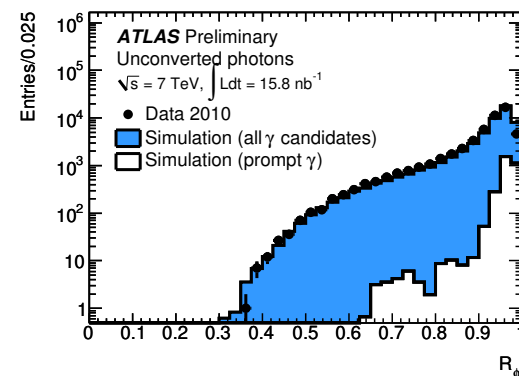
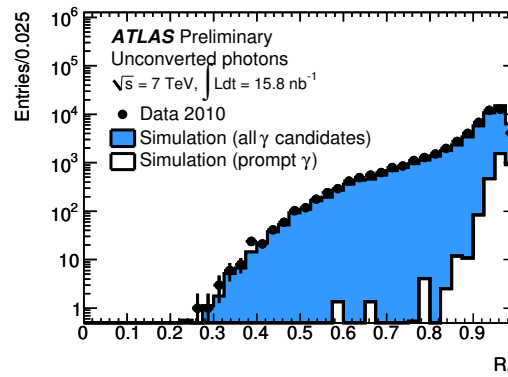
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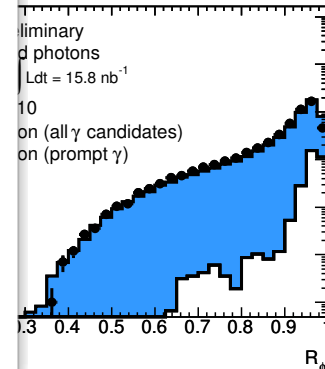
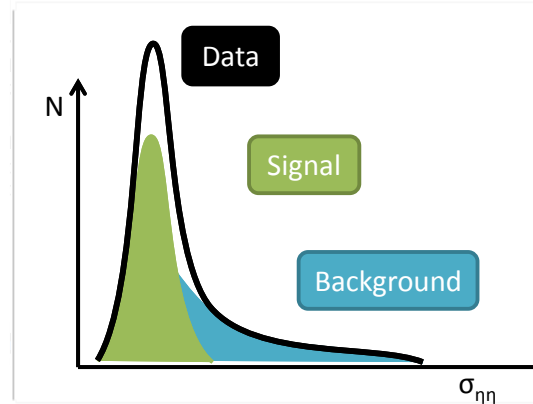
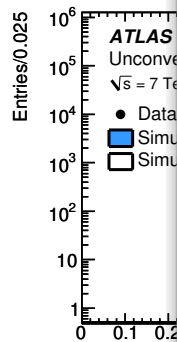
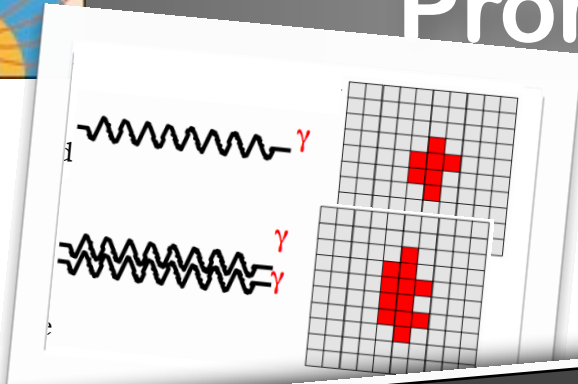
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Shower shape variables

Pre-select γ 's

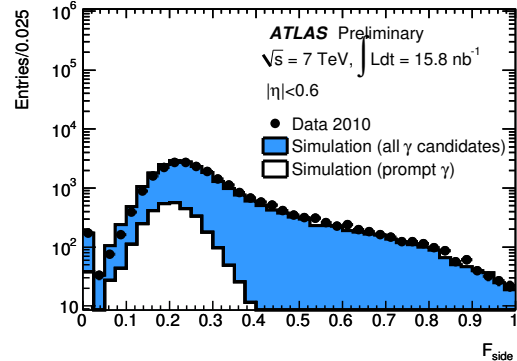
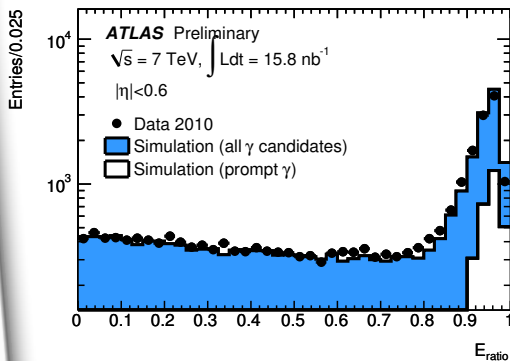
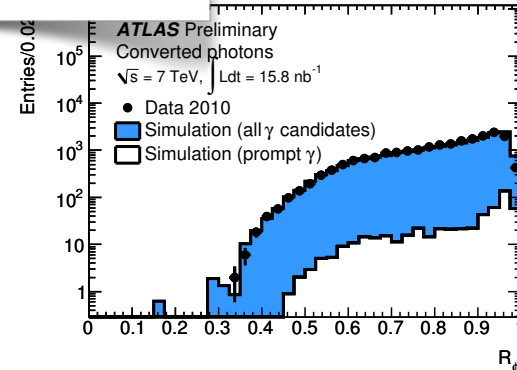
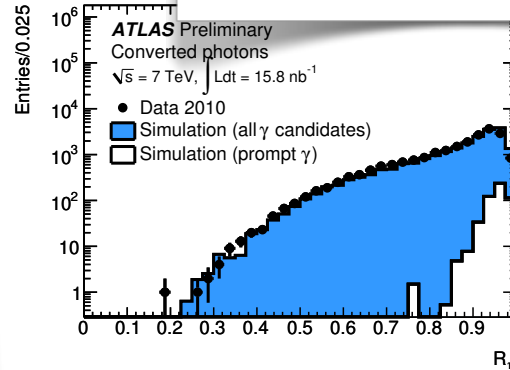
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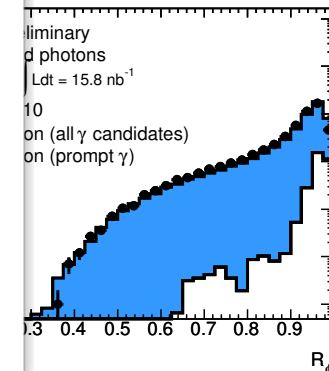
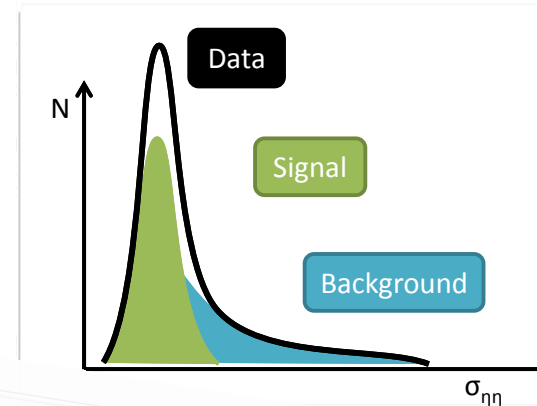
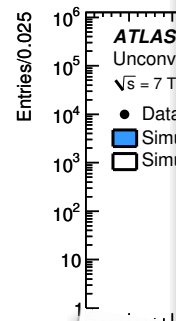
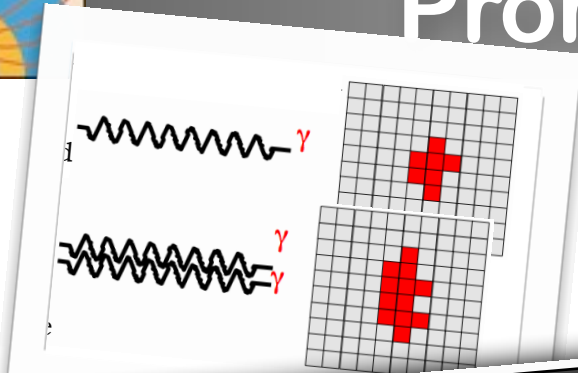


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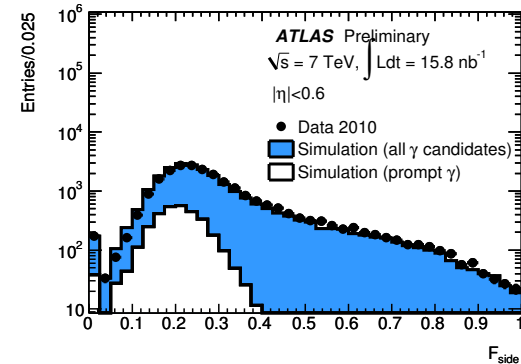
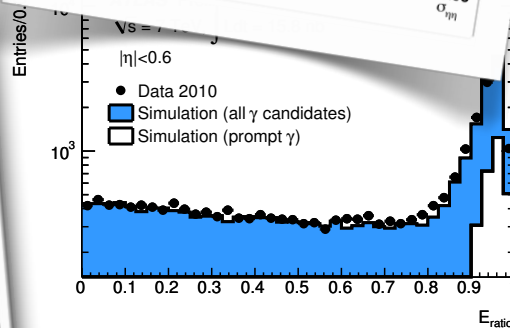
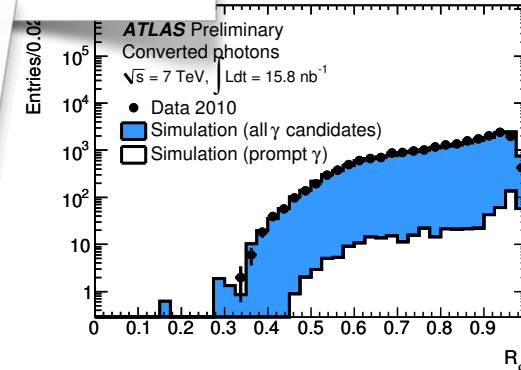
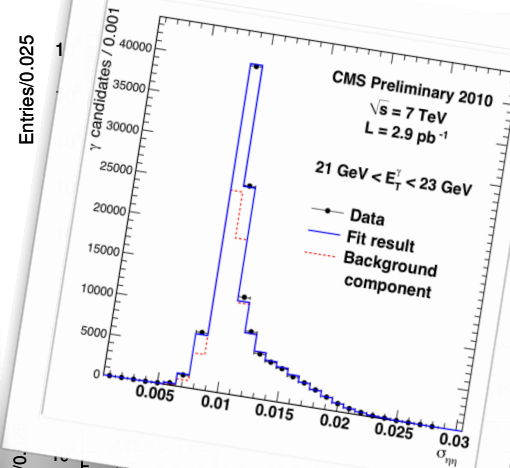
Prompt Photon ID



Shower shape variables

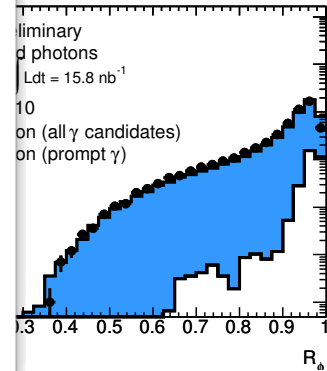
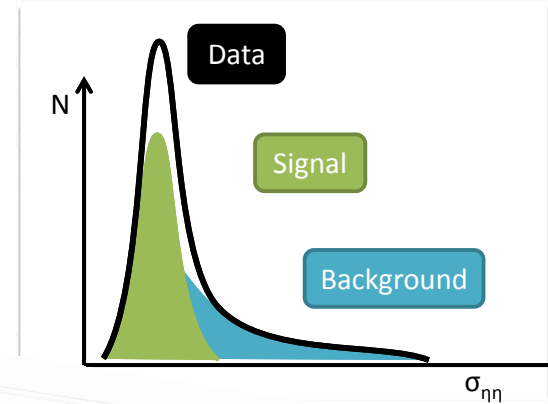
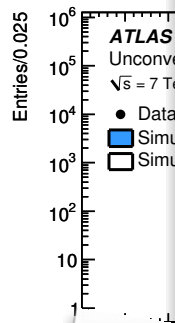
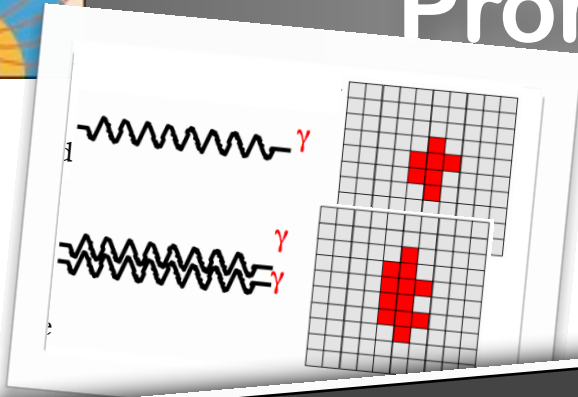
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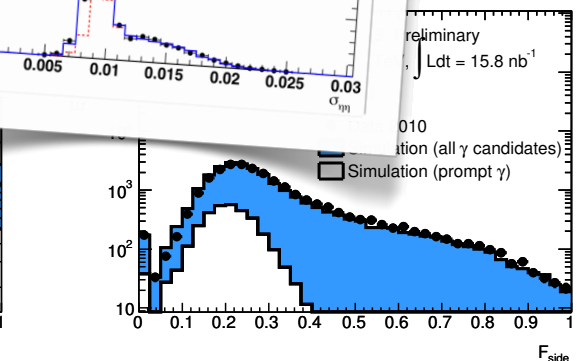
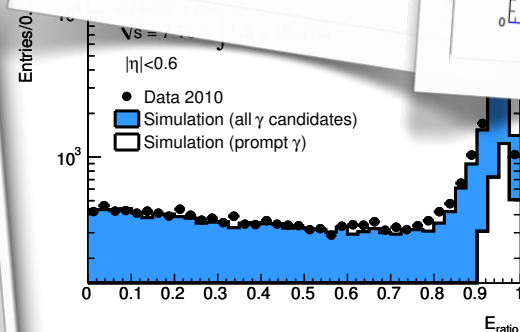
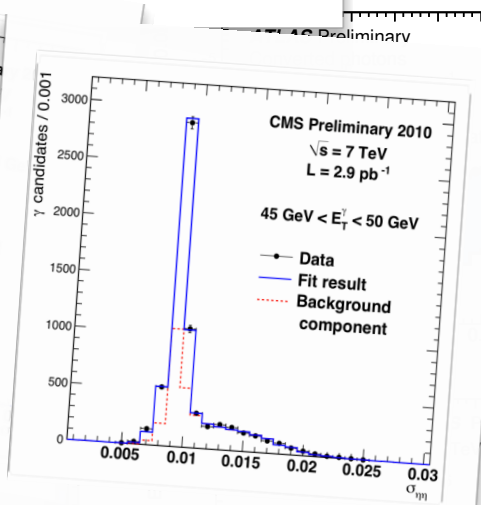
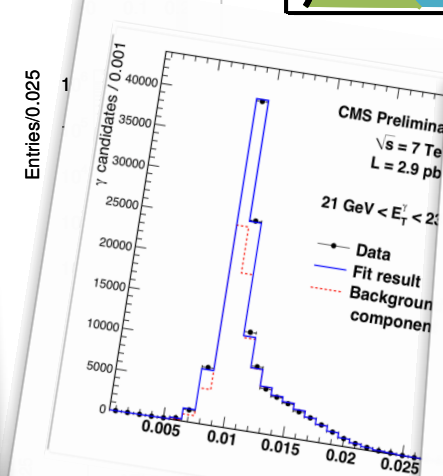
Prompt Photon ID



Shower shape variables

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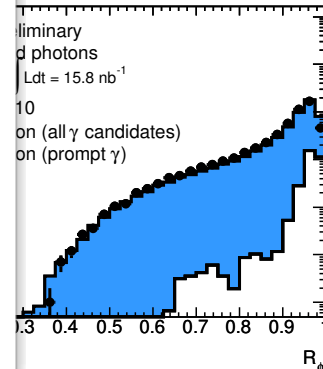
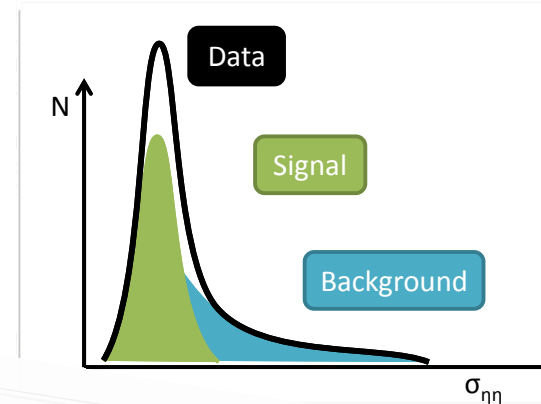
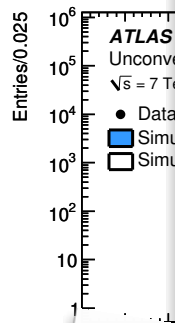
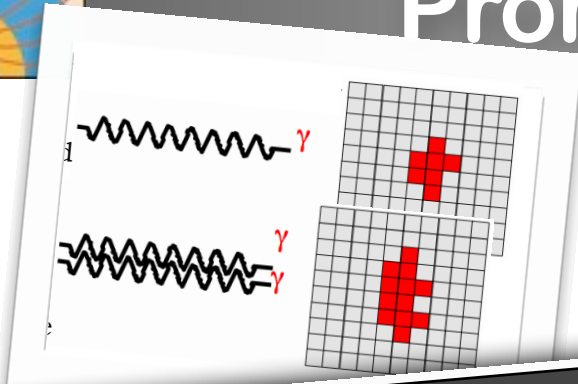


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14-08-2012



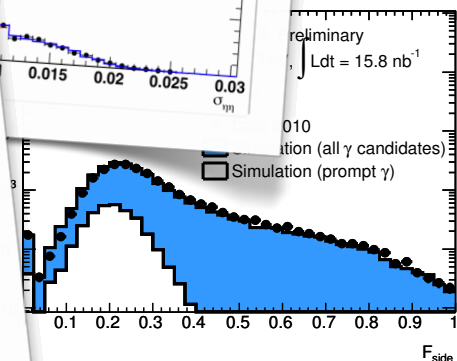
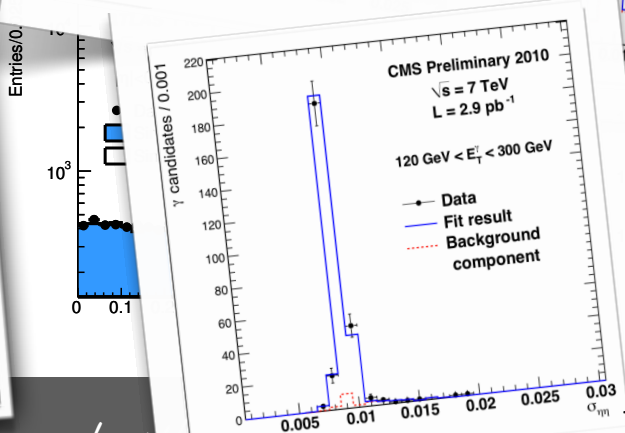
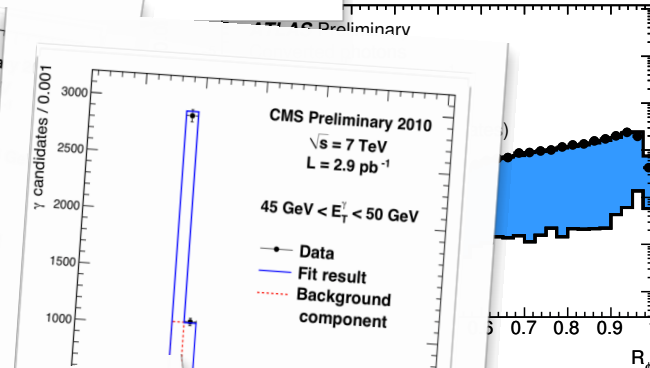
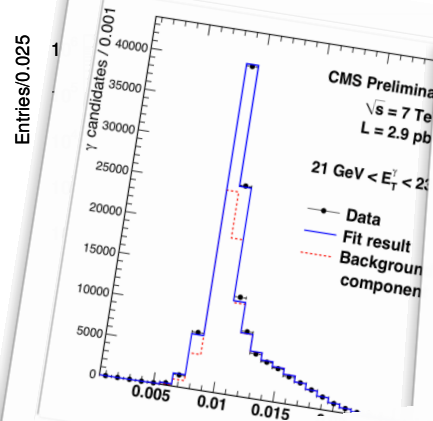
Prompt Photon ID



Shower shape variables

Pre-select γ 's

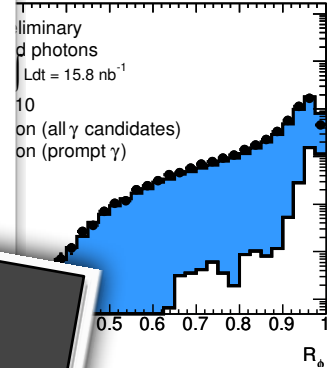
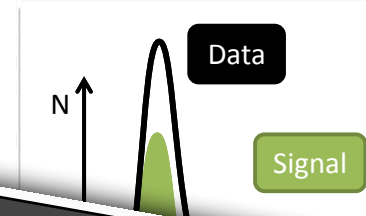
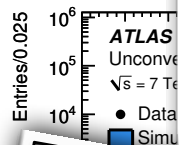
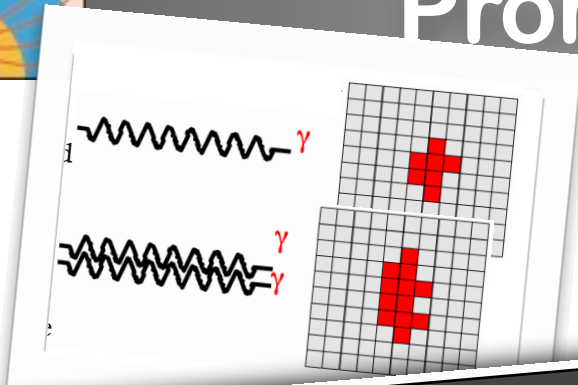
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14-08-2012

Prompt Photon ID



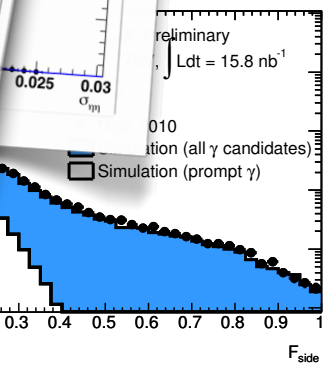
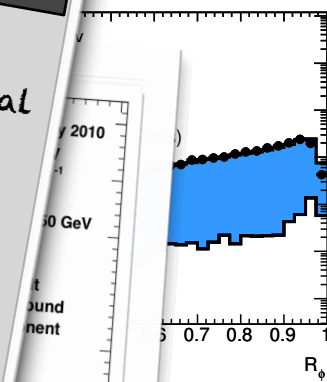
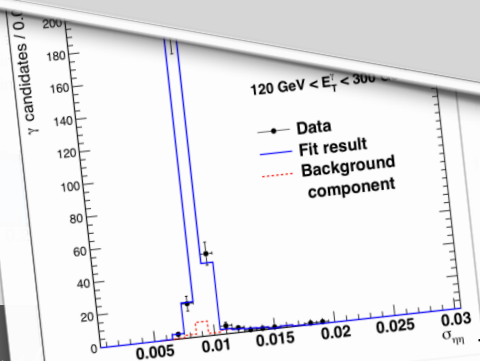
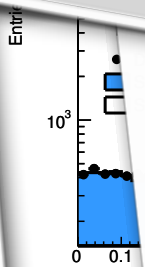
Shower shape variables

Pre-select γ 's

- $R_\eta = 3 \times 7 / 7 \times 7$ ratio of energies in η
- $R_\phi = 3 \times 7 / 7 \times 7$ ratio of energies in ϕ
- $w_\eta =$ RMS width along η of cluster
- $E_T(\text{had})/E_T =$ hadronic E fraction
- ID π^0 giving two sep. showers
- $E_{\text{ratio}} =$ Ratio of 1st & 2nd maxima along η
- $\Delta E =$ diff between 2nd max and min
- ID π^0 giving two merged showers
- $F_{\text{side}} =$ energy in 7 cells centered around 1st max (excluding the core cells)
- $w_{s3} =$ RMS width of core cells

Energy Calibration

- Correct clustered energy for
 - losses due to interactions in material
 - rear leakage
 - due to cracks, punch through, etc
- lateral leakage
 - due to fixed cluster window
- Do this separately for
 - unconverted γ 's
 - converted γ 's



14-08-2012

→ Cavanaugh, ...



Photon Identification



Non-PF Isolation

- cone 0.4
 - centered on γ candidate (include ECAL & HCAL, tracker) e.g.
$$I = \text{Trk Iso} + \text{HCAL Iso} + \text{ECAL Iso}$$
 - ATLAS: exclude 5x7 EM cells around barycenter
 - CMS: exclude "Jurasic" symbol around barycenter
- small lateral leakage subtracted on average
- ambient energy density subtracted on average (UE + PU)
- γ candidates having < 3 GeV isolation considered "isolated"

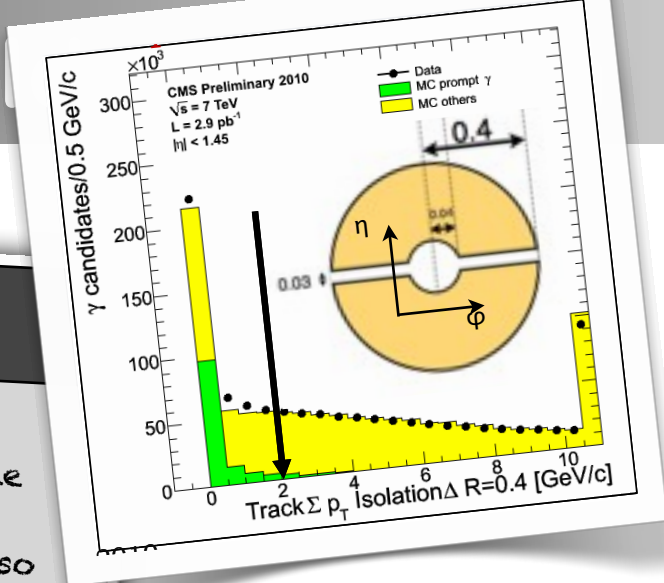


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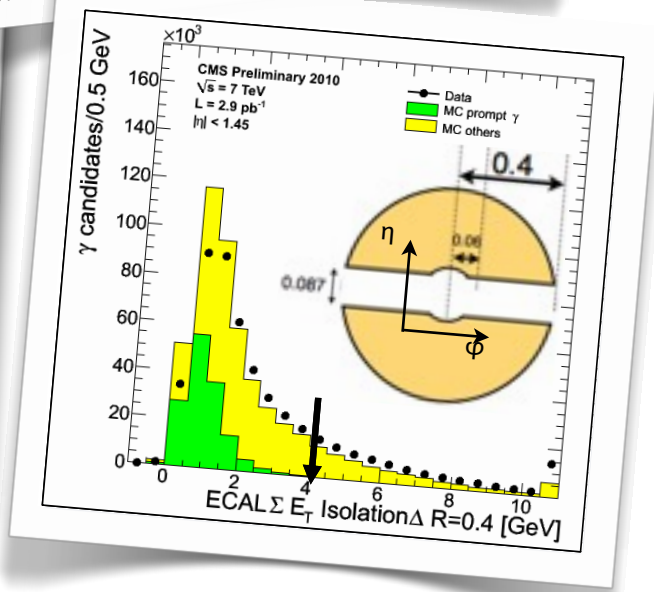
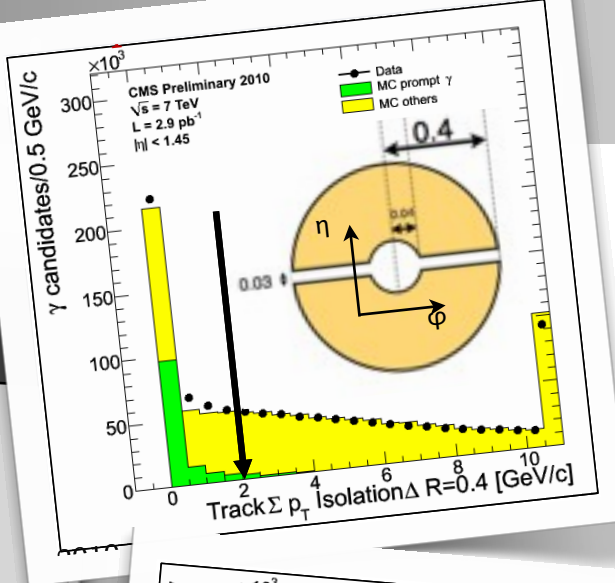


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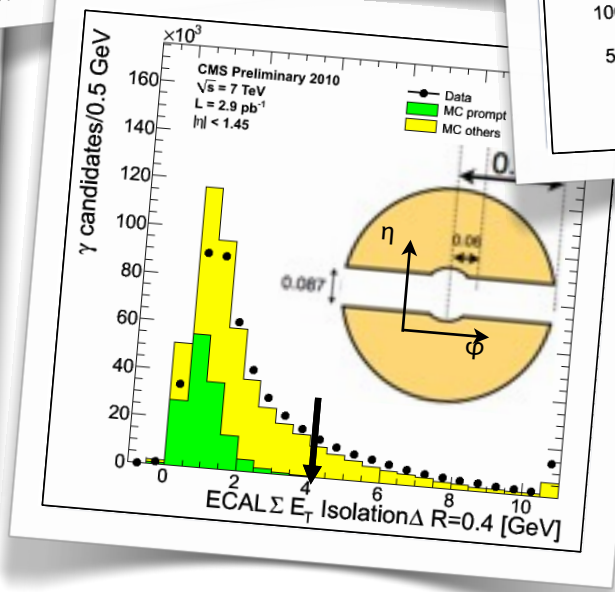
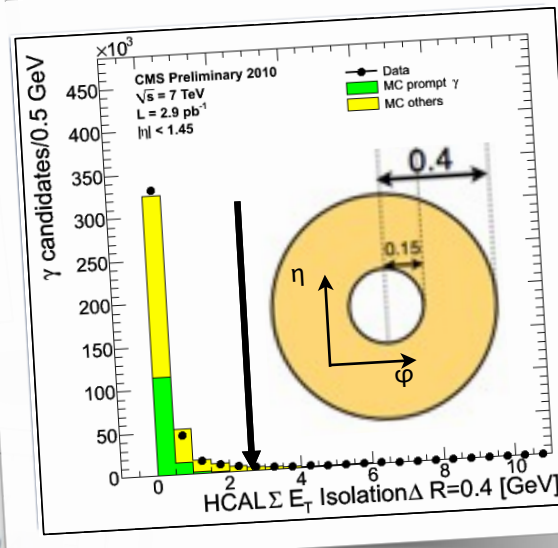
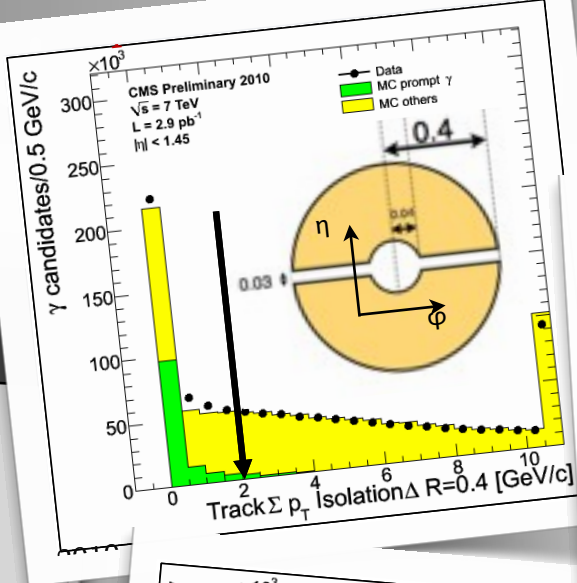


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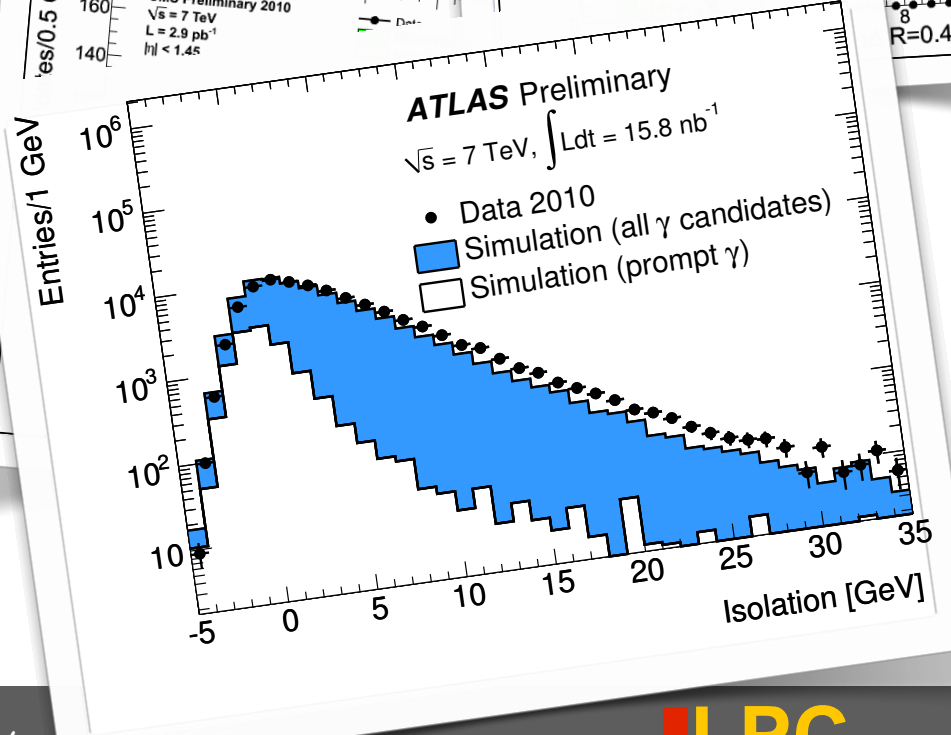
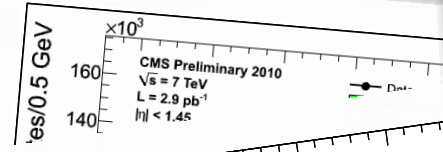
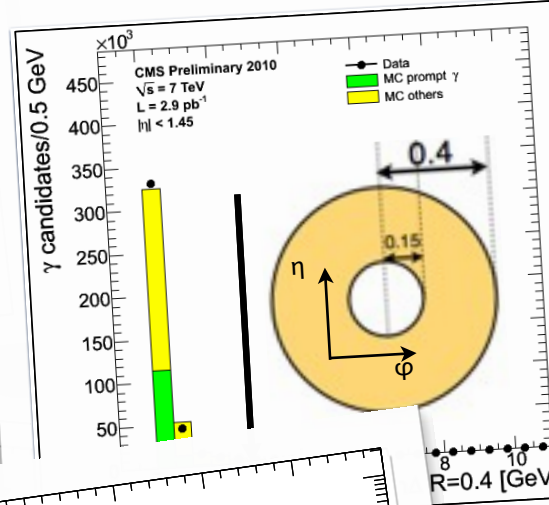
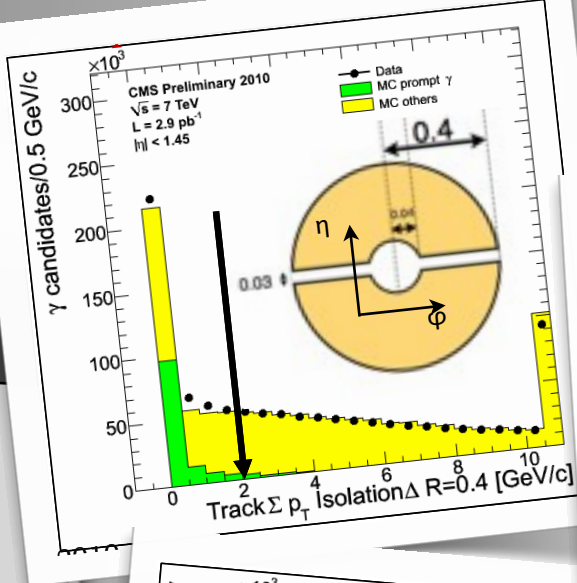


Photon

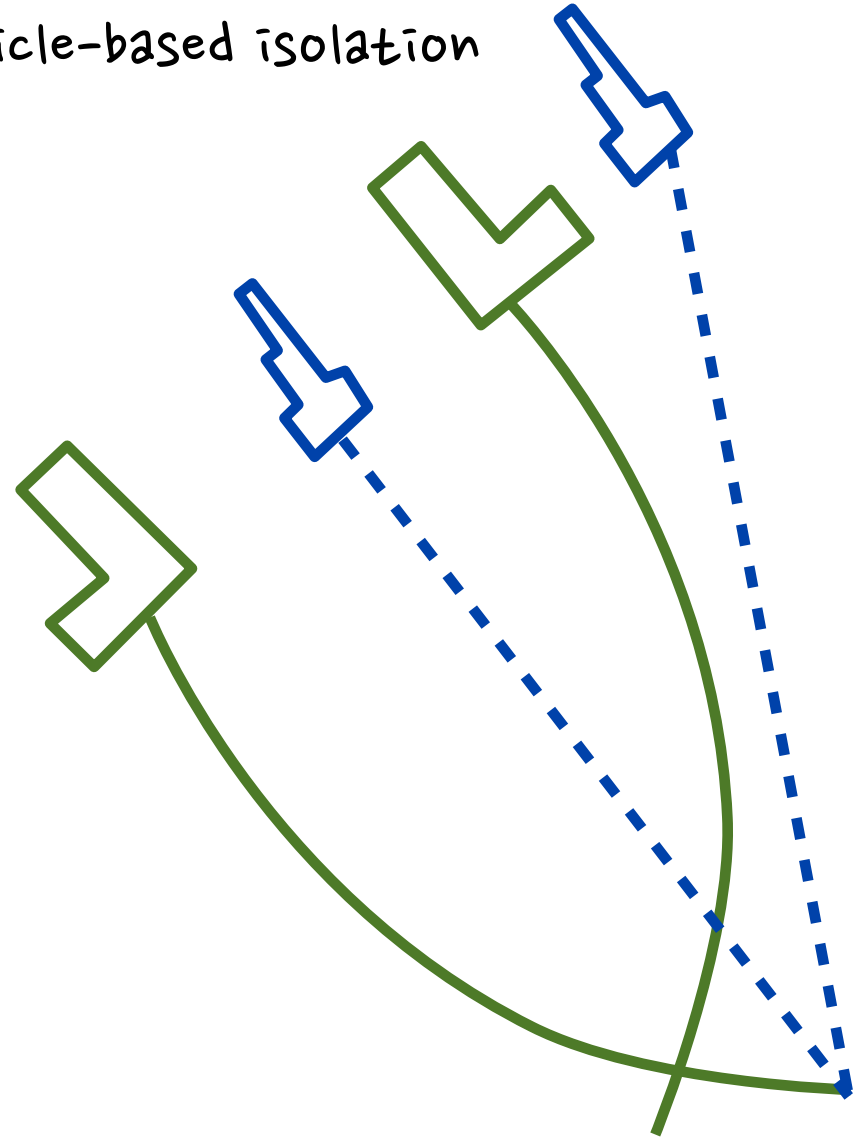
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Particle-based isolation



Isolation from other particles

- Absolute energy in cone:

$$I = \sum p_T(h^\pm) + \sum E_T(\gamma) + \sum E_T(h^0)$$

- Relative energy in cone:

$$I = \frac{\sum p_T(h^\pm) + \sum E_T(\gamma) + \sum E_T(h^0)}{p_T(e^\pm)}$$

- Apply average correction for neutral PU
 - determine from event energy density
- Typically require $I < 15\%$
 - efficiency \approx independent of PU!



Photon Identification

Isolation from other particles

- Use simple ΔR angular distance ≈ 0.3
- between electron and surrounding particles

Isolation from other particles

- Absolute energy in cone:

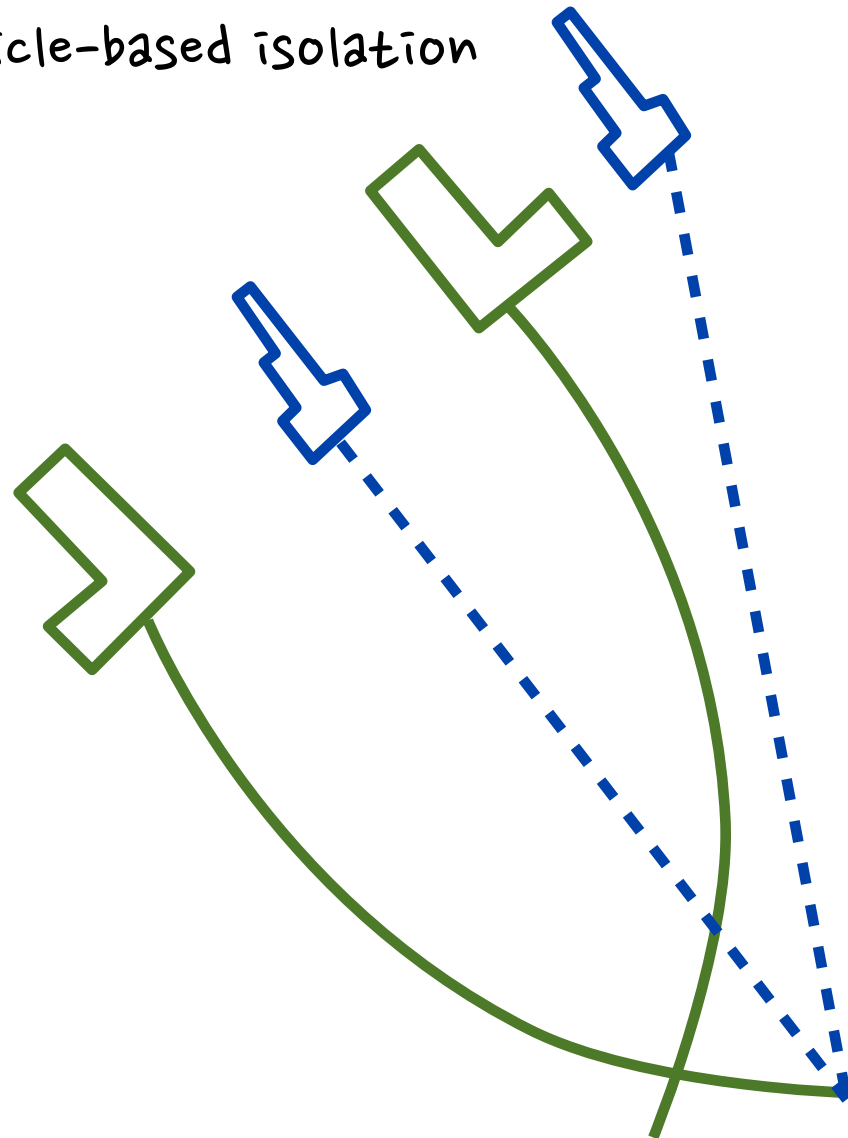
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Particle-based isolation





Photon Identification

Isolation from other particles

- Use simple $\Delta R = \sqrt{(\Delta\phi^2 + \Delta\eta^2)} \approx 0.3$ angular distance
- between electron and surrounding particles

Isolation from other particles

- Absolute energy in cone:

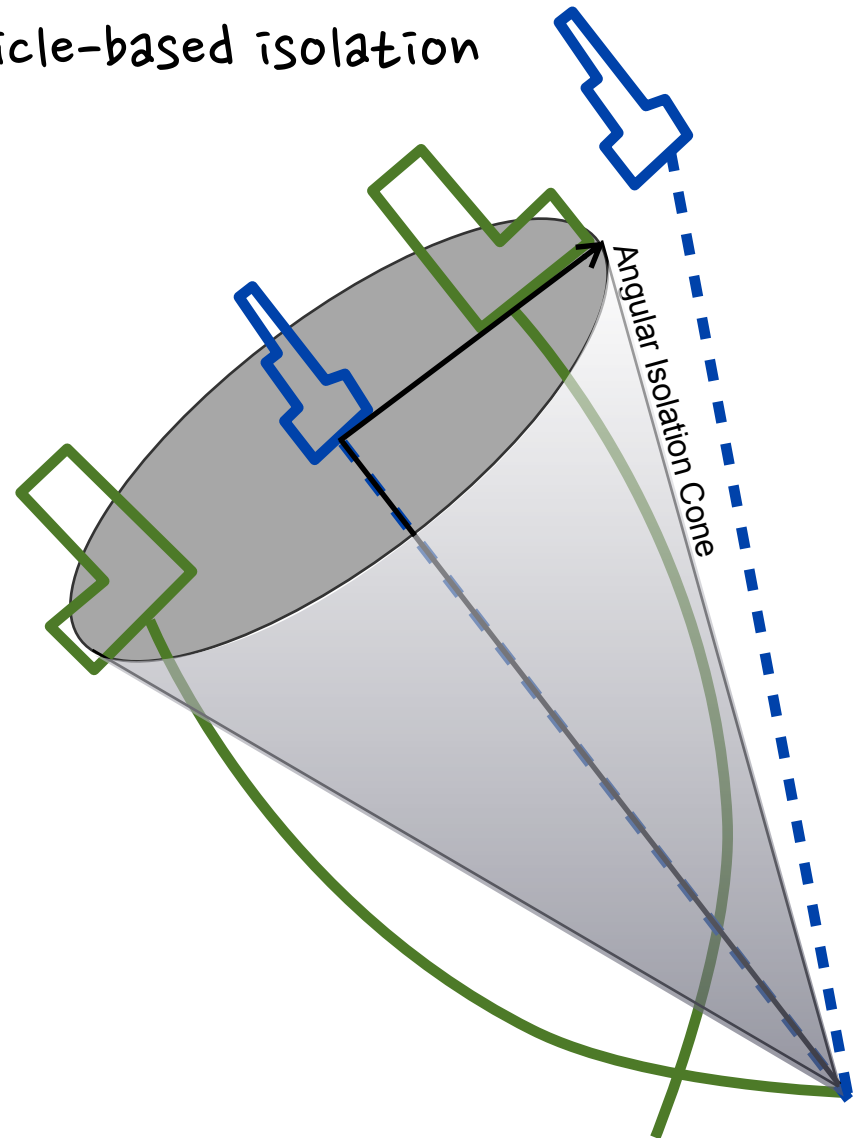
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Photon Identification

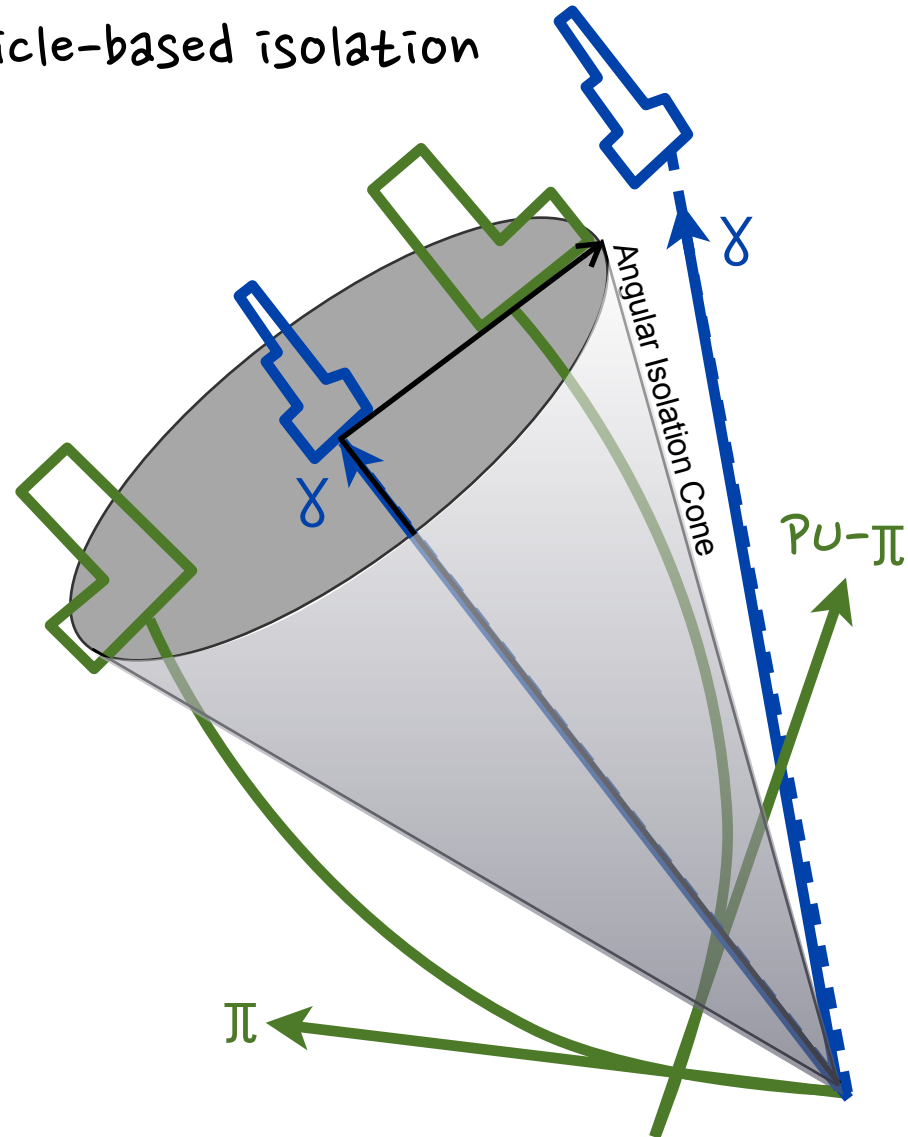
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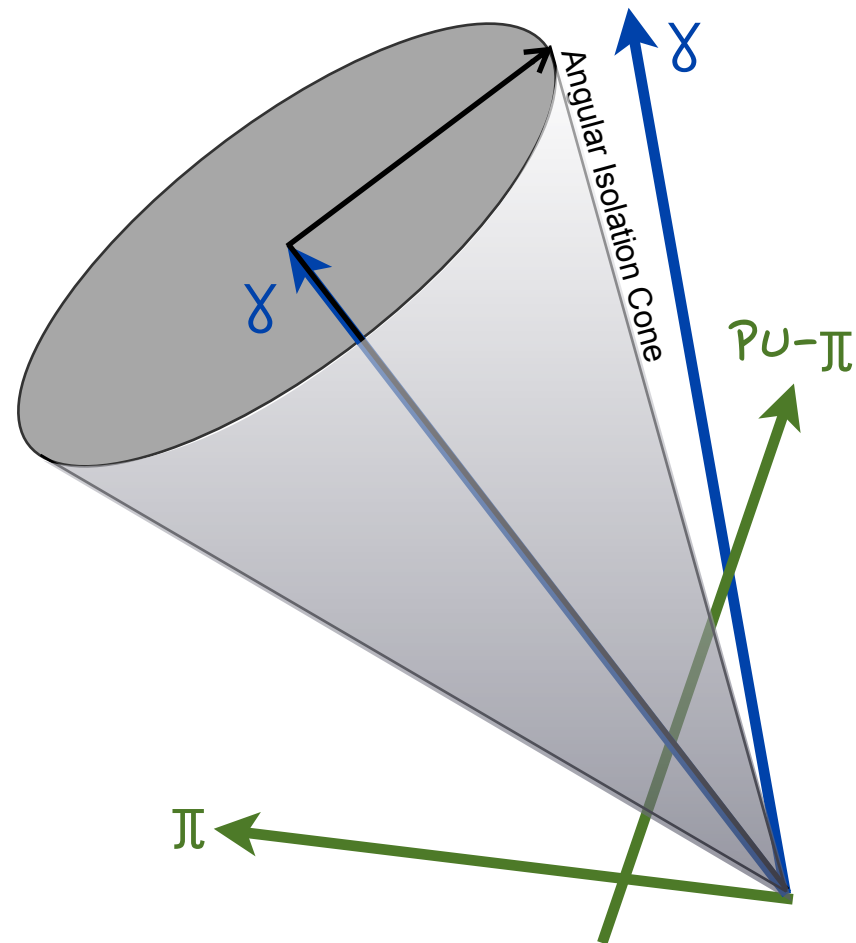
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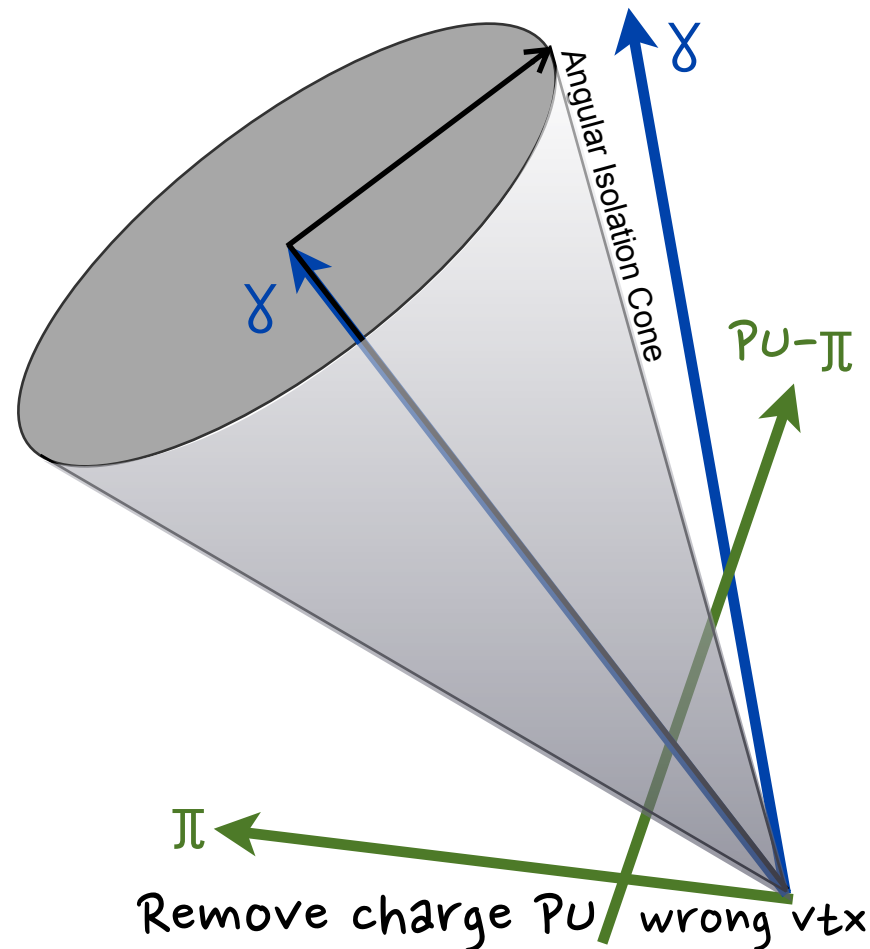
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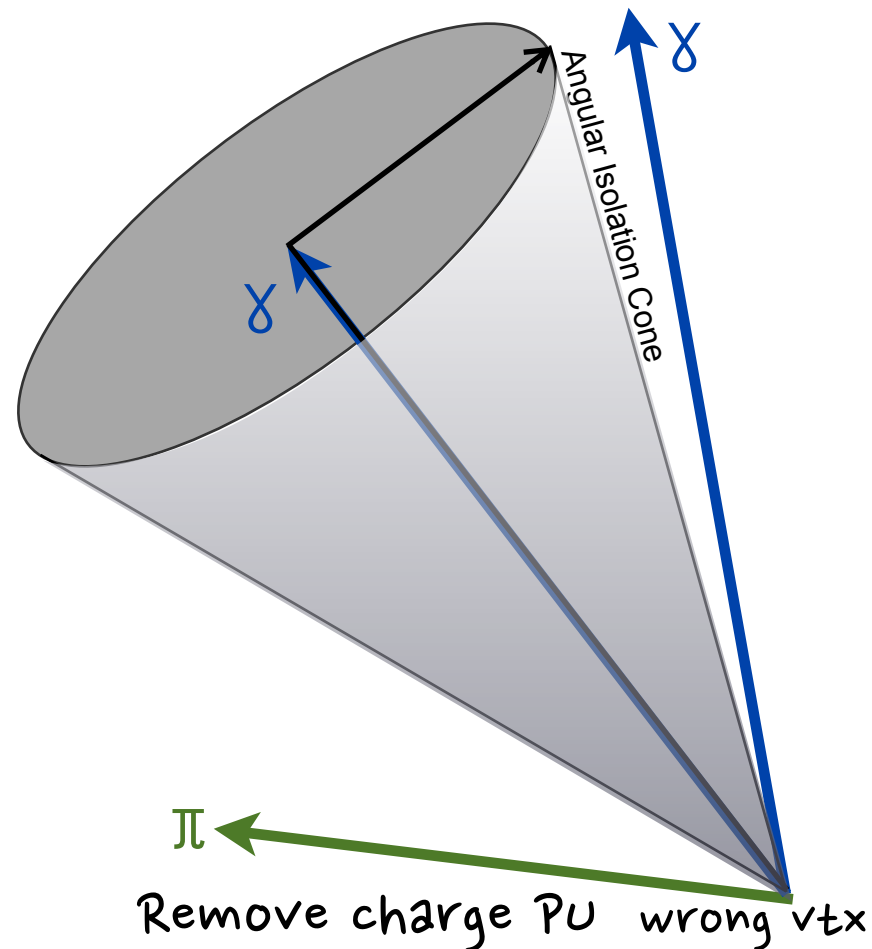
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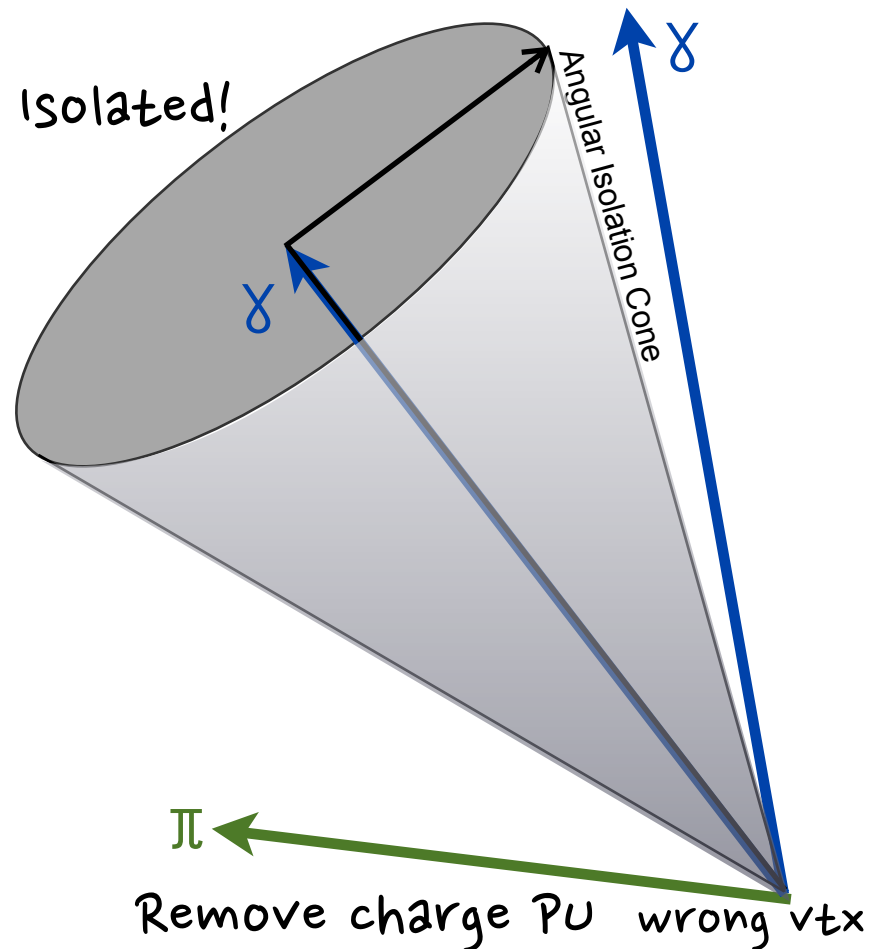
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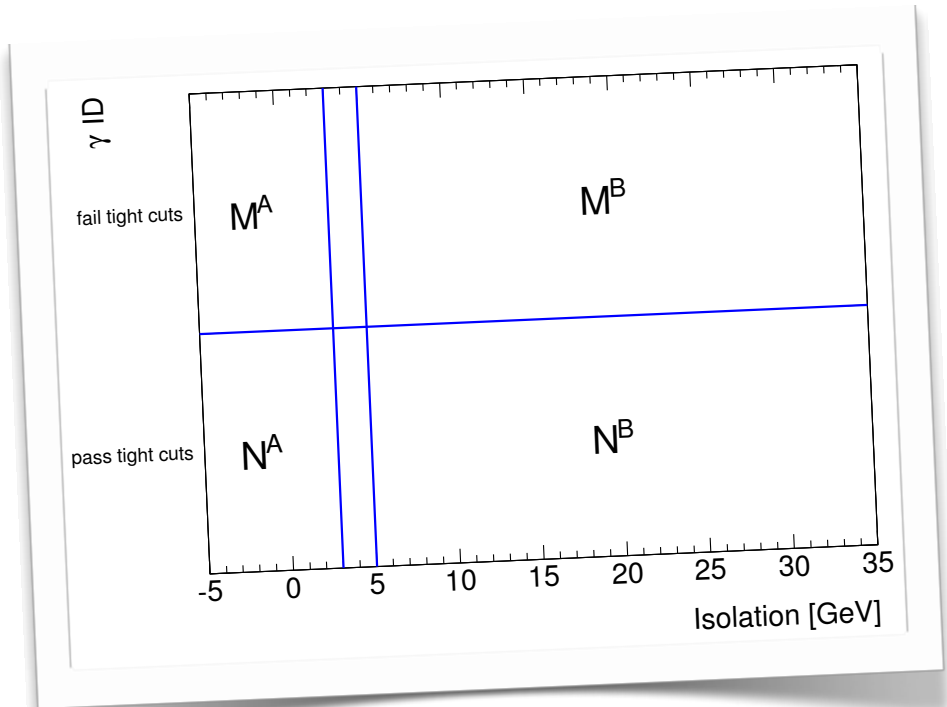




Measuring Efficiency/Purity

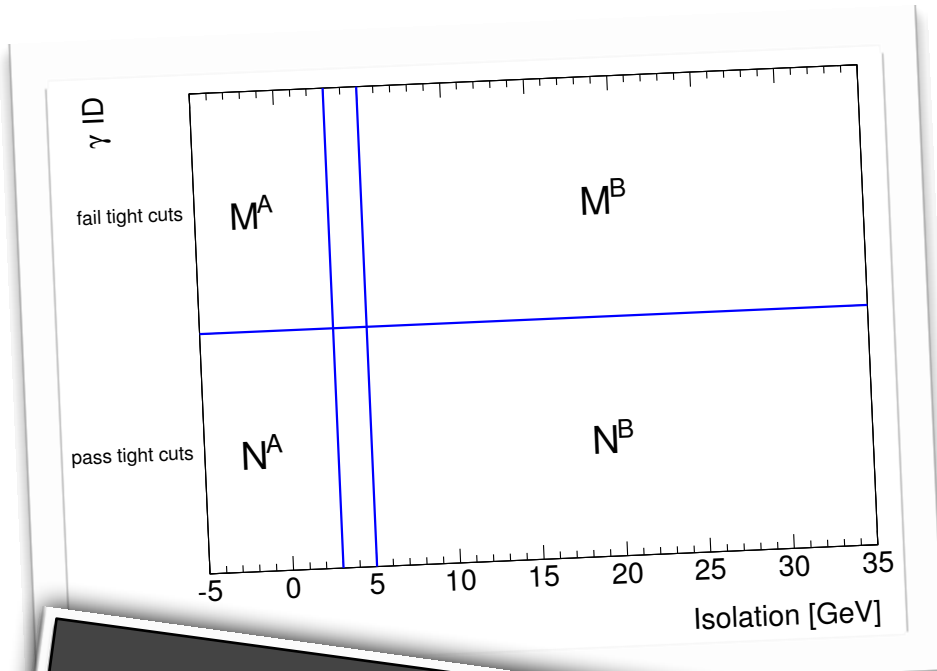


Measuring Efficiency/Purity





Measuring Efficiency/Purity

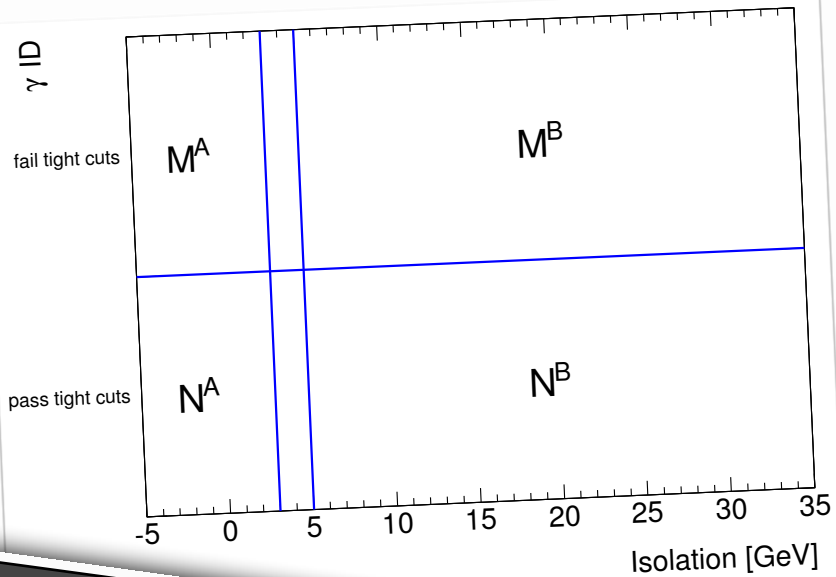


Photon yield and purity

- Assuming that γ ID is uncorrelated with the Isolation, then the number of γ s is:
 - $N(\gamma) = N^A - \text{background}$
 - $\text{background} = N^B M^A/M^B$
- And the γ purity is just:
 - $P = 1 - N^B/N^A M^A/M^B$

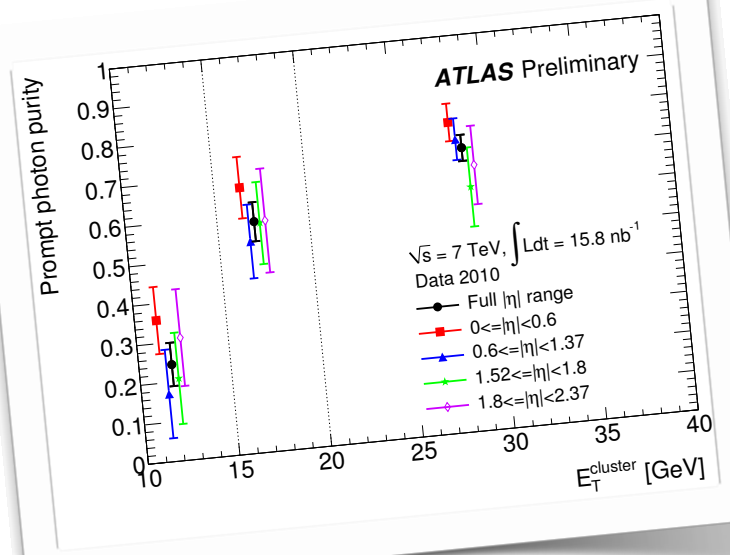


Measuring Efficiency/Purity



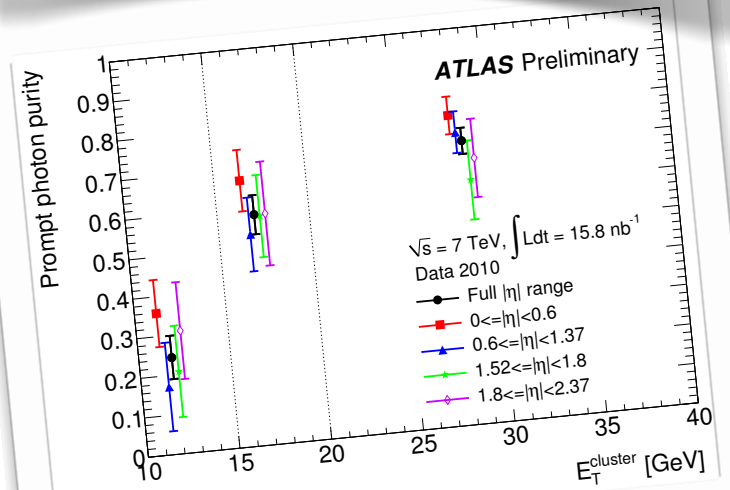
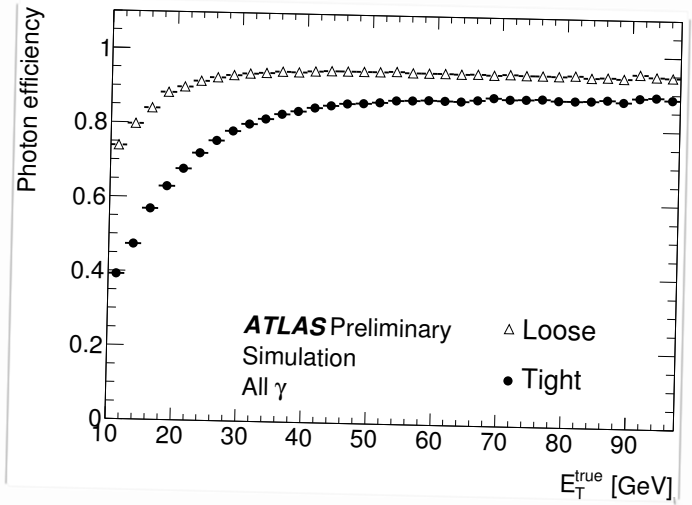
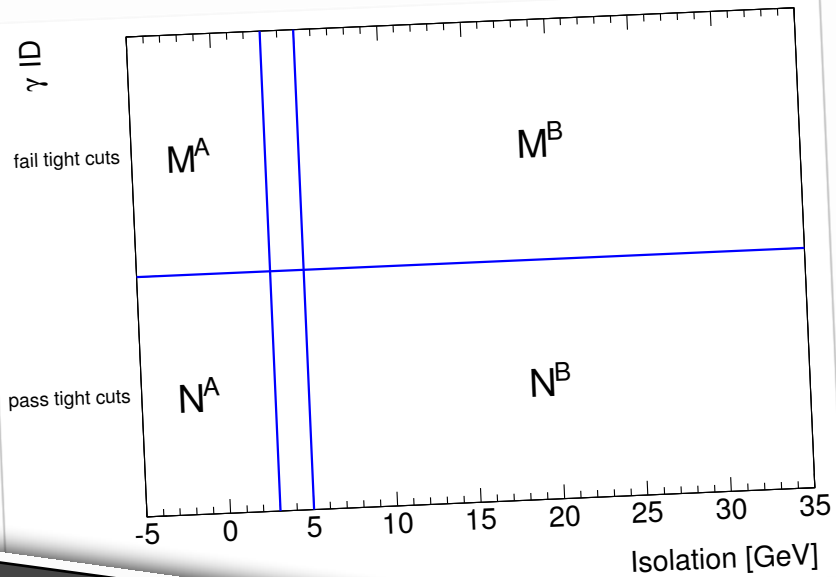
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Measuring Efficiency/Purity



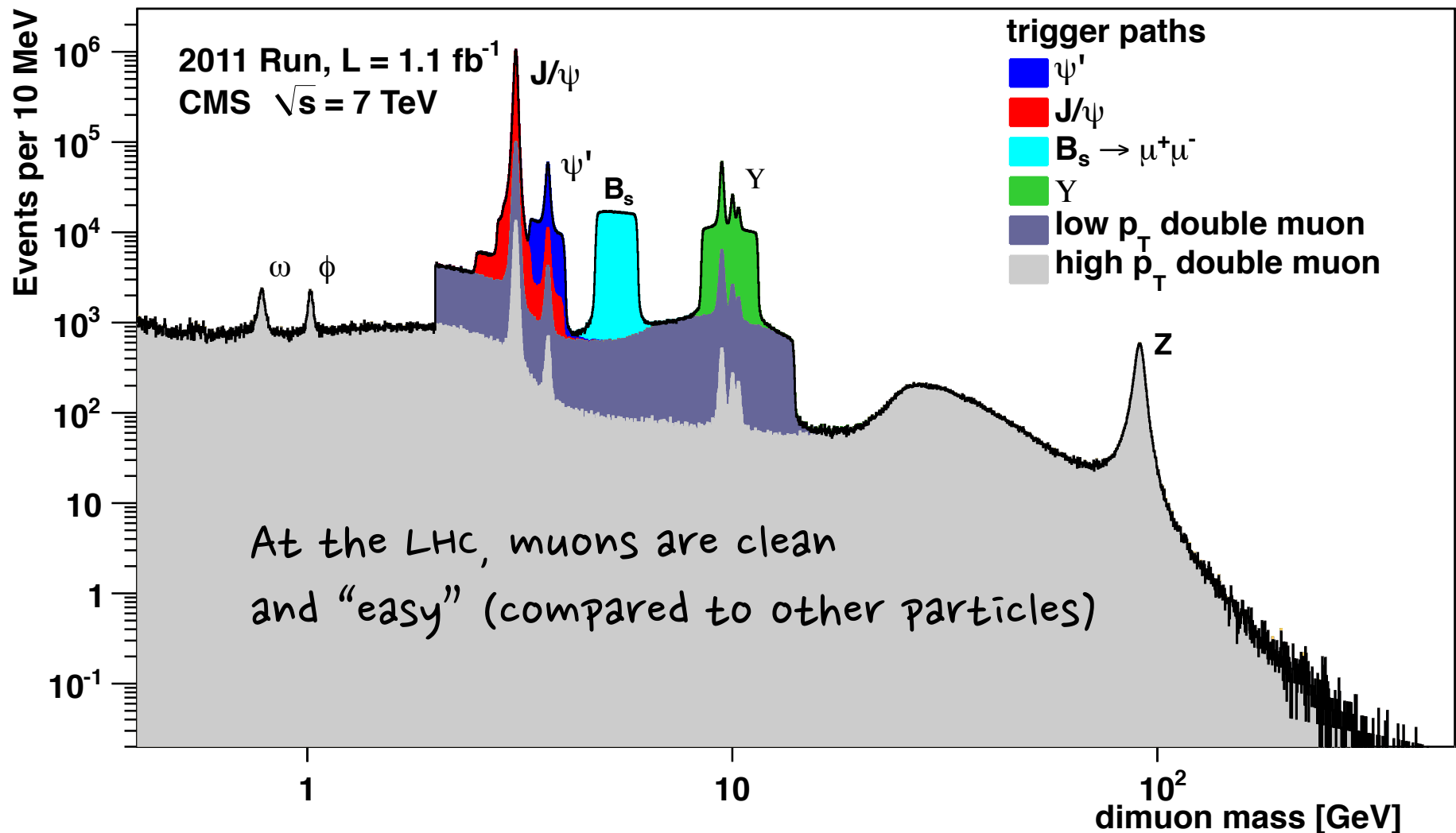
Photon yield and purity

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 - $N(\gamma) = N^A - \text{background}$
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 - $P = 1 - N^B/N^A \cdot M^A/M^B$



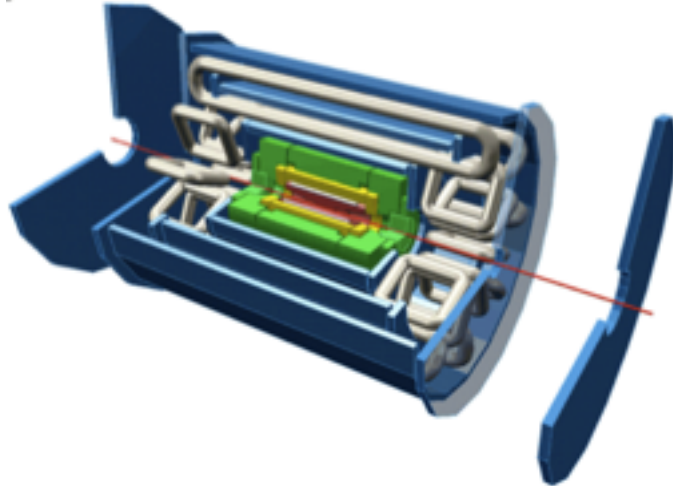
A spectroscopist's delight!

The power of muon identification!

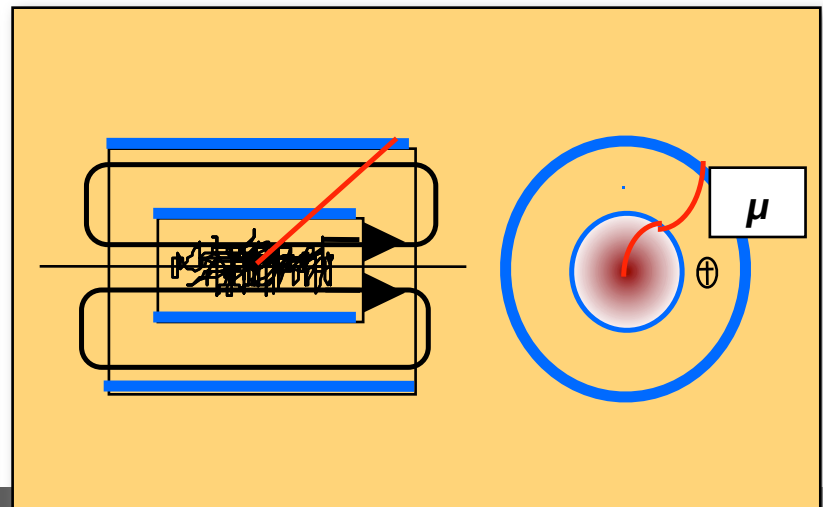
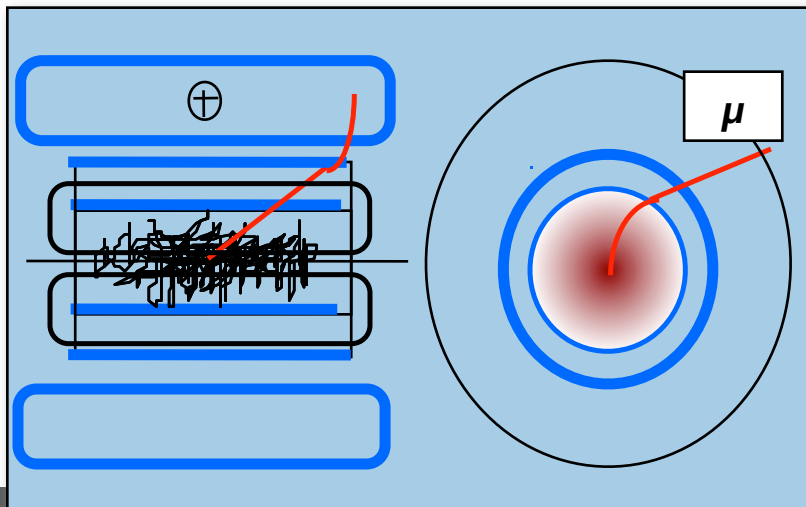
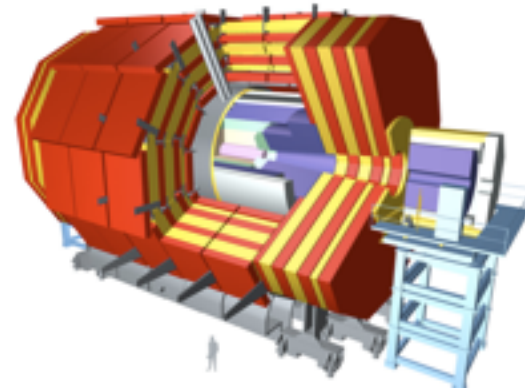


At the LHC, muons are clean and "easy" (compared to other particles)

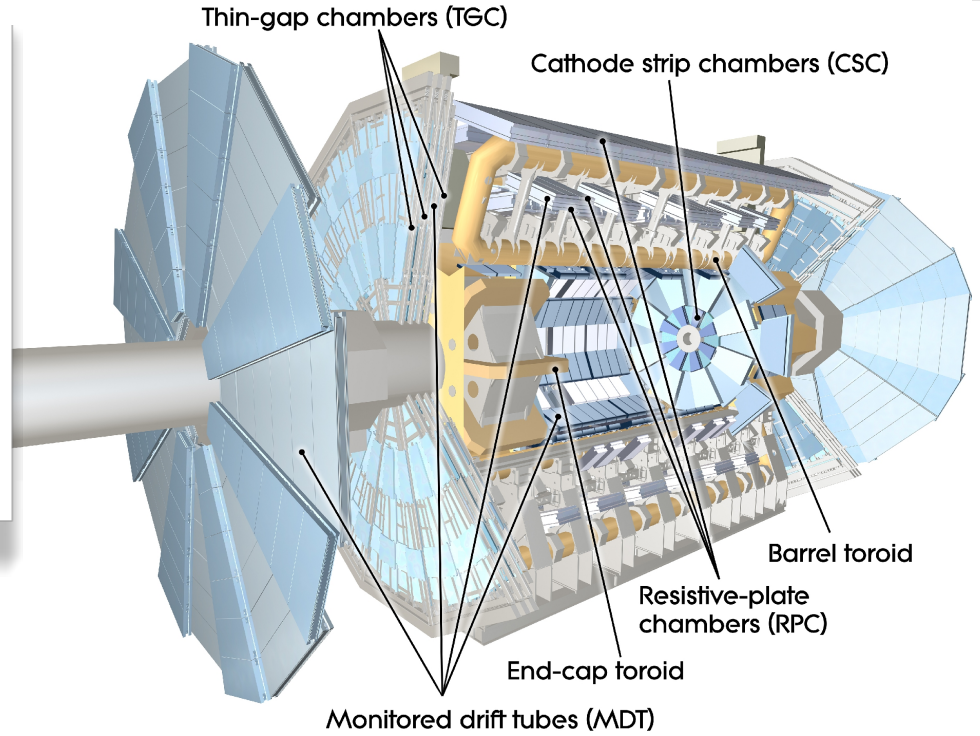
ATLAS



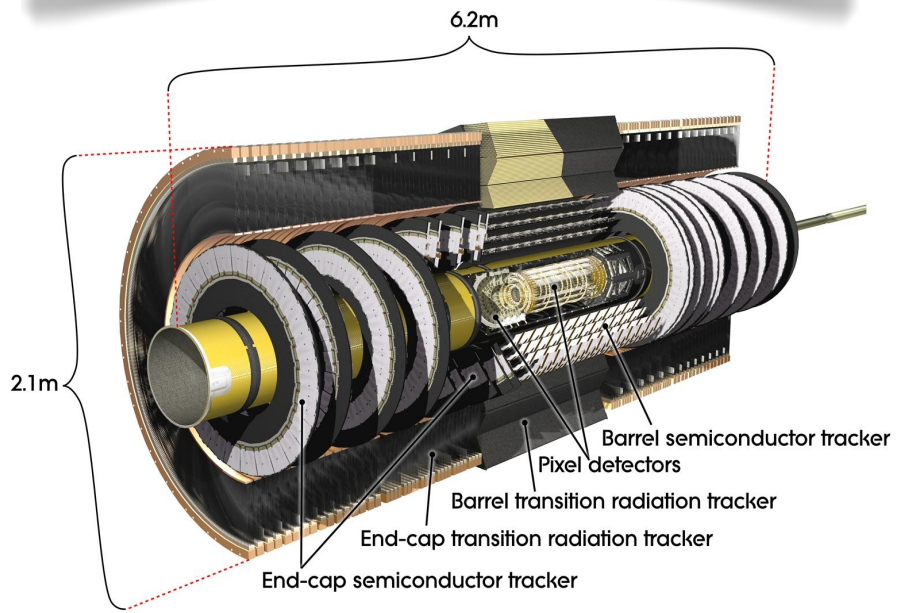
CMS



- Use a toroidal field (about 0.5T)
- **Precision chambers**
 - Monitored Drift Tubes in barrel and endcaps
3 layers for $|\eta| < 2.0$, 2 layers for $2.0 < |\eta| < 2.7$,
resolution of $35 \mu\text{m}$ per chamber
 - Cathode Strip Chambers :
1 layer (inner) for $2.0 < |\eta| < 2.7$,
resolution in precise coordinate of $40 \mu\text{m}$ per station
- **Trigger chambers**
 - Resistive Plate Chambers in barrel ($|\eta| < 1.05$),
1.5 ns of time resolution
 - Thin Gap Chambers in endcaps ($1.05 < |\eta| < 2.7$)



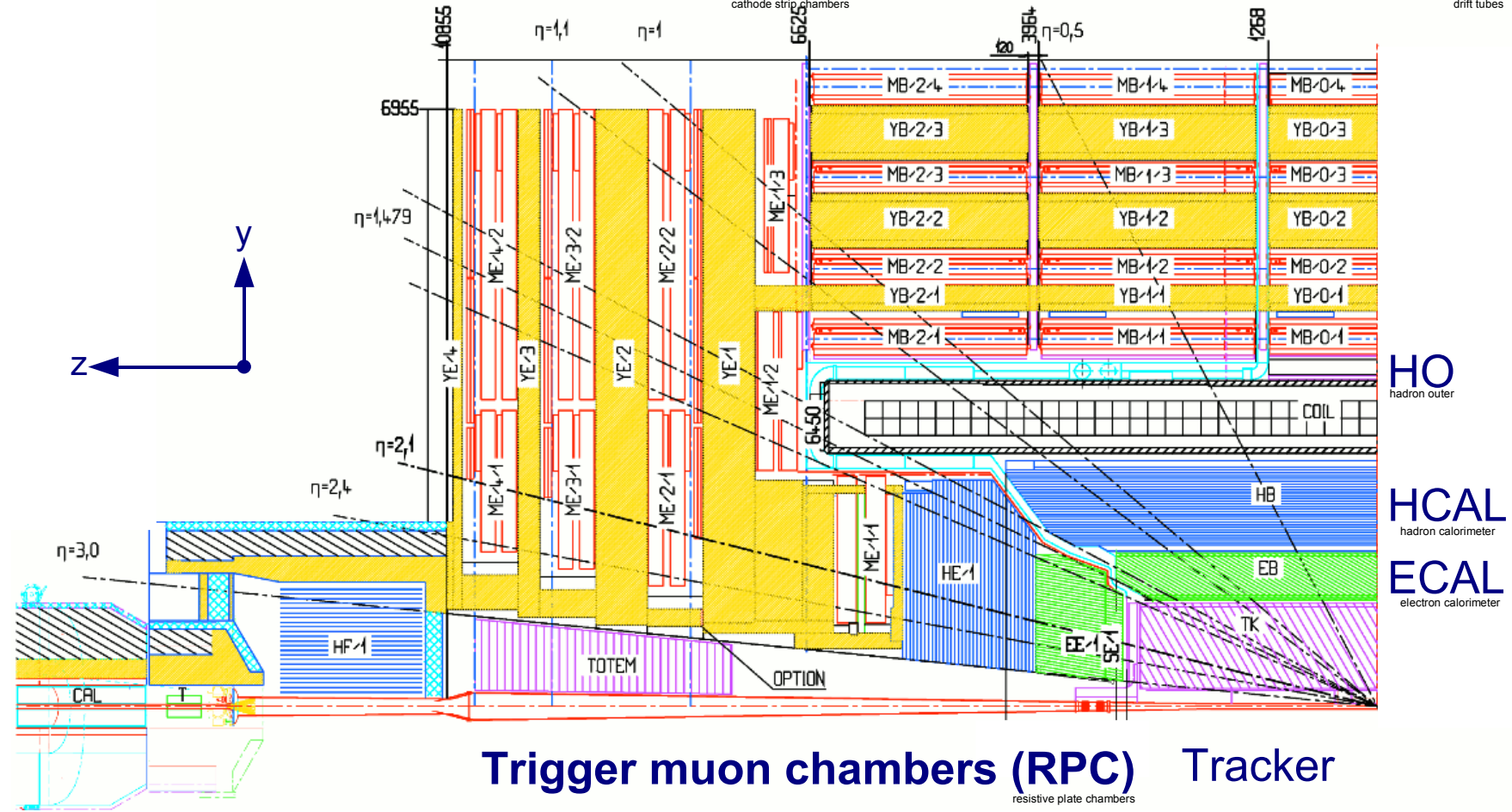
- Inside solenoid (2T)
- Pixels, SemiConductor Tracker, Transition Tradiation Tracker
- Cover $|\eta| \leq 2.5$ region, except TRT ($|\eta| \leq 2$)



Forward muon chambers (CSC) Central muon chambers (DT)

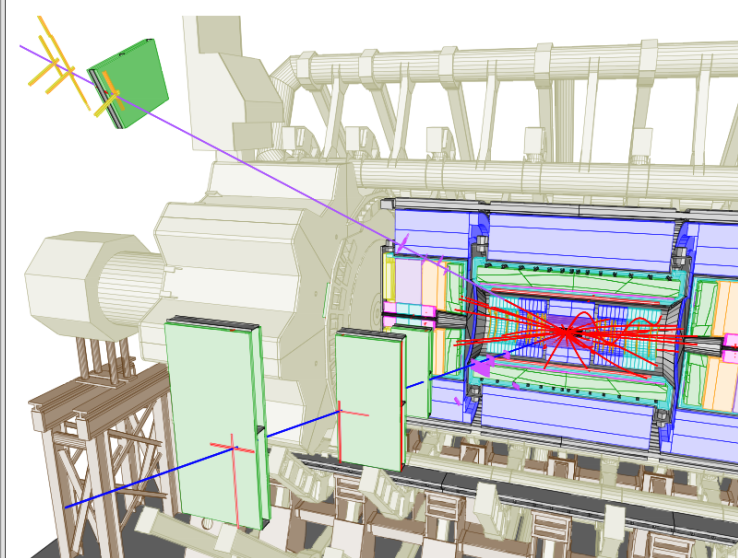
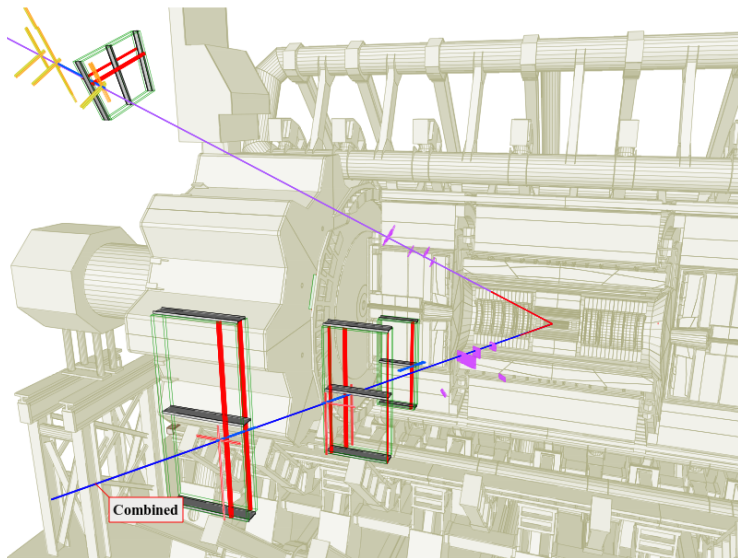
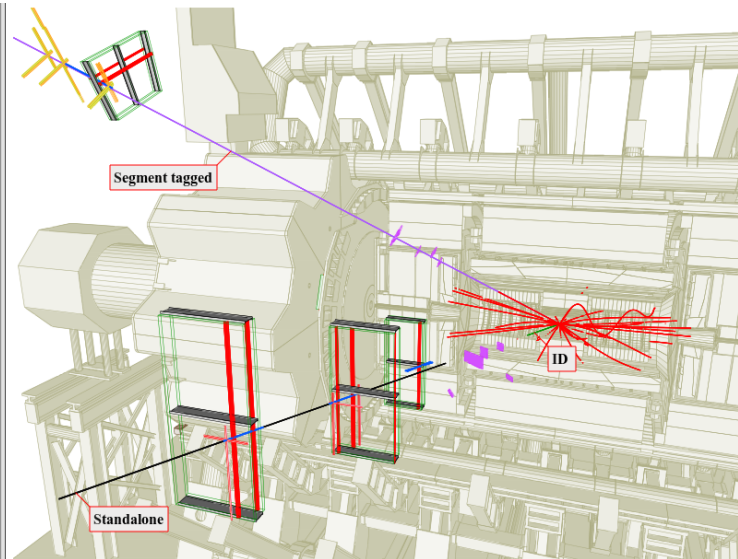
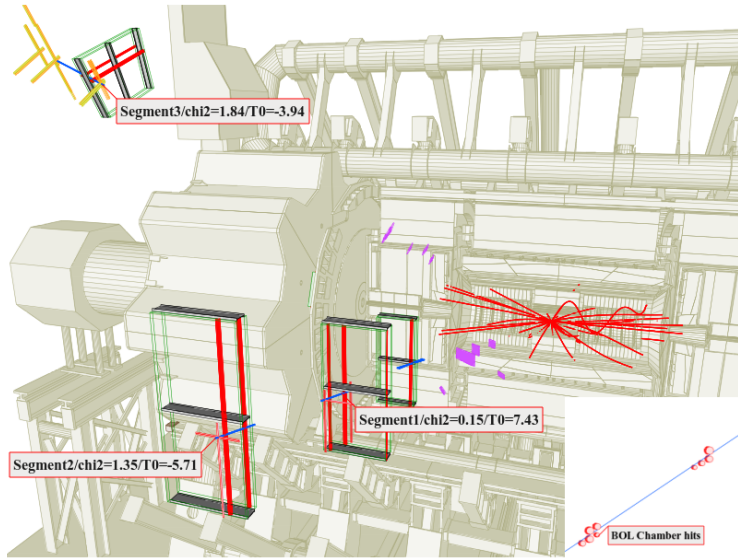
cathode strip chambers

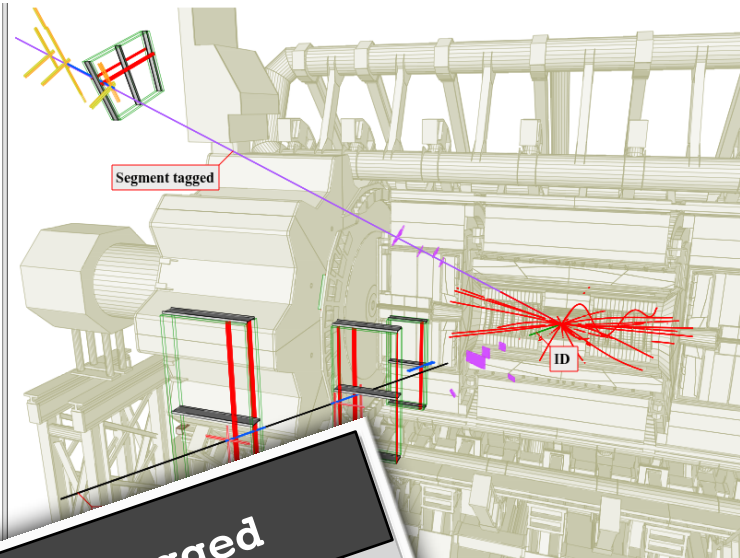
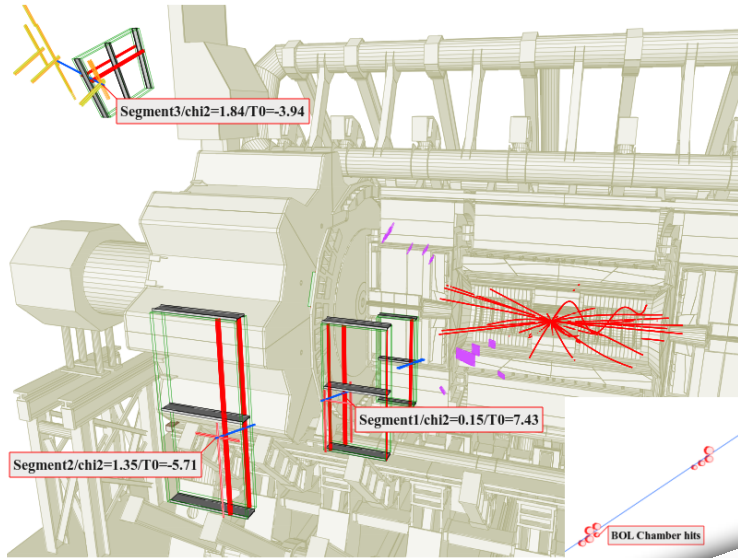
drift tubes



Trigger muon chambers (RPC) Tracker

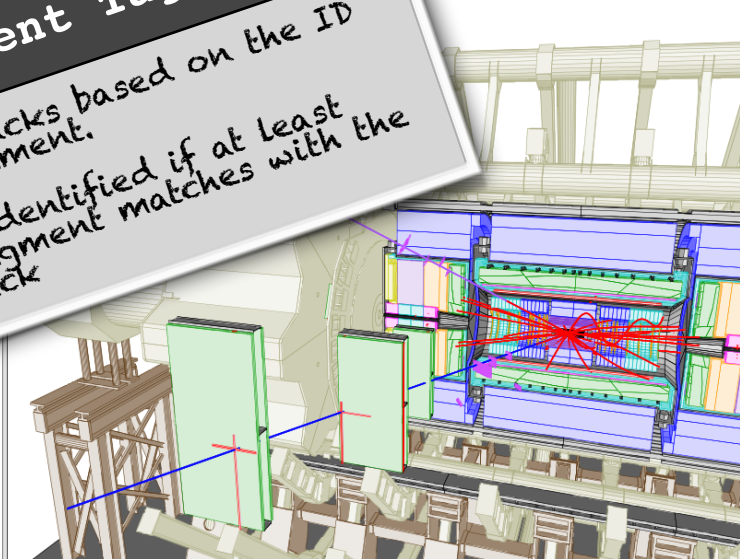
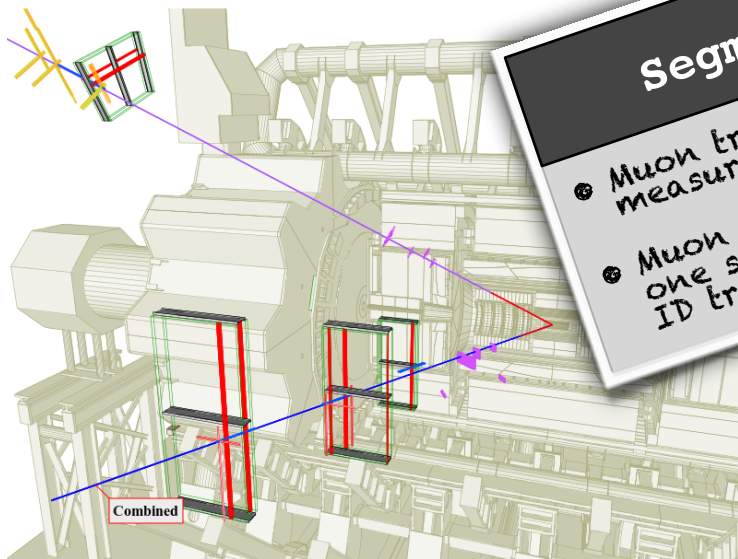
resistive plate chambers

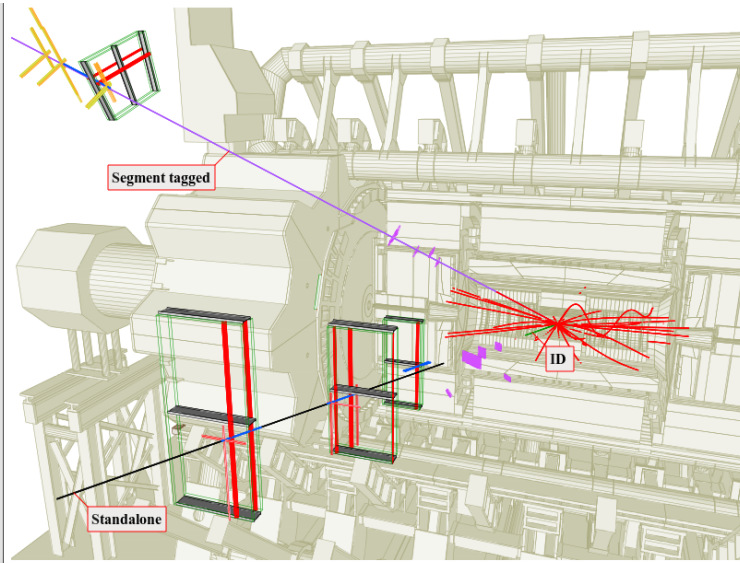
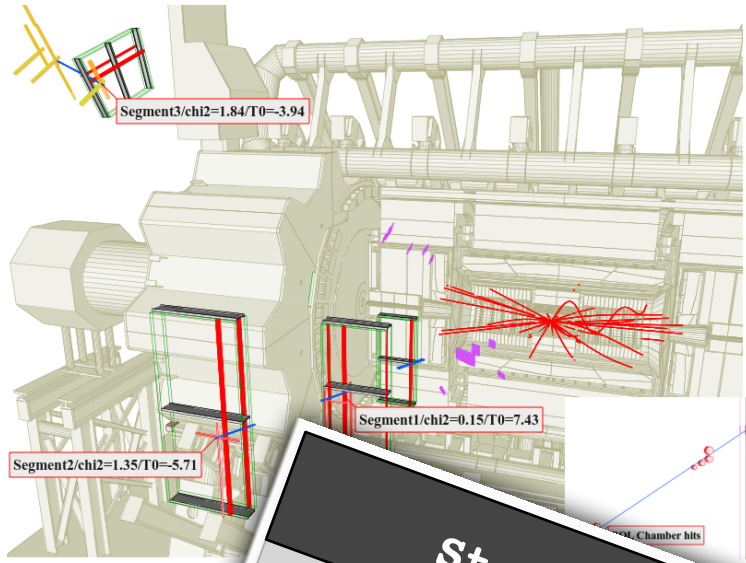




Segment Tagged

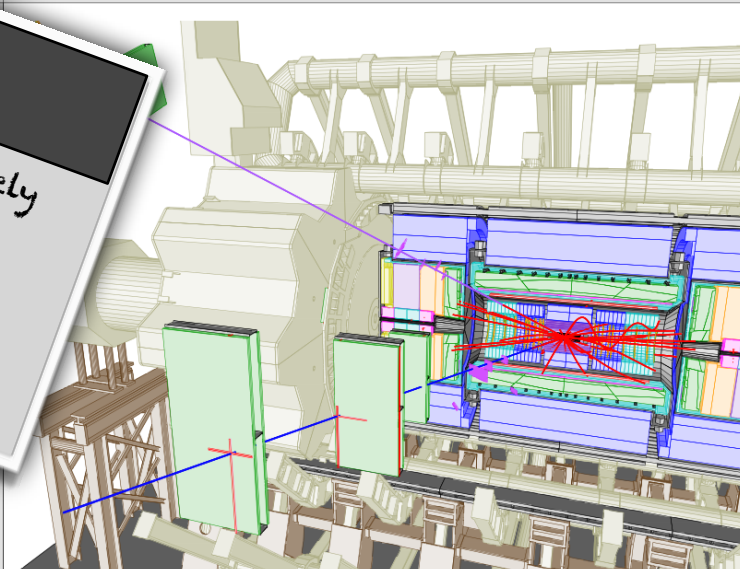
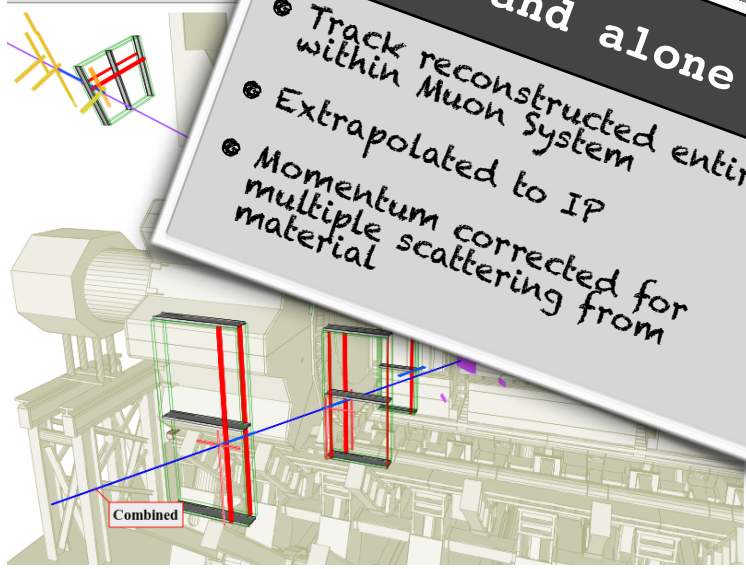
- Muon tracks based on the ID measurement.
- Muon identified if at least one segment matches with the ID track

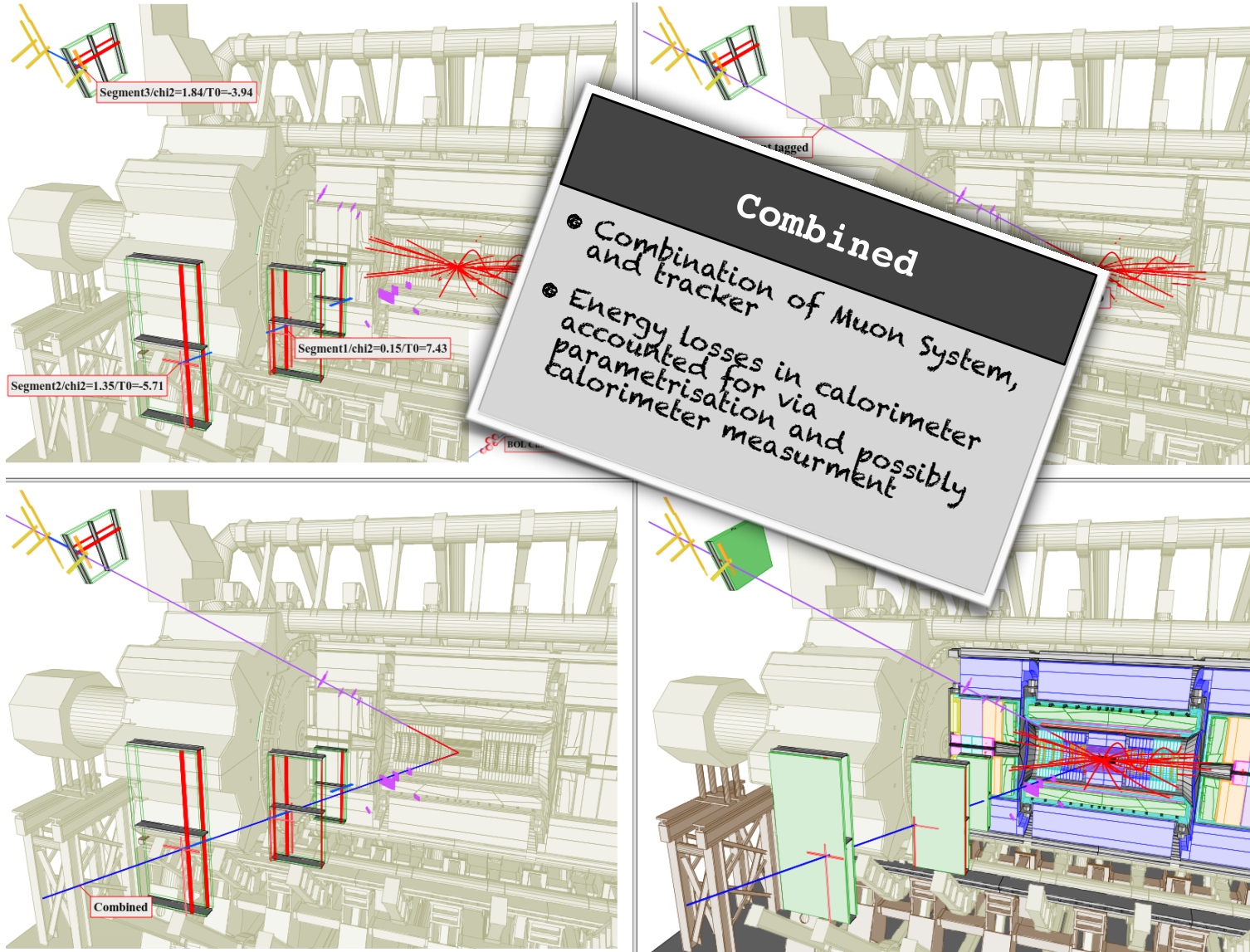


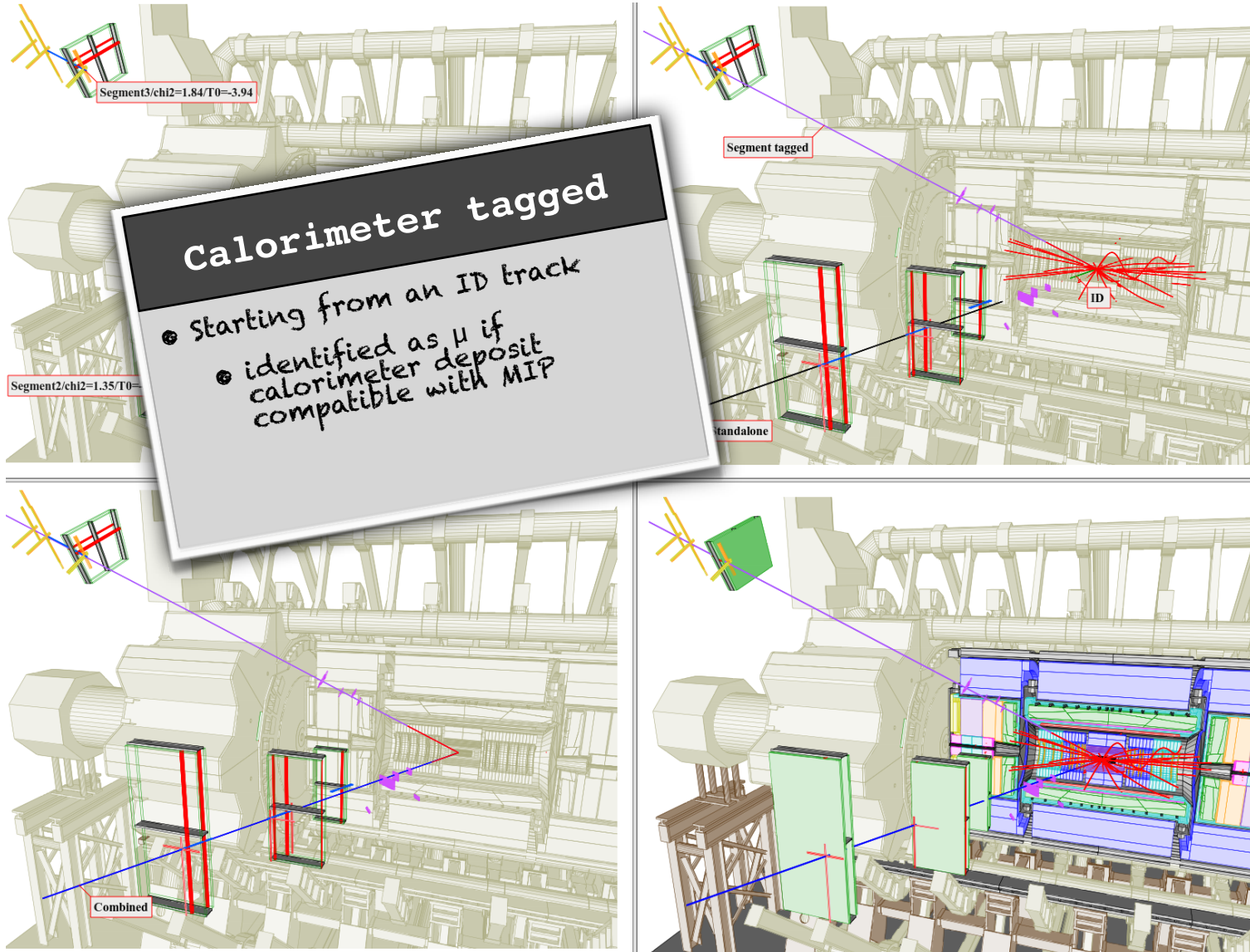


Stand alone

- Track reconstructed entirely within Muon System
- Extrapolated to IP
- Momentum corrected for multiple scattering from material

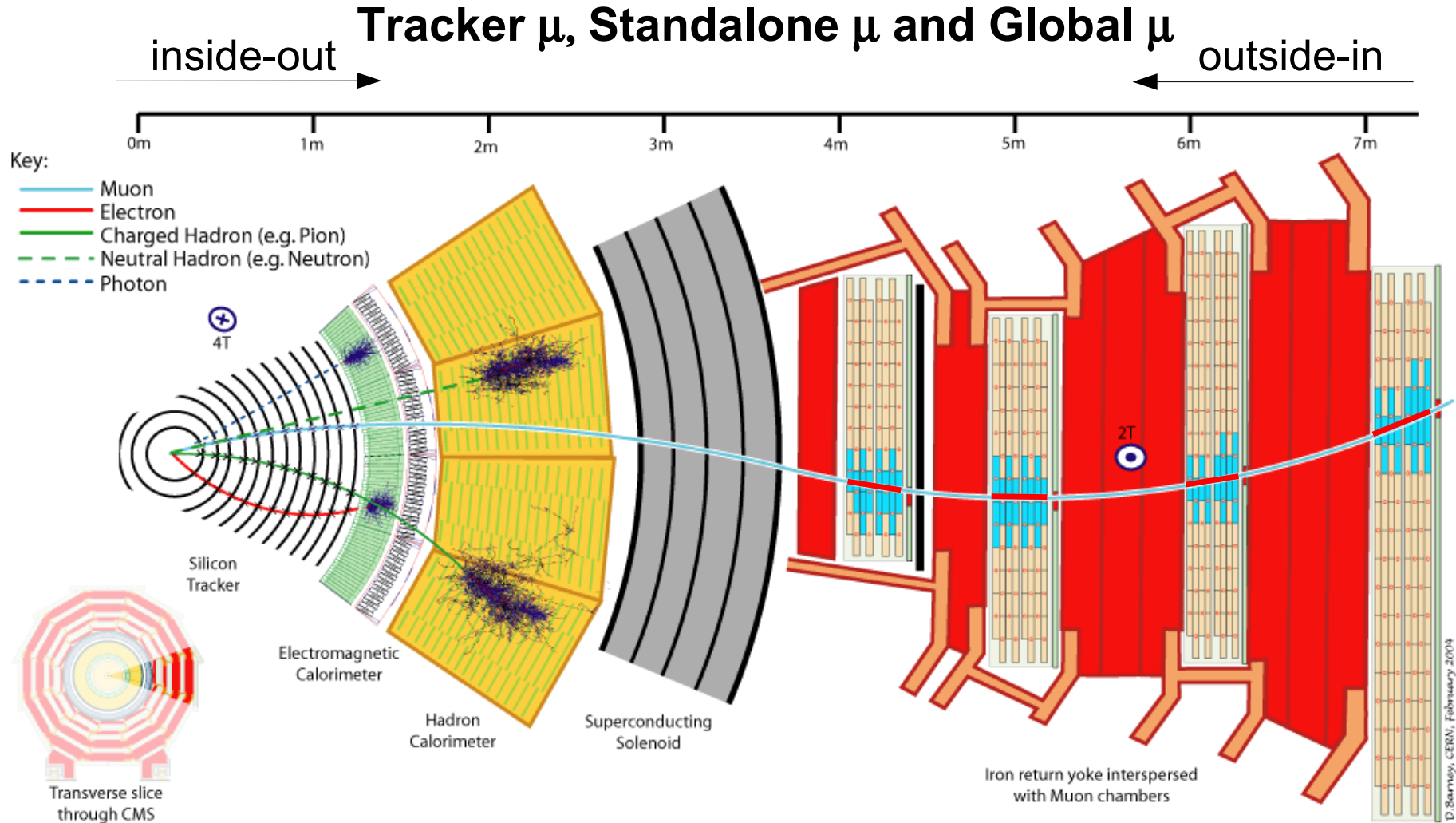


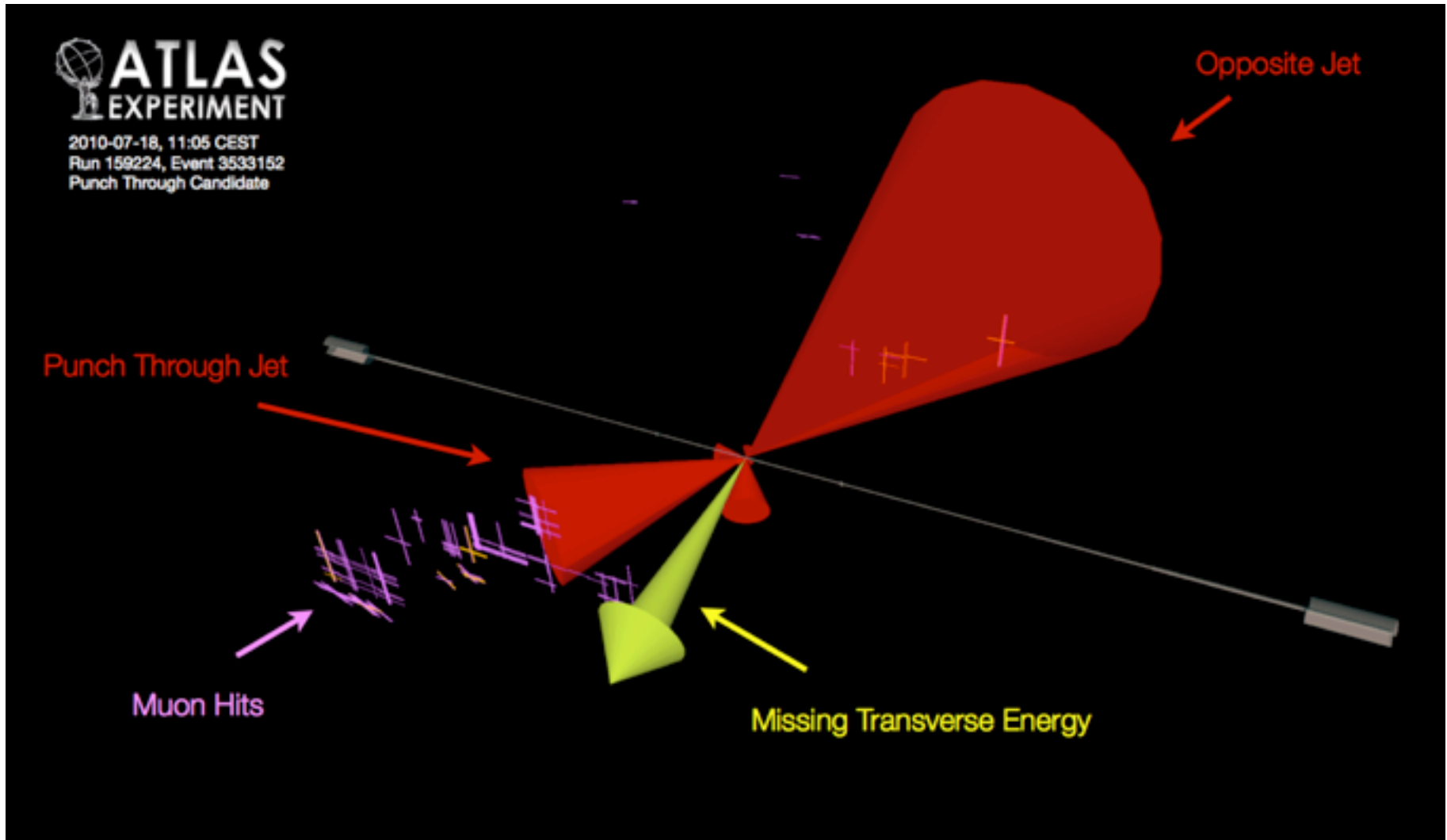






Muon Identification





ATLAS EXPERIMENT

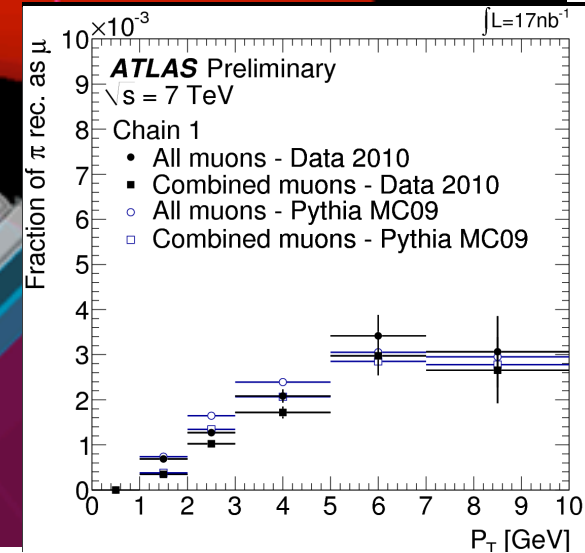
2010-07-18, 11:05 CEST
 Run 159224, Event 3533152
 Punch Through Candidate

Muon Hits

Missing Transverse Energy

Punch Through Jet

Opposite Jet





Summary & Outlook

- We now have a complete list of identified individual particles
 - $e, \gamma, \mu, \pi, K_L^0, PU-\pi$
 - This list of particles describes the entire event
 - all detector hits are used; redundancy exploited;
 - unused energy avoided; double counting of energy avoided
 - Some of these particles can be identified as prompt
 - we discussed electrons, photons, muons
 - pile-up can be removed from isolation consideration
- Next, we will use the above list of particles to identify composite or unstable particles
 - hadronic decays of τ -lepton, quark/gluon jets, b-jets, t-jets, and ν 's
- More tomorrow!