



$H \rightarrow \gamma\gamma$ search and direct photon pair production differential cross section measurement at DØ

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URA thesis award talk

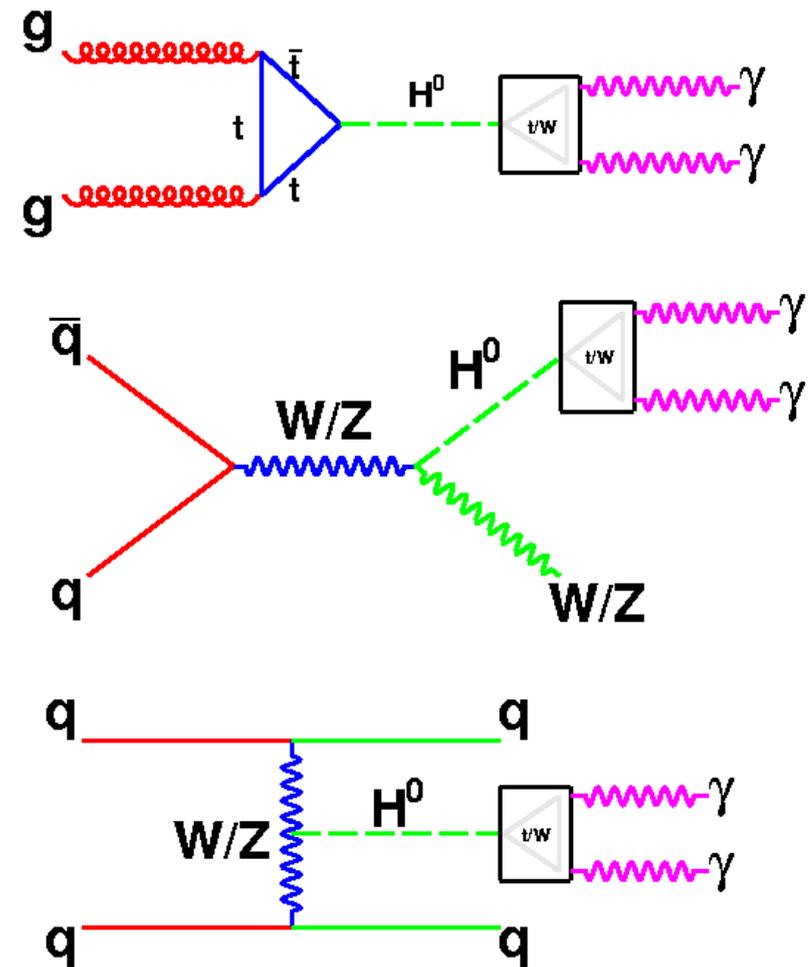


Why search for $H \rightarrow \gamma\gamma$ at Tevatron ?

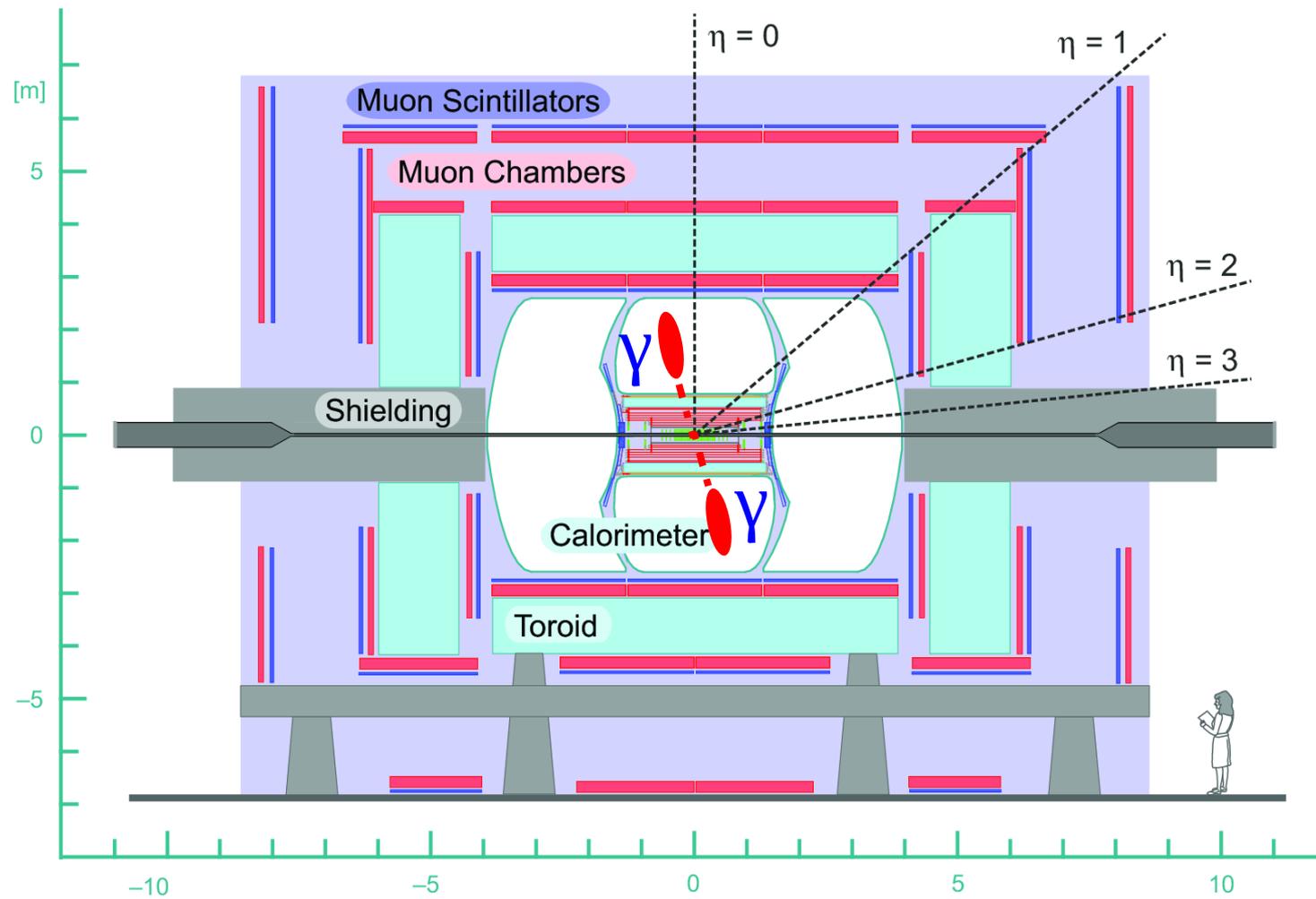
- Higgs boson is undiscovered !
- Clean signature.
- Current direct experimental search + indirect constraint prefer light Higgs boson.
- Contribute to the Tevatron Standard Model Higgs combination, especially in the intermediate mass region around 125 GeV.
- Golden channel for the discovery of Standard Model Higgs at LHC.

$H \rightarrow \gamma\gamma$

- Standard Model Higgs:
 - Gluon fusion
 - Associated production
 - Vector boson fusion
- Examine the inclusive di-photon dataset ($\gamma\gamma+X$) to search for the high mass resonance.

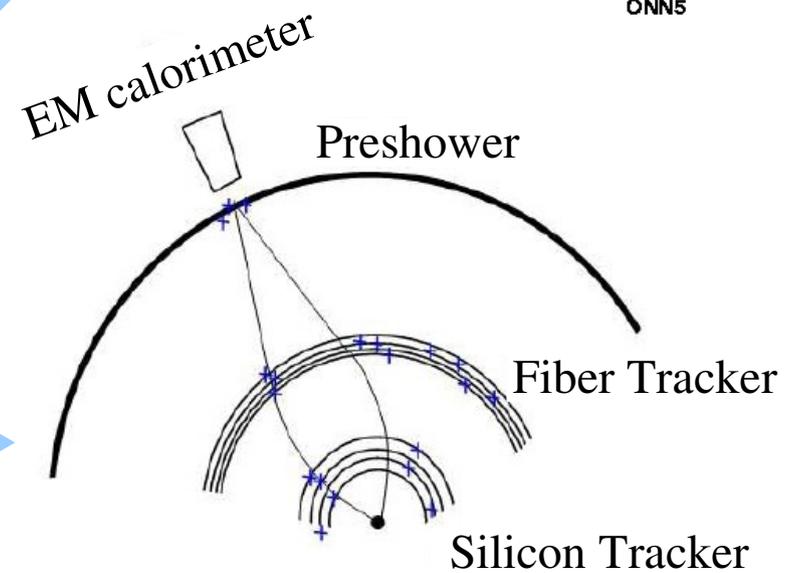
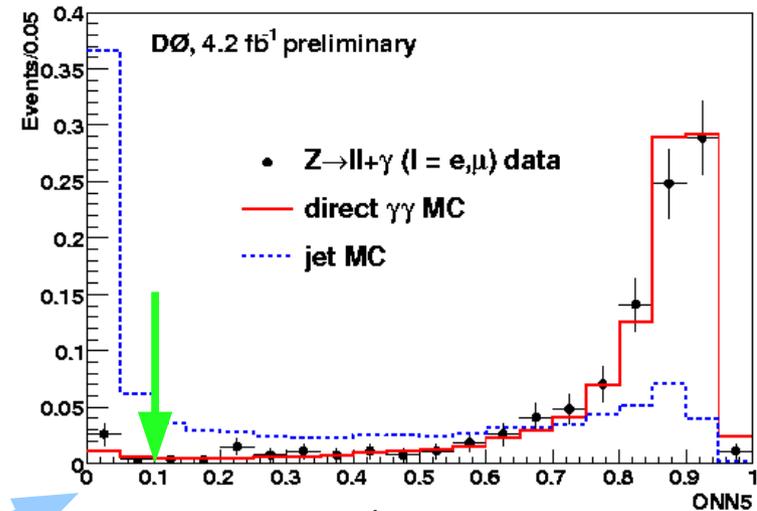


DØ detector



Event selection

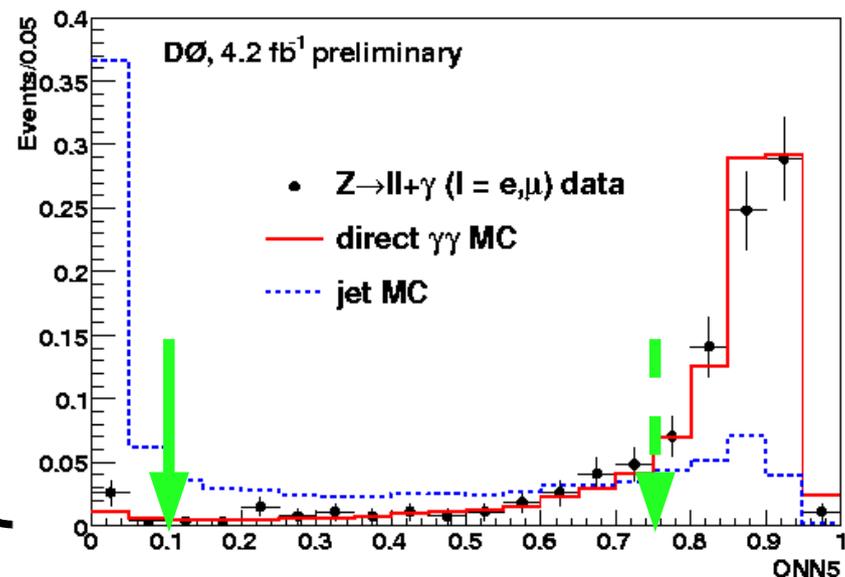
- Select 2 photons
 - Transverse energy $E_T > 25$ GeV
 - Isolated in calorimeter
 - Isolated in tracker
 - Energy shower shape consistent with photon
 - **Neural network discriminant (ONN) > 0.1**
 - No spatially well-matched tracks
 - **No pattern of hits in the tracker in a road be consistent with an electron**



Backgrounds

- Reducible background
 - $Z/\gamma^* \rightarrow ee$, both electrons are misidentified as photons, **estimated with Geant Monte Carlo.**
 - **Non- $\gamma\gamma$** (γ +jet,jet+jet), when the jet(s) is(are) misidentified as photon(s), **estimated from data.**
- Irreducible background – **direct $\gamma\gamma$**
 - **estimated from data.**

- Using $ONN=0.75$ as the boundary to separate the events into 4 categories:
 - N_{pp} : both pass the $ONN>0.75$
 - N_{pf} : first passes, second fails
 - N_{fp} : vice-versa
 - N_{ff} : both fail



$$\begin{pmatrix} N_{ff} \\ N_{fp} \\ N_{pf} \\ N_{pp} \end{pmatrix} = E \times \begin{pmatrix} N_{jj} \\ N_{j\gamma} \\ N_{\gamma j} \\ N_{\gamma\gamma} \end{pmatrix}$$

$\left\{ \begin{array}{l} \epsilon_{j1}, \epsilon_{j2} \text{ are jet ONN}>0.75 \text{ efficiencies,} \\ \text{estimated from MC, validated from data.} \\ \epsilon_{\gamma 1}, \epsilon_{\gamma 2} \text{ are photon ONN}>0.75 \text{ efficiencies,} \\ \text{estimated from MC, corrected from data.} \end{array} \right.$

$$E = \begin{pmatrix} (1 - \epsilon_{j1})(1 - \epsilon_{j2}) & (1 - \epsilon_{j1})(1 - \epsilon_{\gamma 2}) & (1 - \epsilon_{\gamma 1})(1 - \epsilon_{j2}) & (1 - \epsilon_{\gamma 1})(1 - \epsilon_{\gamma 2}) \\ (1 - \epsilon_{j1})\epsilon_{j2} & (1 - \epsilon_{j1})\epsilon_{\gamma 2} & (1 - \epsilon_{\gamma 1})\epsilon_{j2} & (1 - \epsilon_{\gamma 1})\epsilon_{\gamma 2} \\ \epsilon_{j1}(1 - \epsilon_{j2}) & \epsilon_{j1}(1 - \epsilon_{\gamma 2}) & \epsilon_{\gamma 1}(1 - \epsilon_{j2}) & \epsilon_{\gamma 1}(1 - \epsilon_{\gamma 2}) \\ \epsilon_{j1}\epsilon_{j2} & \epsilon_{j1}\epsilon_{\gamma 2} & \epsilon_{\gamma 1}\epsilon_{j2} & \epsilon_{\gamma 1}\epsilon_{\gamma 2} \end{pmatrix}$$

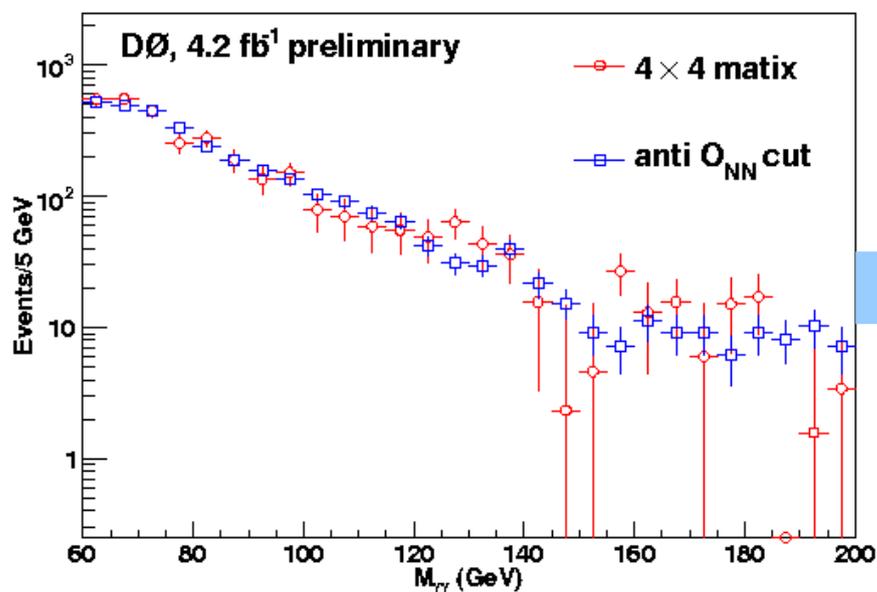
Total	7939
Total - N_{DY}	7722
$N_{\gamma\gamma}$	4538 ± 144
$N_{\gamma j} + N_{j\gamma}$	2189 ± 170
N_{jj}	994 ± 106
non- $\gamma\gamma$	3183 ± 200

➡ Includes potential signal.

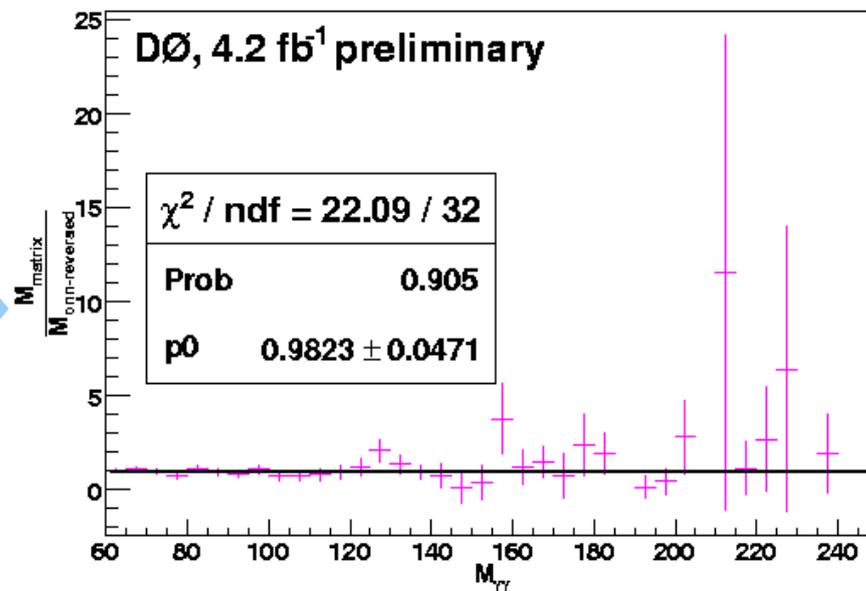
The quoted uncertainties are statistical only.

Non- $\gamma\gamma$ (γ +jet, jet+jet) background

- **Shape** from reversing the ONN cut (0.1) for one photon candidate.
- **Normalization** to the number of events from 4X4 matrix method.

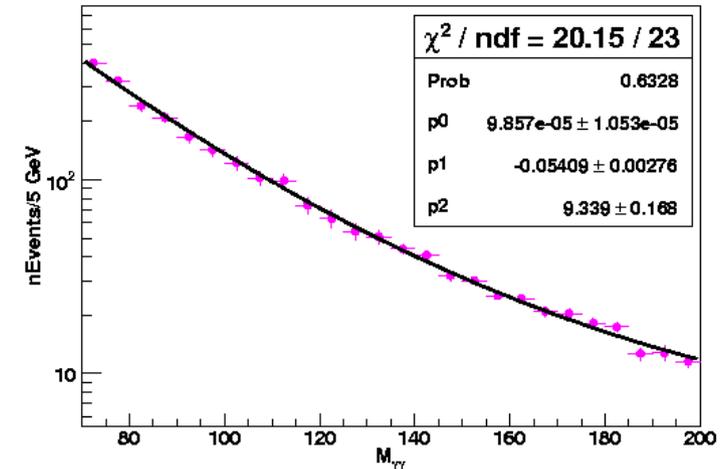


Ratio

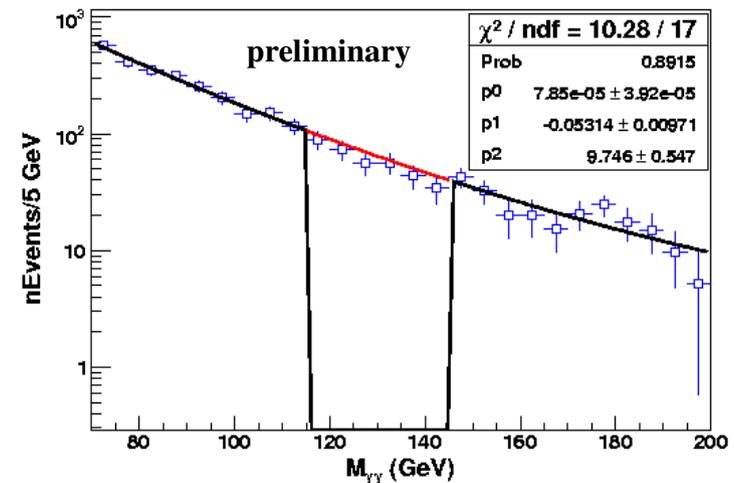


Direct $\gamma\gamma$ background

- Estimated from data using **side-band** fitting after the subtraction of the reducible background ($Z \rightarrow ee$, non- $\gamma\gamma$).
- The exponential fitting function is validated using NLO predictions.

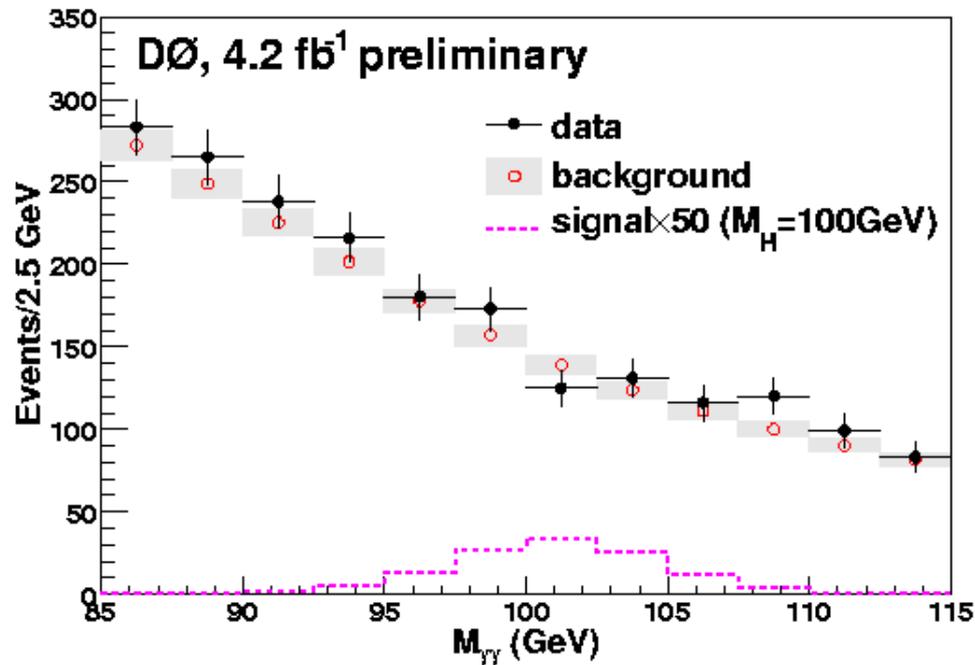


From NLO prediction

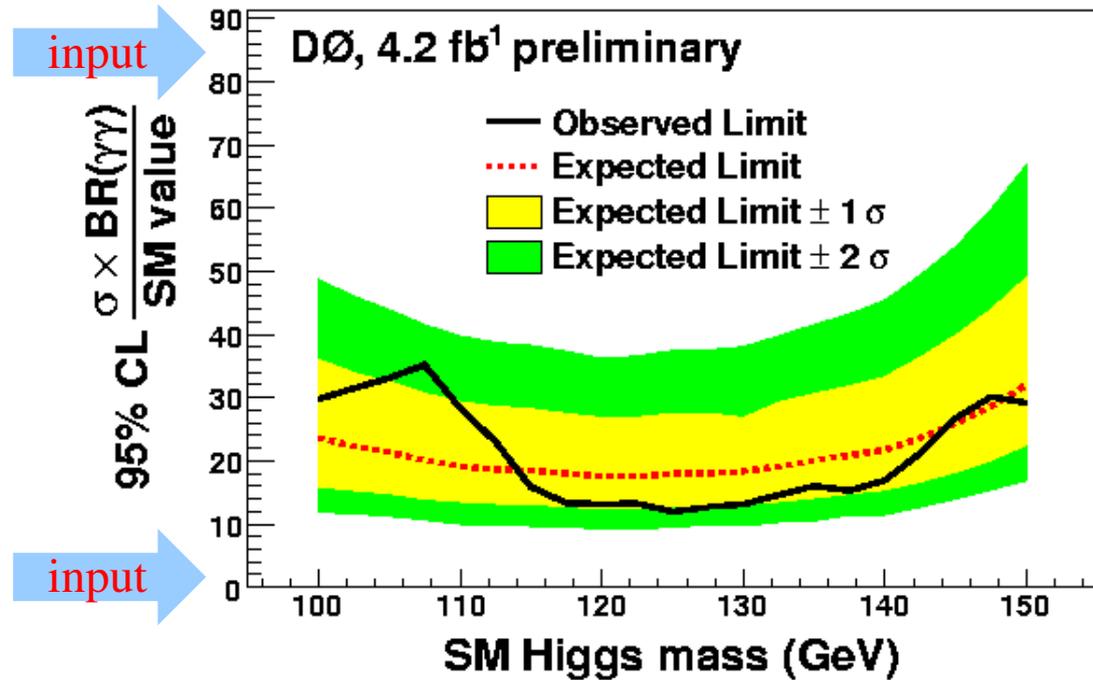


From subtracted data

The results



$M_H = 100 \text{ GeV}$ as an example !

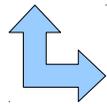




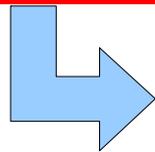
Direct photon pair production cross section

- Check the validity of the predictions of perturbative quantum chromodynamics (pQCD) and soft-gluon resummation methods implemented in theoretical calculations.
- Dominant background for the $H \rightarrow \gamma\gamma$ search
 - The sensitivity could be enhanced if other kinematic information, e.g. the di-photon transverse momentum, included.

$$\frac{d\sigma}{dX} = \frac{N_{data} - N_{bkg}}{L \cdot Acc \cdot \epsilon_{trigger} \cdot \epsilon_{sel} \cdot \Delta}$$

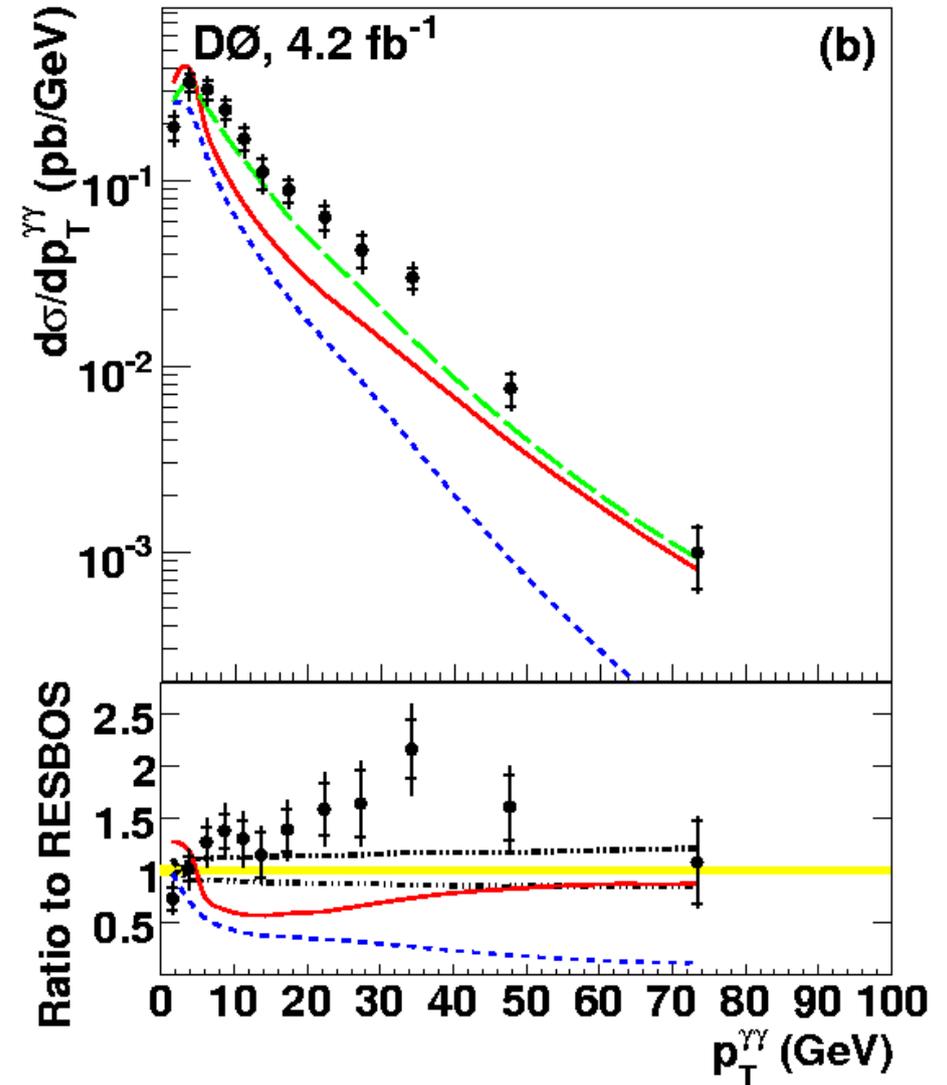
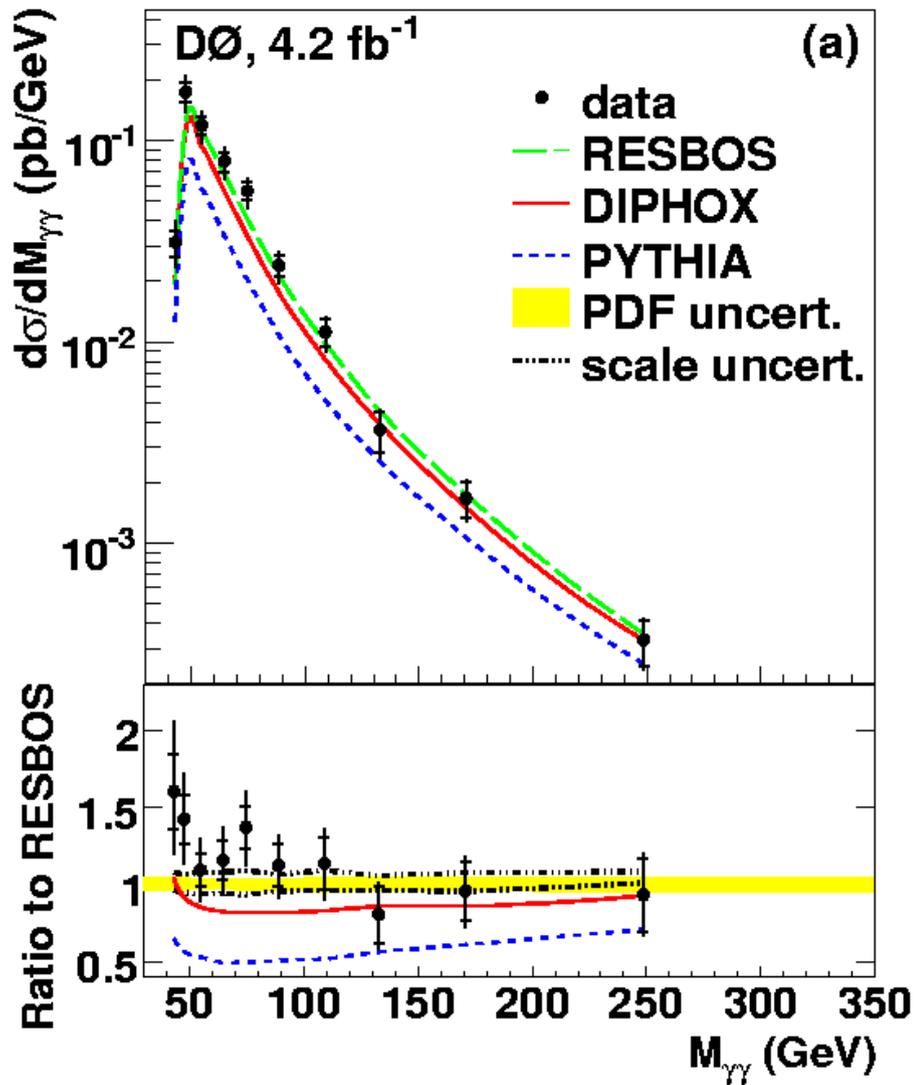


$$\underline{M_{\gamma\gamma}, p_T^{\gamma\gamma}, \Delta\phi_{\gamma\gamma}, \cos\theta_{\gamma\gamma}}$$

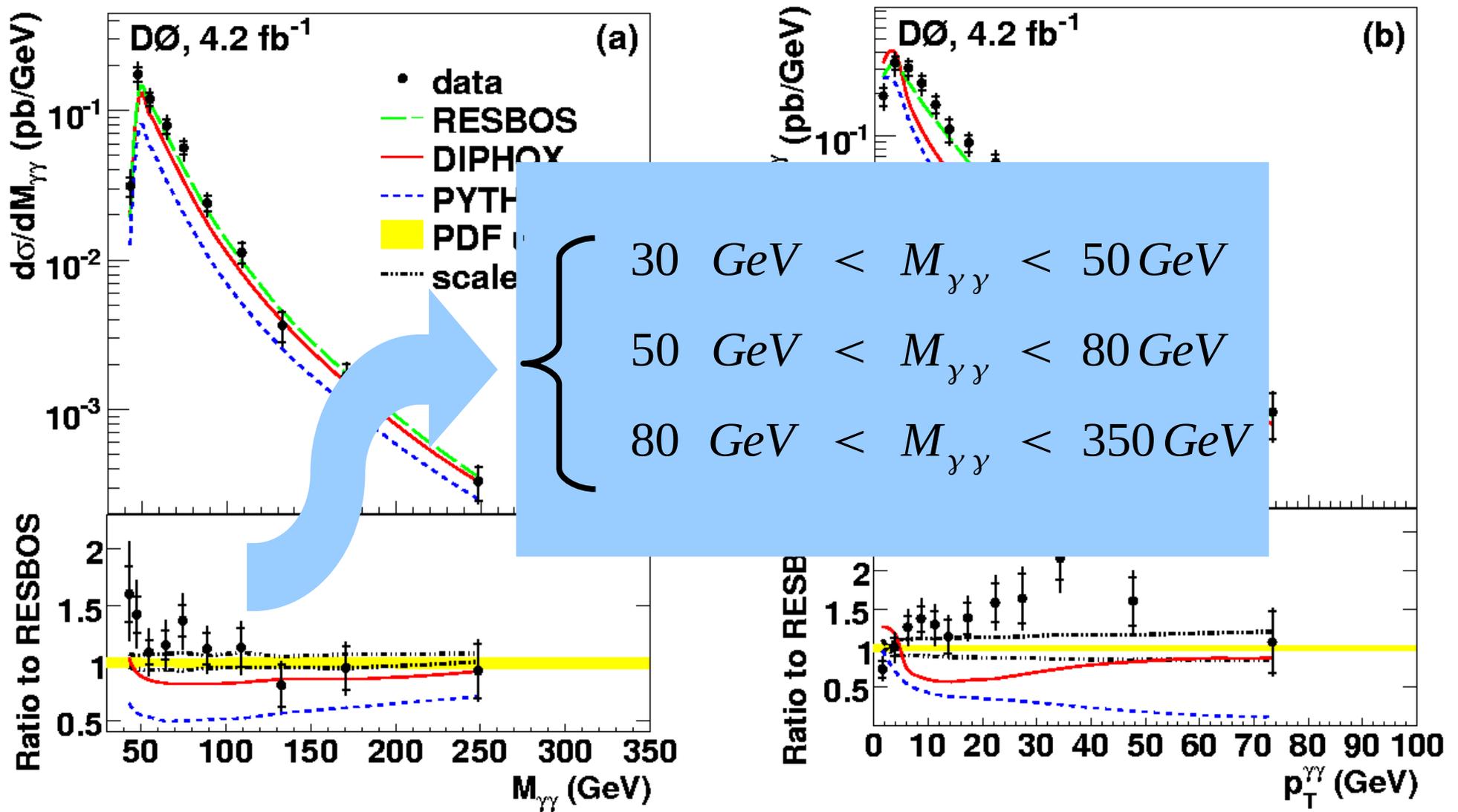


data vs. **resbos**, **diphox** and **pythia**

Differential cross sections



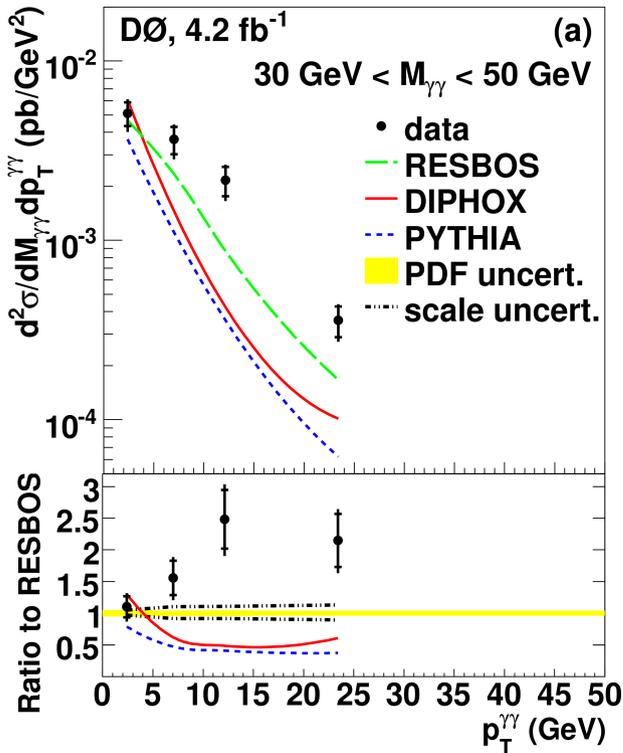
Differential cross sections



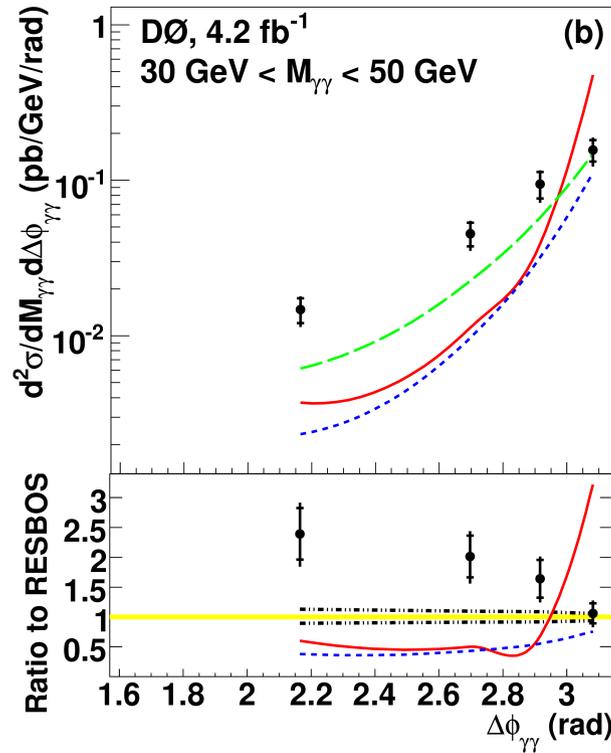
Differential cross sections in

$$30 \text{ GeV} < M_{\gamma\gamma} < 50 \text{ GeV}$$

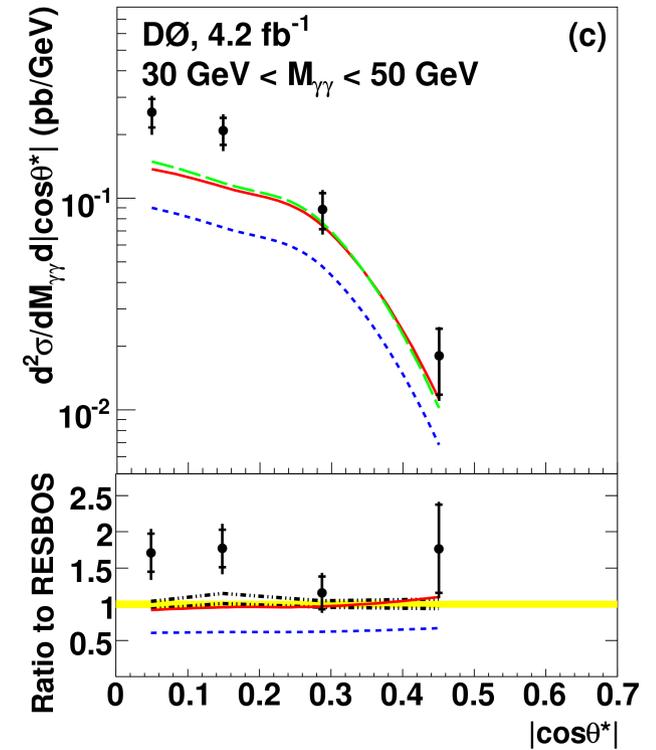
$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot dp_T^{\gamma\gamma}}$$



$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot d\Delta\phi_{\gamma\gamma}}$$



$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot d\cos\theta^*}$$

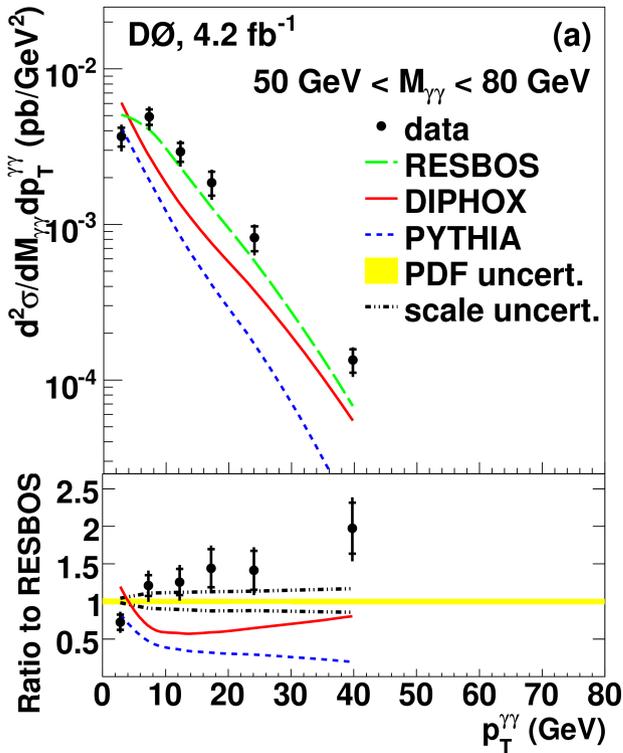


Differential cross sections in

$$50 \text{ GeV} < M_{\gamma\gamma} < 80 \text{ GeV}$$

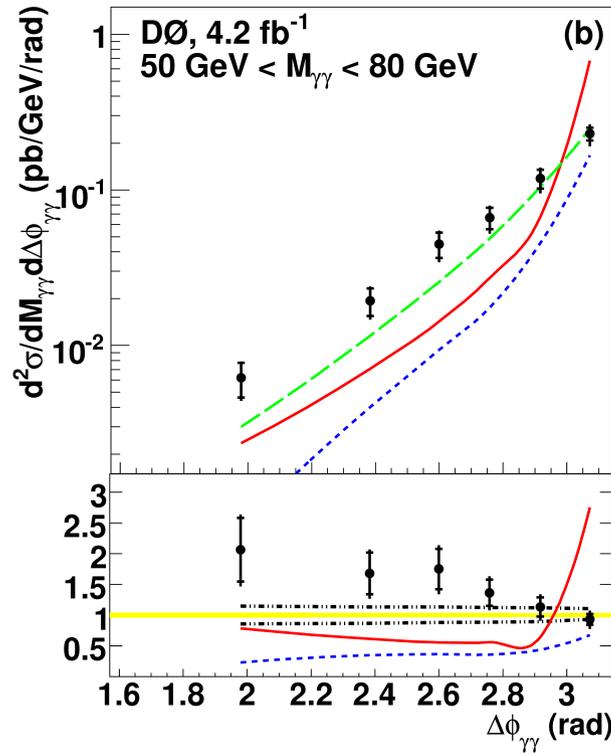
$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot dp_T^{\gamma\gamma}}$$

$$dM_{\gamma\gamma} \cdot dp_T^{\gamma\gamma}$$



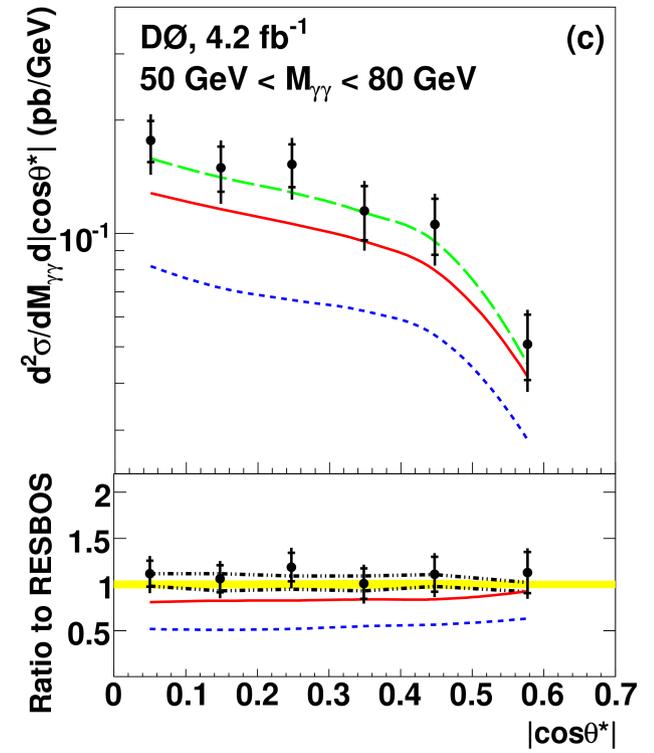
$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot d\Delta\phi_{\gamma\gamma}}$$

$$dM_{\gamma\gamma} \cdot d\Delta\phi_{\gamma\gamma}$$



$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot d\cos\theta}$$

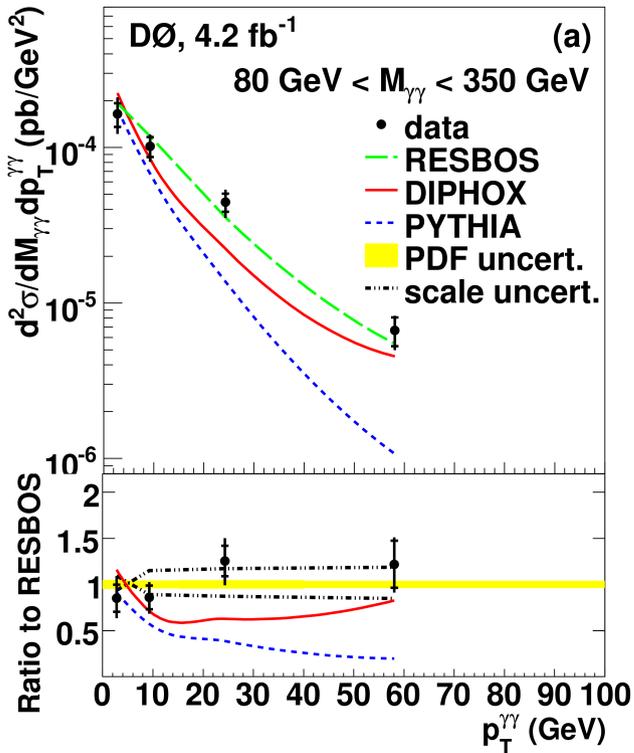
$$dM_{\gamma\gamma} \cdot d\cos\theta$$



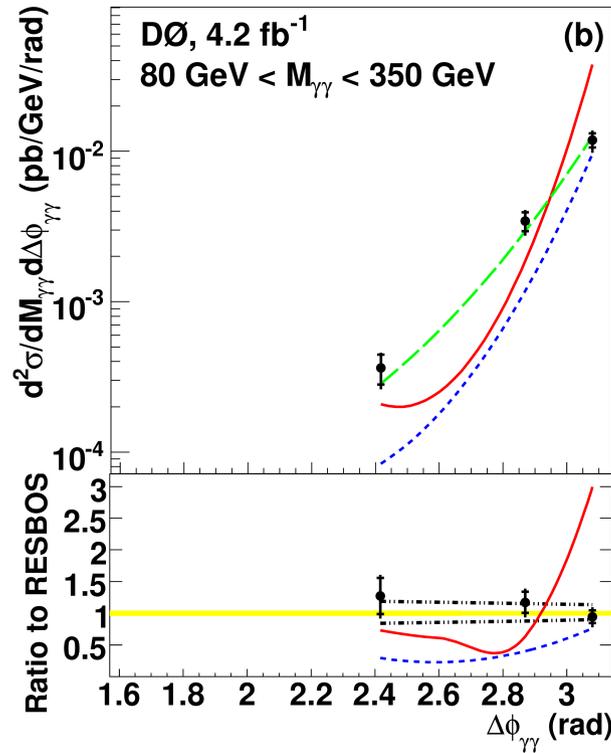
Differential cross sections in

$$80 \text{ GeV} < M_{\gamma\gamma} < 350 \text{ GeV}$$

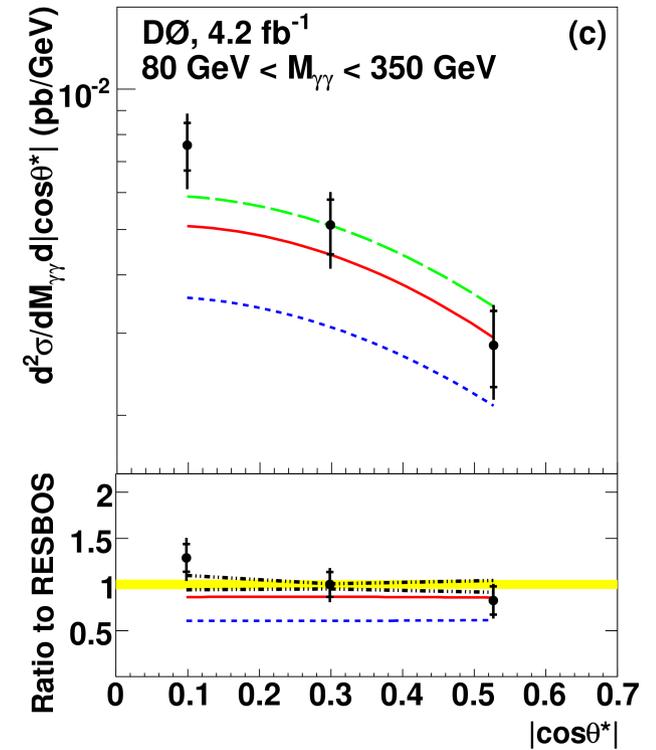
$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot dp_T^{\gamma\gamma}}$$



$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot d\Delta\phi_{\gamma\gamma}}$$



$$\frac{d^2\sigma}{dM_{\gamma\gamma} \cdot d\cos\theta}$$



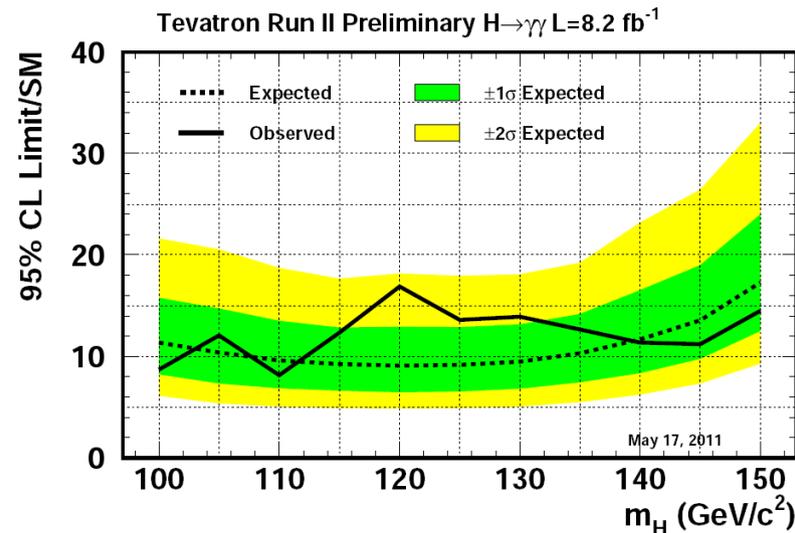


Conclusions

- We presented the measurements of single and double differential cross sections for direct photon pair production:
 - Show necessity to include higher order corrections beyond NLO and resummation to all orders of soft and collinear initial state gluons to describe di-photon production.

Conclusions

- We presented the measurements of single and double differential cross sections for direct photon pair production:
 - Show necessity to include higher order corrections beyond NLO and resummation to all orders of soft and collinear initial state gluons to describe di-photon production.
- We performed the first Higgs $\rightarrow \gamma\gamma$ search at Tevatron. Setting limits a factor of ~ 15 above Standard Model prediction at 130 GeV.
- Recent Tevatron Higgs $\rightarrow \gamma\gamma$ search results:



back-up

SM Higgs at Tevatron

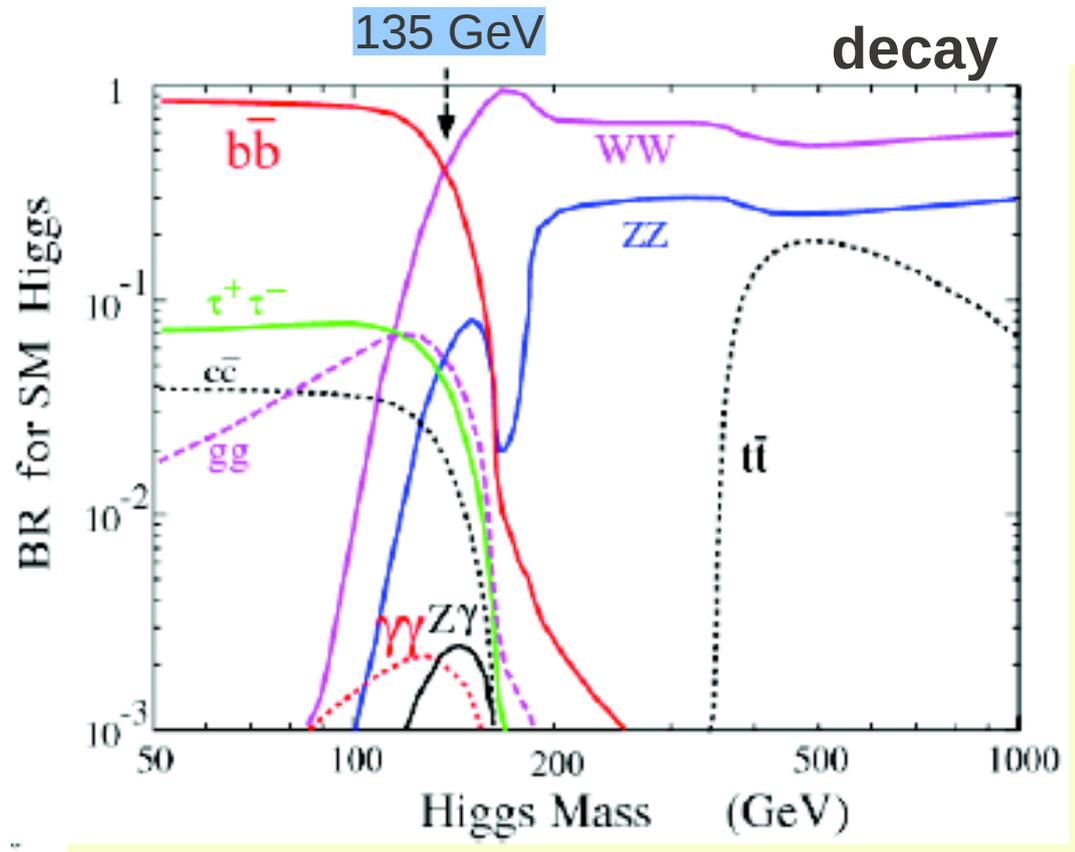
➤ Major decay channels

$H \rightarrow b\bar{b}$ @ $M_H < 135$ GeV

$H \rightarrow WW$ @ $M_H > 135$ GeV

➤ Current results (limits @95% CL)

- SM LEP direct search: $M_H > 114.4$ GeV
- SM indirect constraint + LEP direct search: $M_H < 185$ GeV

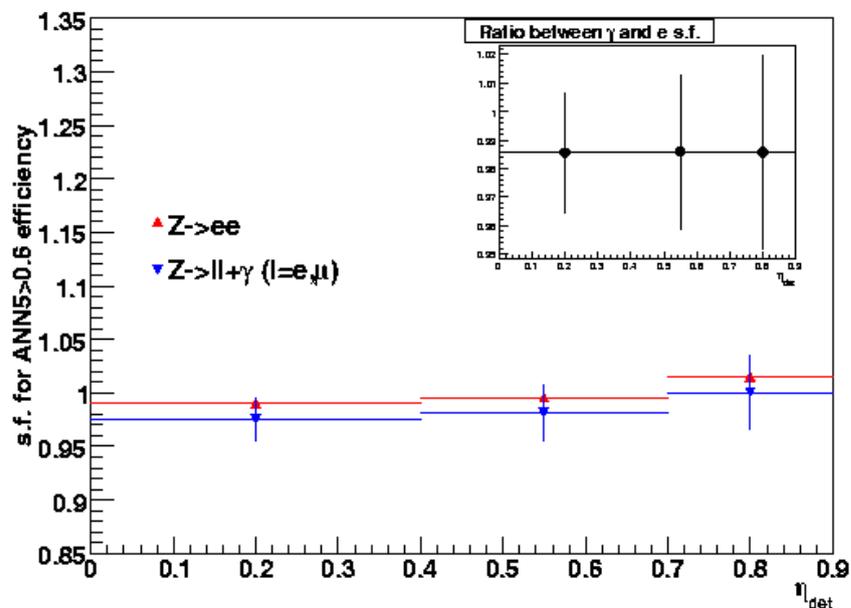
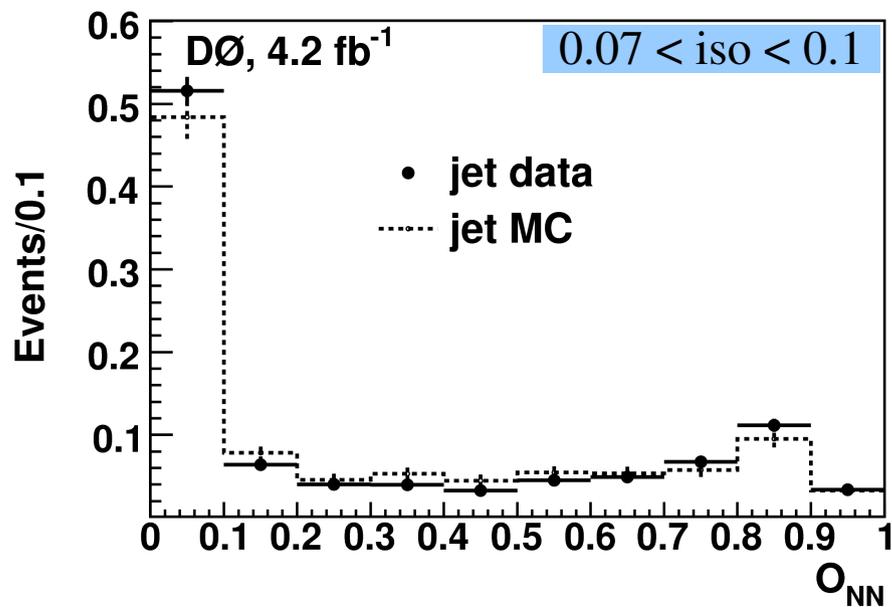


★ Tevatron is sensitive over the whole “interesting” mass range.

$$\begin{pmatrix} N_{ff} \\ N_{fp} \\ N_{pf} \\ N_{pp} \end{pmatrix} = E \times \begin{pmatrix} N_{jj} \\ N_{j\gamma} \\ N_{\gamma j} \\ N_{\gamma\gamma} \end{pmatrix}$$

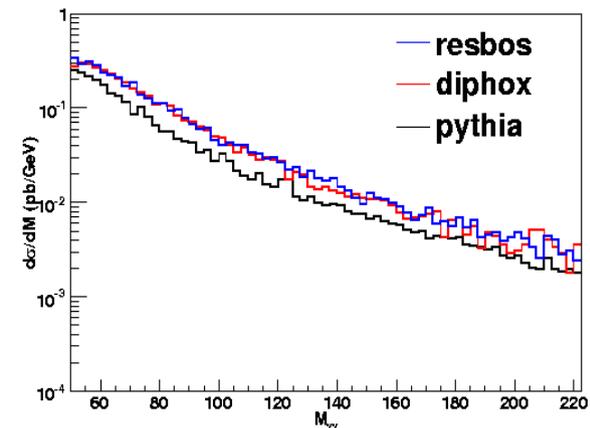
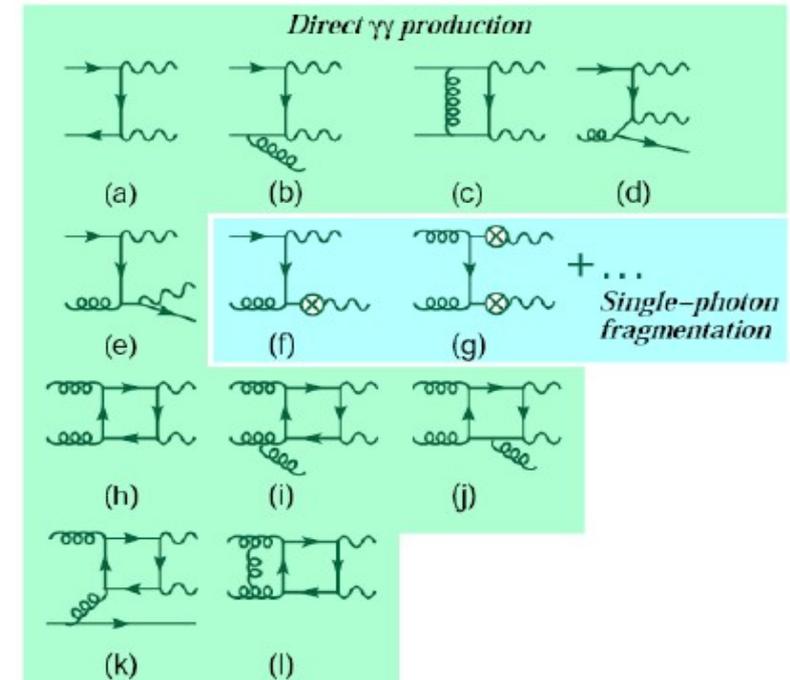


$$E = \begin{pmatrix} (1 - \epsilon_{j1})(1 - \epsilon_{j2}) & (1 - \epsilon_{j1})(1 - \epsilon_{\gamma 2}) & (1 - \epsilon_{\gamma 1})(1 - \epsilon_{j2}) & (1 - \epsilon_{\gamma 1})(1 - \epsilon_{\gamma 2}) \\ (1 - \epsilon_{j1})\epsilon_{j2} & (1 - \epsilon_{j1})\epsilon_{\gamma 2} & (1 - \epsilon_{\gamma 1})\epsilon_{j2} & (1 - \epsilon_{\gamma 1})\epsilon_{\gamma 2} \\ \epsilon_{j1}(1 - \epsilon_{j2}) & \epsilon_{j1}(1 - \epsilon_{\gamma 2}) & \epsilon_{\gamma 1}(1 - \epsilon_{j2}) & \epsilon_{\gamma 1}(1 - \epsilon_{\gamma 2}) \\ \epsilon_{j1}\epsilon_{j2} & \epsilon_{j1}\epsilon_{\gamma 2} & \epsilon_{\gamma 1}\epsilon_{j2} & \epsilon_{\gamma 1}\epsilon_{\gamma 2} \end{pmatrix}$$



Direct $\gamma\gamma$ production (DDP) – irreducible background

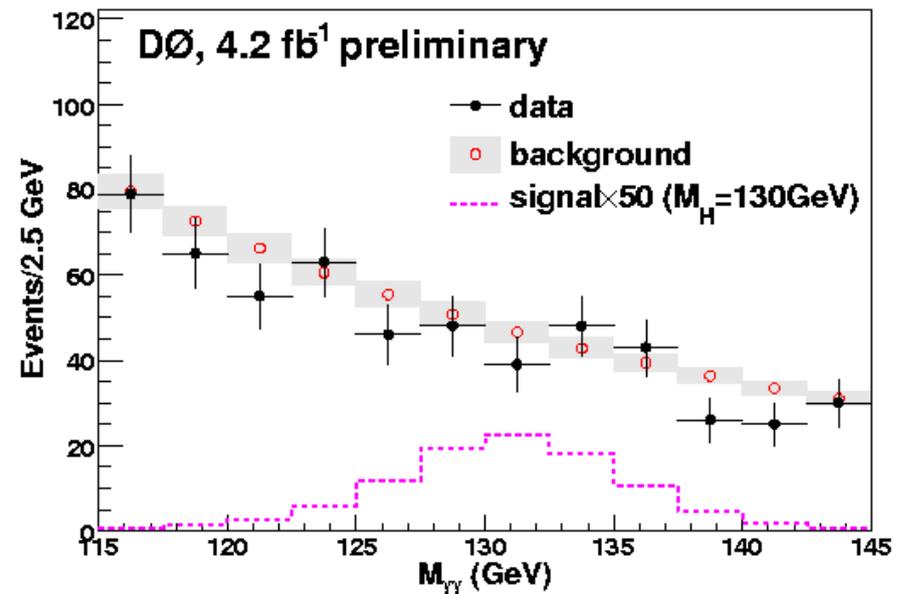
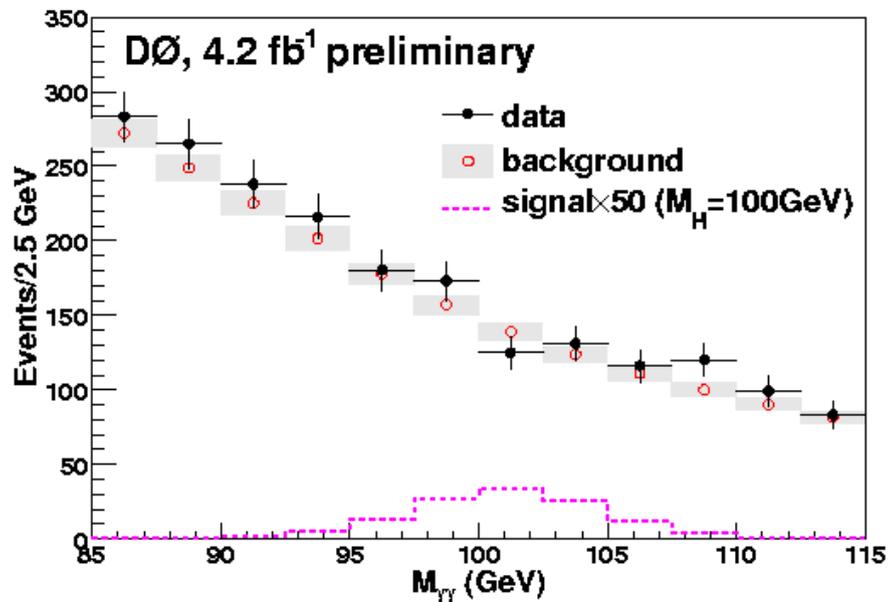
- Challenge to model by theoretical prediction so far, each generator has its own deficit
 - pythia: LO + parton shower
 - **diphox**: NLO, except the gg box contribution (LO)
 - **resbos**: NLO + Resummation of collinear/soft ISR gluons + \sim NLO single fragmentation
- However, in the high mass region ($M_{\gamma\gamma} > 50\text{GeV}$), $M_{\gamma\gamma}$ spectrum agrees reasonably very well between **diphox** and **resbos**.



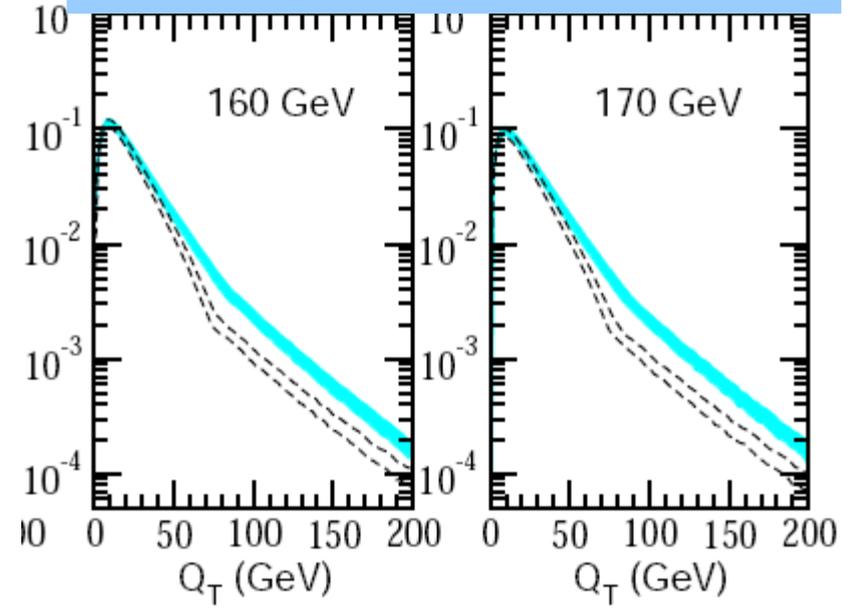
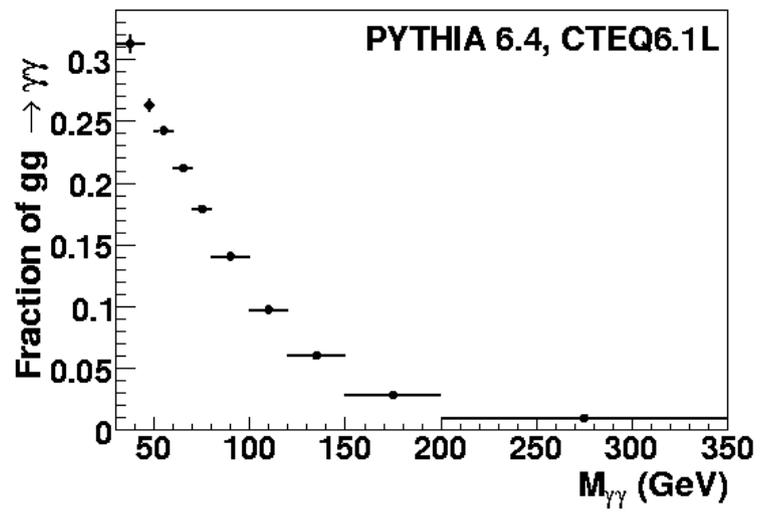
Event yield

Good agreement between data and SM background prediction:

	100 GeV	110 GeV	120 GeV	130 GeV	140 GeV	150 GeV
$Z/\gamma^* \rightarrow ee$	134 ± 27	53 ± 12	17 ± 5	9 ± 3	5 ± 2	3 ± 2
$\gamma j + jj$	712 ± 102	455 ± 65	299 ± 43	202 ± 29	140 ± 20	100 ± 14
QCD $\gamma\gamma$	1080 ± 96	764 ± 62	539 ± 41	404 ± 28	280 ± 19	207 ± 14
total background	1926 ± 35	1272 ± 21	855 ± 14	615 ± 10	425 ± 7	310 ± 5
data	2029	1289	861	567	412	295
signal	2.53 ± 0.18	2.53 ± 0.18	2.38 ± 0.17	2.01 ± 0.14	1.45 ± 0.10	0.87 ± 0.06



$gg \rightarrow H \rightarrow WW$ (arXiv: 0909.2305)



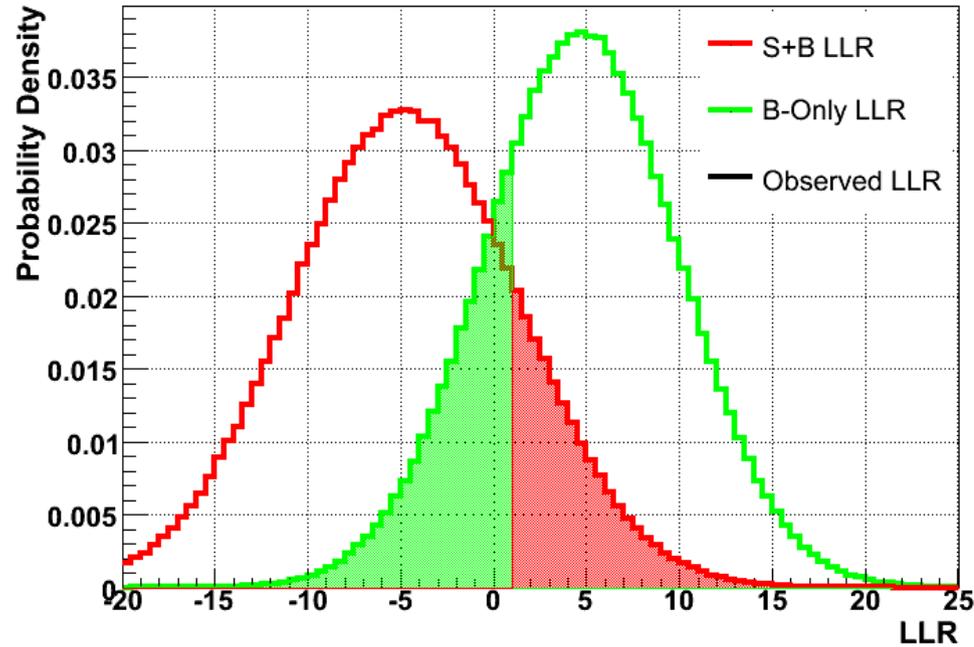
CL Limits calculation

$$\begin{aligned} Q_i &= \frac{\frac{e^{-(s_i+b_i)} \cdot (s_i+b_i)^{d_i}}{d_i!}}{\frac{e^{-b_i} \cdot b_i^{d_i}}{d_i!}} \\ &= e^{-s_i} \cdot \left(1 + \frac{s_i}{b_i}\right)^{d_i} \end{aligned} \quad (7.15)$$

where s_i and b_i are the numbers of signal and background events per bin of invariant mass distribution; and d_i is the number of selected events per bin in data. Poisson sampling is performed for d_i in each bin with the poisson mean set to b_i and $s_i + b_i$ respectively for the “background-only” and “signal+ background” hypotheses. And the likelihood ratio is computed as

$$Q = \prod_{\text{all bins}} Q_i \quad (7.16)$$

Probability distribution function (p.d.f) of the $-2\ln Q$ variable under the background-only and signal+background hypotheses can be obtained from the Poisson sampling.



where the two histograms are the p.d.f of the $-2\ln Q$ statistic under background-only hypothesis and signal+background hypothesis and the vertical line stands for the value of data. The integral of the signal+background (background-only) p.d.f from the vertical data line to $+\infty$ is called CL_{s+b} (CL_b). If $CL_{s+b} < 0.05$, exclusion can be claimed at 95% C.L. A more conservative exclusion estimator is defined in place of CL_{s+b} as

$$CL_s = \frac{CL_{s+b}}{CL_b}$$

