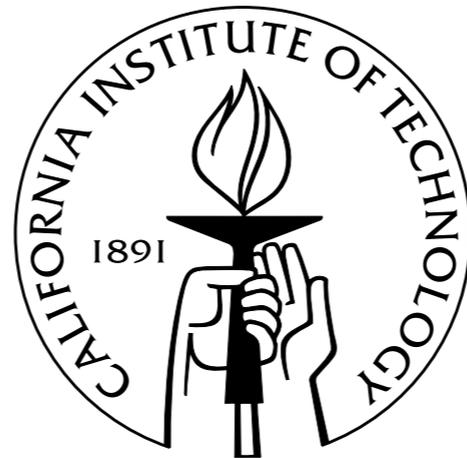


Naturalness and the Weak Gravity Conjecture

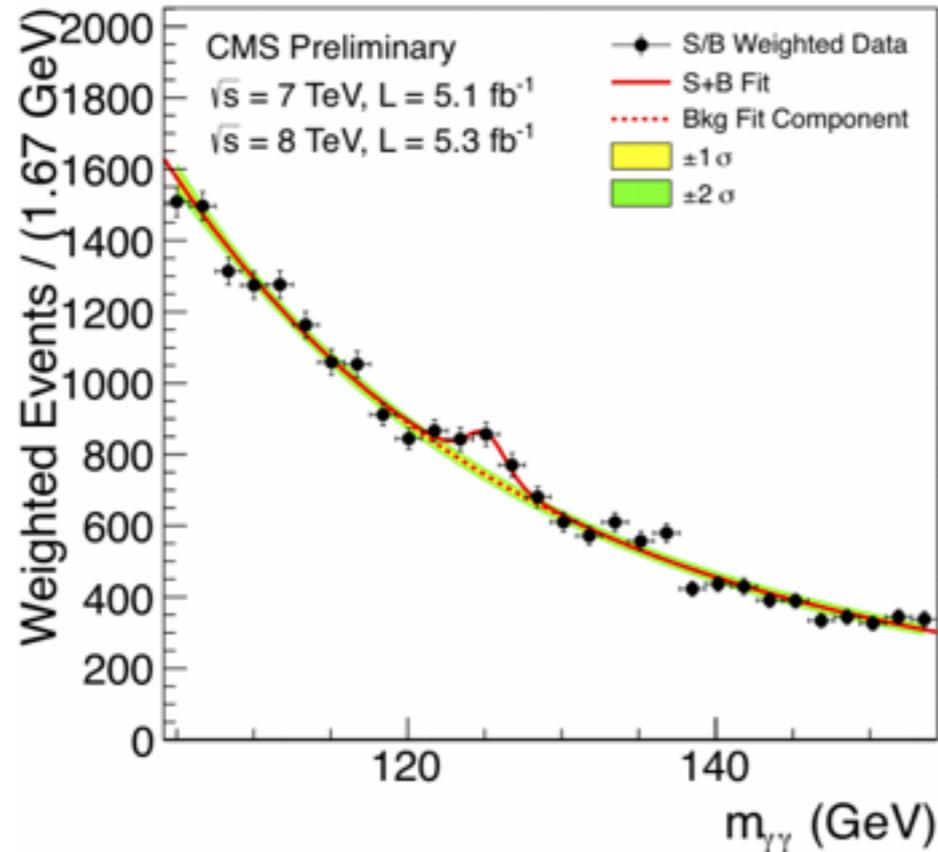
Clifford Cheung



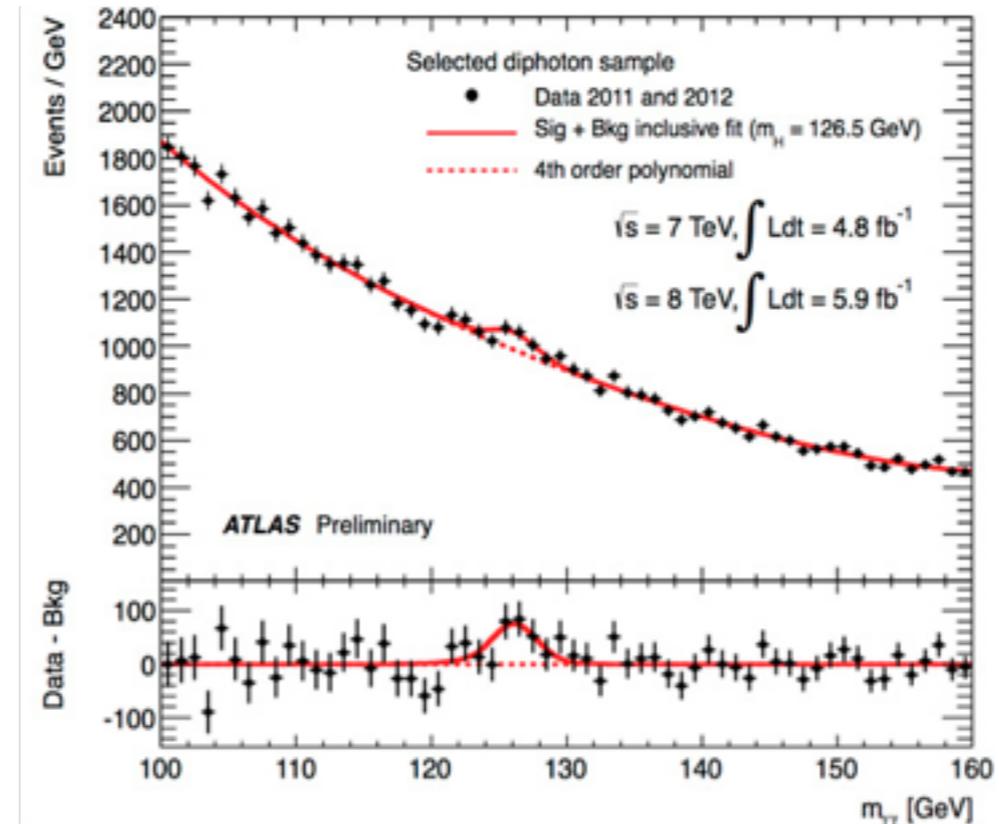
C. Cheung, G. N. Remmen (1402.2287, 1407.7865)

Nature exhibits a fundamental scalar...

CMS

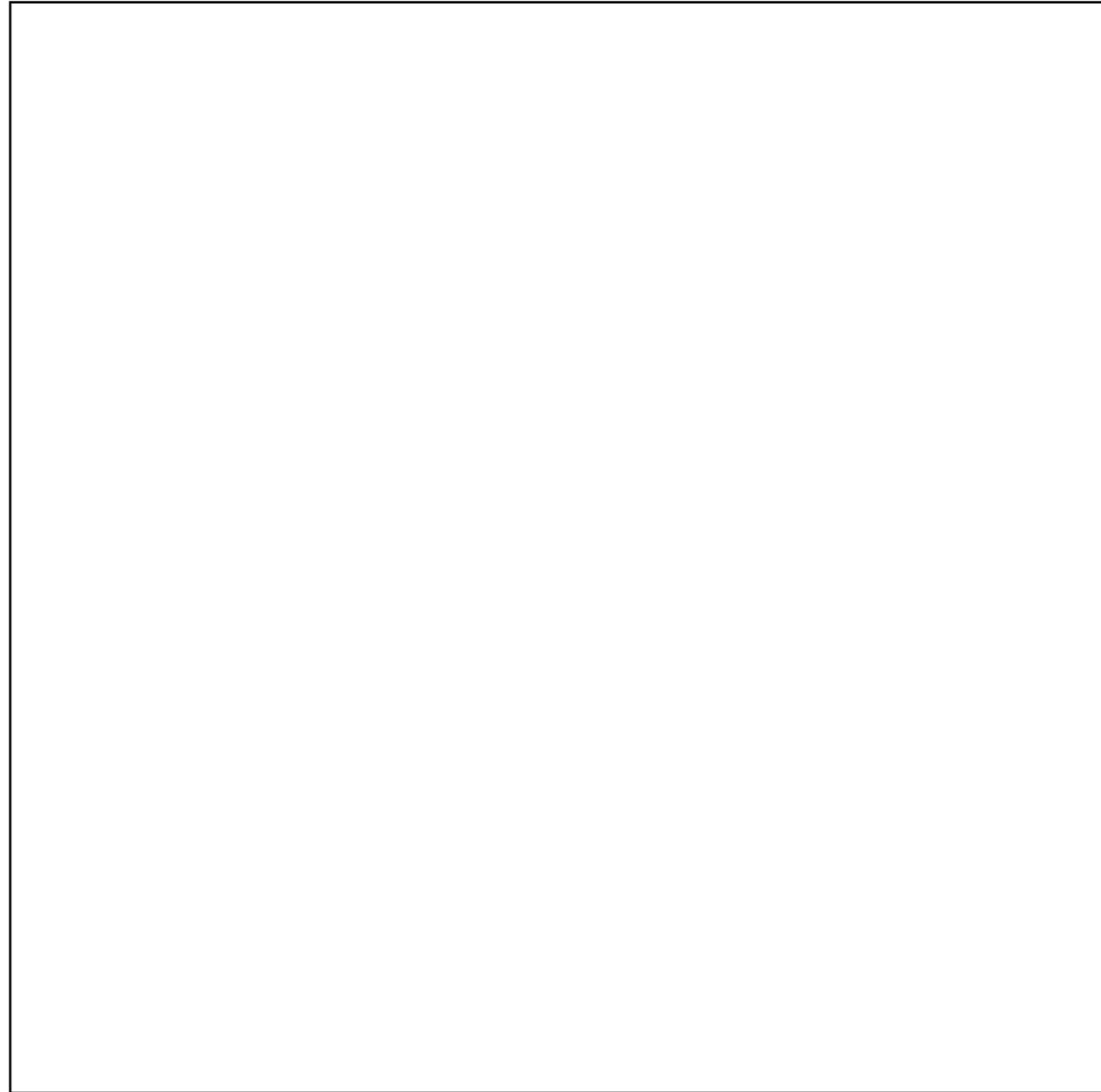


ATLAS



...but alas, no “naturalons,” yet.

theory space



theory space

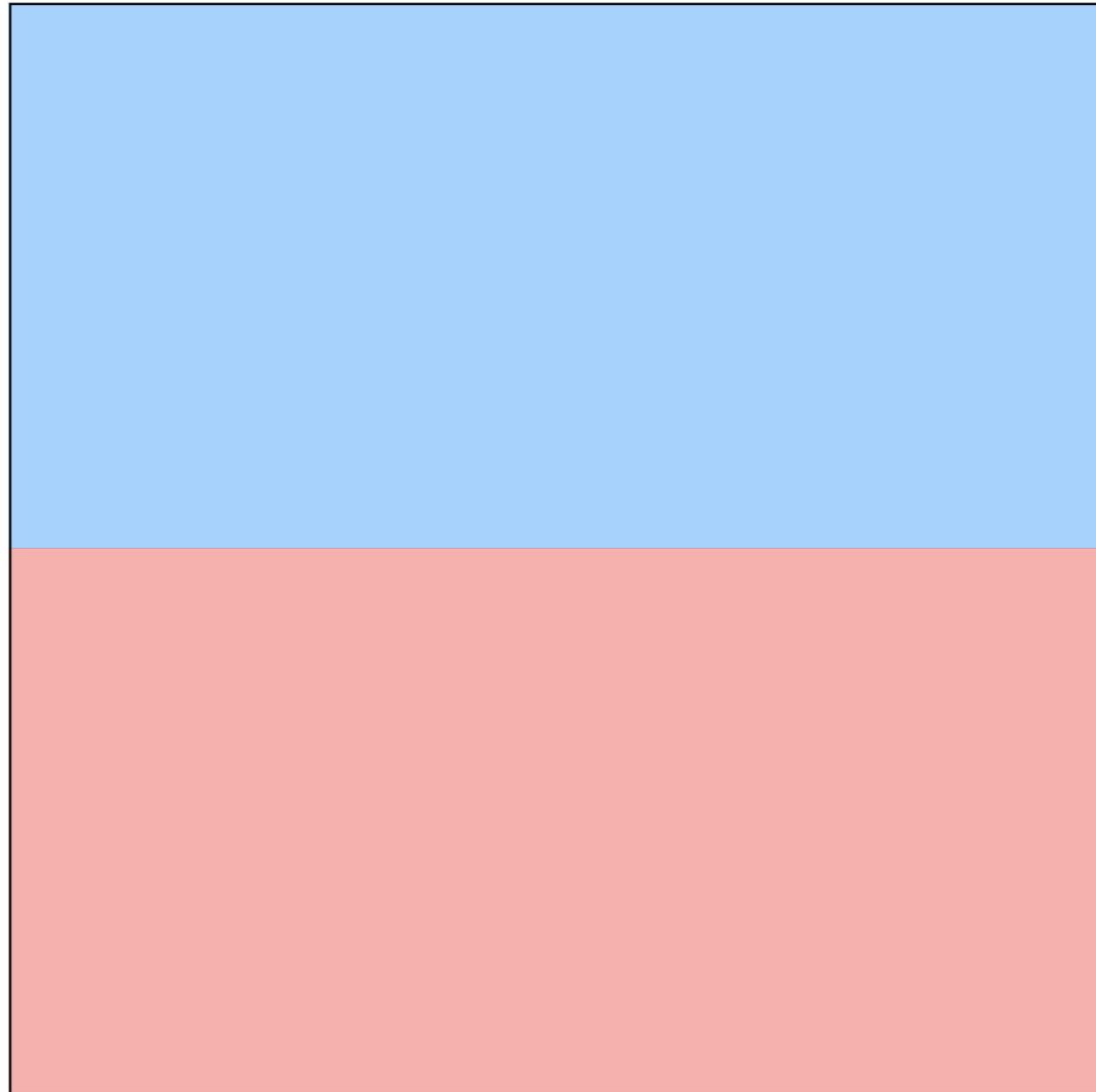
natural



theory space

natural

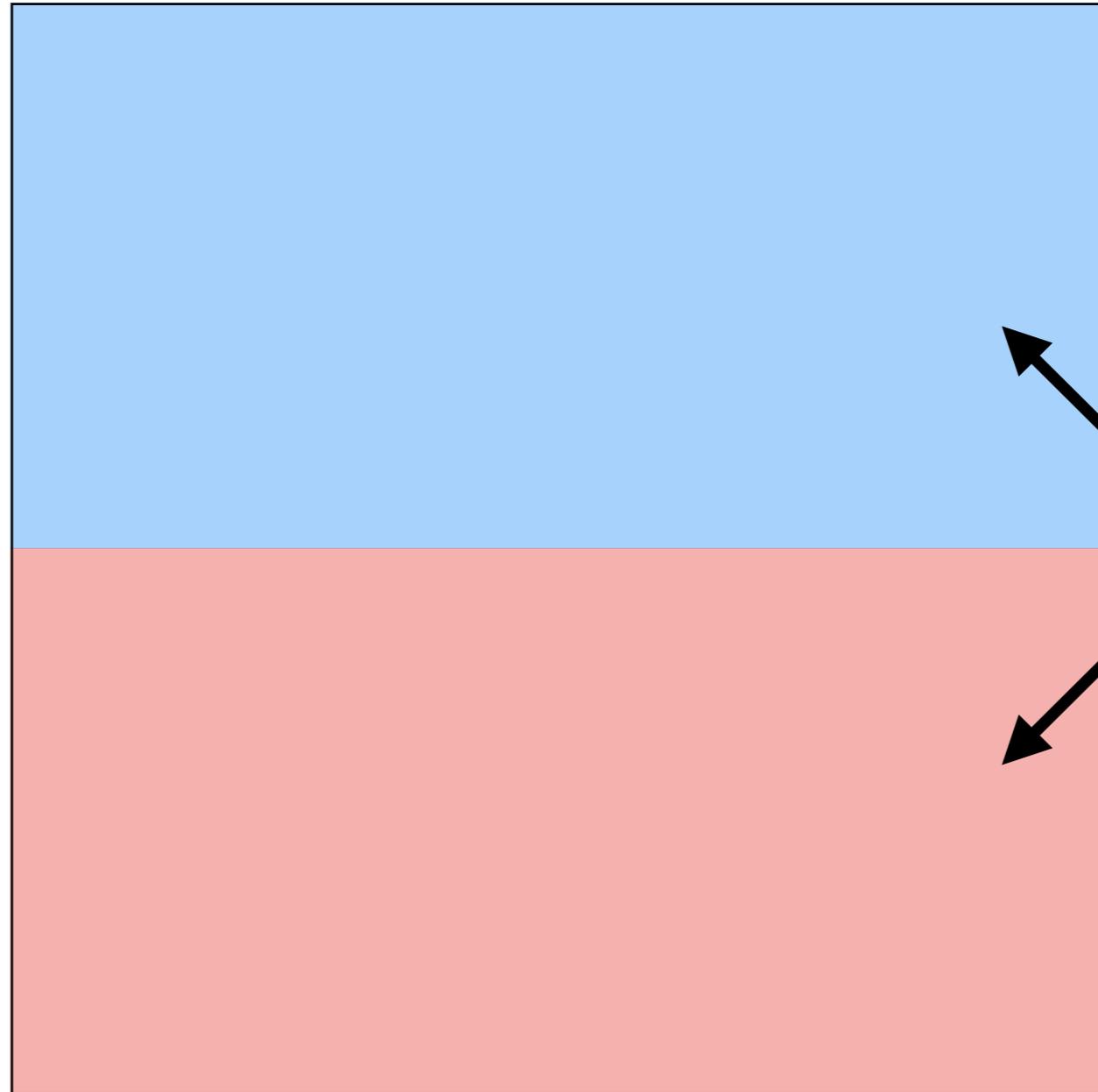
anthropic
selection



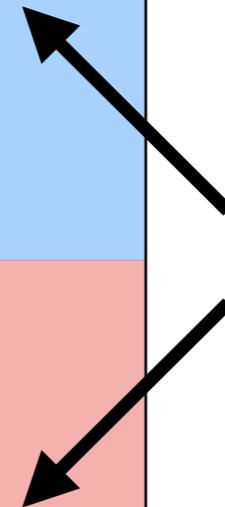
theory space

natural

anthropic
selection



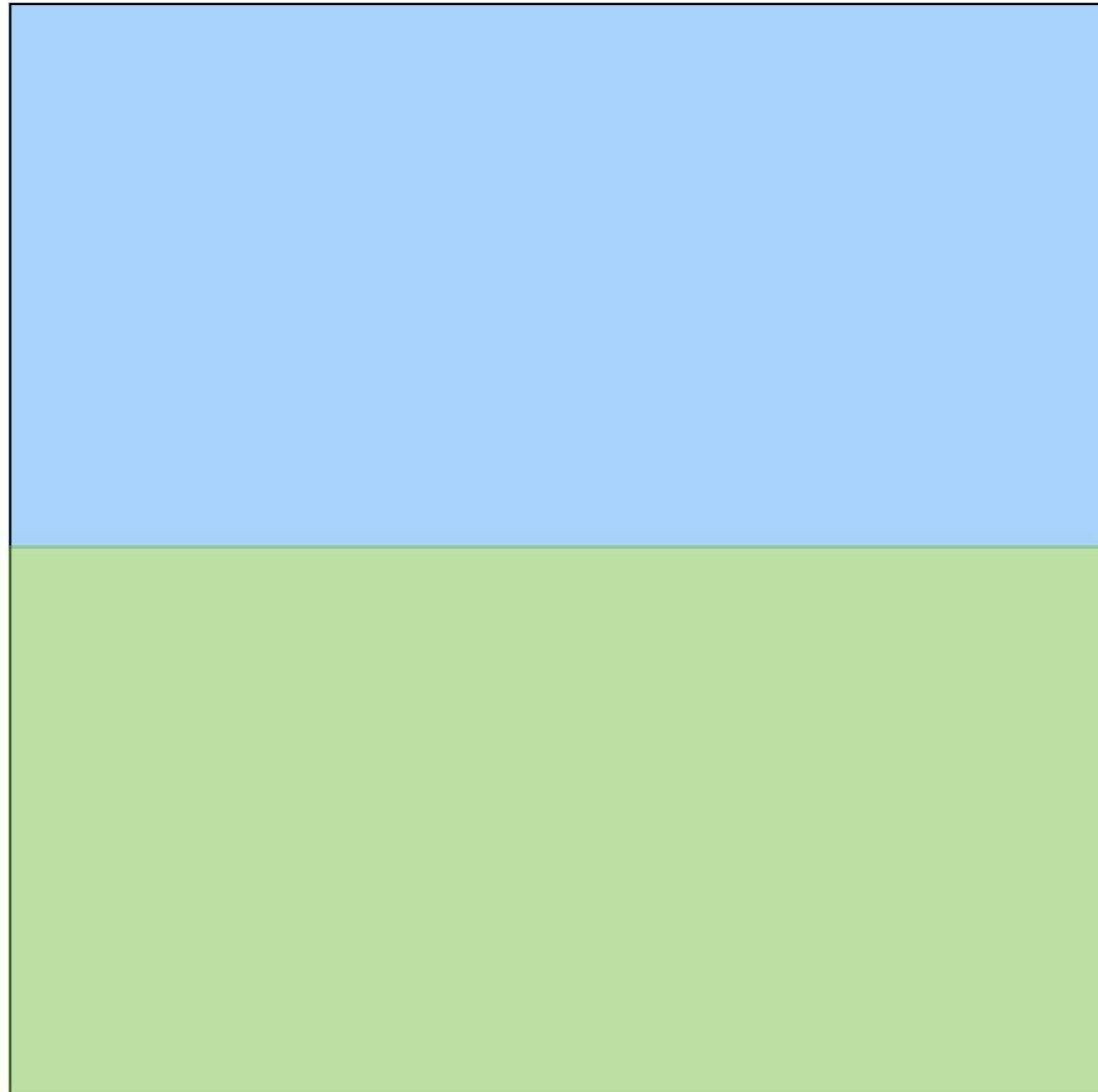
false
dichotomy



theory space

natural

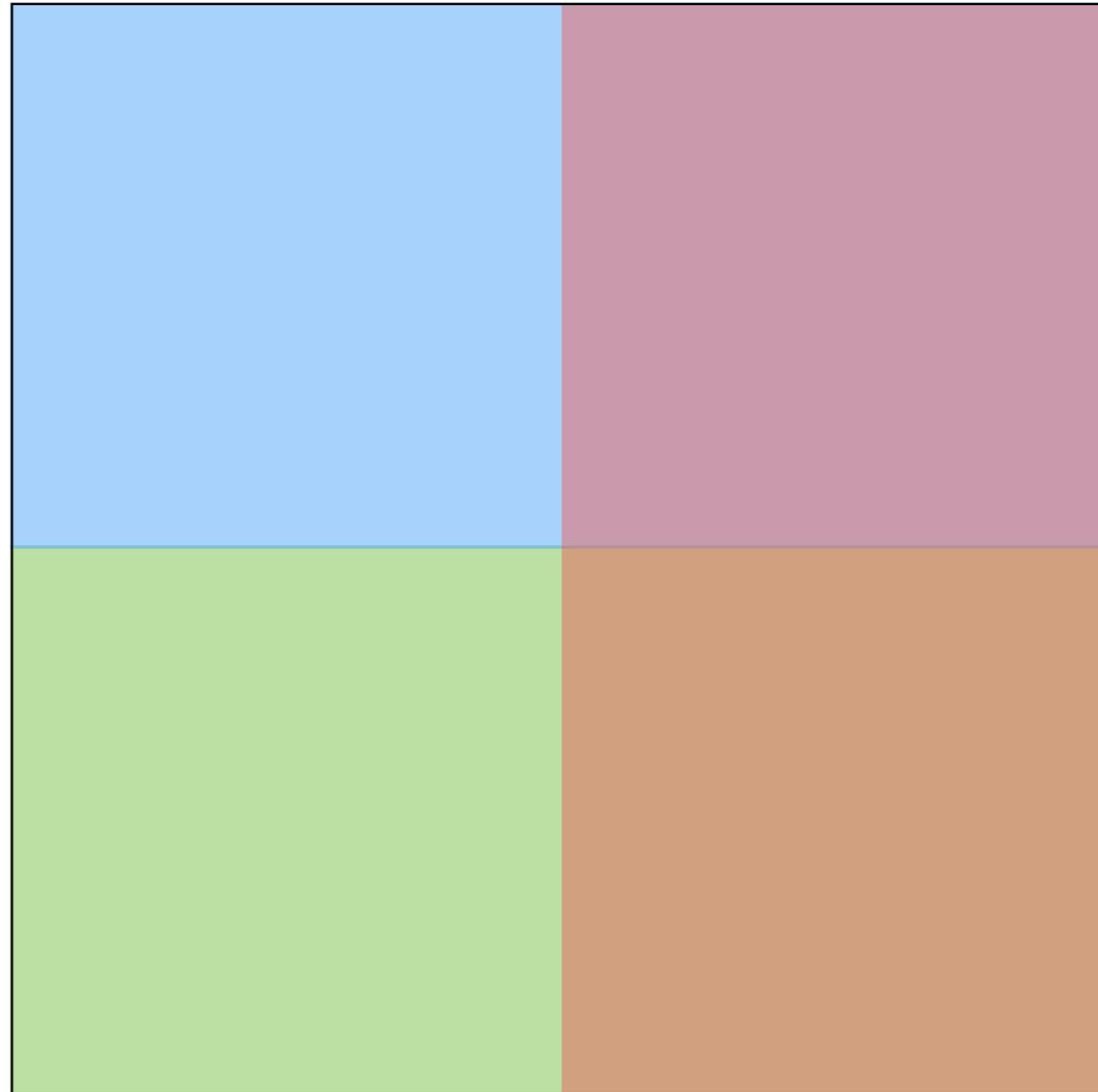
un-natural



theory space

natural

un-natural



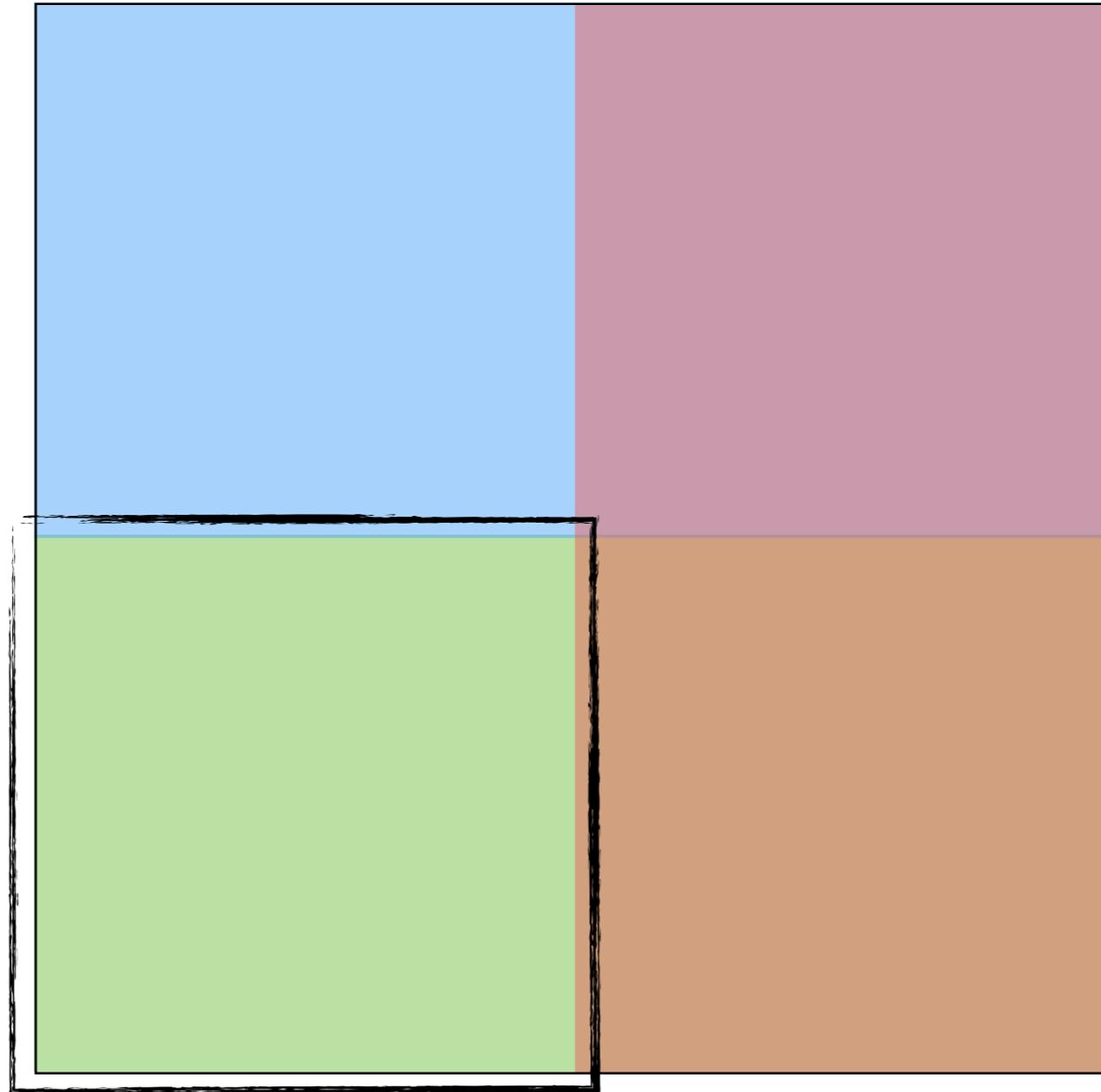
anthropic
selection

theory space

natural

anthropic
selection

un-natural



“The 3rd Way”

Imagine a “desert” above the weak scale.

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$$m_H \sim 10^{-16} m_{\text{Pl}} \qquad V_0^{1/4} \sim 10^{-30} m_{\text{Pl}}$$

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Possible lessons?

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$$m_H \sim 10^{-16} m_{\text{Pl}} \quad V_0^{1/4} \sim 10^{-30} m_{\text{Pl}}$$

Possible lessons?

- We suck (at computing symmetry unprotected dimensionful parameters).
- Don't modify gravity, understand tuning.

What is the invariant meaning of naturalness?

Consider an EFT for the hierarchy problem.

$$\begin{aligned}\mathcal{L} = & \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_0^2 \phi^2 \quad \leftarrow \text{(light state)} \\ & + \frac{1}{2} \partial_\mu \chi \partial^\mu \chi - \frac{1}{2} M_0^2 \chi^2 \quad \leftarrow \text{(heavy state)} \\ & - \frac{1}{4} \lambda_0 \phi^2 \chi^2 + \dots\end{aligned}$$

Claim: even pulling every dirty trick in the book, there's still an irreducible tuning.

Computing with a hard cutoff ...

$$\delta m_0^2 \sim \frac{\lambda_0}{16\pi^2} (\Lambda^2 + M_0^2 \log M_0^2 / \Lambda^2 + M_0^2 + \dots)$$

... or dimensional regularization...

$$\delta m_0^2 \sim \frac{\lambda_0}{16\pi^2} \left(\frac{M_0^2}{\epsilon} + M_0^2 \log M_0^2 / \mu^2 + M_0^2 + \dots \right)$$

... the problem persists.

Regulator abracadabra won't fix ubiquitous tree-level hierarchy problems.

ultraviolet
symmetries



dangerous symmetry
breaking parameters

The issue is a generalization of the doublet-triplet splitting problem.

$$\delta m_0^2 \sim \lambda_0 \langle \chi \rangle^2$$



GUT scale,
PQ scale

Let's re-examine the terms.

$$\delta m_0^2 \sim \frac{\lambda_0}{16\pi^2} \left(\frac{M_0^2}{\epsilon} + \underset{\substack{\uparrow \\ \text{scheme} \\ \text{independent}}}{M_0^2} \log M_0^2 / \mu^2 + \underset{\substack{\uparrow \\ \text{scheme} \\ \text{dependent}}}{M_0^2} + \dots \right)$$

Finite and $1/\epsilon$ contributions are unobservable and can be absorbed entirely into pole mass.

However, the running is physical.

criteria of “running naturalness”

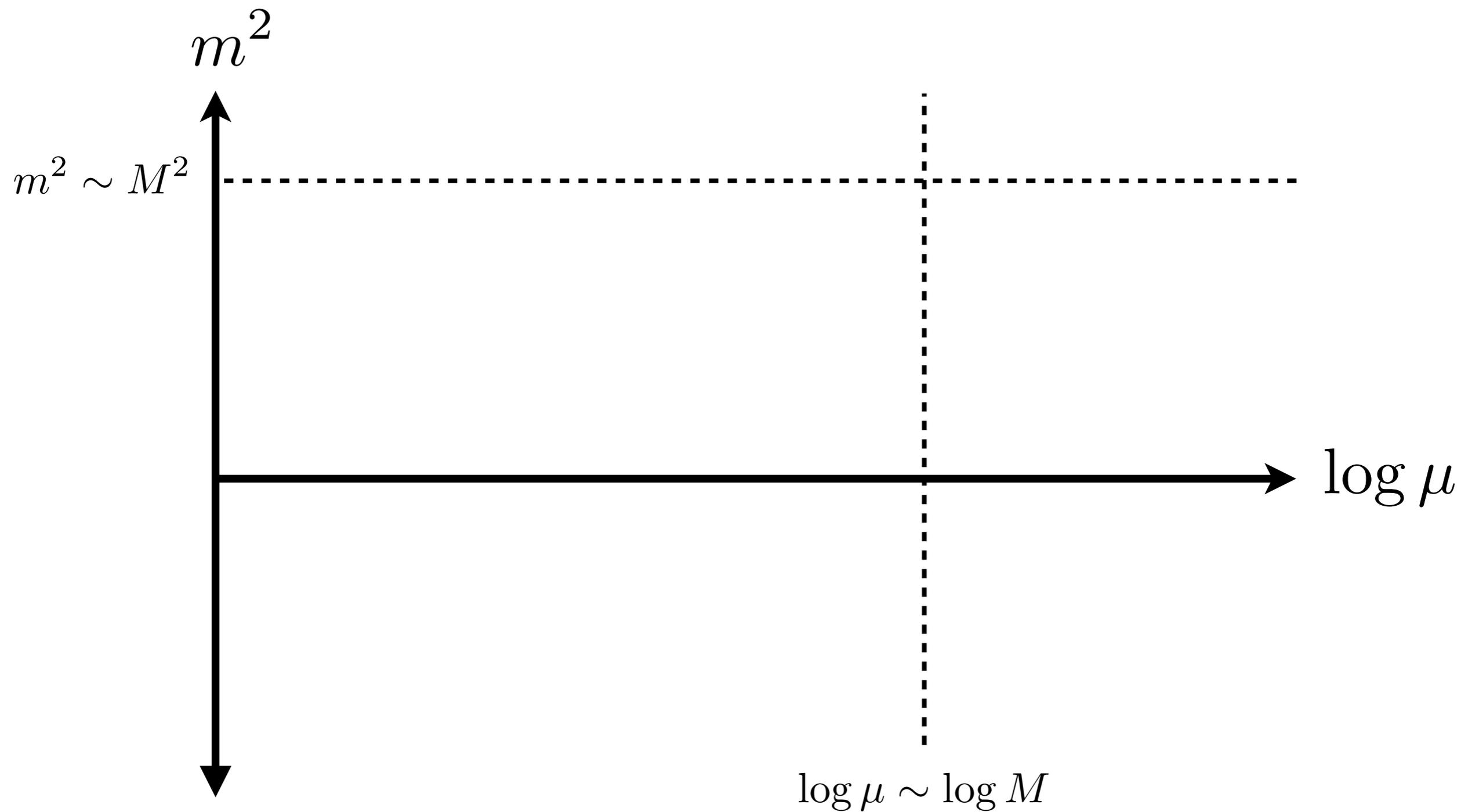
Theory is fine-tuned if the mass changes by orders of magnitude in an e-fold of running.

$$\frac{dm^2}{d \log \mu} = \frac{\lambda M^2}{16\pi^2}$$

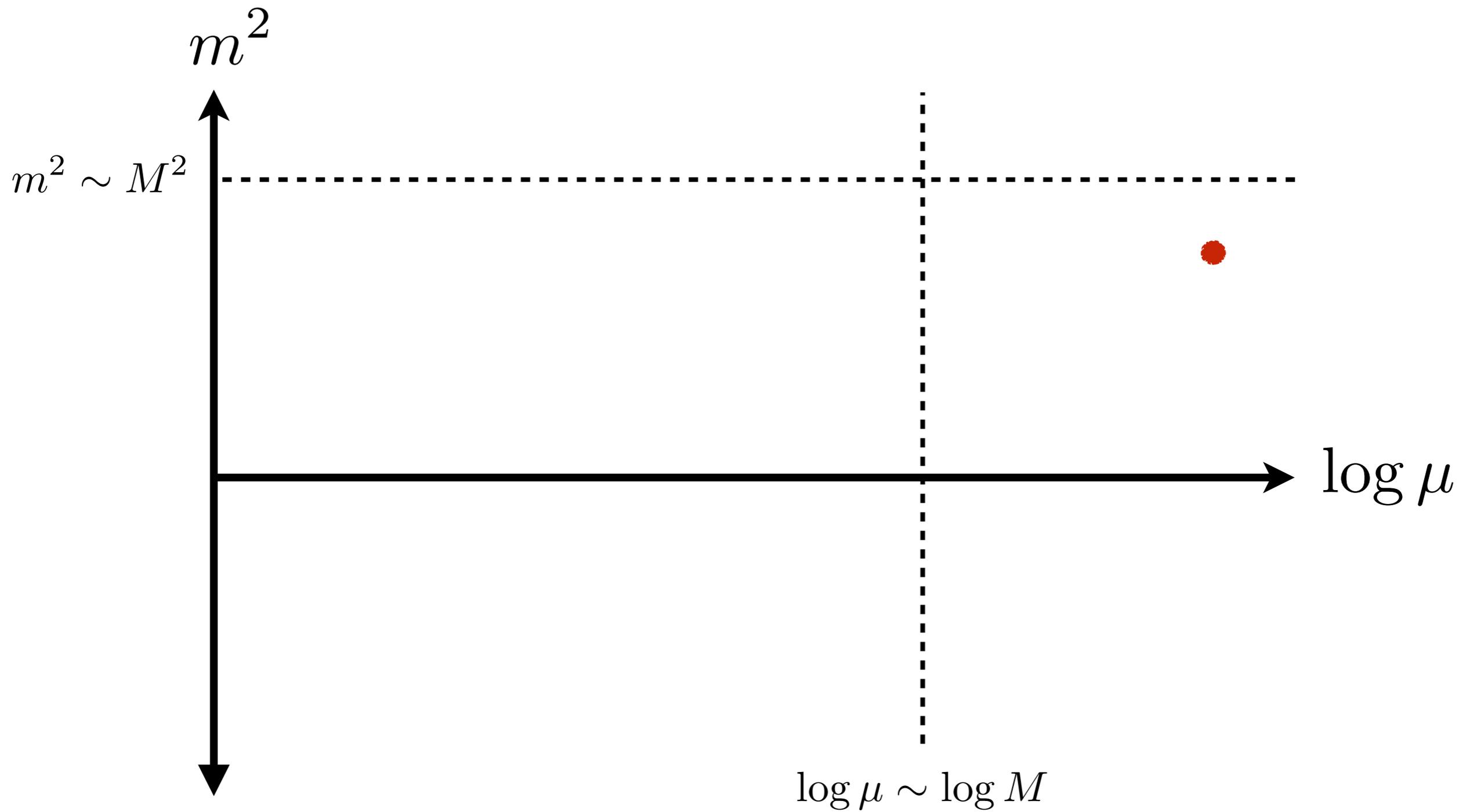
This is a part of the hierarchy problem that won't ever go away.



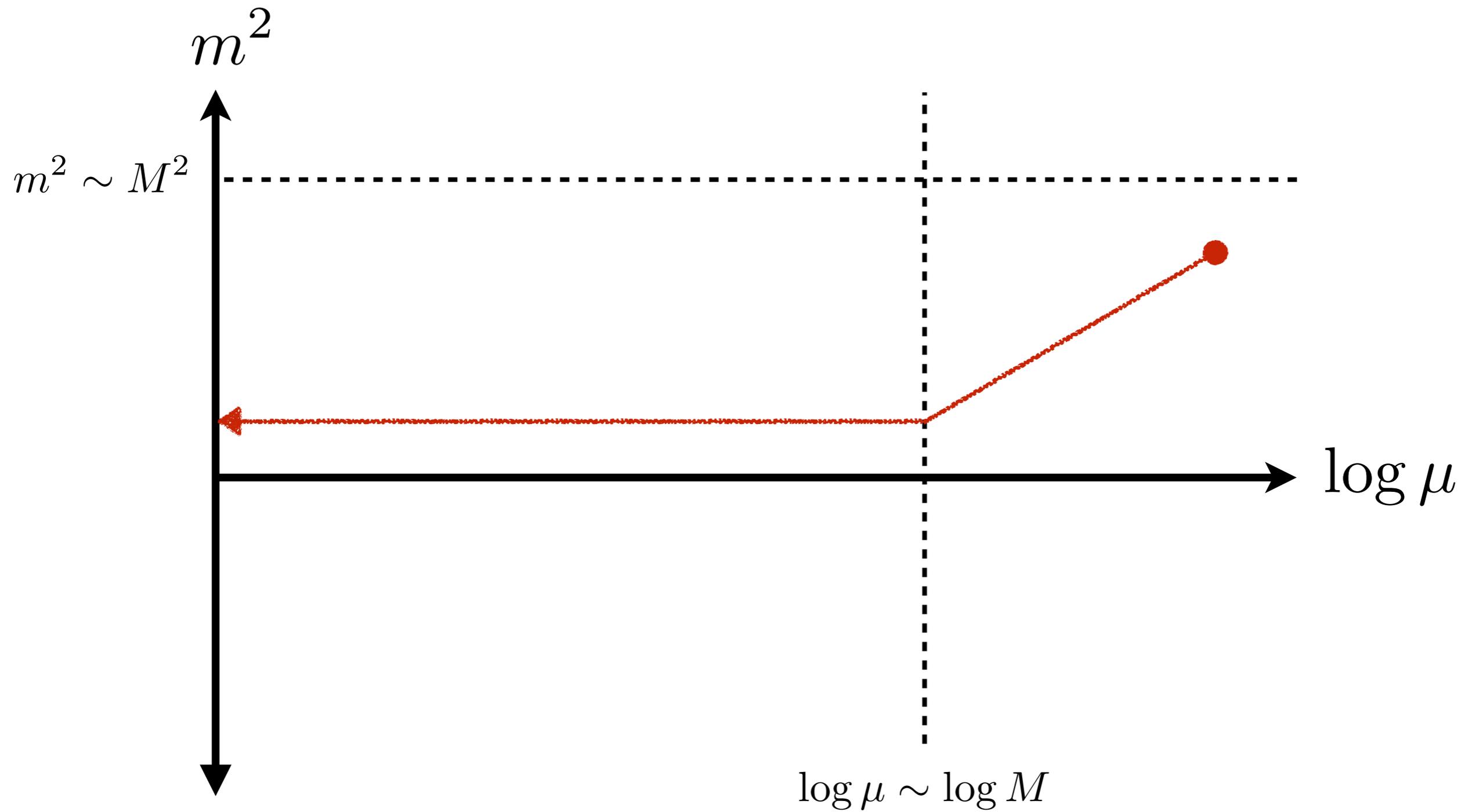
Consider the running of the light mass.



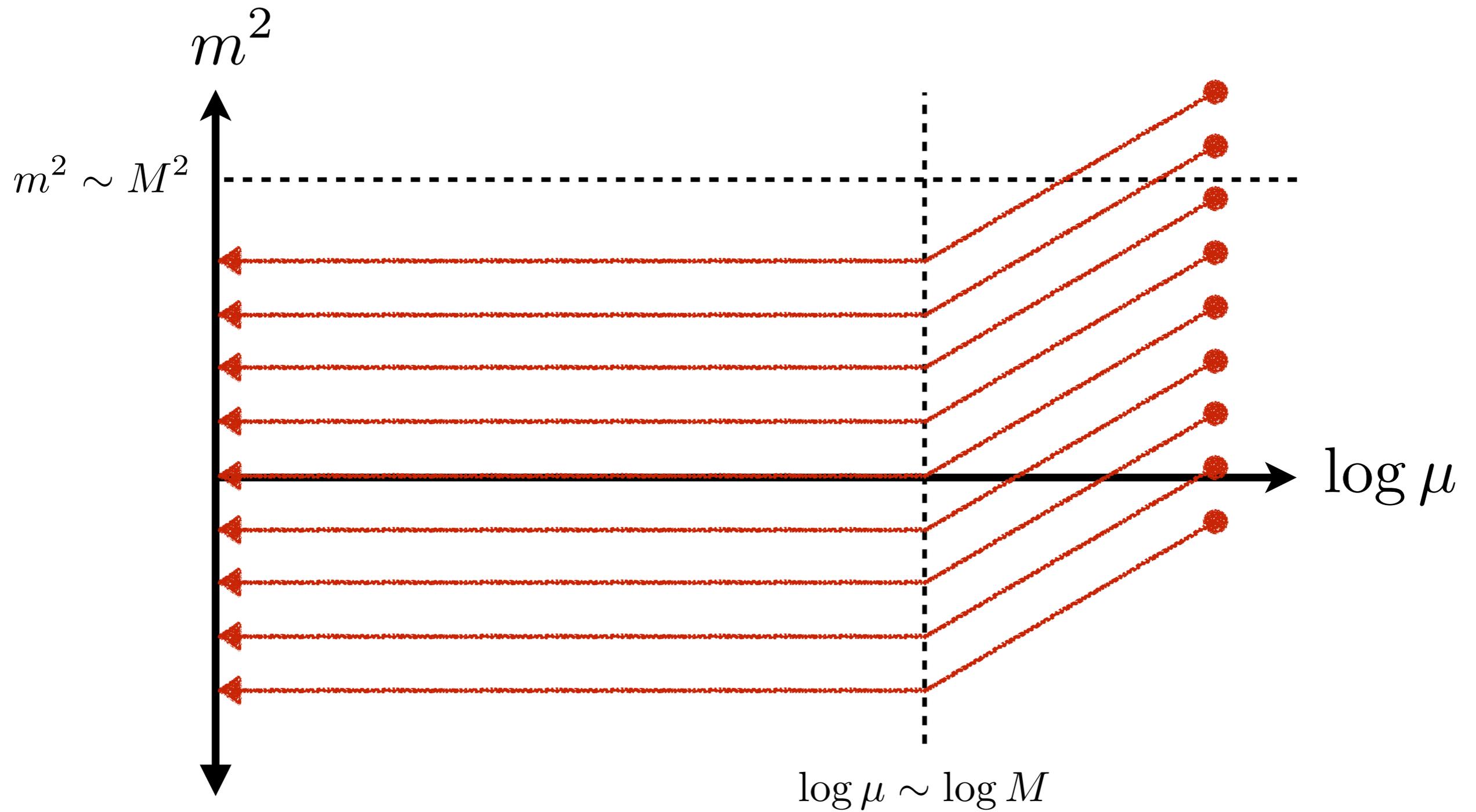
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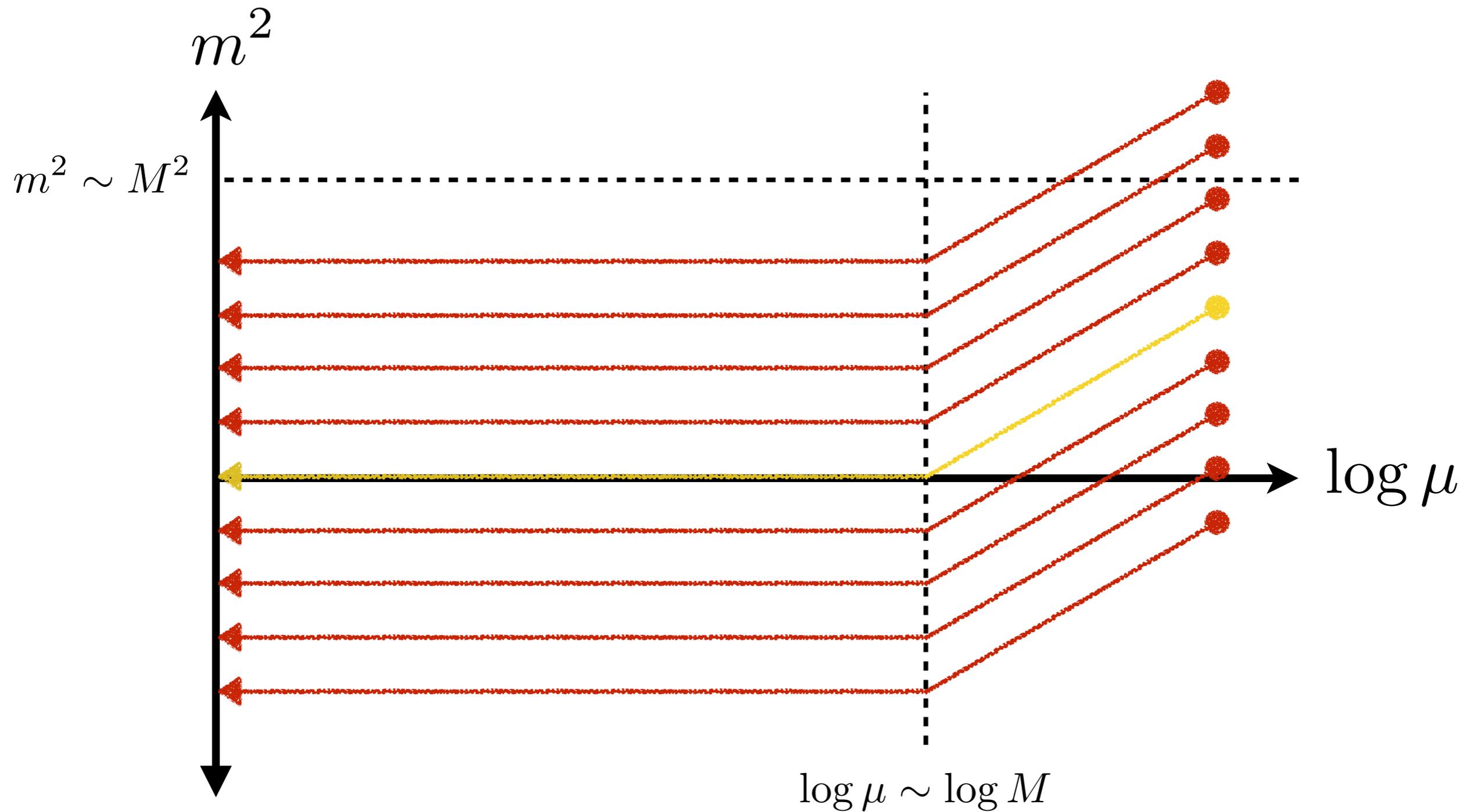
Fix a UV boundary condition.



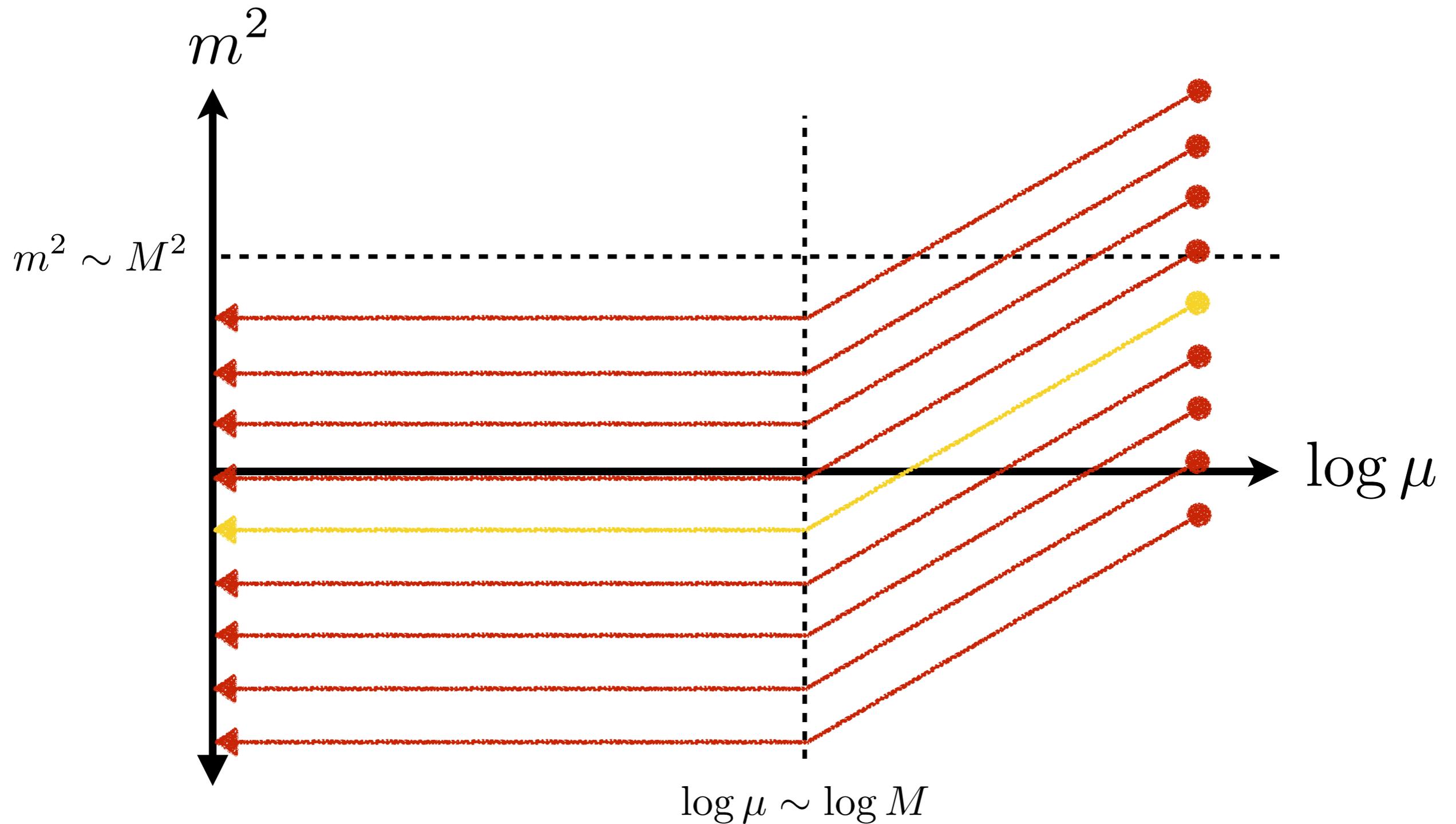
Run down to obtain pole mass.



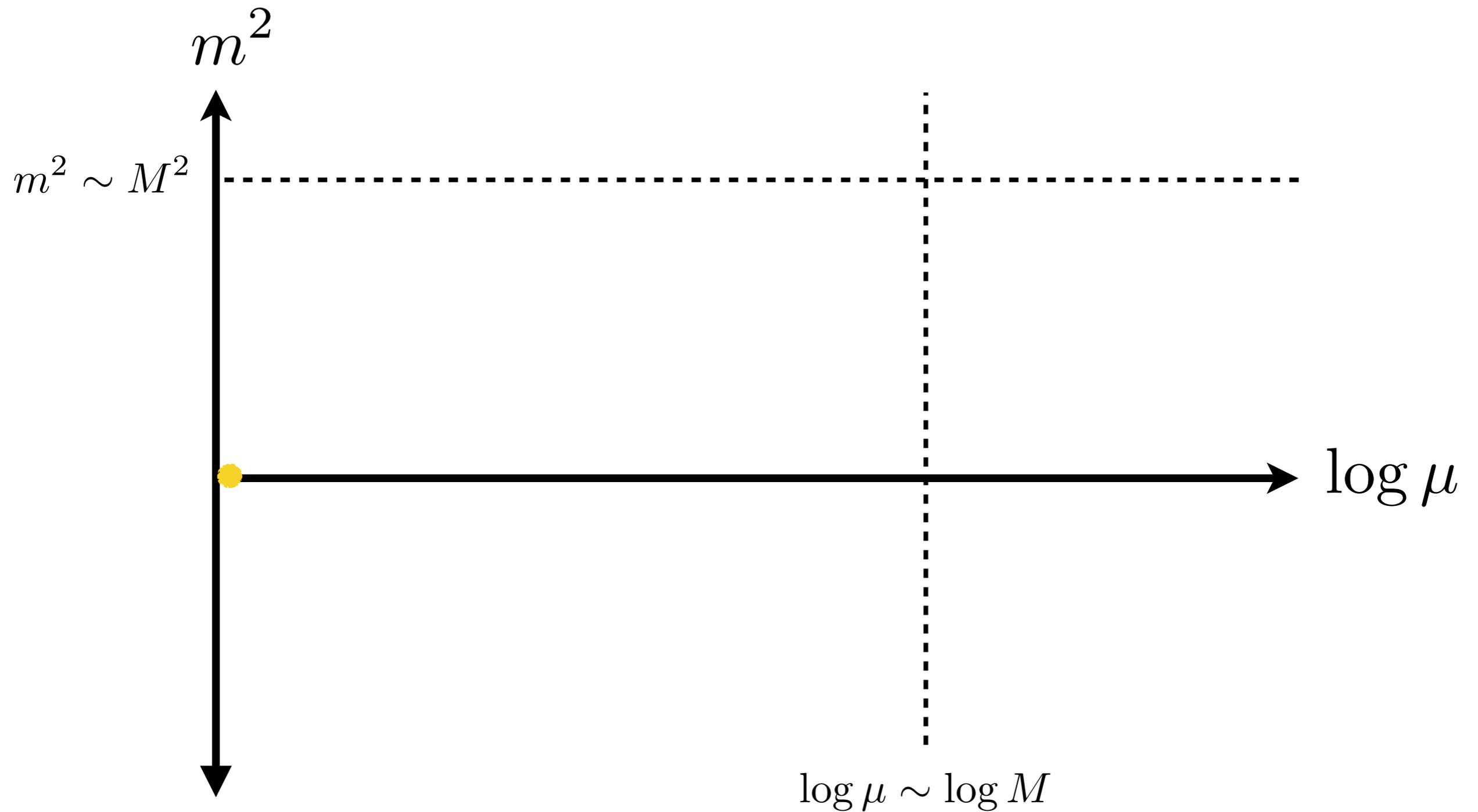
Now, throw darts in the UV.



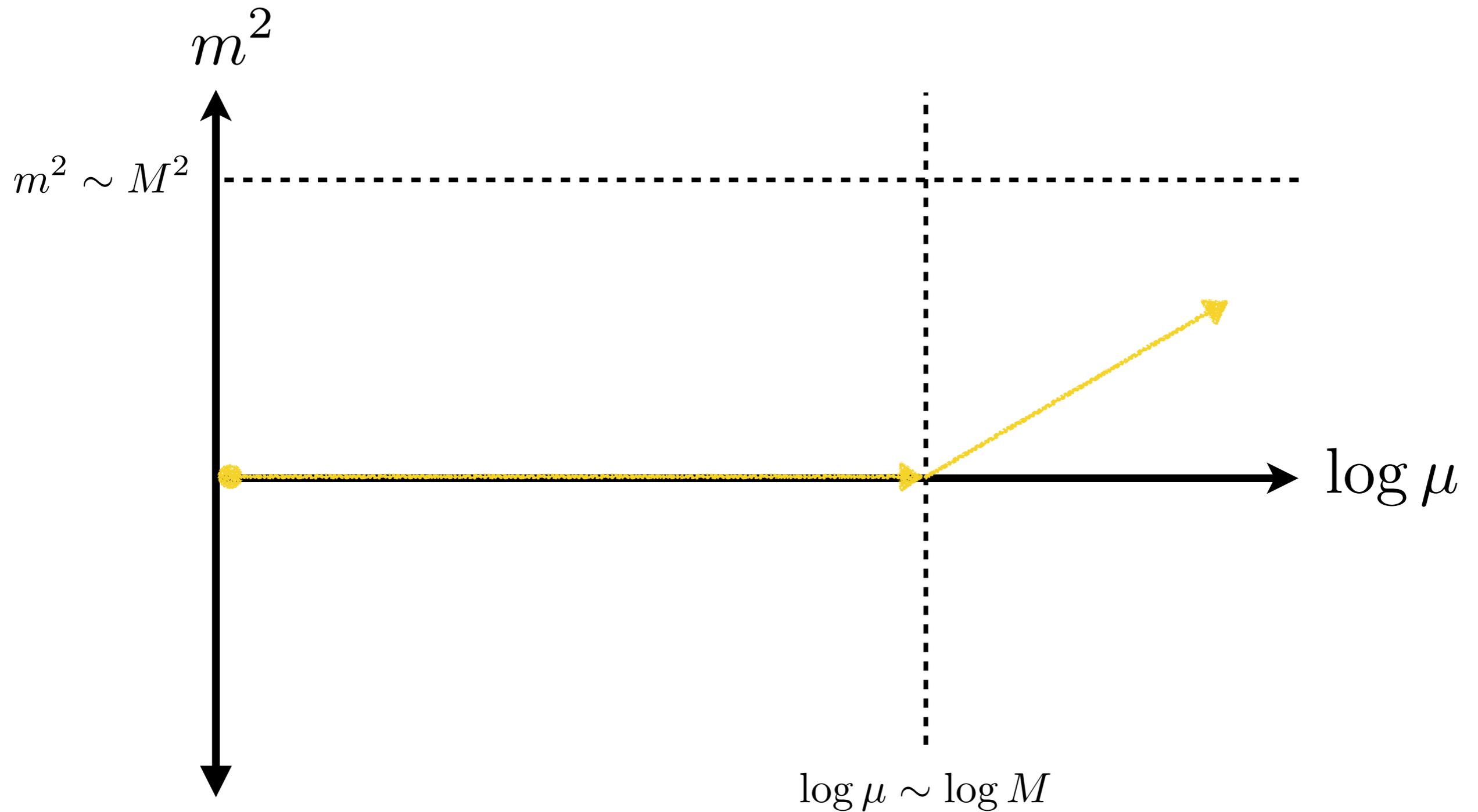
A light pole mass is exceedingly rare.



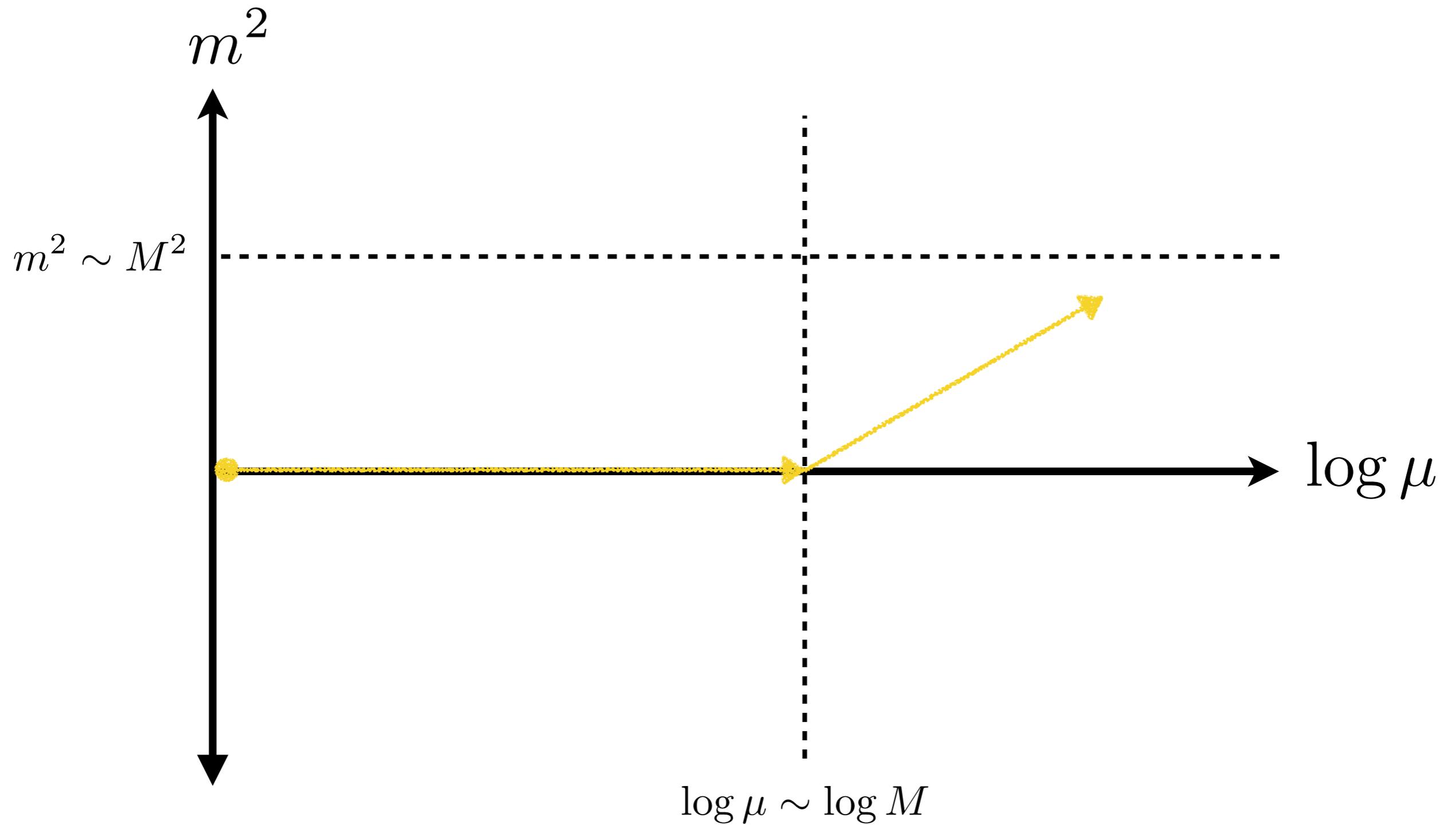
Moreover, it is unstable to variations of M .



What if pole mass is the boundary condition?



What if pole mass is the boundary condition?



Then varying M doesn't do much.

Fine-tuning comes from throwing darts evenly in the UV. Possibilities for “3rd Way”:

- Maybe the more “fundamental” boundary condition is **on-shell**, not **UV**. Indeed, S-matrix is only observable in quantum gravity.

- Maybe the darts aren’t thrown evenly.

There are constraints on theories mandated by consistency.

weak gravity conjecture

weak gravity conjecture (WGC)

(Arkani-Hamed, Motl, Nicolis, Vafa)

A long-range U(1) coupled consistently to gravity *requires* a state with

$$q > m/m_{\text{Pl}}$$

which is a non-perturbative, highly non-trivial criterion for healthy theories.

“Gravity is the weakest force.”

evidence #1

The WGC is satisfied by a litany of healthy field and string theories.

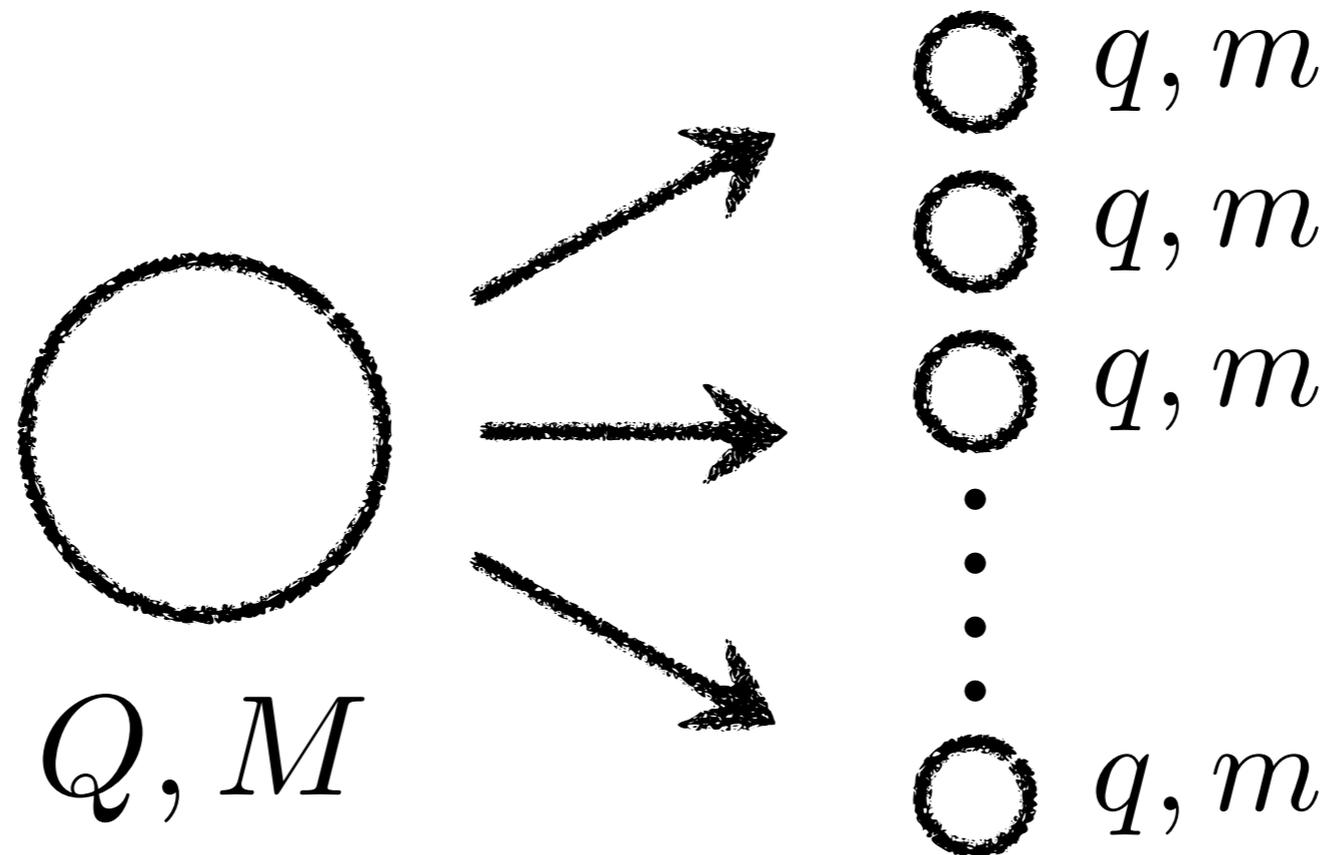
For example, for $SU(2) \rightarrow U(1)$ gauge theory,

$$g > m_W / m_{P1} \xrightarrow{(m_W = gv)} m_{P1} > v$$

and similarly for the monopoles.

evidence #2

The authors of the WGC justified it with a Gedanken experiment with black holes:



number of particles
in final state $= Q/q$ conservation
of charge

total rest mass
in final state $= mQ/q < M$ conservation
of energy

For an extremal black hole, $Q = M/m_{\text{Pl}}$, so

$$q > m/m_{\text{Pl}}$$

When the WGC criterion fails, extremal black holes are exactly stable.

In such a theory there will be a huge number of stable black hole remnants.

This yields serious pathologies:

- thermodynamic catastrophes
- tension with holography

But we have ignored a crucial effect, which is that charges and masses are renormalized!

$$q(\mu) > m(\mu)/m_{\text{Pl}}$$


renormalized
quantities

We should evaluate quantities at pole mass.

Note: WGC can bound a radiatively *unstable* quantity (mass) by a *stable* one (charge).

scalar QED

Take the very simplest case of a U(1) charged particle with a hierarchy problem:

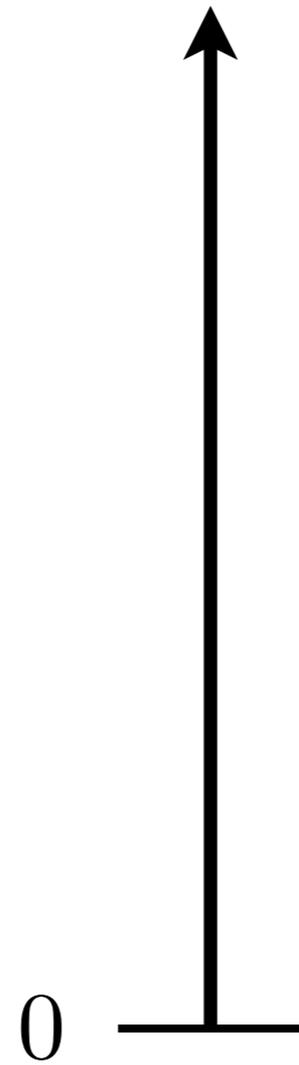
$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^2 + |D_{\mu}\phi|^2 - m^2|\phi|^2 - \frac{\lambda}{4}|\phi|^4$$

where the “selectron” has charge q :

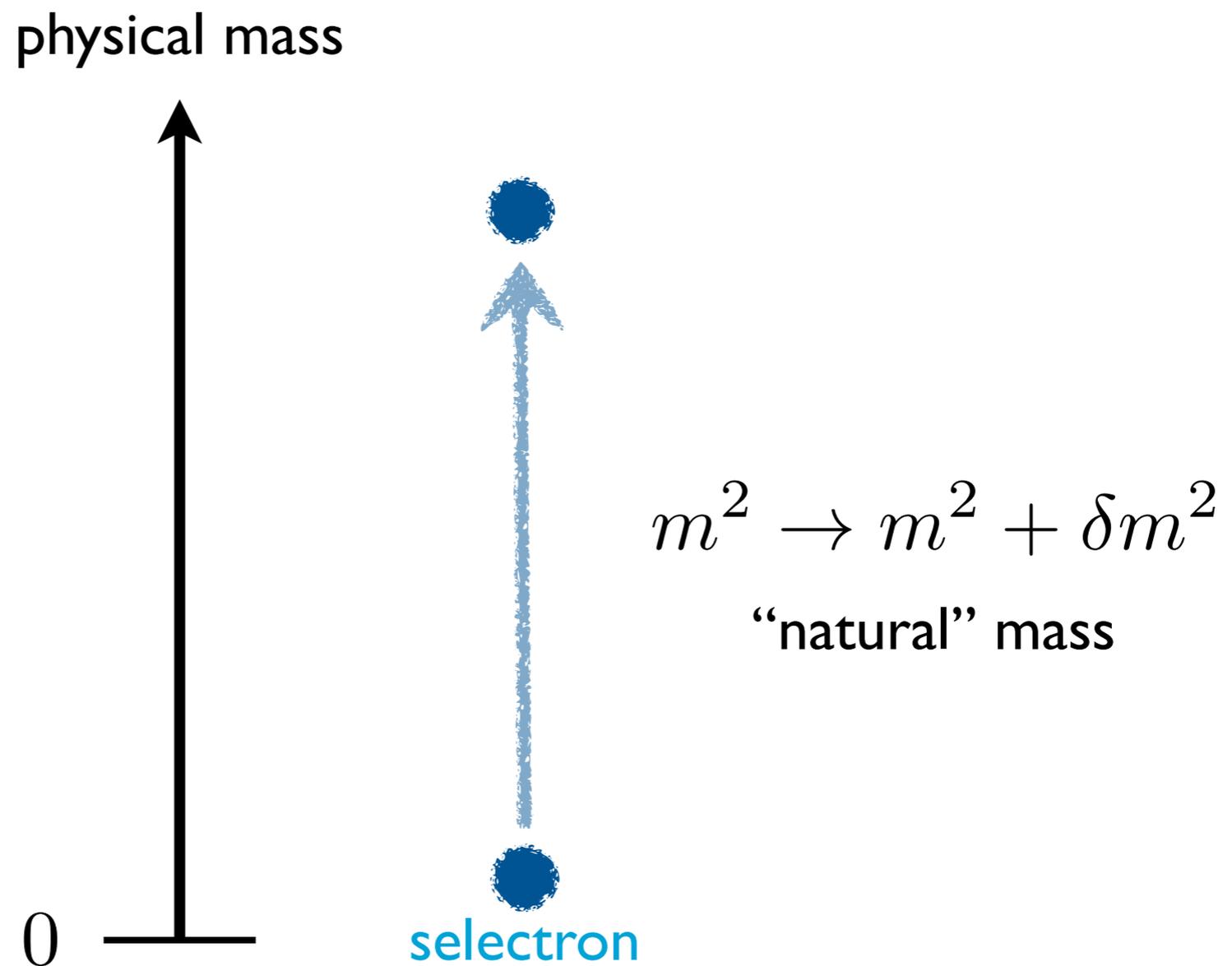
$$D_{\mu} = \partial_{\mu} + iqA_{\mu}$$

There is fundamental tension between naturalness and the WGC.

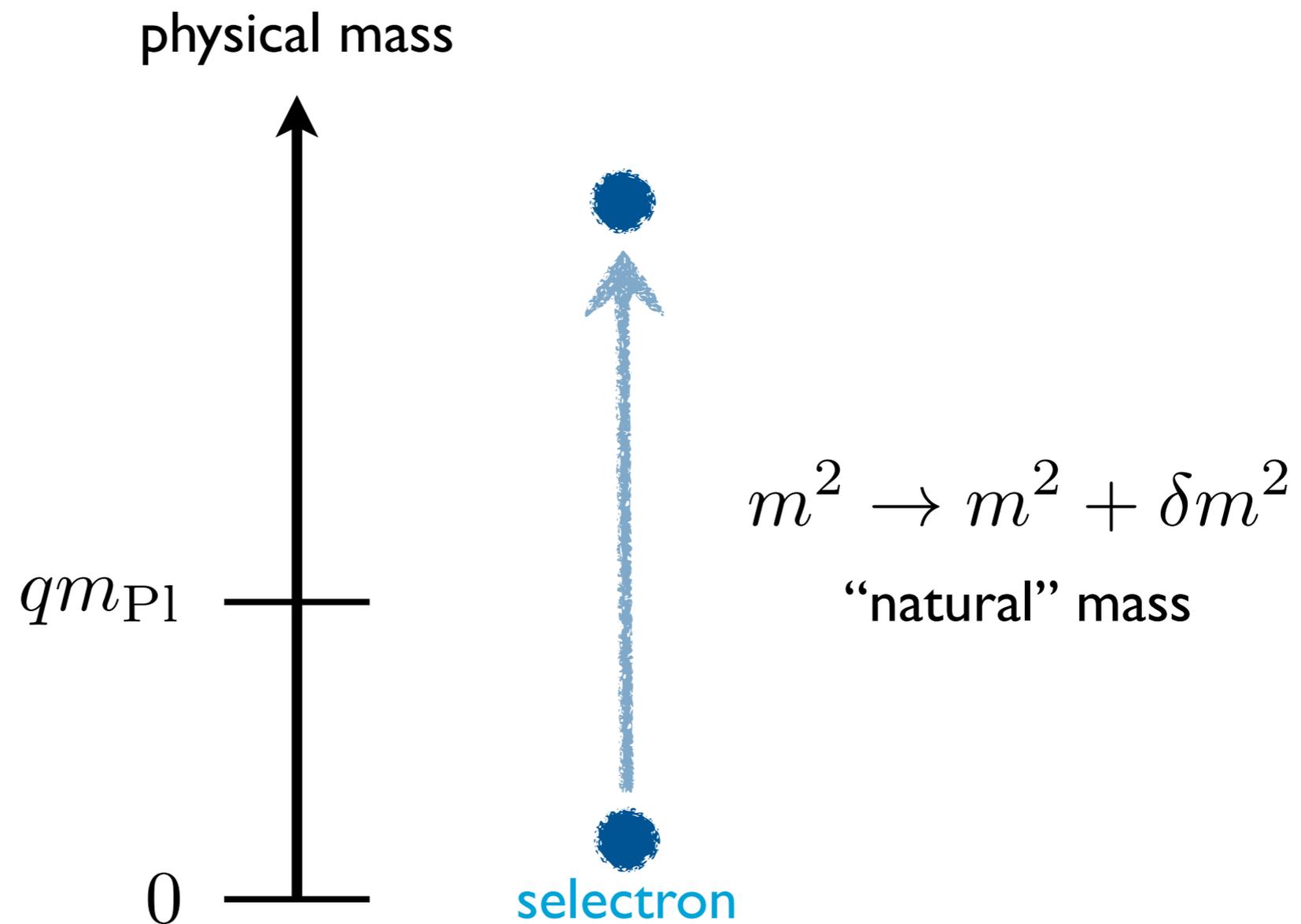
physical mass



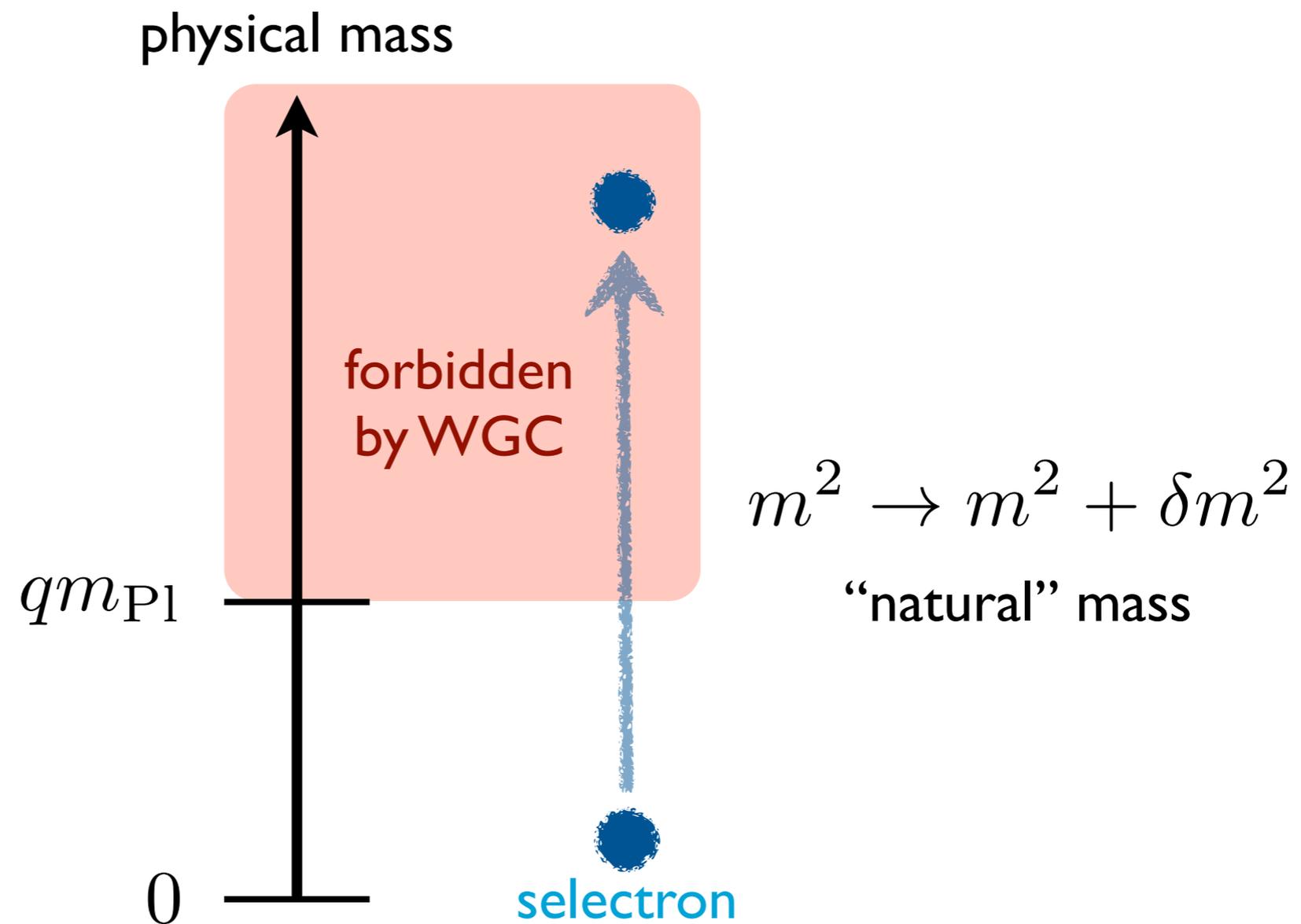
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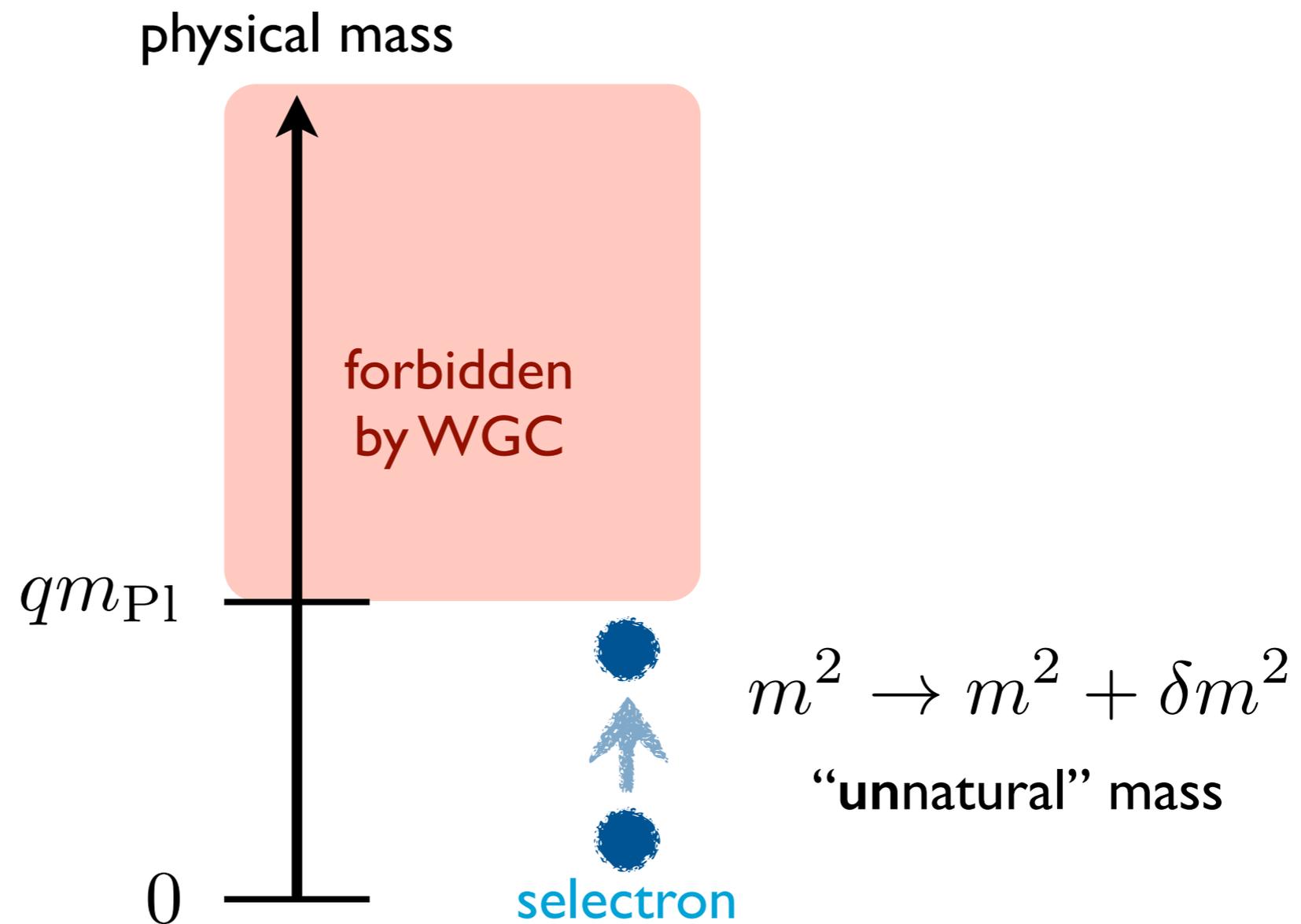
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There is fundamental tension between naturalness and the WGC.



There is fundamental tension between naturalness and the WGC.



Let's quantify the tension.

$$m^2 \rightarrow m^2 + \delta m^2$$

$$\delta m^2 = \frac{\Lambda^2}{16\pi^2} (aq^2 + b\lambda)$$

incalculable coefficients



Naturalness principle: absent symmetries, the physical mass squared is $\sim \delta m^2$, so a, b are $\mathcal{O}(1)$ coefficients.

Setting the physical mass equal to its natural value yields a charge to mass ratio

$$z = qm_{\text{P1}}/m$$
$$= \frac{4\pi m_{\text{P1}}}{\Lambda} \frac{1}{\sqrt{a + b\lambda/q^2}}$$

Since the charged scalar is the only state in the spectrum, the WGC implies

$$z > 1$$

So, the loop cutoff is bounded from above.

$$\Lambda < \frac{4\pi m_{\text{Pl}}}{\sqrt{a}} \quad q^2 \gg \lambda$$

$$\Lambda < 4\pi m_{\text{Pl}} \sqrt{\frac{q^2}{b\lambda}} \quad q^2 \ll \lambda$$

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reasonable: cutoff
below Planck

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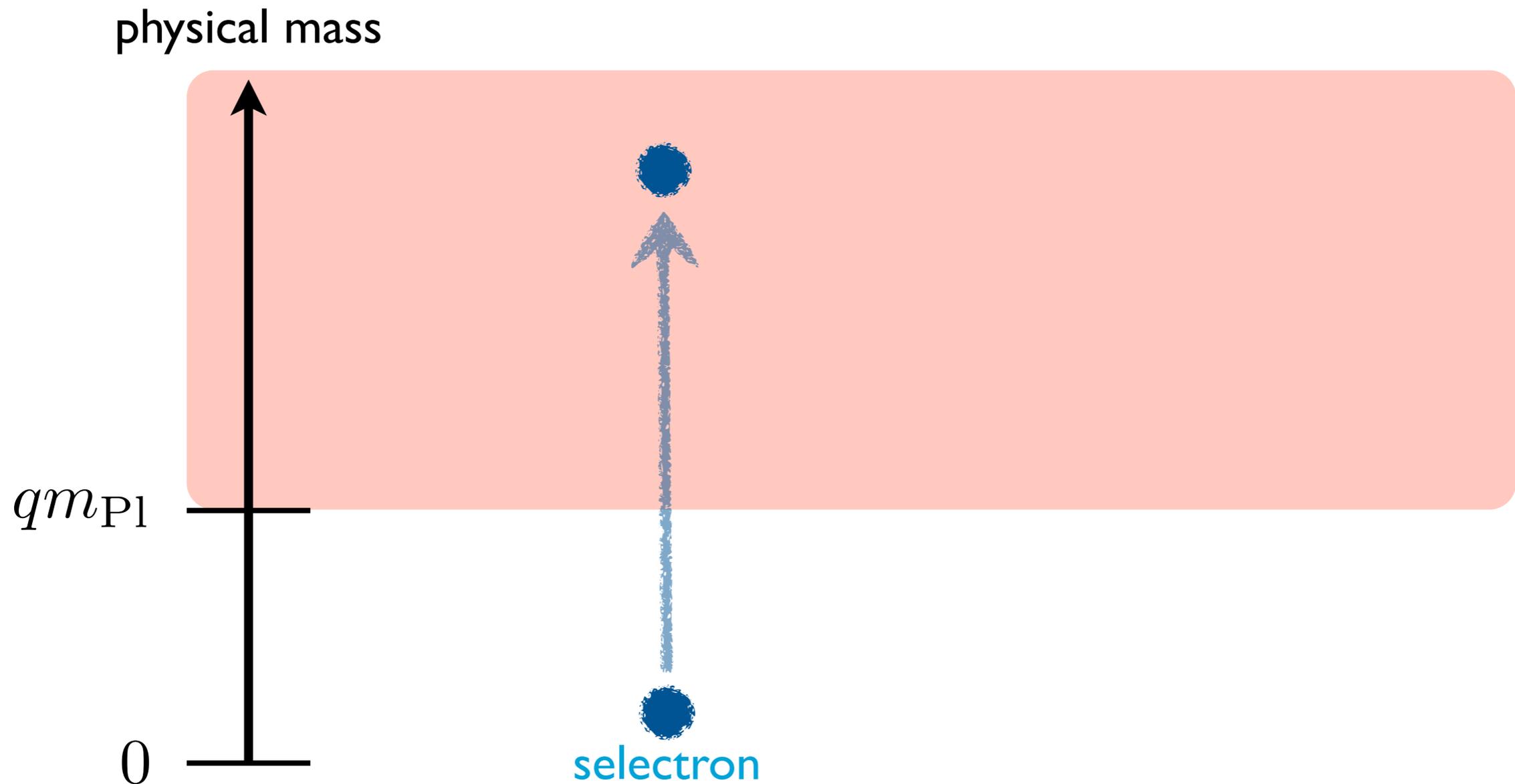
$$\Lambda < 4\pi m_{\text{Pl}} \sqrt{\frac{q^2}{b\lambda}} \quad q^2 \ll \lambda$$

parametrically low cutoff!
(conjectured in original paper)

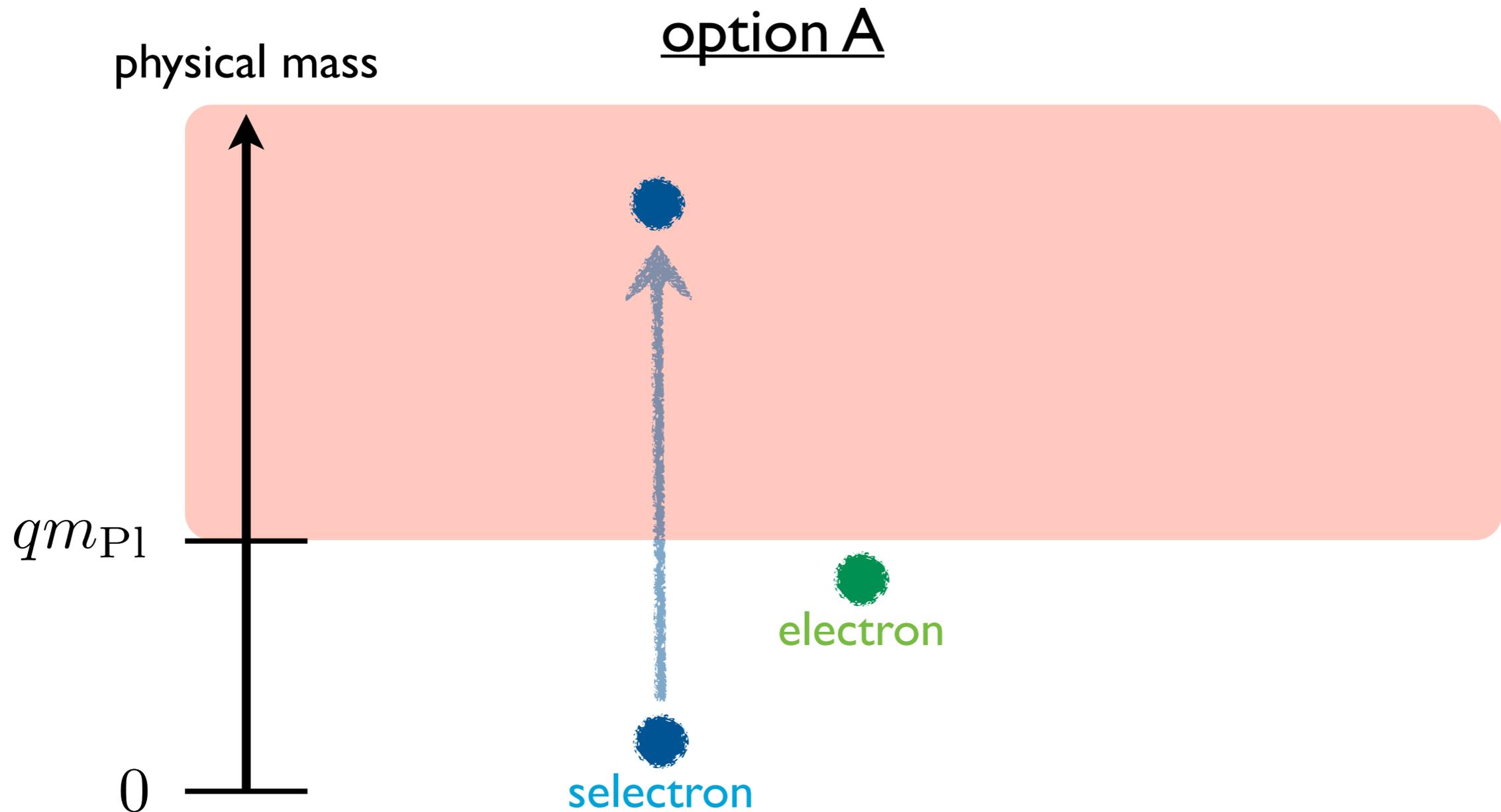
Naturalness and WGC *can* be reconciled if we revisit and *modify* our premises.

There is an obvious strategy.

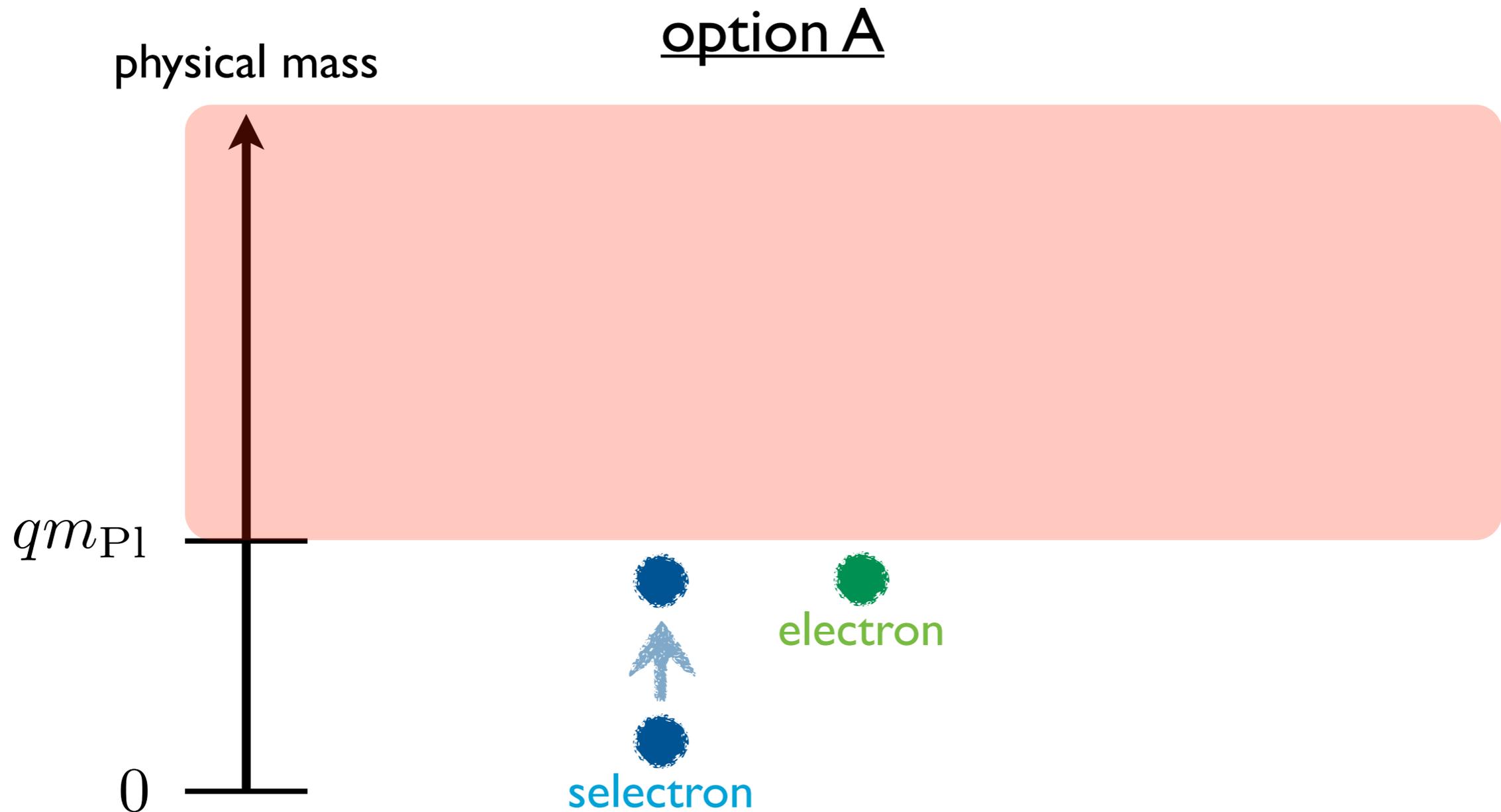
Add new degrees of freedom below cutoff.



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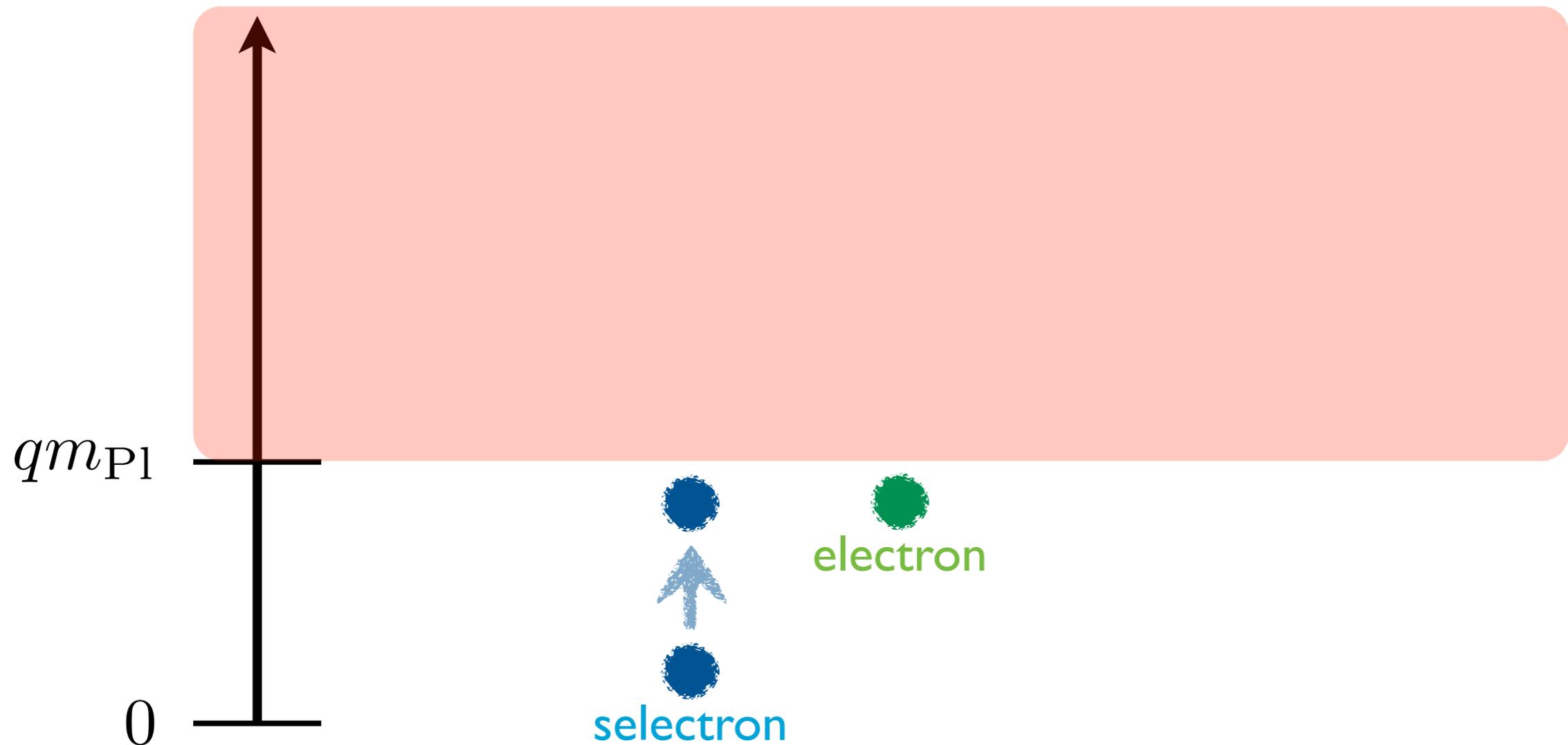


Add new degrees of freedom below cutoff.

(uninteresting)

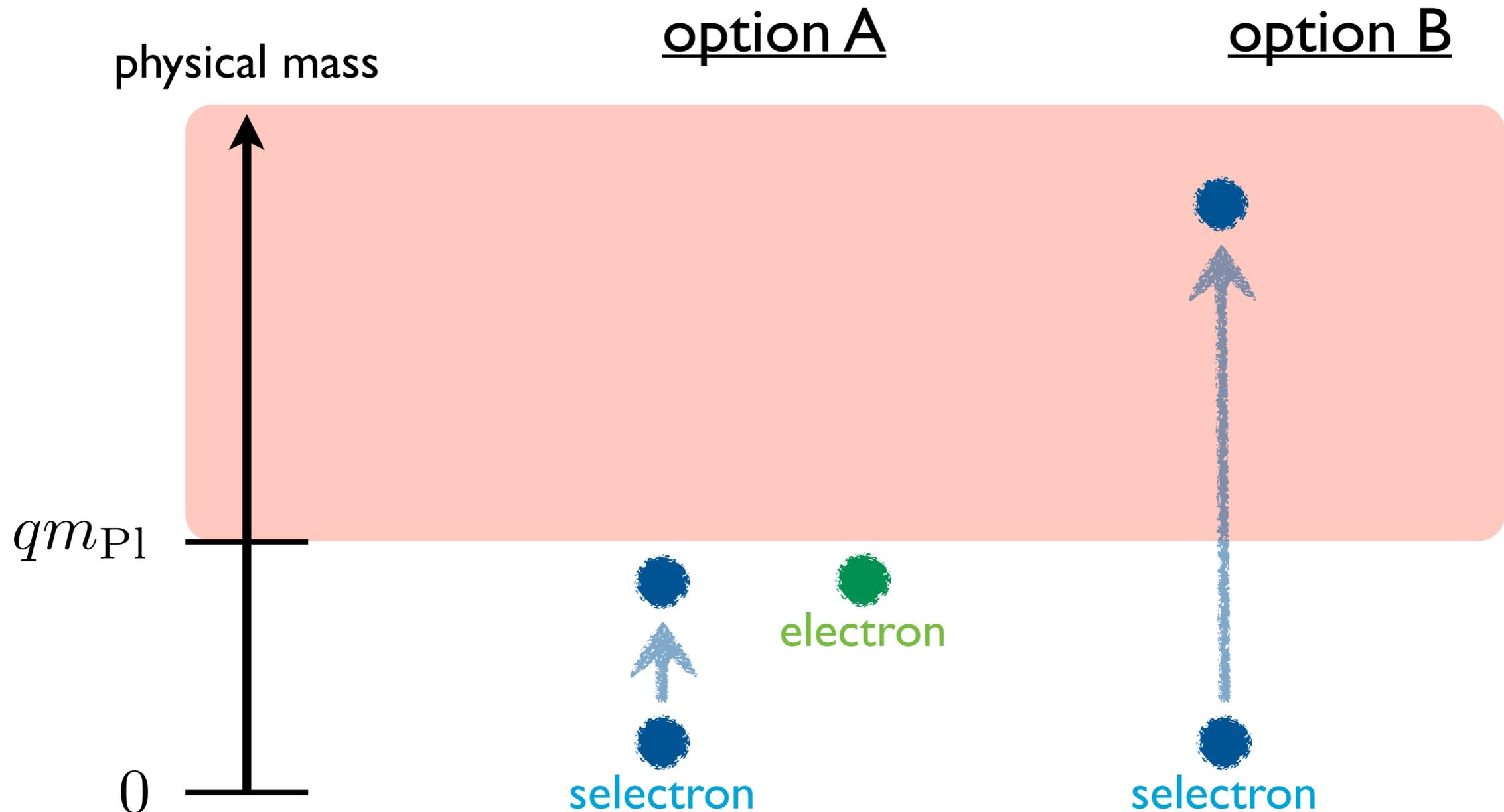
option A

physical mass



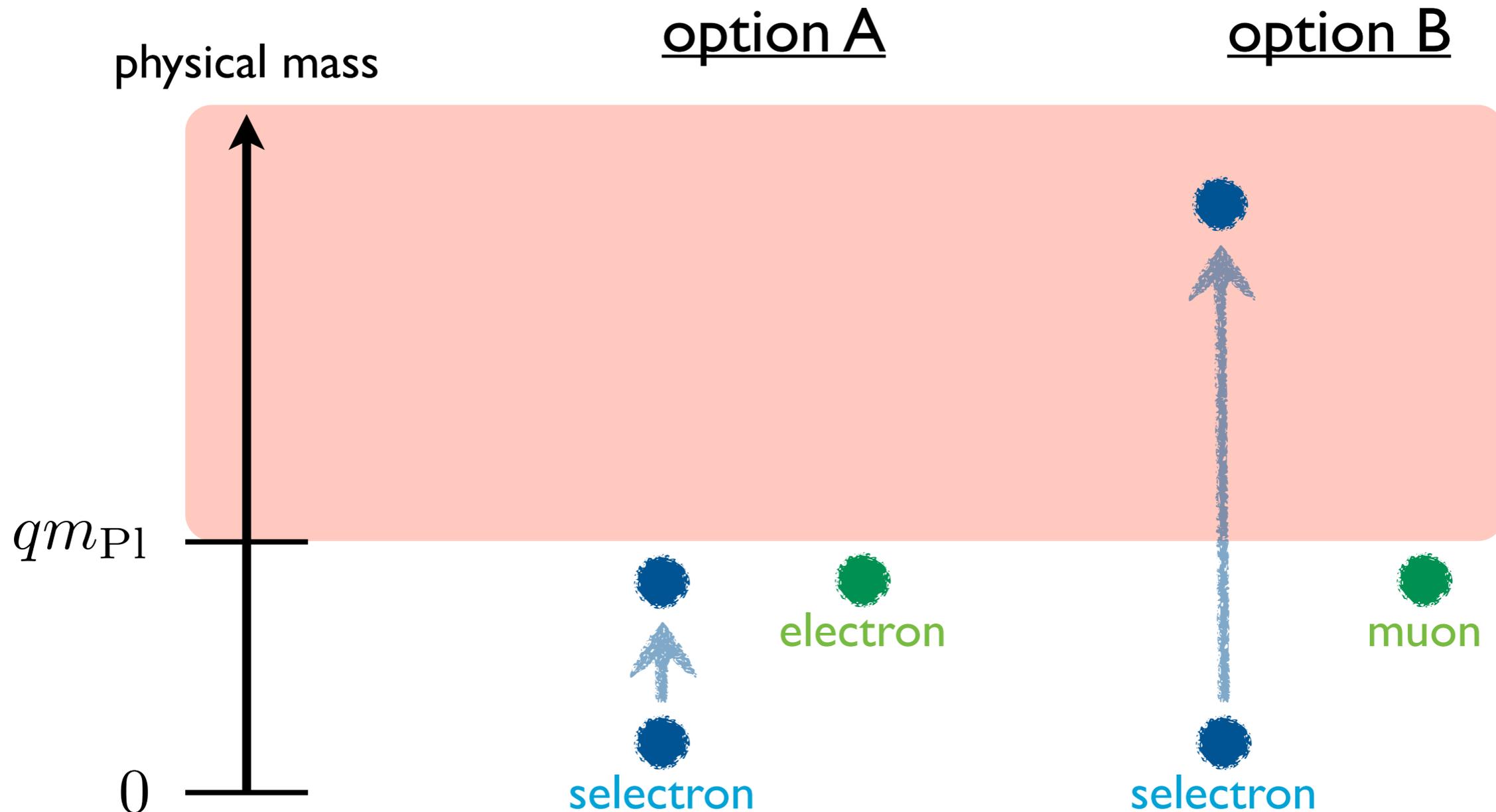
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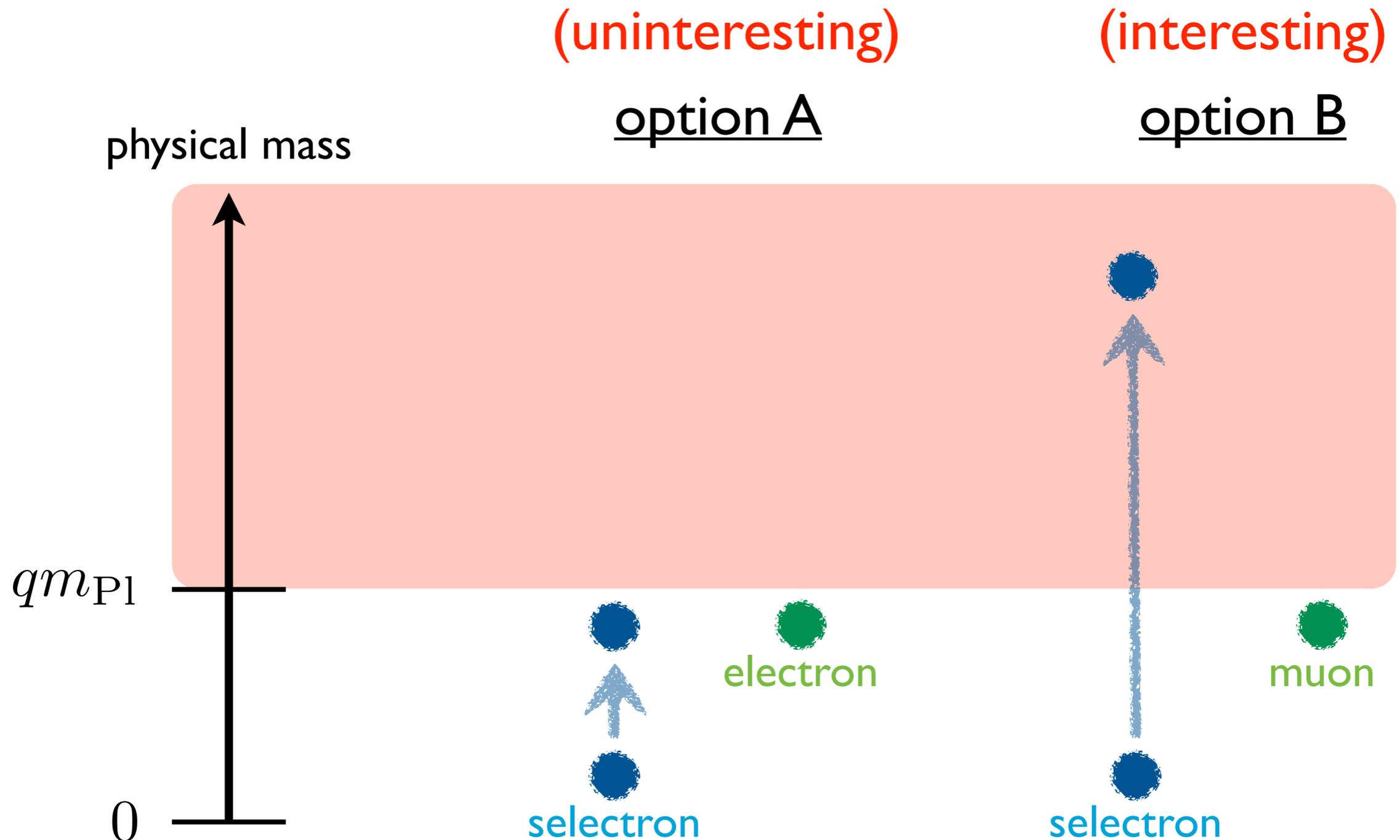


Add new degrees of freedom below cutoff.

(uninteresting)



Add new degrees of freedom below cutoff.



(technically, another option : Higgs phase)

$$\delta m^2 < 0$$

The WGC is ambiguous in the Higgs phase because $[q, m] \neq 0$. Whose mass, charge?

More importantly, black holes do not have Higgsed U(1) hair. No justification for WGC!

Some lessons from scalar QED:

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say we observe
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$!(\text{low cutoff}) \rightarrow \overset{\text{False}}{!(\text{small charge})} \parallel !(\text{naturalness})$


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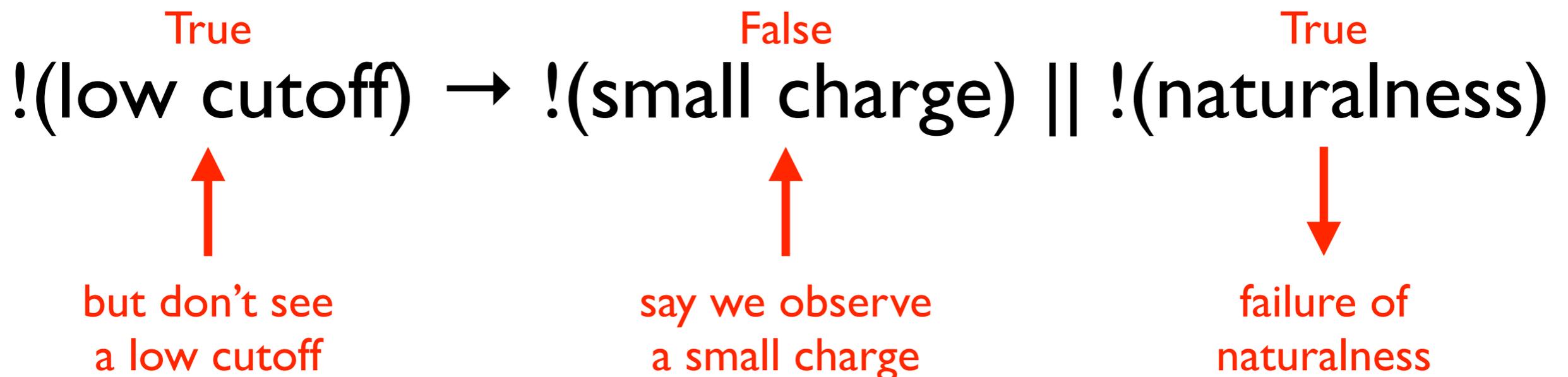
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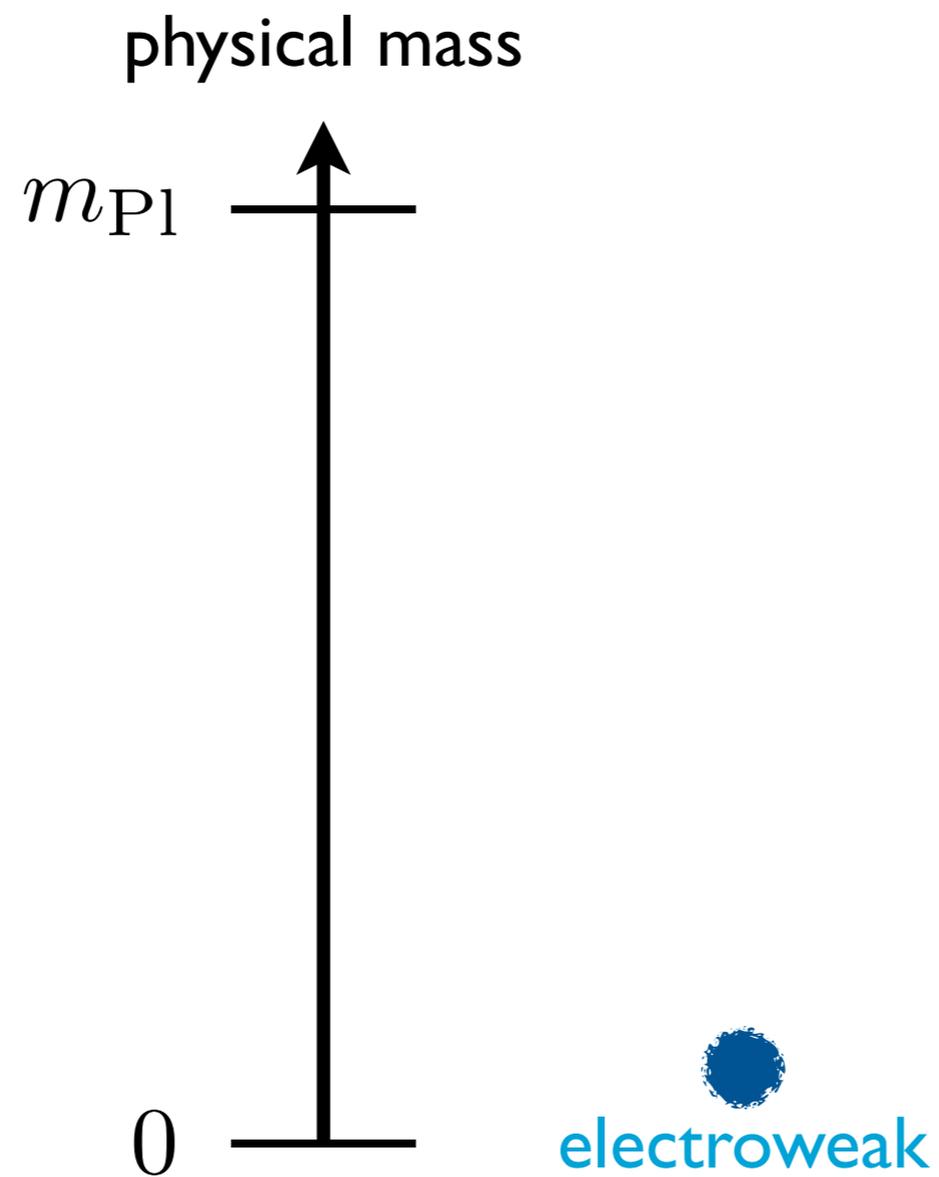
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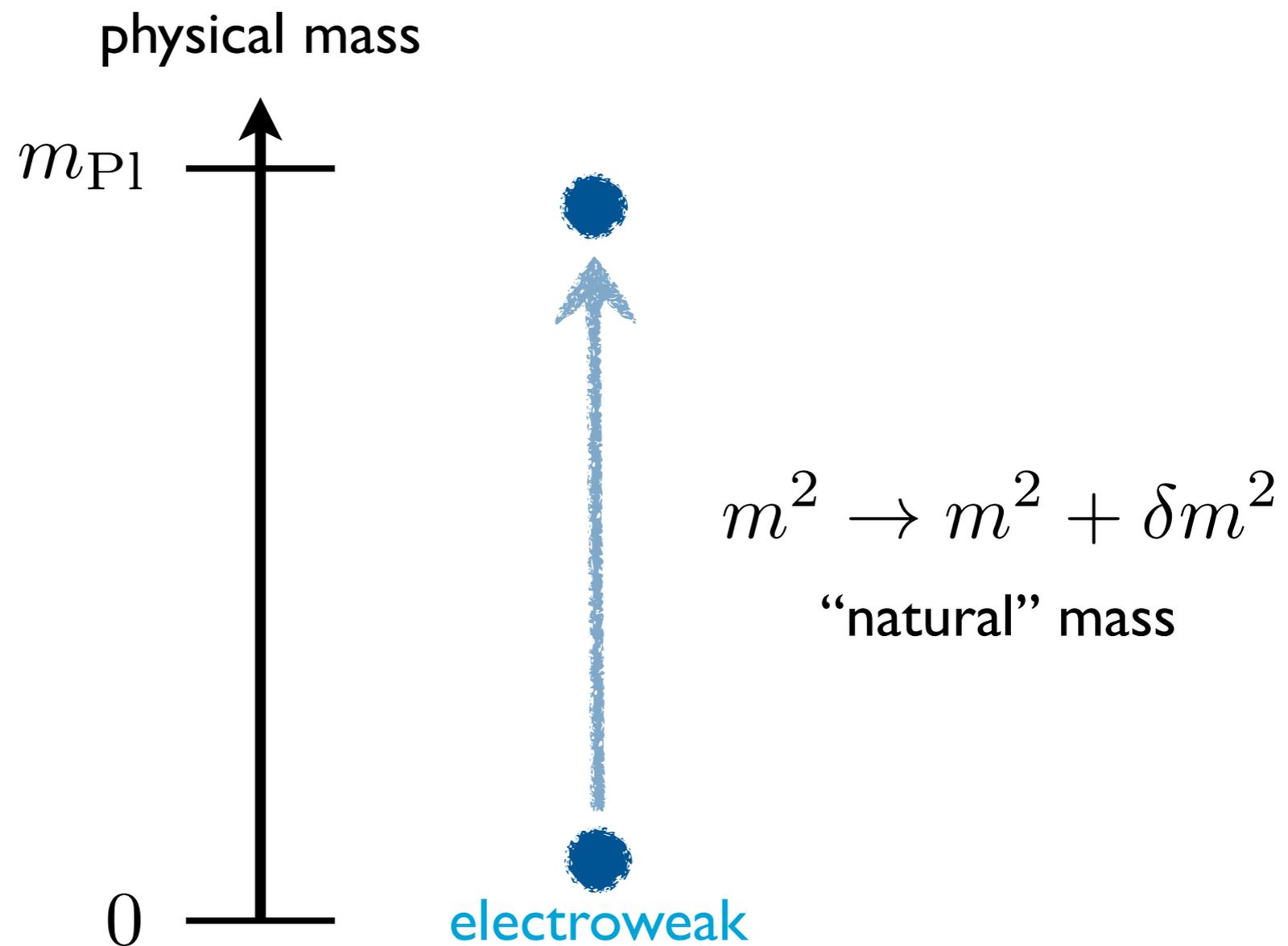
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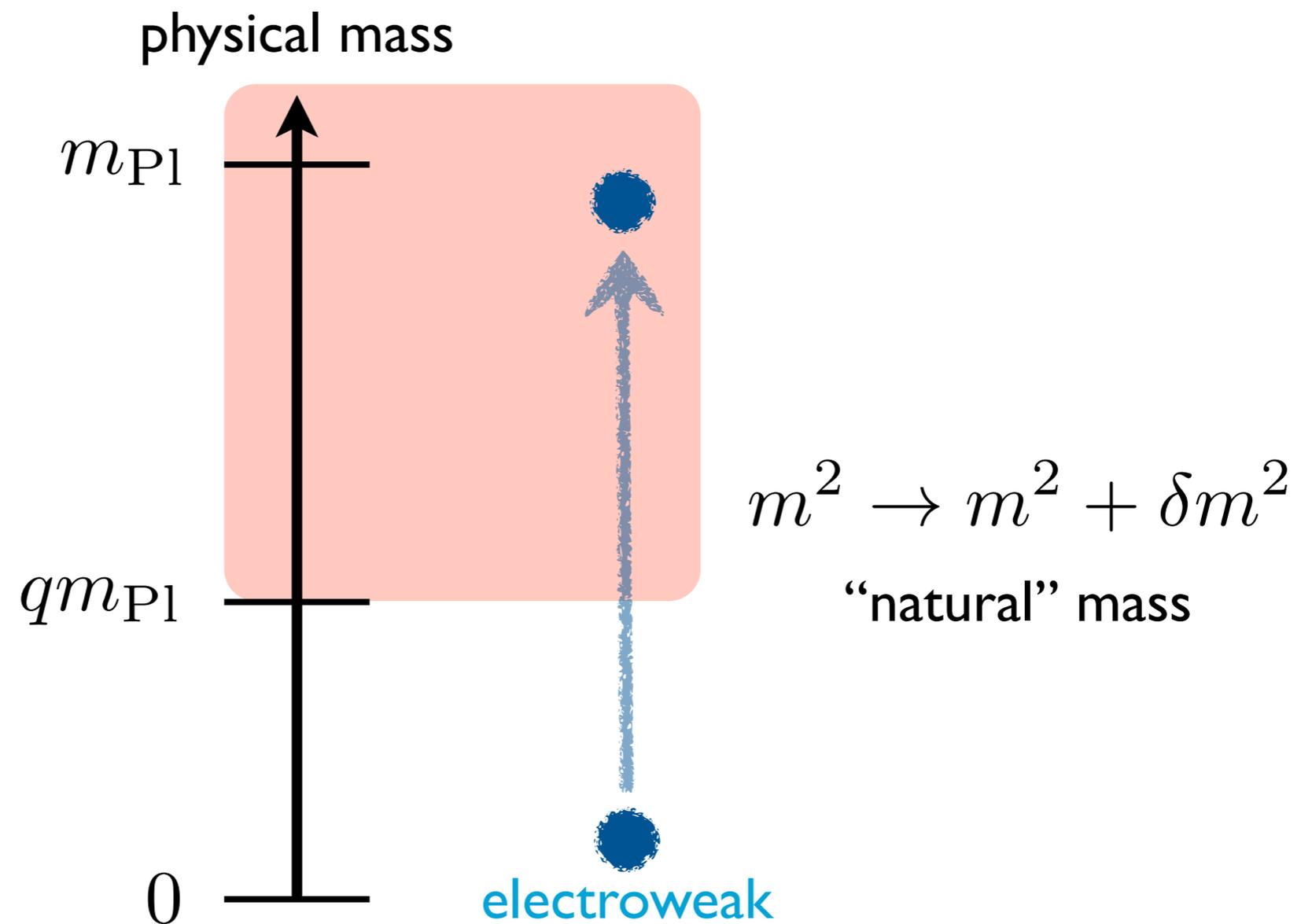
Let's apply this to the hierarchy problem.



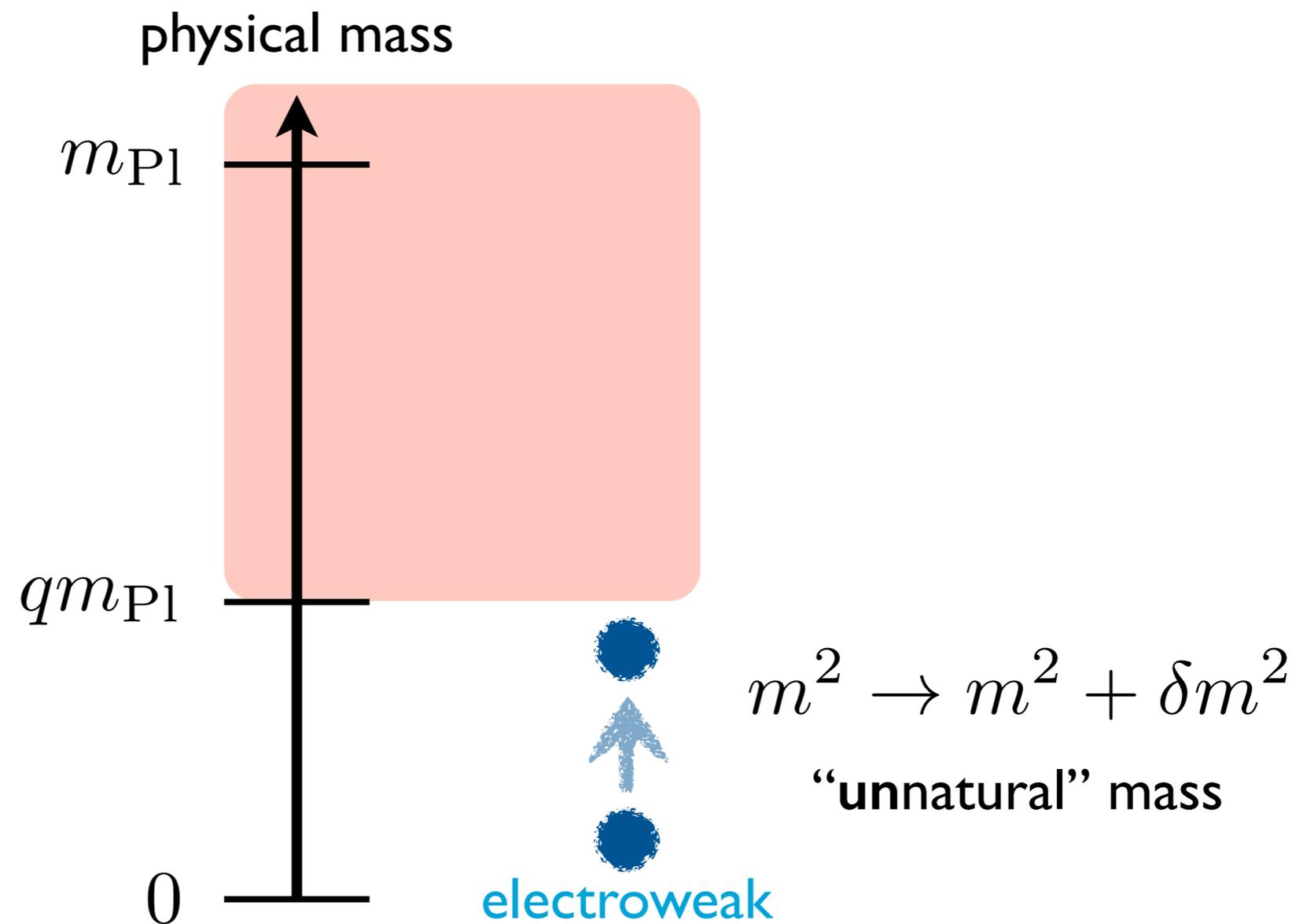
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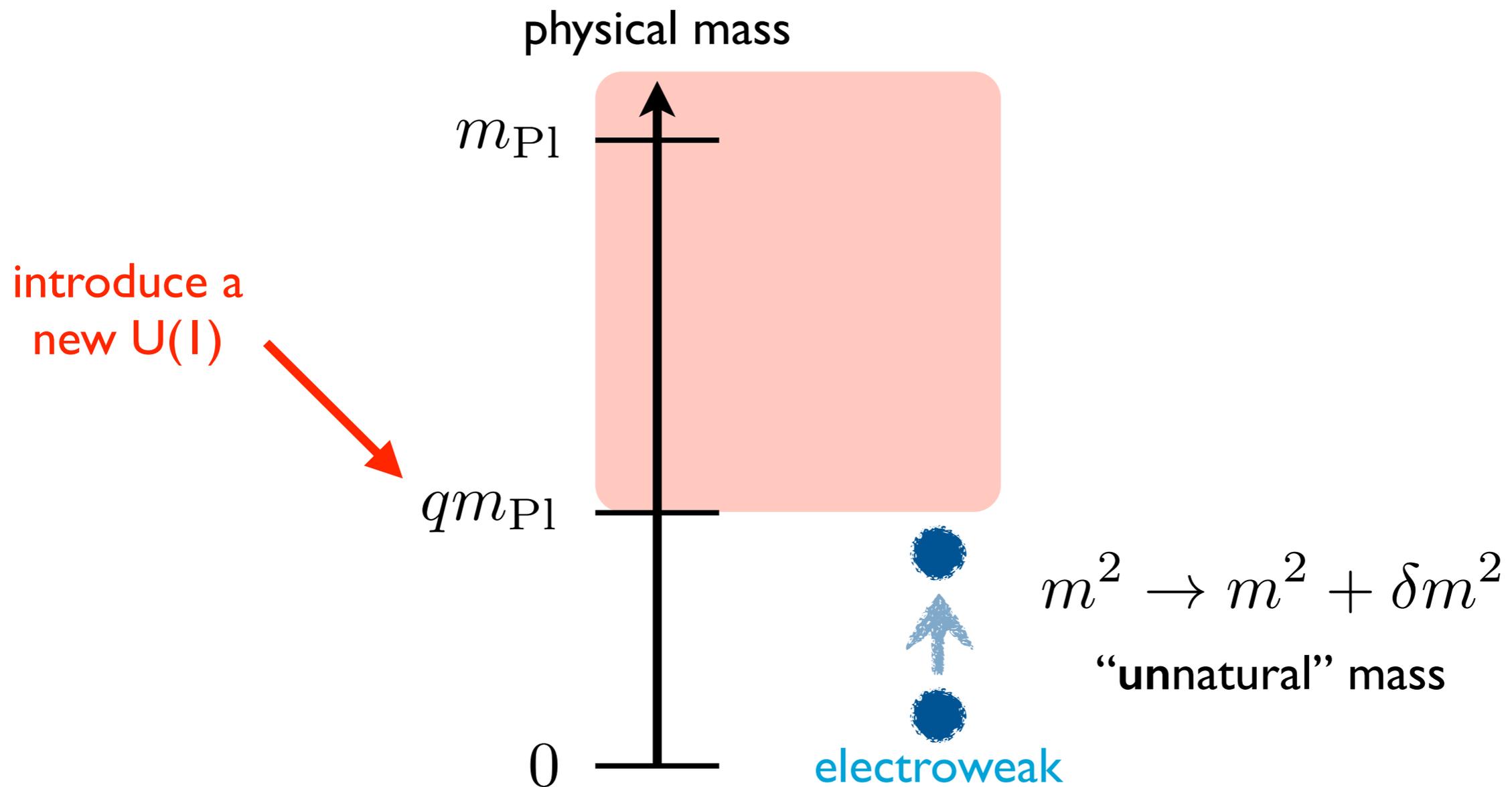
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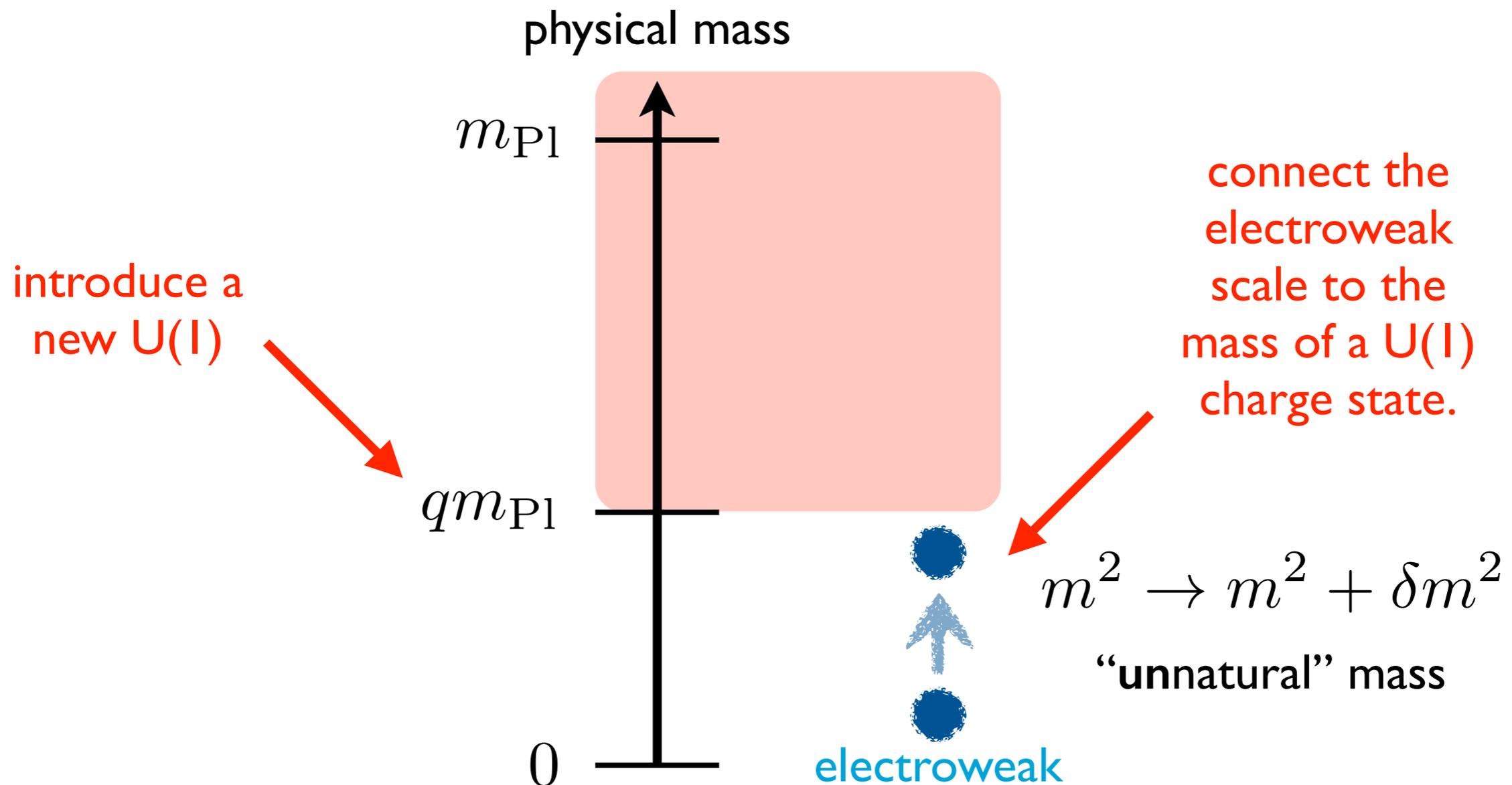
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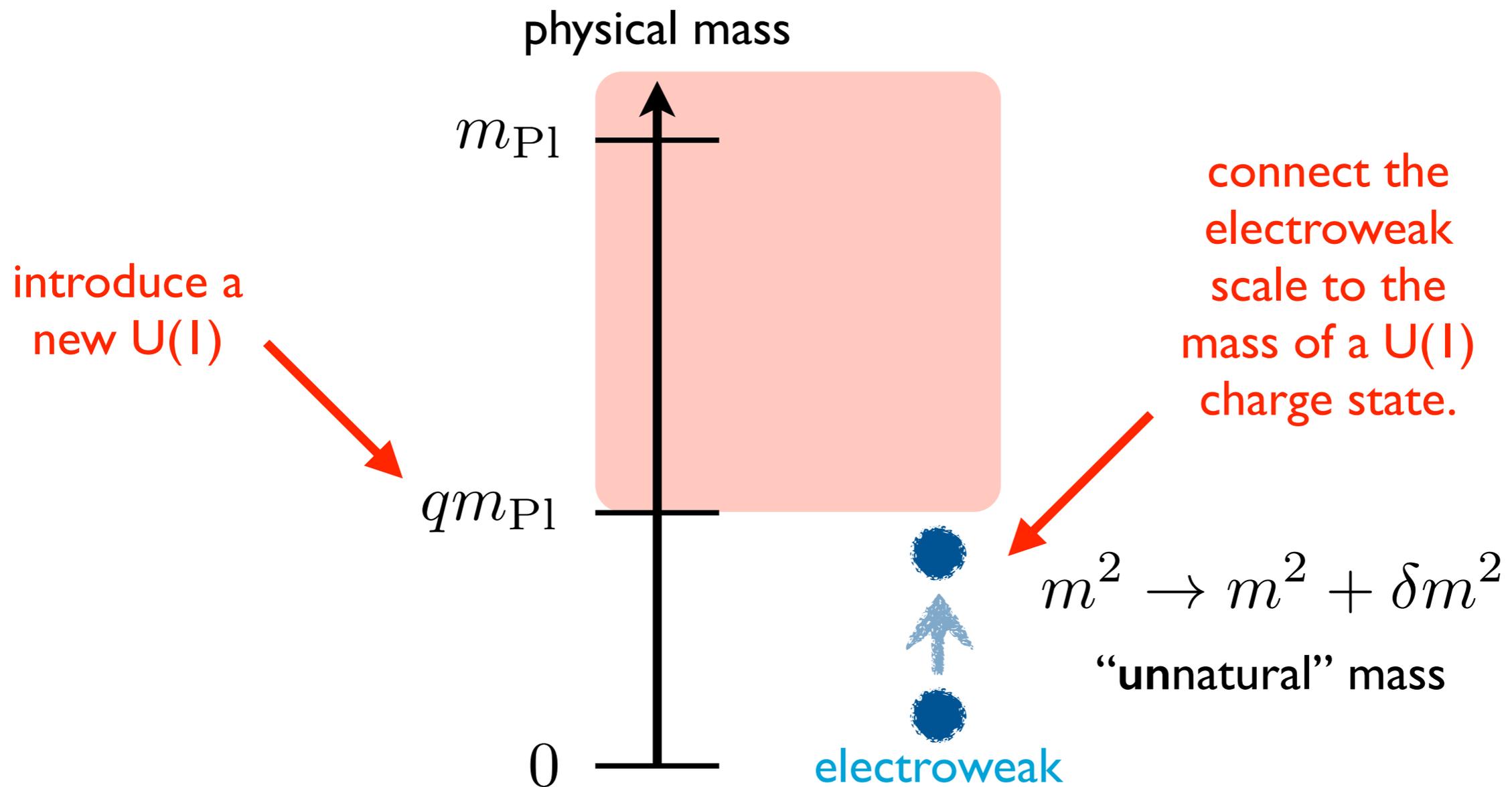
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Let's apply this to the hierarchy problem.



The electroweak scale is unnatural, but only because a natural value is forbidden!



a simple model

Weakly gauge $U(1)_{B-L}$ with Dirac neutrinos.

$$-\mathcal{L} = m_\nu \bar{\nu}_L \nu_R + \text{h.c.} \quad m_\nu \sim y_\nu v$$

Assuming that $m_\nu \sim 0.1$ eV, we fix

$$q \sim 10^{-29} \quad (\sim m_\nu / m_{\text{Pl}})$$

so that the WGC is marginally satisfied.

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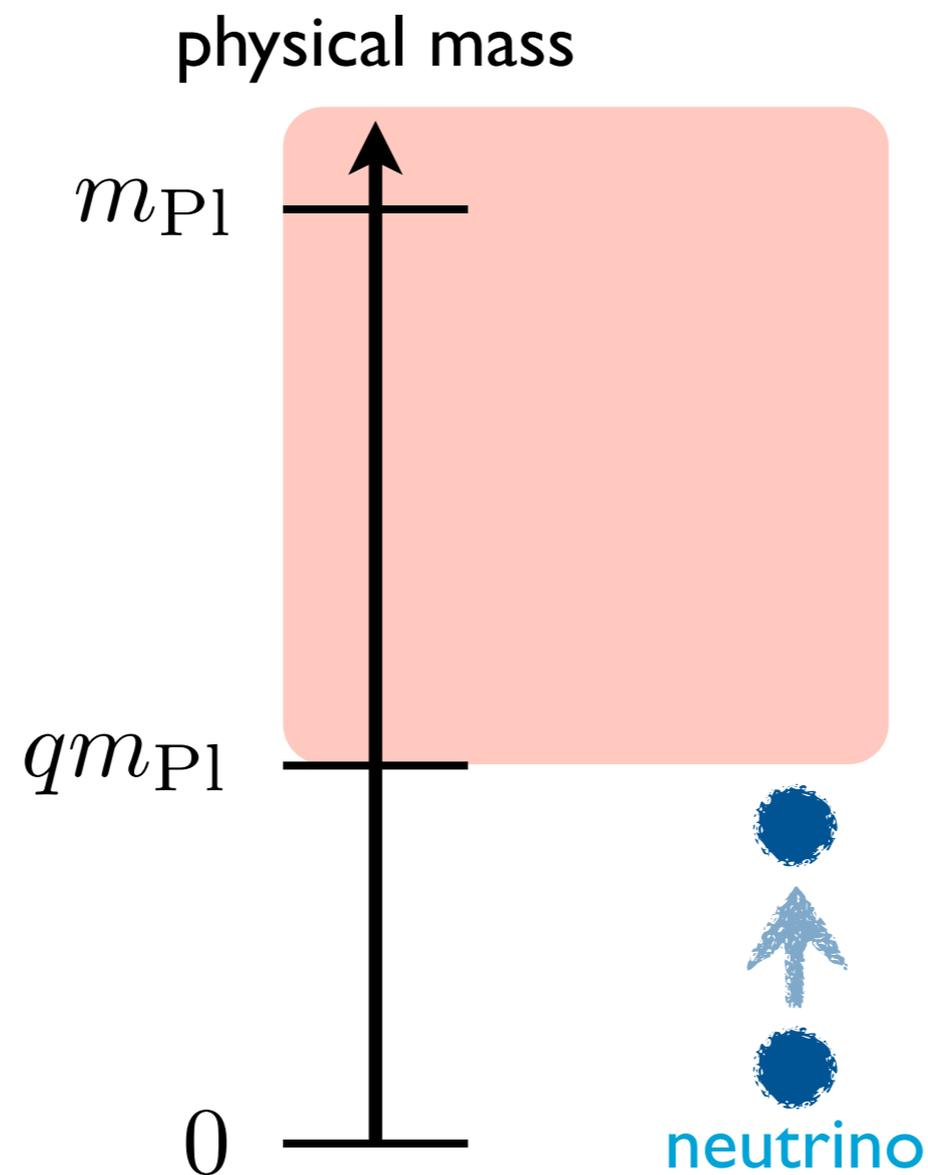
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Assuming that $m_\nu \sim 0.1 \text{ eV}$, we fix

(technically natural) $q \sim 10^{-29} \quad (\sim m_\nu / m_{\text{Pl}})$

so that the WGC is marginally satisfied.

Fixing couplings, were the electroweak scale larger, then the WGC condition would fail.



The model is a proof of concept but it has a prediction: a massless gauge boson.

There are very stringent limits of fifth forces and violation of equivalence principle:

$$q \lesssim 10^{-24} \quad (\text{torsion balance})$$

An incredibly small charge!

But, naturalness & small charge \rightarrow low cutoff.

$$q \lesssim 10^{-24} \quad \longrightarrow \quad \Lambda \lesssim q m_{\text{P}1} \lesssim \text{keV}$$

Such an extremely low cutoff is not there!

Hence, literally any fifth force observation will exclude the low cutoff conjecture and invalidate the argument from naturalness.

conclusions

- “Running Naturalness” is an irreducible notion of fine-tuning.
- Still, fine-tuning presumes evenly thrown darts in the UV as boundary cond.
- Perhaps “on-shell” is somehow more fundamental than “UV” (cf. S-matrix).
- Perhaps there is UV selection. WGC forbids certain natural theories.
- Milli-charges offer an experimental test.

thanks!