

Recent progress of HCC D&S

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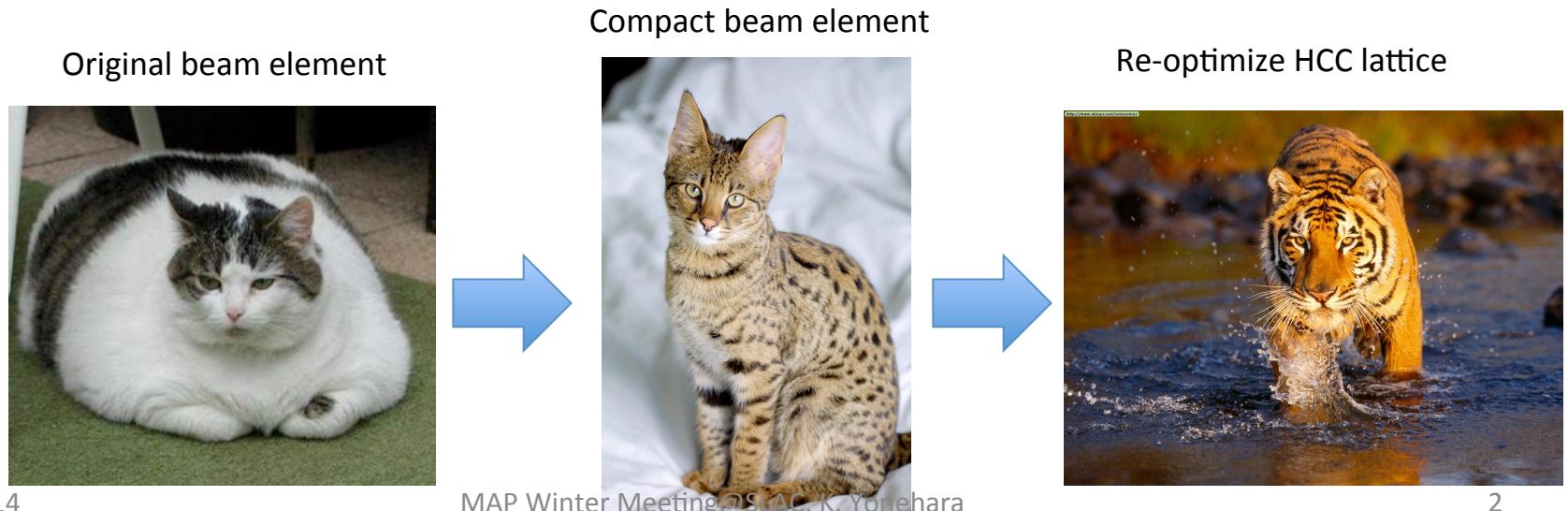


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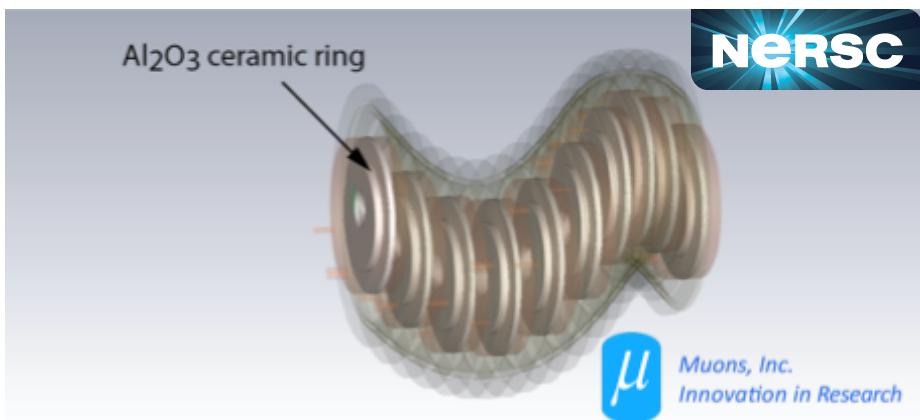
- R&D of HCC beam element
 - How to mitigate the geometry constraint?
- Re-optimizing HCC lattice
- Summary



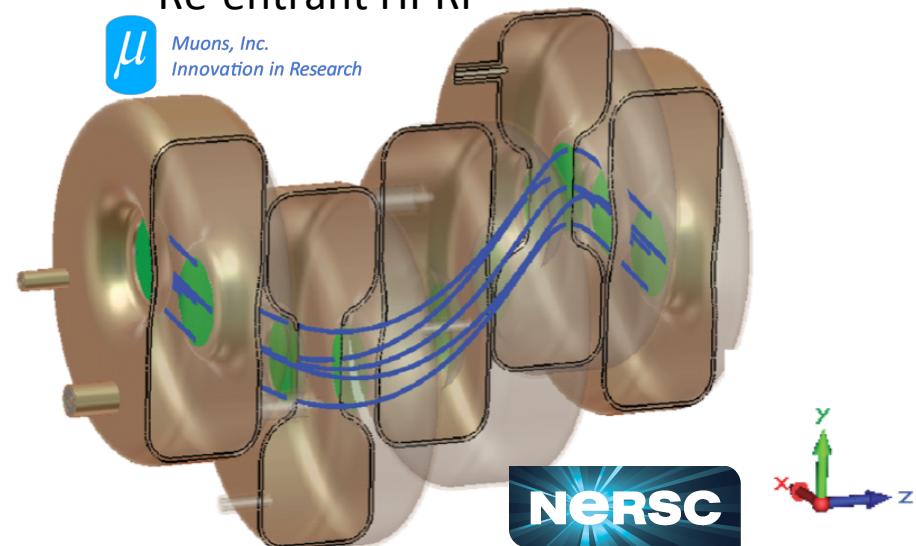


HCC RF system

Dielectric loaded HPRF

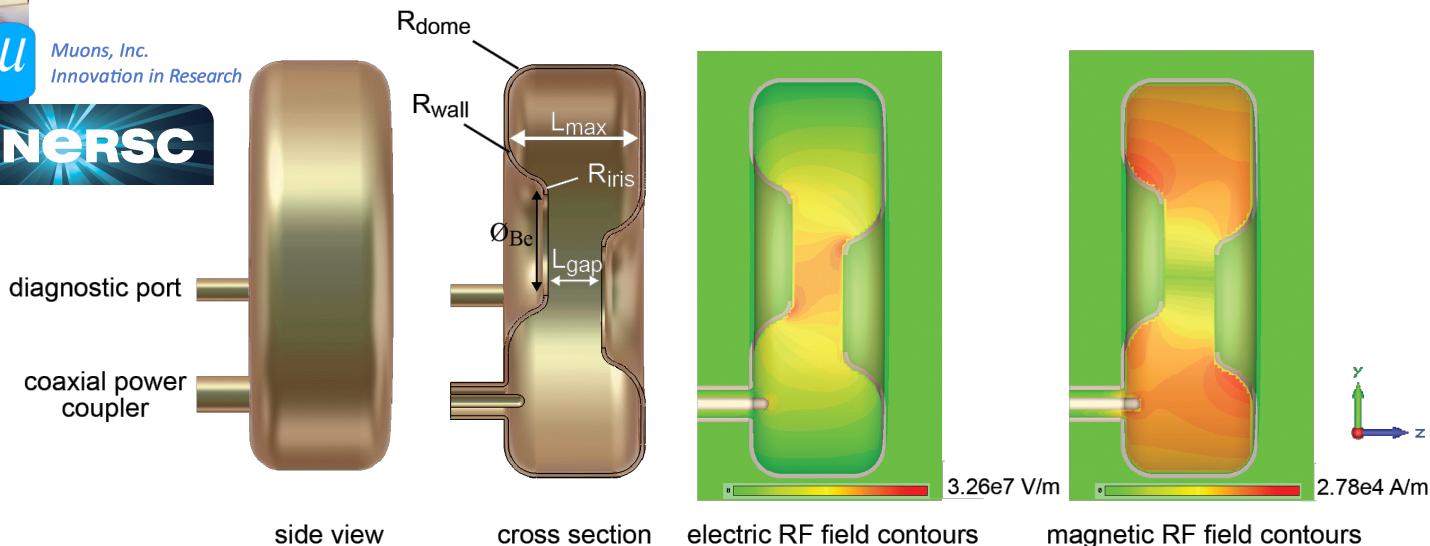


Re-entrant HPRF



Re-entrant HPRF cavity

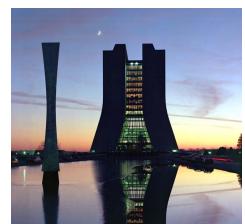
μ Muons, Inc.
Innovation in Research



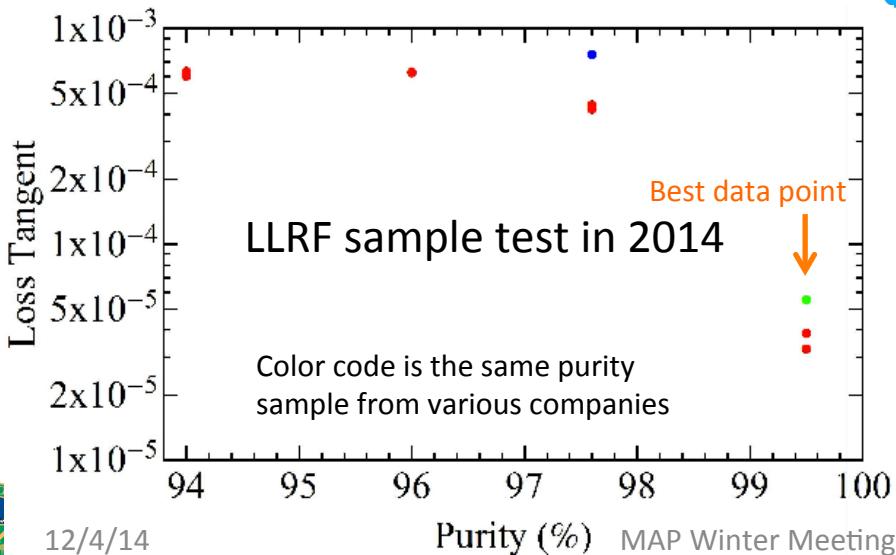
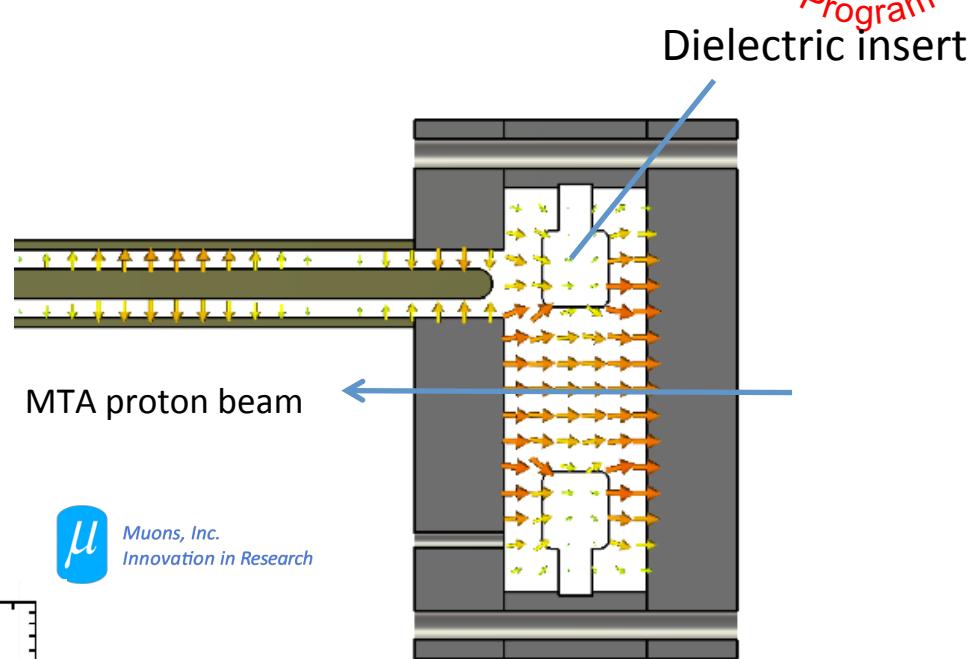
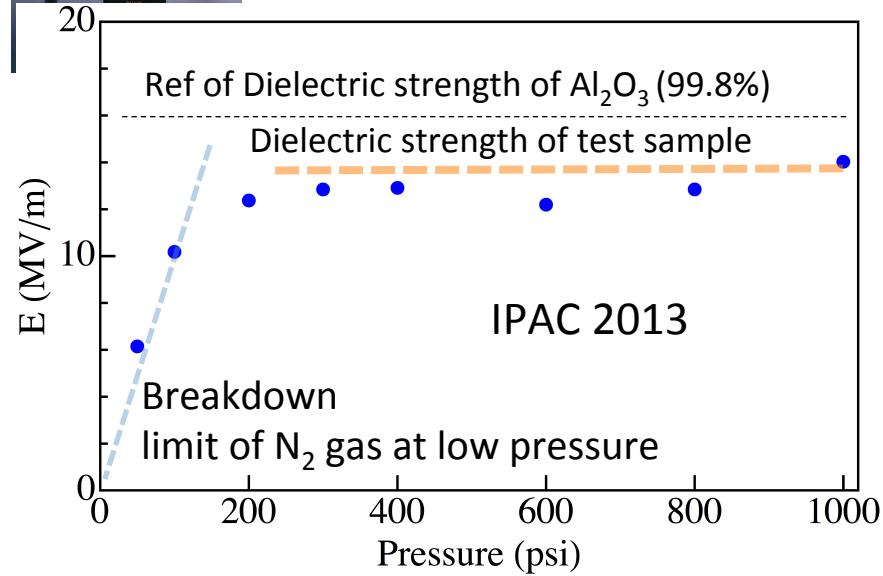
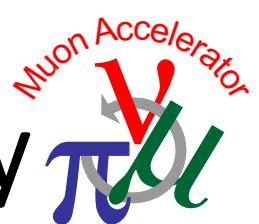
- Since huge RF power (see magnetic field above plot) can be stored in a dome the required peak power is almost factor 10 smaller than the dielectric loaded RF (4.4 MW/m @ 300 K, 325 MHz)
- However, the peak RF gradient should be higher than the DLRF
 - ▶ Accelerating gap in the re-entrant RF is 1/2 of the DLRF
 - ▶ Thus, the plasma loading in the re-entrant RF cavity becomes significant
- Real EM fields of the re-entrant RF was tested in g4bl and observed cooling
- Longitudinal matching is the current issue
 - ▶ Too strong longitudinal kick to hold particles in the RF bucket

Require more lattice tuning & test cooling simulation

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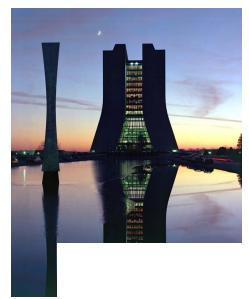
Dielectric Loaded HPRF cavity



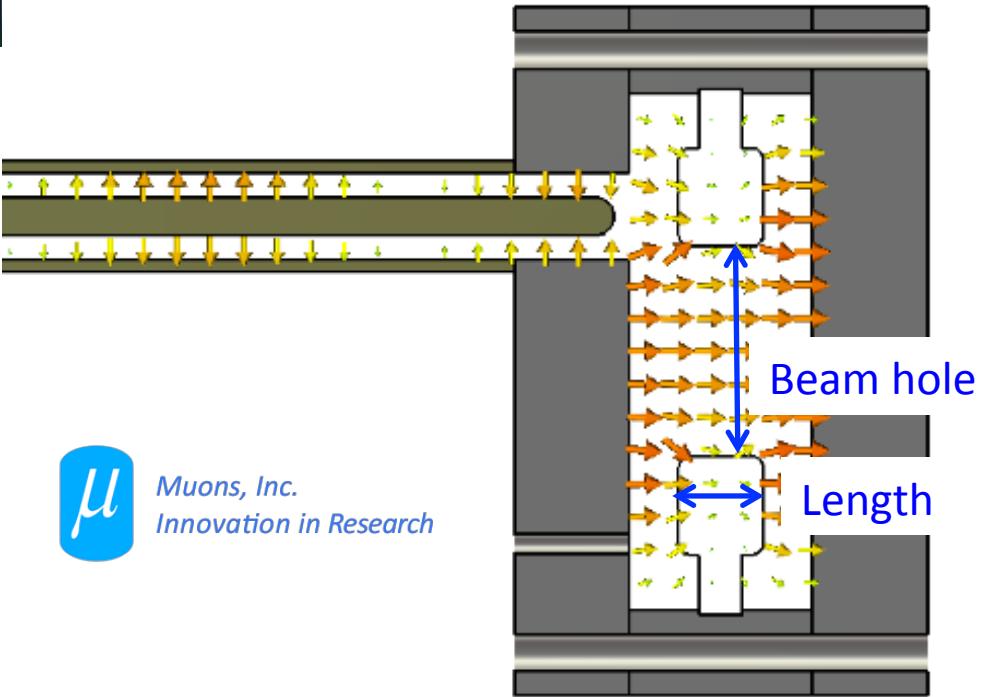
Proposed dielectric loaded HPRF test cell
Tests will be made in FY15 & 16 to observe
the loss tangent and the dielectric strength
at high RF power with intense beam

Test, Test, Test!





Strategy of DL HPRF test for HCC



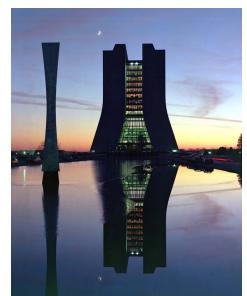
If the high power test shows that the dielectric strength of the 99.5 % sample is very high ($> 20 \text{ MV/m}$) we can re-evaluate the insert with **smaller beam hole and longer** than the test sample



Allow to design a very compact cavity

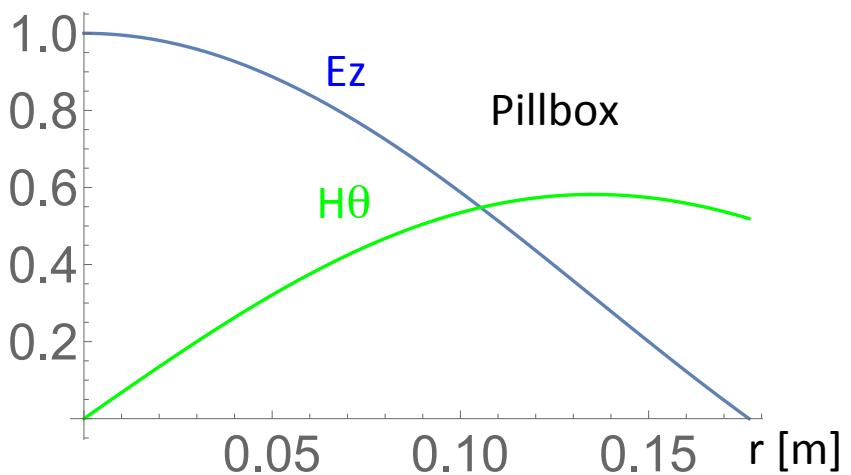
- Original design was made to reduce the risk of permanent damage of ceramics
 - We assumed to use a 99.8 % ring
 - Big tolerance for operating gradients
- From recent tests,
- Obtained loss tangent of 99.5 % is twice lower than our original design
 - Company spec shows that the 99.5 % is twice higher dielectric strength than 99.8 %
 - We consider a 99.5 % ring is better than the 99.8 % one





Considerable advantages in DL HPRF

$f = 650 \text{ MHz}$, $E = 20 \text{ MV/m}$, $\tan \delta = 0.4e-4$, $\epsilon = 9.6$

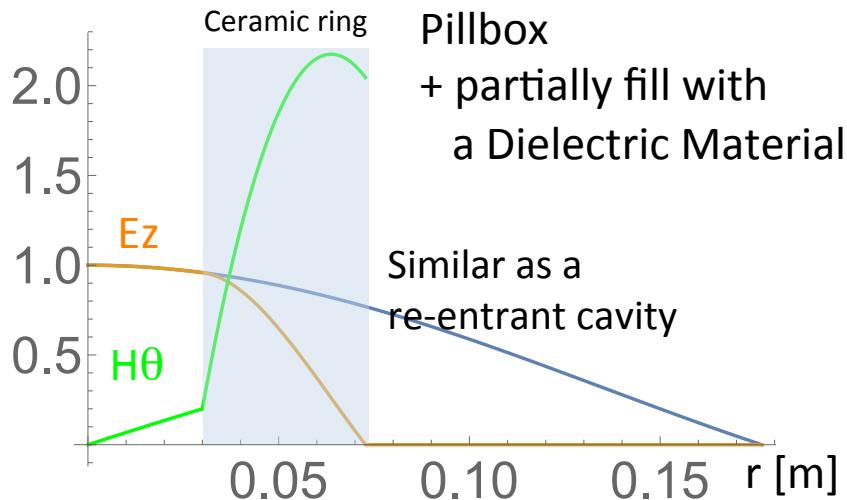


$$\underline{IR = 0.177 \text{ m}}$$

$$Q = 7,518$$

$$\underline{U = 2.34 \text{ J}}$$

$$P = 1.27 \text{ MW}$$



$$\underline{IR = 0.073 \text{ m}}$$

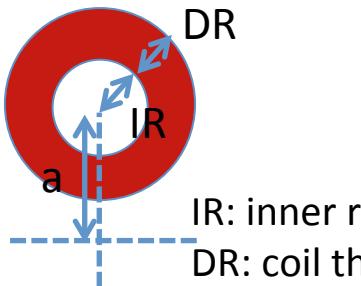
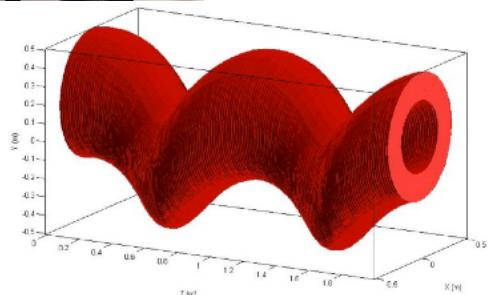
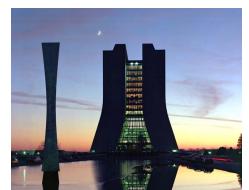
$$Q = 5,073$$

$$\underline{U = 4.17 \text{ J}}$$

$$P = 3.36 \text{ MW}$$

- Cavity radius of the DL HPRF is significantly small
- DL HPRF has almost twice larger stored energy than the pillbox
- Peak power is still high but getting lower at smaller beam hole

Constraint of HCC magnet

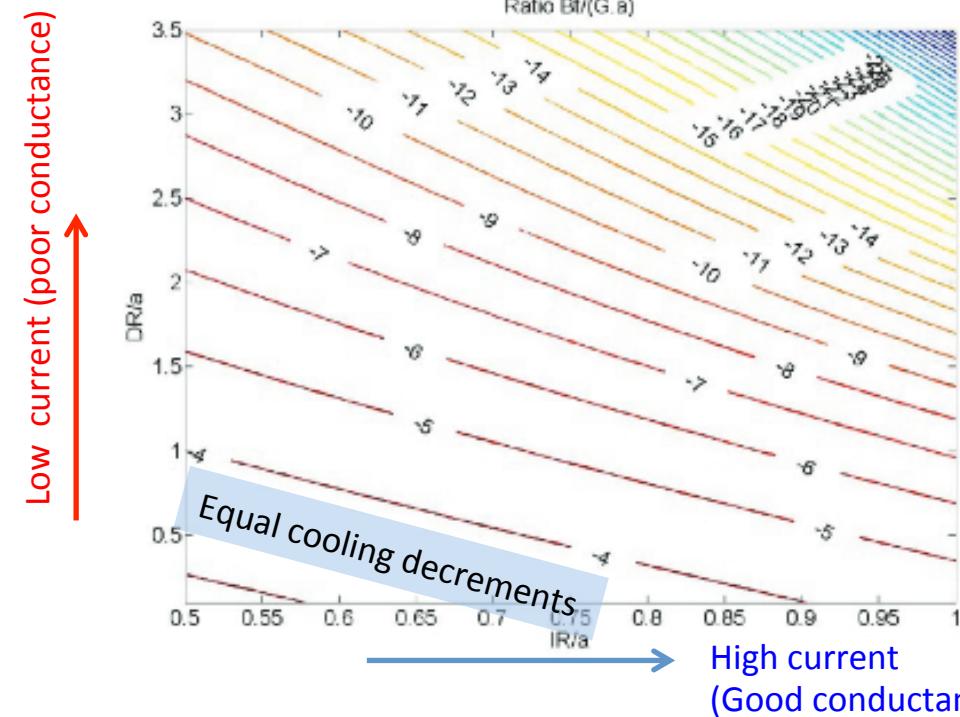


IR: inner radius

DR: coil thickness

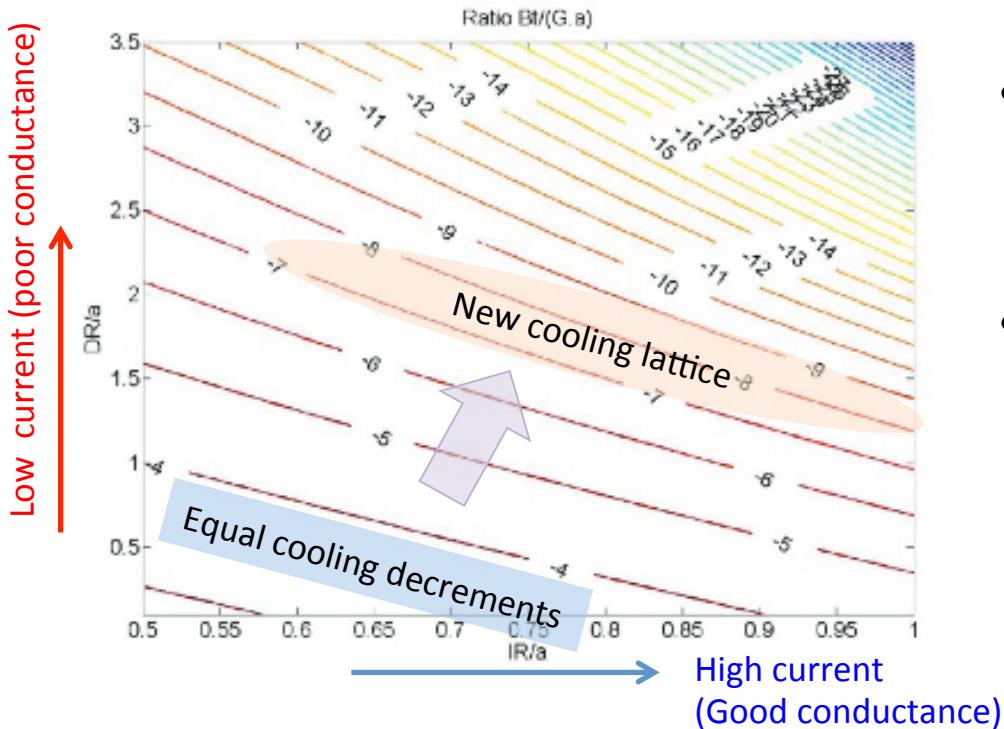
a: Helical radius = $\kappa \lambda / 2 \pi$

- M. Lopes made a field calculation code SolCalc written in Matlab (IPAC'14)
- He demonstrated the limit of IR at the equal cooling decrements



- Left plot shows the correlation between the HS coil geometry and the ratio of helical dipole and field gradient
- Dispersion is determined from the strength of helical dipole and the field gradient (see next slide)
- Therefore, the cooling decrements in the HCC is determined by the coil geometry
- Maximum allowable current (conductance) is another parameters to determine the coil geometry
- **Coil geometry becomes tighter (smaller IR) if higher current is needed**

Propose new configuration of HCC magnet



- Relax the constraint of HCC magnet!
Larger IR & thinner DR
→ Helical dipole stronger
→ Field gradient smaller
- What happens cooling performance in the new lattice configuration?

Emittance evolution in cooling channel

Generic emittance evolution formulae

$$\frac{d\epsilon_n}{ds} = \epsilon \frac{d\beta\gamma}{ds} + \beta\gamma \frac{d\epsilon}{ds} = \epsilon_n \frac{1}{\beta^2 E} \frac{dE}{ds} + \frac{\beta\gamma}{2} (\hat{\beta}_\xi \sigma_\xi^2)$$

$$\epsilon_n(s) = (\epsilon_{n,0} - \epsilon_{n,eq}) e^{-\Lambda_\xi s} + \epsilon_{n,eq}$$

$\epsilon_{n,0}$; initial emittance

$$\epsilon_{n,eq} = \frac{\beta\gamma}{2} (\hat{\beta}_\xi \sigma_\xi^2) / \left(\frac{1}{\beta^2 E} \frac{dE}{ds} \right) ; \text{ equilibrium emittance}$$

$$\Lambda_\xi = \left(\frac{g_\xi}{\beta^2 E} \frac{dE}{ds} \right) ; \text{ cooling decrement}$$

$$\varphi = \int \frac{ds}{\hat{\beta}_\xi} \rightarrow \hat{\beta}_\xi = \frac{1}{\omega \hat{Q}_\xi} ; \text{ beta function in HCC} \\ (\text{beta is constant})$$

σ_ξ^2 ; rms fraction due to stochastic process

g_ξ ; partition of cooling

Exact formulae for emittance evolution in HCC are given in Slava & Rol's PRSTAB and IPAC'14

In case of HCC,

- These formulae are still good approximation
- ϵ_n has a strong coupling among three orthogonal components (horizontal, vertical & longitudinal)
- The coupling is given by the dispersion

$$g_L \rightarrow g_{L,0} + \delta g_L$$

$$g_T \rightarrow 1 - \frac{\delta g_L}{2}$$

In the HCC, δg_L is given (D. Neuffer)

$$\delta g_L = \frac{\kappa^2}{1+\kappa^2} \frac{\partial \rho}{\rho} \frac{p}{\partial p} = \frac{\kappa^2}{1+\kappa^2} \hat{D}$$

Equal cooling decrements

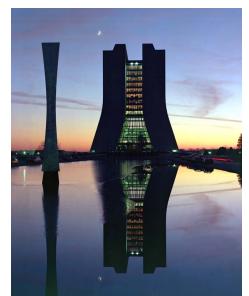
$$\rightarrow g_L = g_T$$

→ Need strong focusing (high field gradient)

In case of low field gradient...

→ $g_L > g_T$ Longitudinal enhance cooling





Equilibrium Emittance in various cooling decrements



Equal cooling decrements

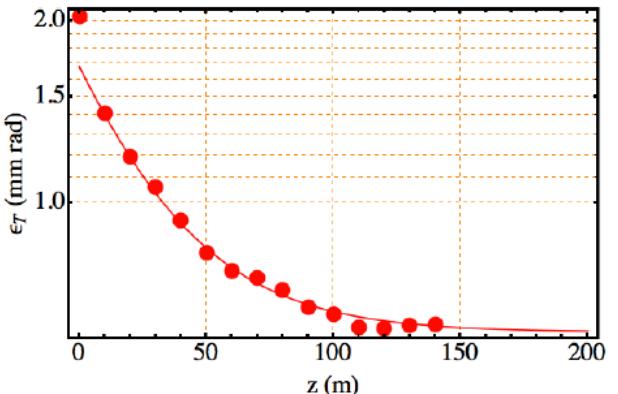
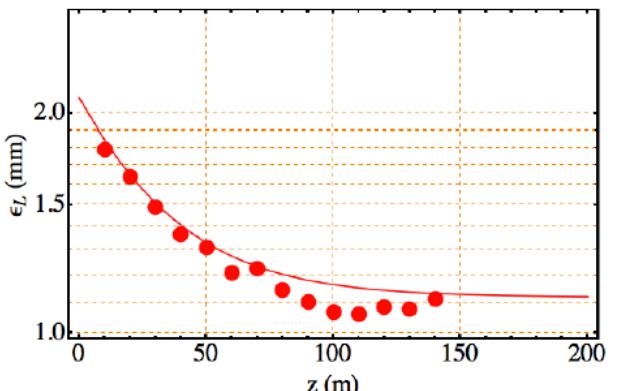
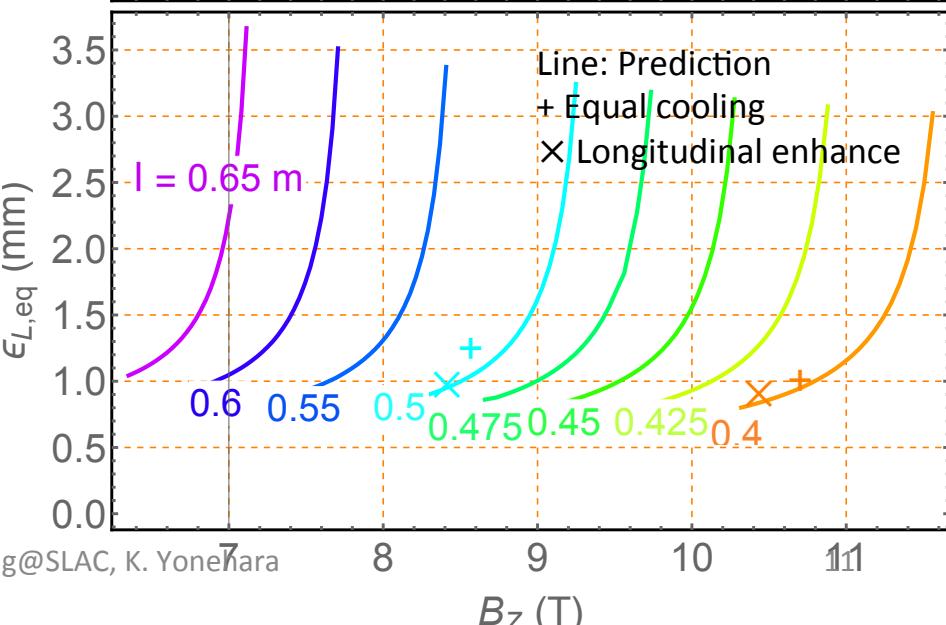
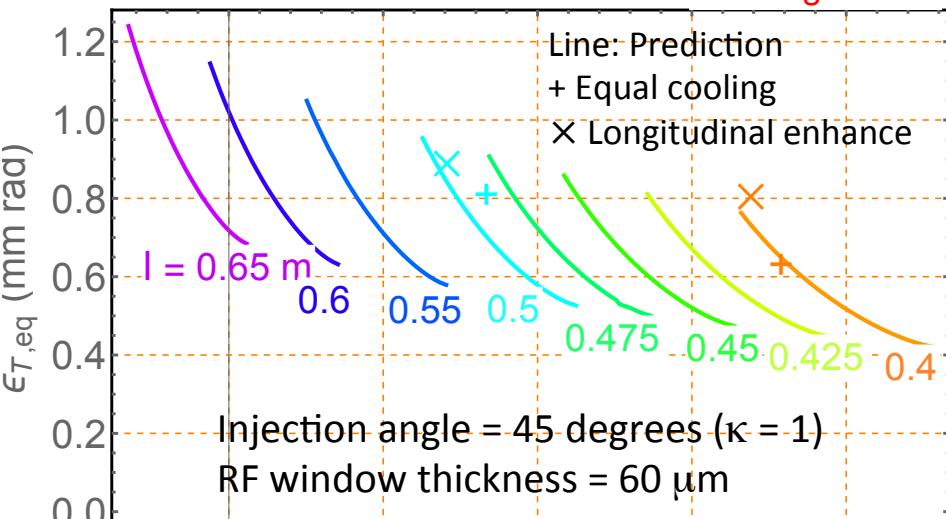


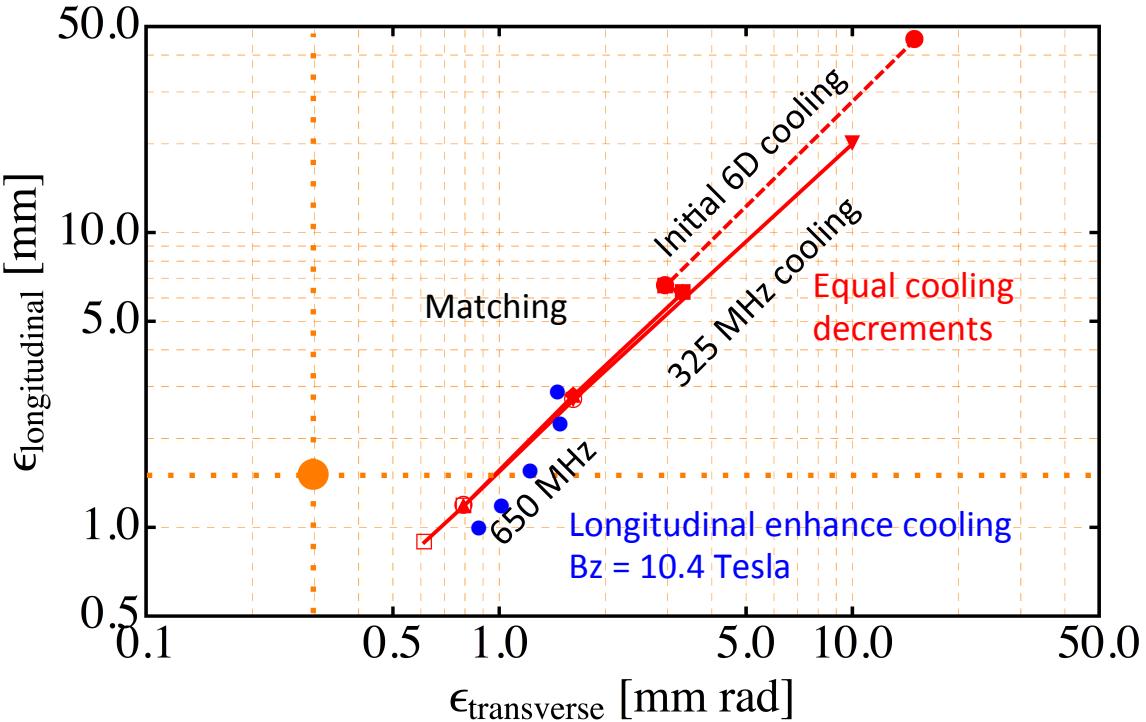
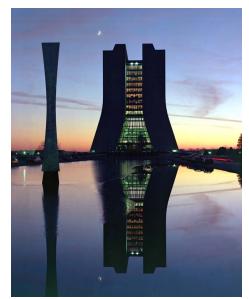
Figure 5: Estimated transverse emittance evolution. Red line is the predicted emittance evolution curve, $\epsilon_T(s) = (1.68 - 0.62)e^{-0.0307s} + 0.62$.



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Recent cooling result



Matching

Transmission: 56 % → **72 %**

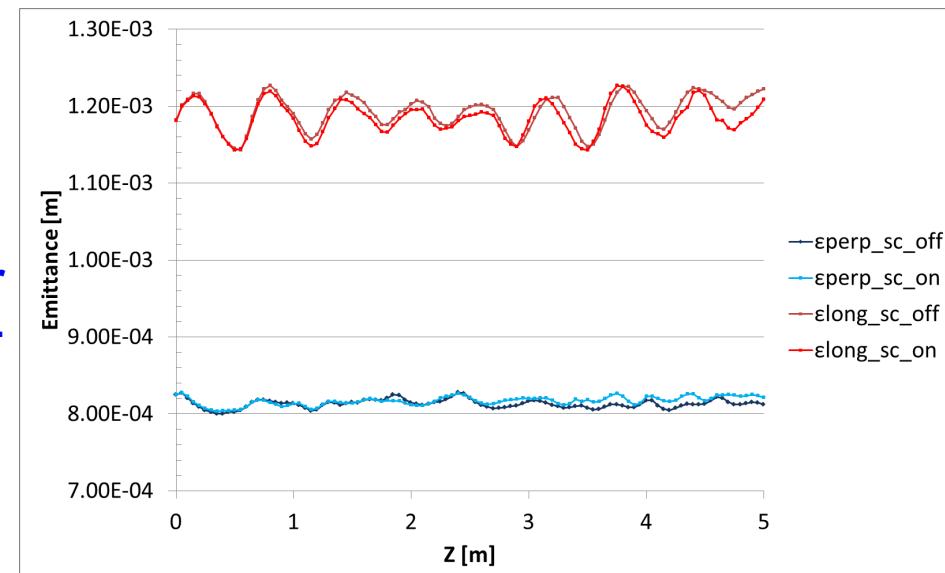
6D HCC

- RF parameter
 $E = 20 \text{ MV/m}$
 $v = 325 \text{ & } 650 \text{ MHz}$
- Gas pressure
 160 atm at 300 K
 43 atm at 80 K
- Magnetic fields
 $B_z = 4 - 12 \text{ Tesla}$
- Equilibrium emittance
 $\epsilon_T = 0.6 \text{ mm} \text{ (goal: } 0.3 \text{ mm)}$
 $\epsilon_L = 0.9 \text{ mm} \text{ (goal: } 1.5 \text{ mm)}$
- Transmission (no bunch merge)
 ~ 60 %
- Channel length
 $380 \text{ m} \rightarrow \text{280 m}$

Space Charge in the HCC

- High bunch charge ($\sim 10^{13}$ muons per bunch) after the bunch merge may affect emittance evolution within the HCC due to the space charge field from beam-induced plasma in the HPRF cavities
- G4Beamline space charge capabilities are currently limited and will require development to fully simulate effects of high-pressure hydrogen gas environment
- Current results:
 - Existing *spacecharge/w* module in G4Beamline indicates no effect on emittance evolution over short HCC sections
- In progress:
 - Extension of simulations to longer HCC channel sections
 - G4Beamline code development to better model space charge fields from a beam-induced plasma
 - Increasing bunch charge to find the bunch charge limit in the HCC

Emittance evolution in L=5m HCC section with G4Beamline *spacecharge/w* module on and off
Bunch charge $1e13 \mu^+$ per bunch



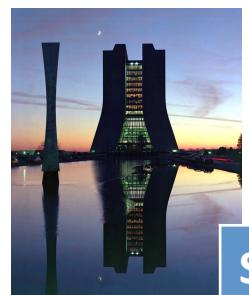


Power estimation

Section	Units	1	2	3	4
λ , period	m	1.0	0.8	0.5	0.4
β_T/β_L (equal)	m	0.20/2.1	0.16/2.1	0.098/1.5	0.079/1.5
Length	m	80	50	100	50
HS Coil R_{in}	mm	275	217	105	100
HS Coil R_{out}	mm	300	247	195	190
Coils/period		10	10	10	10
Coil Length	mm	100	80	50	40
J, curr density	A/mm ²	258.2	271.4	149.1	188.8
Anti solenoid	T	-2.49	-3.21	-8.78	-7.15
Stored B	MJ/m	4.84*	4.7	10.7	10.5
RF gradient	MV/m	20	20	20	20
RF frequency	MHz	325	325	650	650
Peak power	MJ/m	11 (36)**	11 (36)	13 (50)	13 (50)

*HS coil
+Anti Sol

**77 K
operation
(300 K)





Summary

- Progress on HCC machine design
 - Re-evaluation & Re-optimization based on **recent experimental results**
- Longitudinal enhance cooling is considerable option for later 6D cooling stage
 - Significantly reduce the geometry constraint
 - HCC does not have any limit due to longitudinal space charge
- We should evaluate hydrogen safety if the gas filled channel is the solution