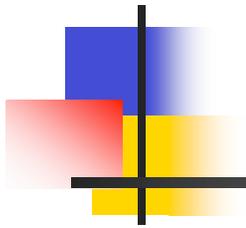


Cosmology For Particle Physicists



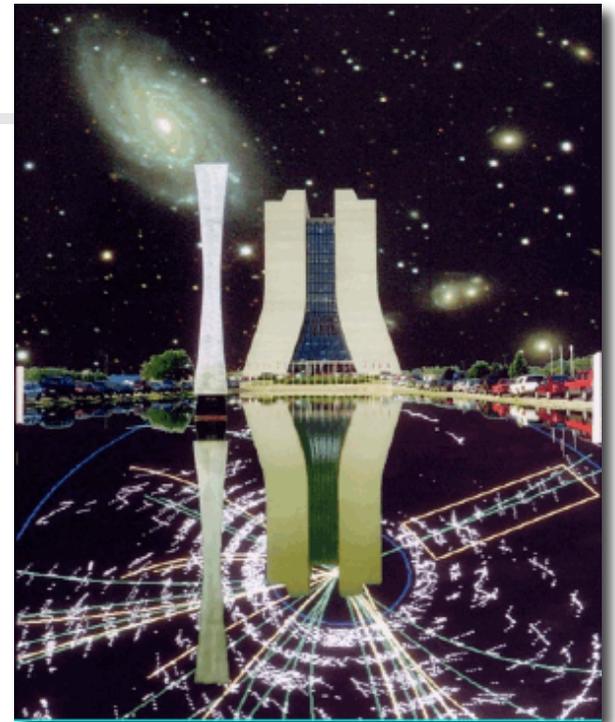
Dan Hooper

Fermilab/University of Chicago

Hadron Collider Physics Summer School

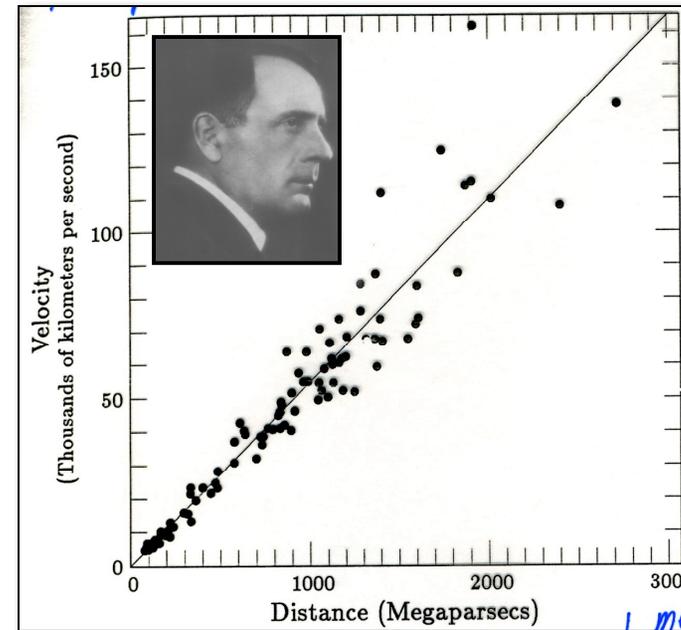
Fermilab

August 23, 2010



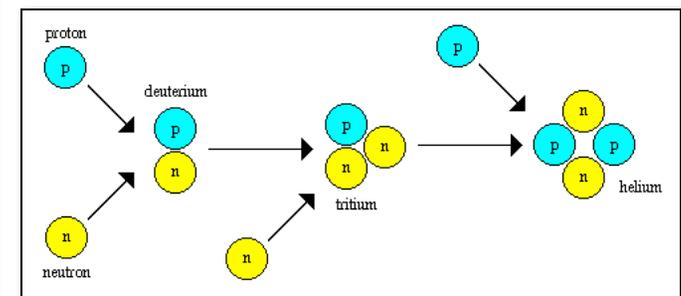
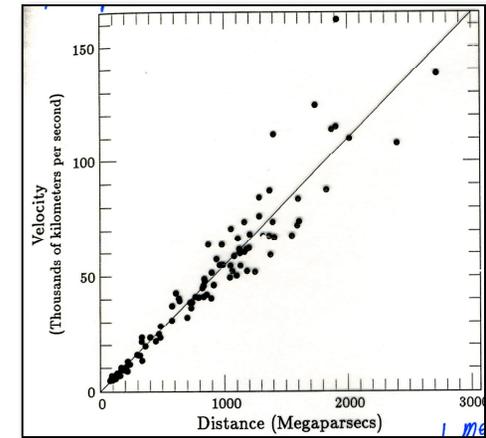
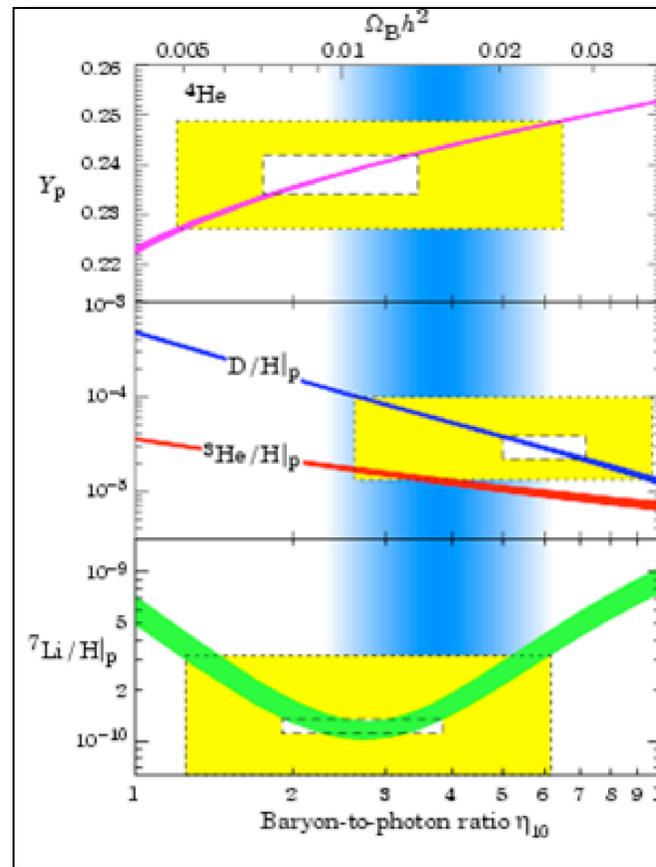
The Pillars Of Physical Cosmology

- Hubble Expansion



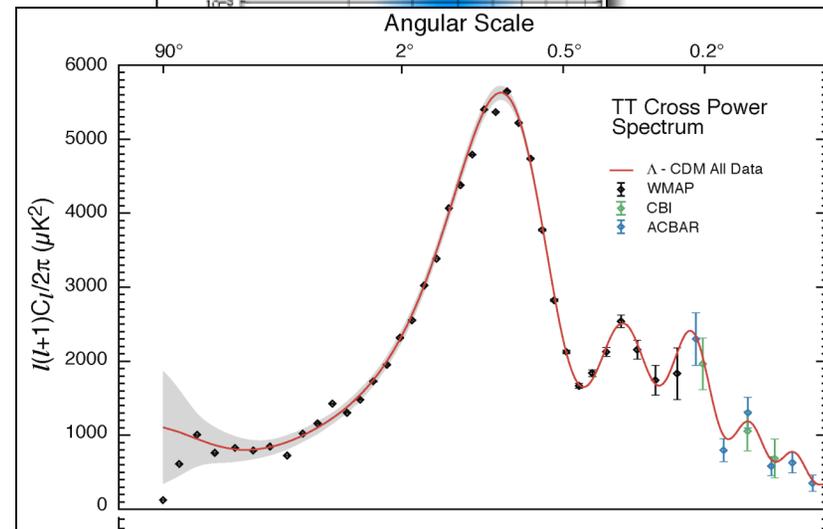
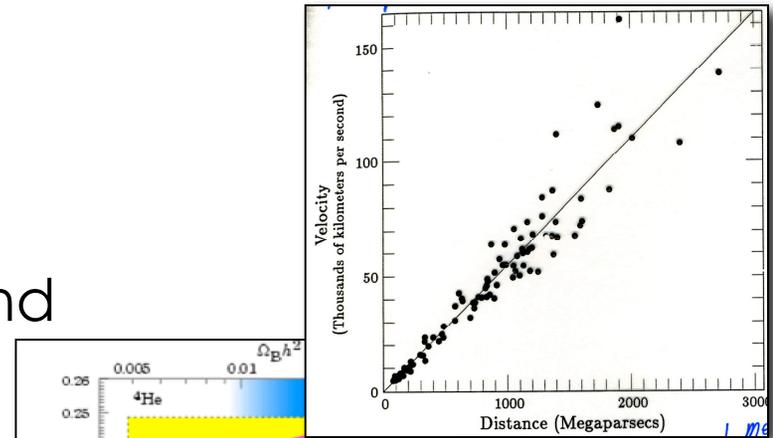
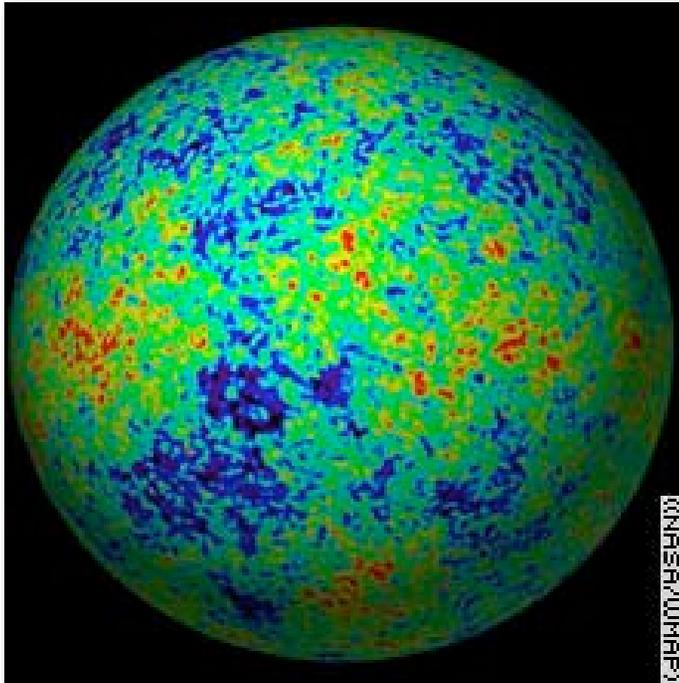
The Pillars Of Physical Cosmology

- Hubble Expansion
- The Light Element Abundances



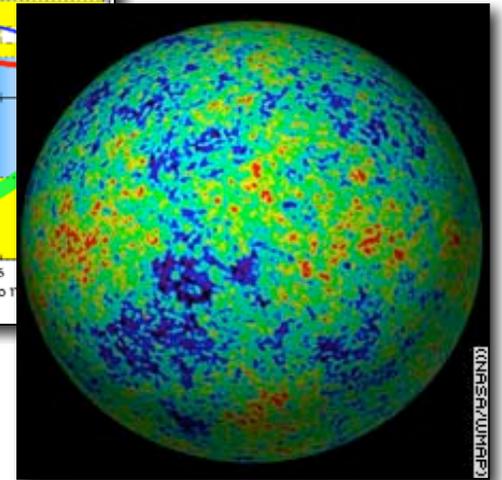
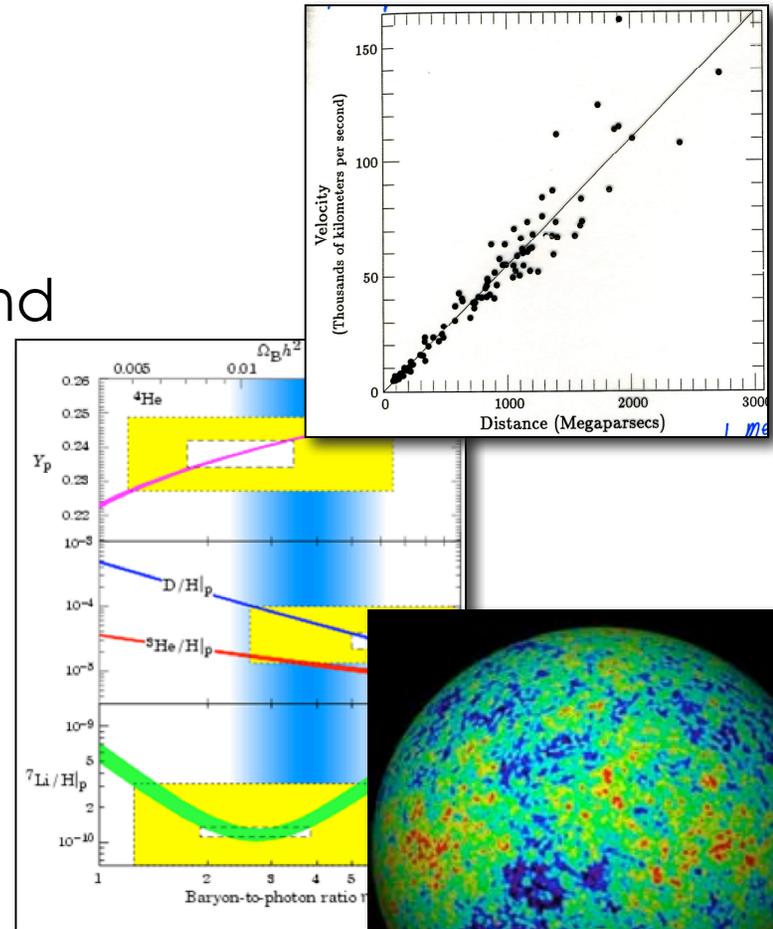
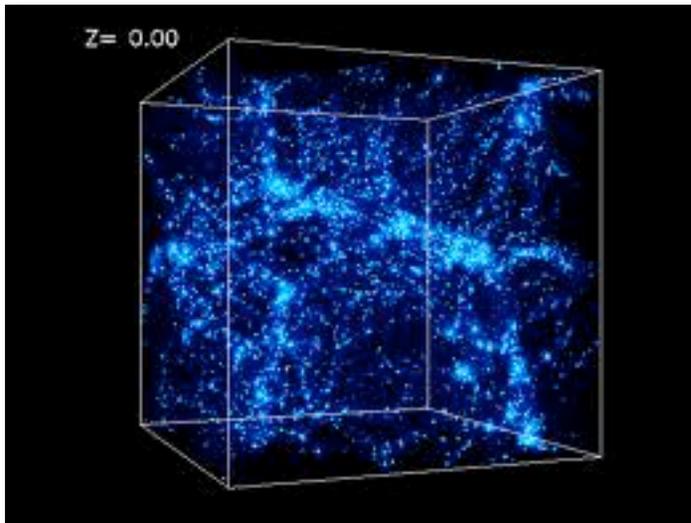
The Pillars Of Physical Cosmology

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- The Cosmic Microwave Background



The Pillars Of Physical Cosmology

- Hubble Expansion
- The Light Element Abundances
- The Cosmic Microwave Background
- Large Scale Structure

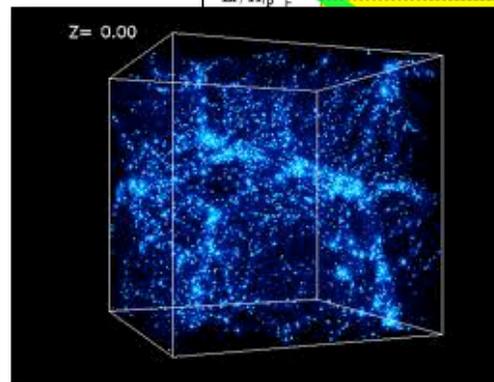
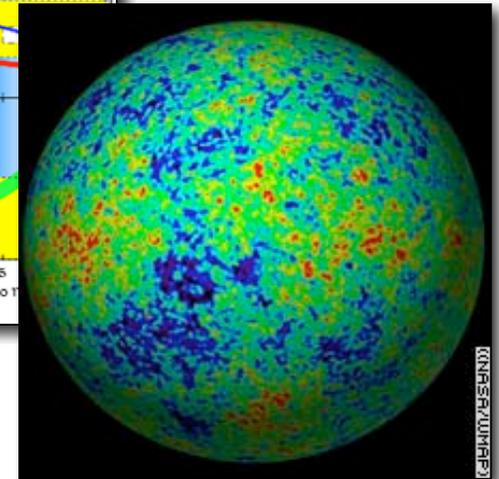
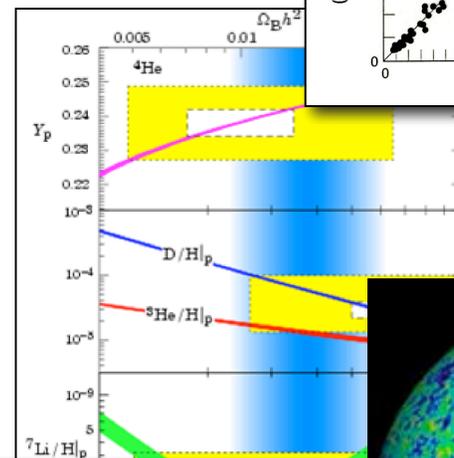
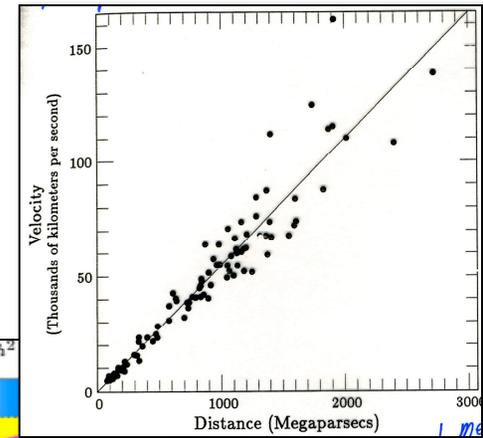


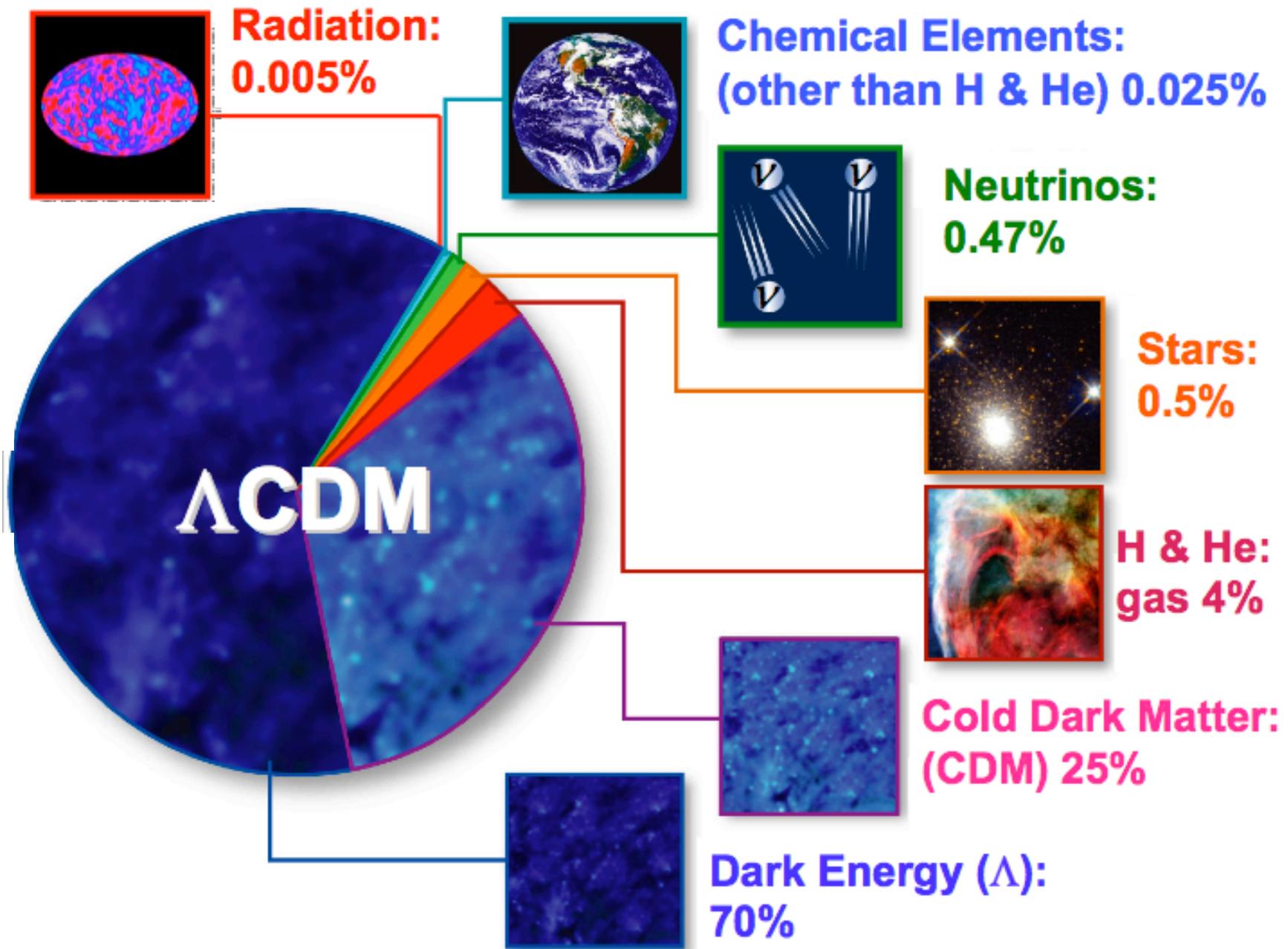
The Pillars Of Physical Cosmology

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The Big Bang Theory, although remarkably successful, requires us to introduce:

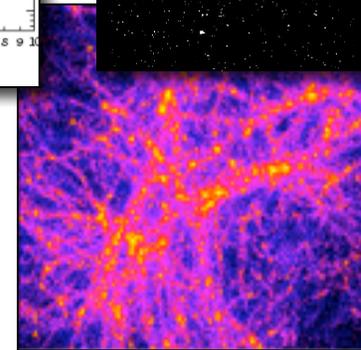
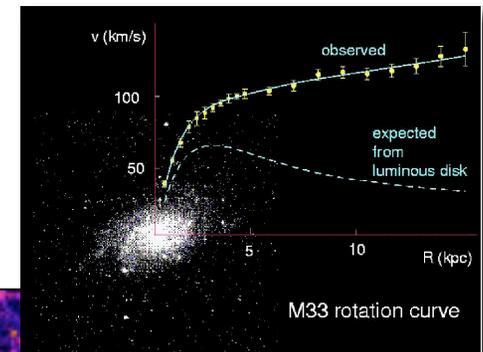
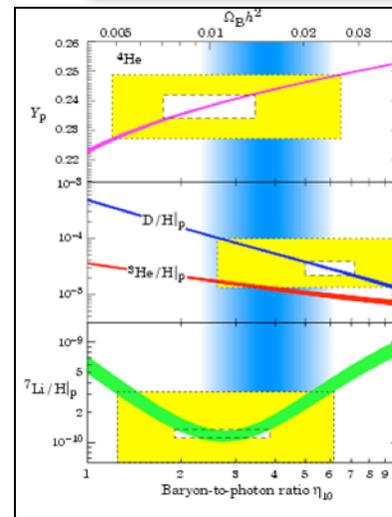
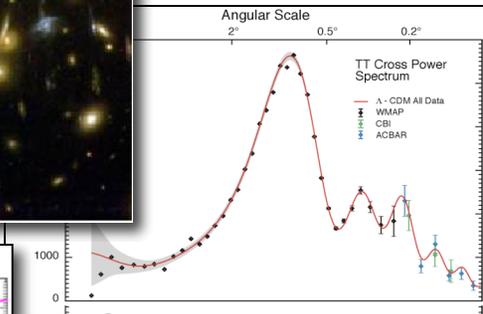
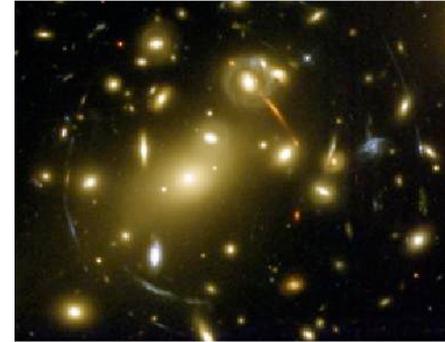
- 1) something that drives accelerated expansion in the early universe (inflation)
- 2) dark matter
- 3) something that drives accelerated expansion in the current era (dark energy)





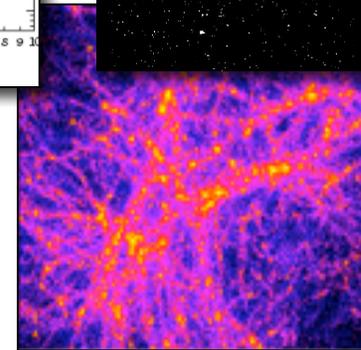
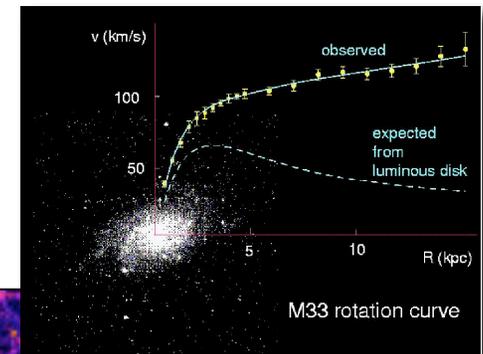
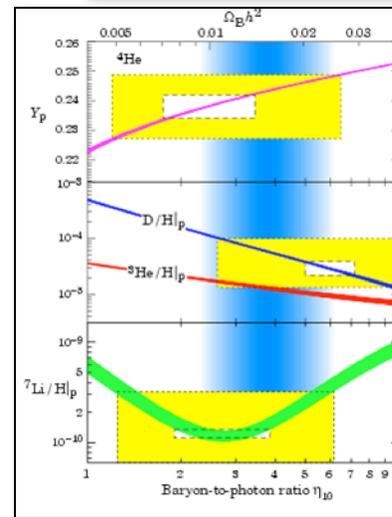
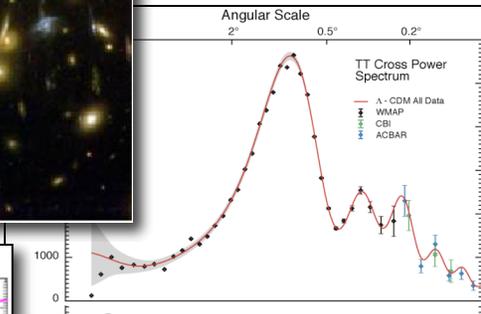
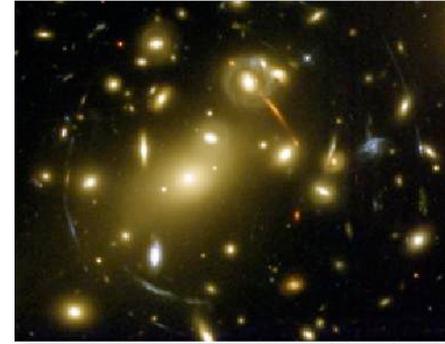
Evidence For Dark Matter

- Galactic rotation curves
- Gravitational lensing
- Light element abundances
- Cosmic microwave background anisotropies
- Large scale structure



Evidence For Dark Matter

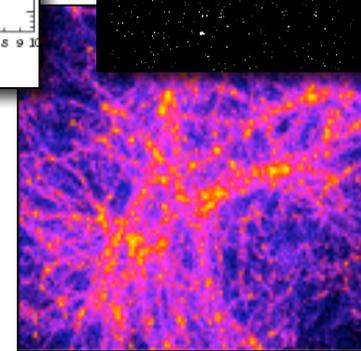
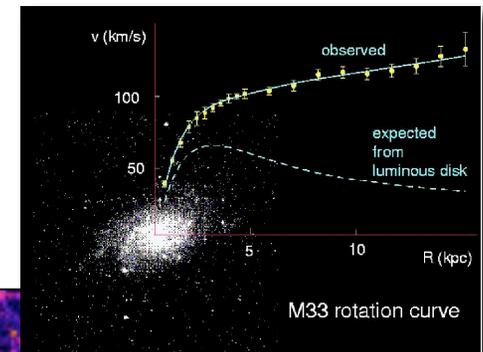
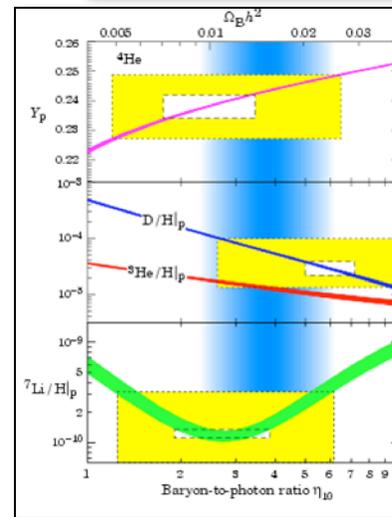
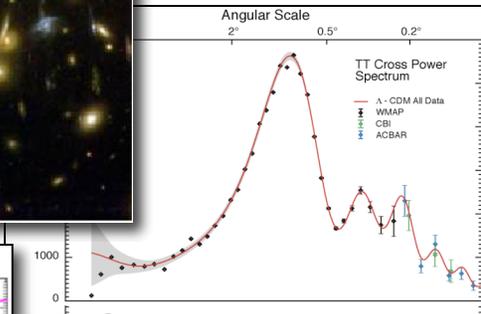
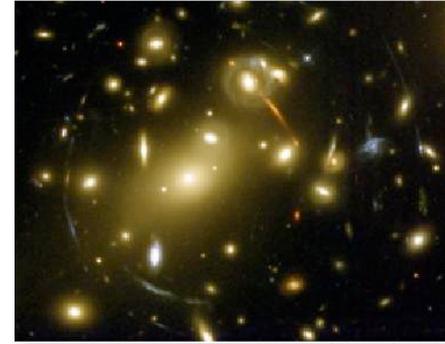
- There exists a wide variety of independent indications that dark matter exists
- Each of these observations infer dark matter's presence through its gravitational influence
- Still no observations of dark matter's electroweak interactions (or other non-gravitational interactions)



Evidence For Dark Matter

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Instead of dark matter, might we not understand gravity?



Modified Newtonian Dynamics (MOND)

- Begin by modifying Newtonian dynamics as follows:

$$F = ma \longrightarrow F = ma \times \mu(a)$$

where $\mu \approx 1$, except at small accelerations, at which $\mu = a/a_0$

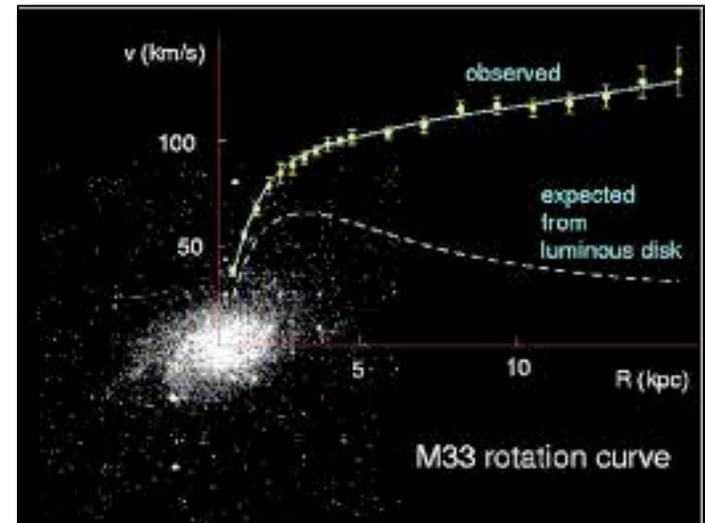
- For a circular orbit,

$$F = \frac{GMm}{r^2} = ma\mu.$$

Which in the low-acceleration limit yields:

$$a = \frac{\sqrt{GMa_0}}{r} = \frac{v^2}{r} \implies v = (GMa_0)^{1/4}$$

Rotational velocity independent of galactic radius (flat rotation curve)



Modified Newtonian Dynamics (MOND)

- MOND has been quite successful in explaining galactic dynamics of galaxies, and provides an explanation for the Tully-Fisher relationship
- Galaxy clusters have been less well described by MOND
- MOND cannot be applied to questions of cosmology
(toy theory – modifies Newton, not Einstein)

MONDs and Dark Matters Throughout History

- Case 1: Astronomical tables of Uranus' orbit deviated from observations. It was suggested that either another planet was perturbing Uranus' orbit (Dark Matter) or that Newton's laws had broken down (MOND).

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-One Point For Dark Matter

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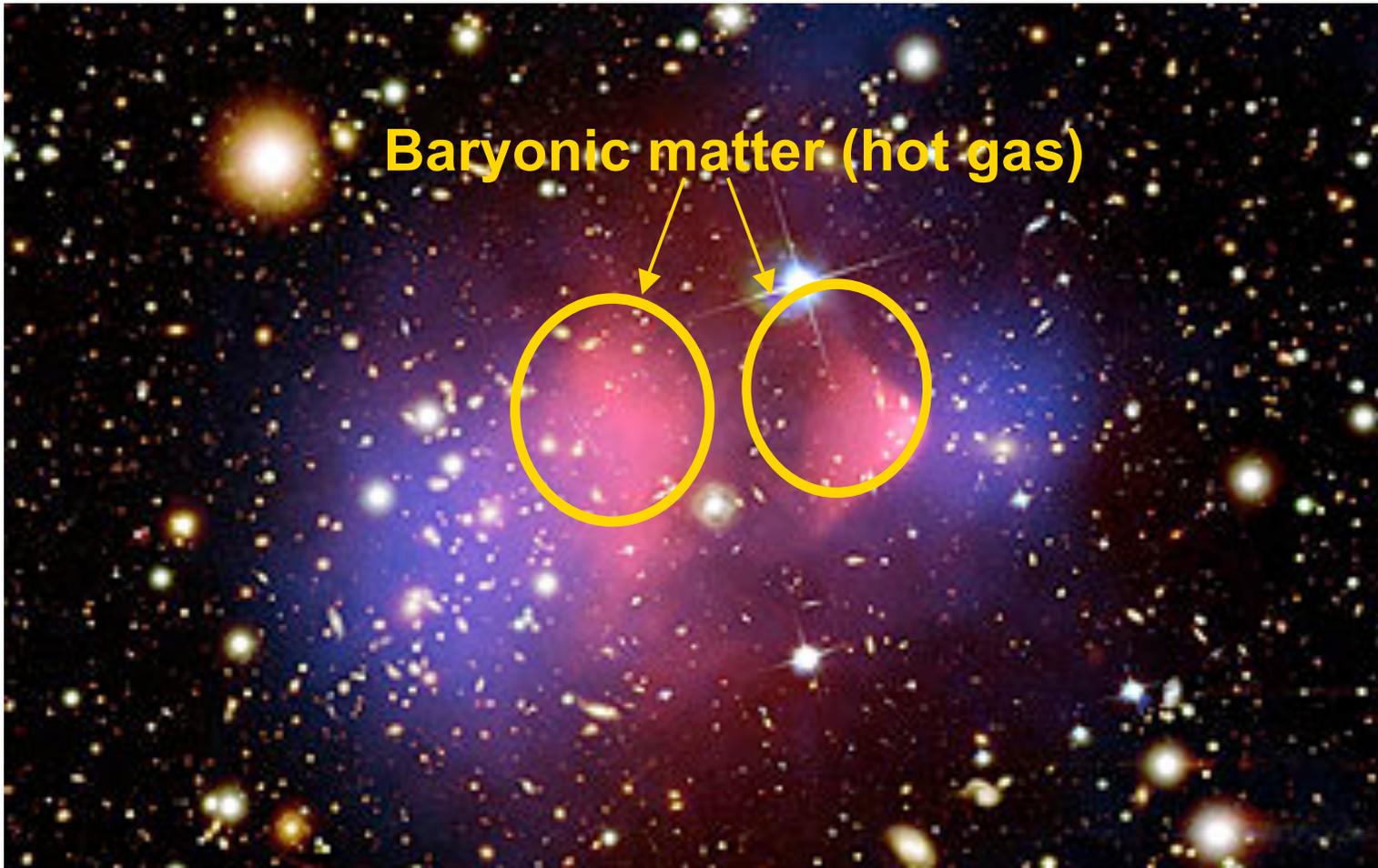
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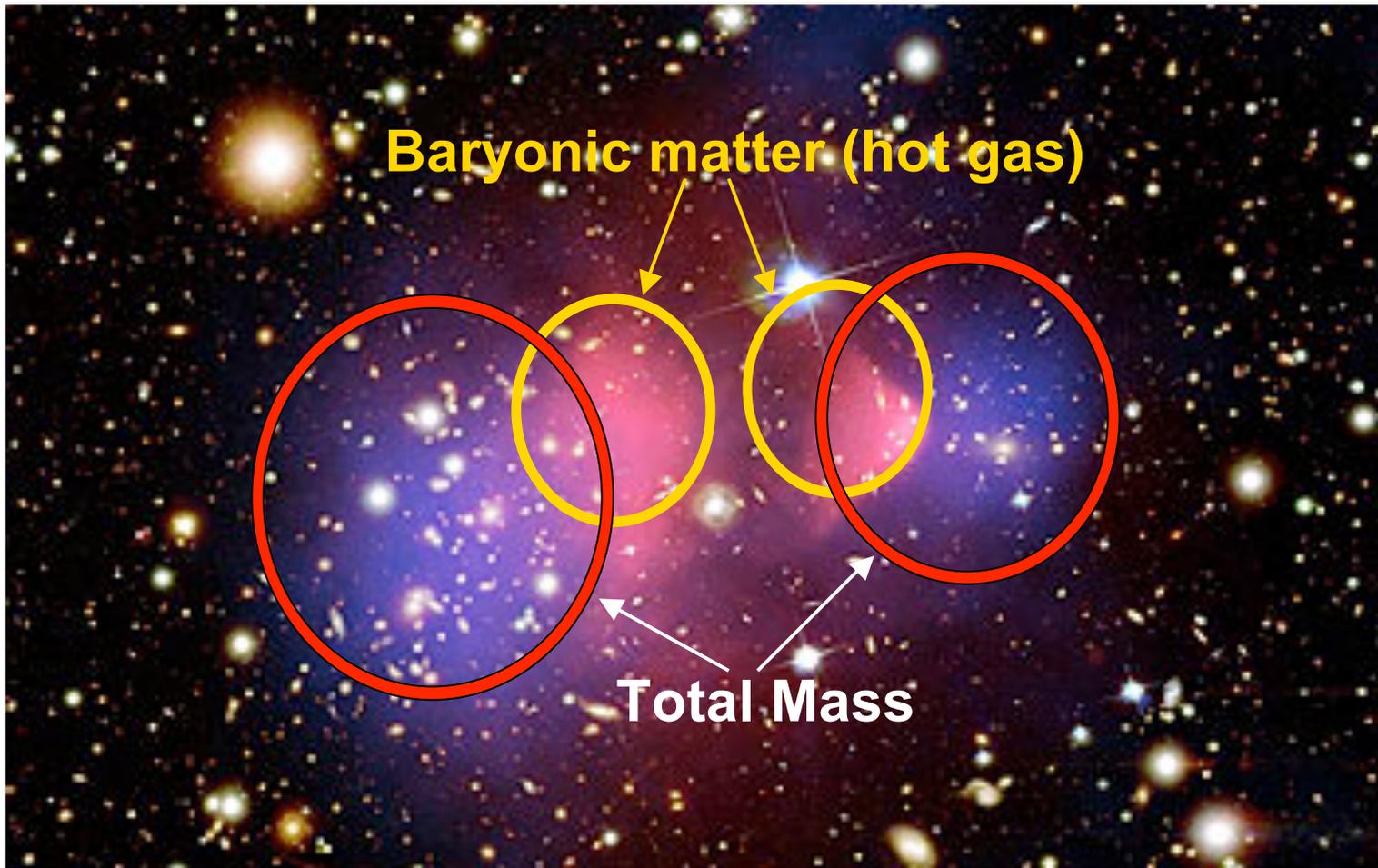
-One Point For MOND



**NASA/Chandra Press Release,
August 21, 2006**



**NASA/Chandra Press Release,
August 21, 2006**



**NASA/Chandra Press Release,
August 21, 2006**

MOND Takes Another Bullet



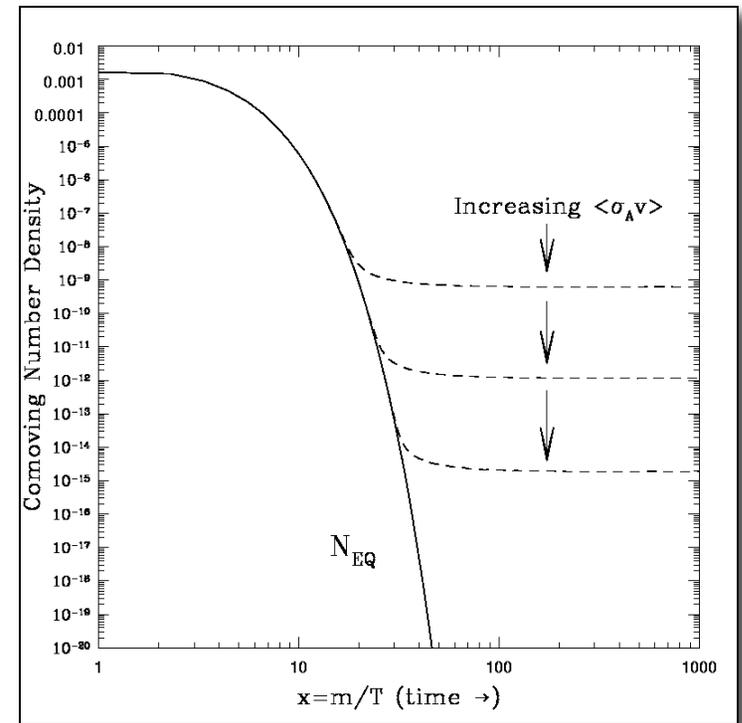
MACS J0025.4-1222,
R. Massey et al, August 27, 2008

Why WIMPs?

The thermal abundance of a WIMP

- $T \gg M_X$, WIMPs in thermal equilibrium
- $T < M_X$, number density becomes Boltzmann suppressed

$$n_{X,eq} = g_X \left(\frac{m_X T}{2\pi} \right)^{3/2} e^{-m_X/T}$$



Why WIMPs?

The thermal abundance of a WIMP

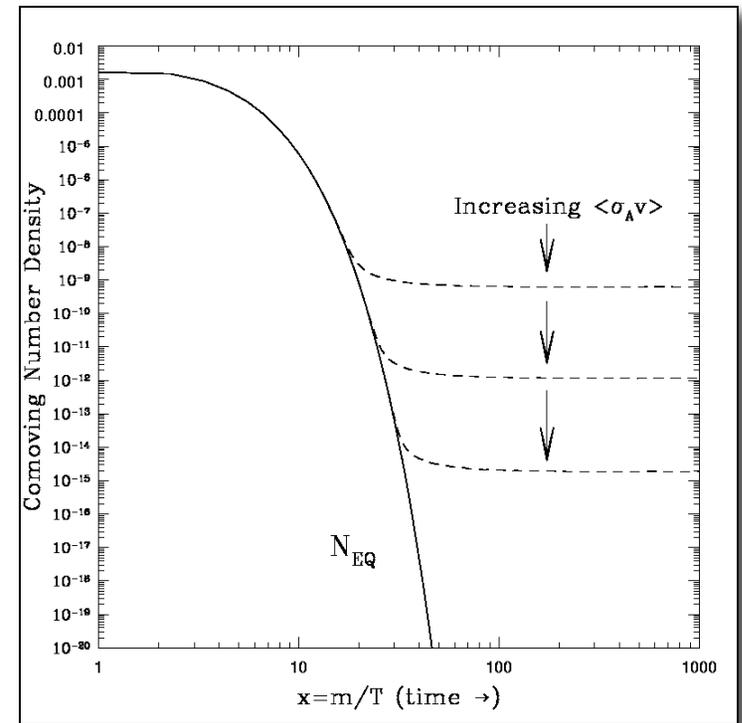
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- Including the effects of expansion pulls the density away from its equilibrium value:

$$\frac{dn_X}{dt} + 3Hn_X = - \langle \sigma_{X\bar{X}} |v| \rangle (n_X^2 - n_{X,eq}^2)$$

\swarrow
 $H \equiv \dot{R}/R = (8\pi^3 \rho / 3M_{Pl})^{1/2}$



Why WIMPs?

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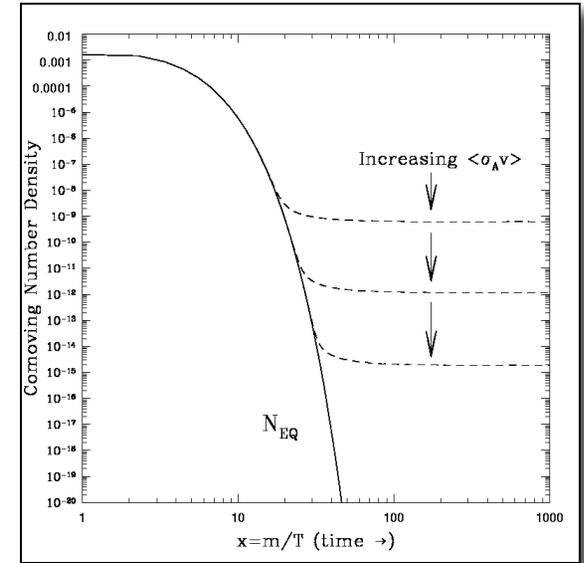
- The departure from equilibrium occurs at a temperature given by:

$$x_{FO} \equiv \frac{m_X}{T_{FO}} \approx \ln \left[c(c+2) \sqrt{\frac{45}{8} \frac{g_X}{2\pi^3} \frac{m_X M_{Pl} (a + 6b/x_{FO})}{g_*^{1/2} x_{FO}^{1/2}}} \right]$$

order one factor

of external degrees of freedom

$$\langle \sigma_{X\bar{X}} |v| \rangle = a + b \langle v^2 \rangle + \mathcal{O}(v^4)$$



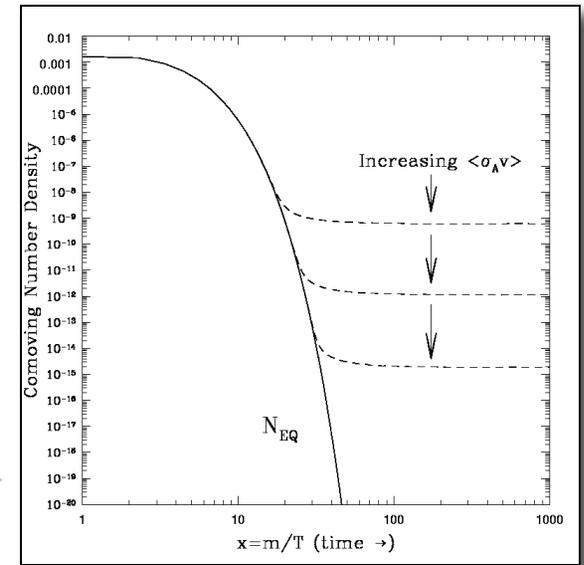
- For a weak-scale mass and cross section, $T_{FO} \sim M_X/20$

Why WIMPs?

The thermal abundance of a WIMP

- Numerically, this yields a thermal relic abundance given by:

$$\Omega_X h^2 \approx 0.1 \left(\frac{x_{\text{FO}}}{20} \right) \left(\frac{g_\star}{80} \right)^{-1/2} \left(\frac{a + 3b/x_{\text{FO}}}{3 \times 10^{-26} \text{ cm}^3/\text{s}} \right)^{-1}$$



- For comparison,

$$\alpha^2 / (100 \text{ GeV})^2 \sim 3 \times 10^{-26} \text{ cm}^3/\text{s}.$$

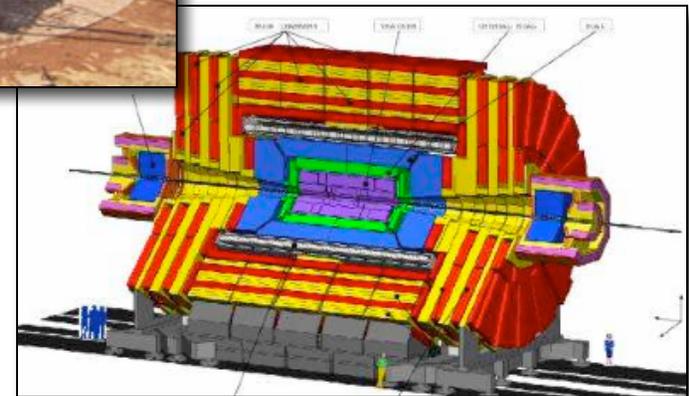
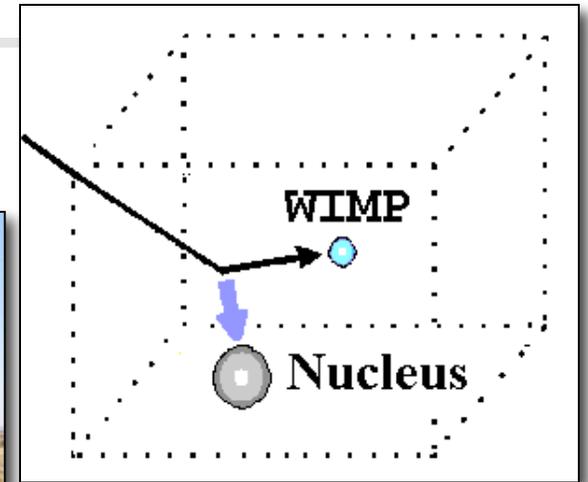
⇒ Weak scale interaction yields approximately the Observed dark matter abundance

$$\langle \sigma_{X\bar{X}} |v| \rangle = a + b \langle v^2 \rangle + \mathcal{O}(v^4)$$

Approximately ~1 pb

WIMP Hunting

- Direct Detection
- Indirect Detection
- Collider Searches

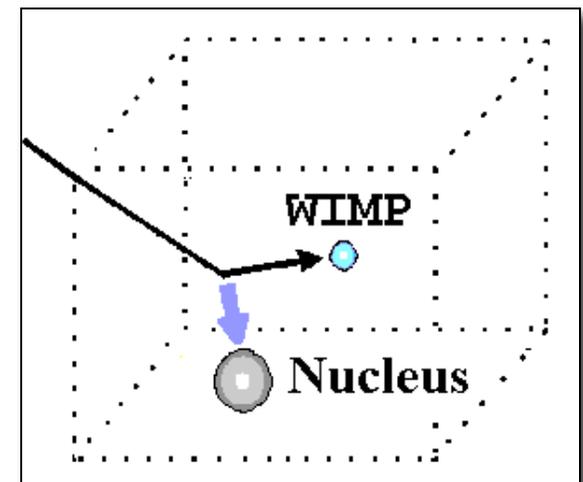


Direct Detection

- A WIMP striking a nucleus will impart a recoil of energy:

$$E_{\text{recoil}} = \frac{|\vec{q}|^2}{2M_{\text{nucleus}}} = \frac{2\mu^2 v^2 (1 - \cos \theta)}{2M_{\text{nucleus}}} = \frac{m_X^2 M_{\text{nucleus}} v^2 (1 - \cos \theta)}{(m_X + M_{\text{nucleus}})^2}$$

- For $m_X \gg M_{\text{nucleus}}$ and $v \sim 300$ km/s, this corresponds to $E_{\text{recoil}} \sim 1-100$ keV



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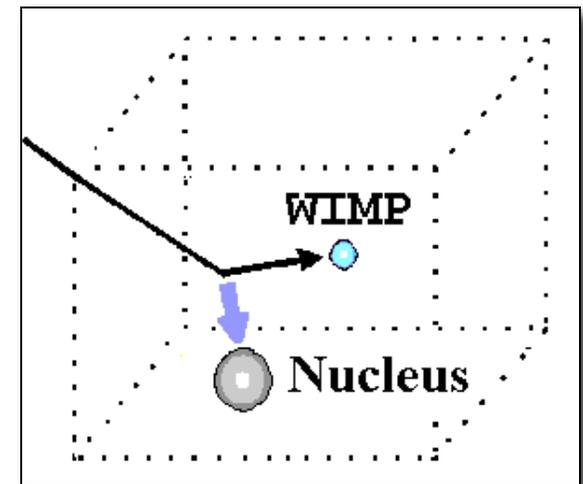
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- The rate of WIMP-nucleus scattering is given by:

$$R \approx \int_{E_{\text{min}}}^{E_{\text{max}}} \int_{v_{\text{min}}}^{v_{\text{max}}} \frac{2\rho}{m_X} \frac{d\sigma}{d|\vec{q}|} v f(v) dv dE_{\text{recoil}}$$

DM velocity distribution



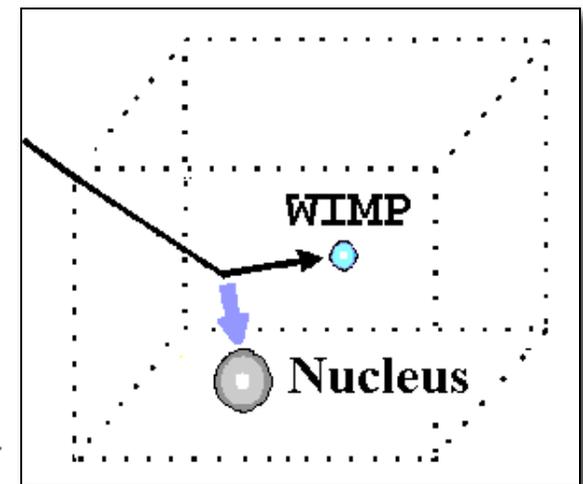
Direct Detection

- In the non-relativistic limit, all interaction forms can be described as spin-independent (coherent) scattering or spin-dependent scattering
- For spin-independent scattering:

$$\sigma \approx \frac{4m_{\chi^0}^2 m_{\text{nucleus}}^2}{\pi(m_{\chi^0} + m_{\text{nucleus}})^2} [Z f_p + (A - Z) f_n]^2$$

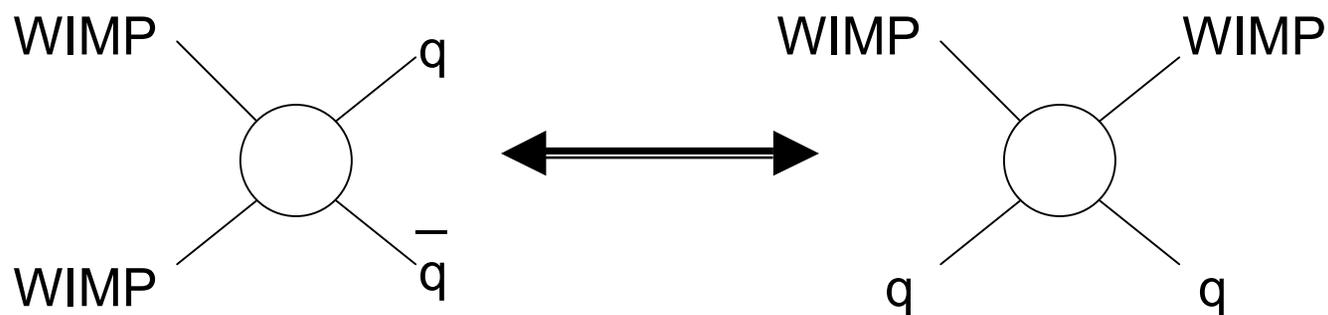
where

$$f_{p,n} = \sum_{q=u,d,s} f_{T_q}^{(p,n)} a_q \frac{m_{p,n}}{m_q} + \frac{2}{27} f_{TG}^{(p,n)} \sum_{q=c,b,t} a_q \frac{m_{p,n}}{m_q}$$



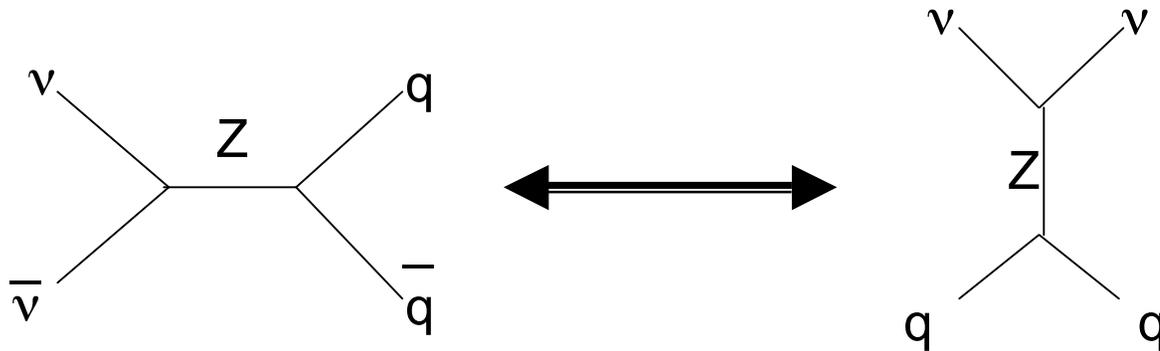
WIMP Annihilation and Elastic Scattering

- From the observed density of dark matter, we can estimate the couplings of the WIMP (in lieu of resonances, coannihilations, etc.)
- Annihilation rate to quarks can be used to estimate the elastic scattering cross section with nuclei



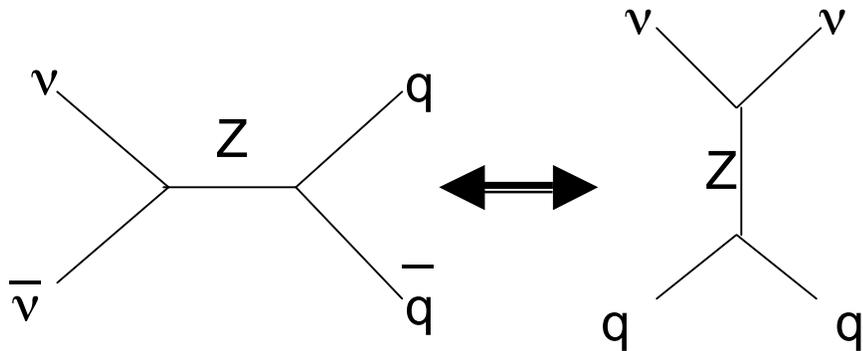
WIMP Annihilation and Elastic Scattering

- Case Example: Dirac Fermion or Scalar WIMP with vector-like interactions (such as a heavy 4th generation neutrino, or a sneutrino)



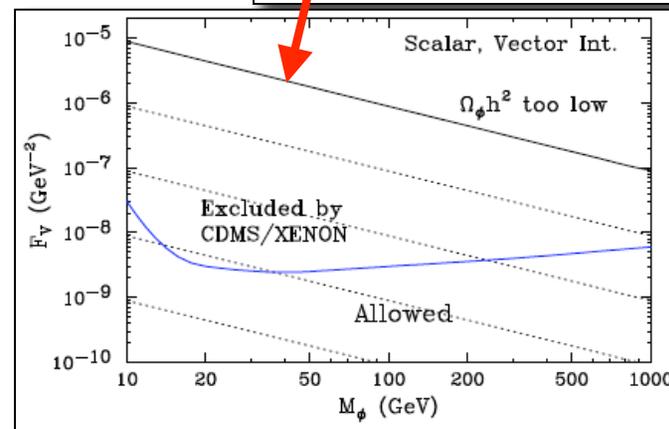
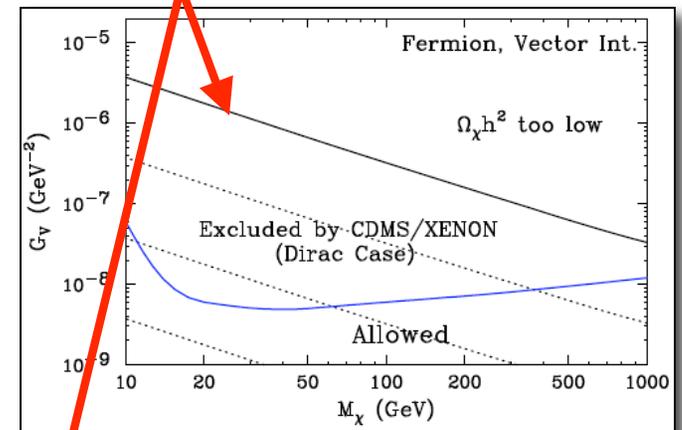
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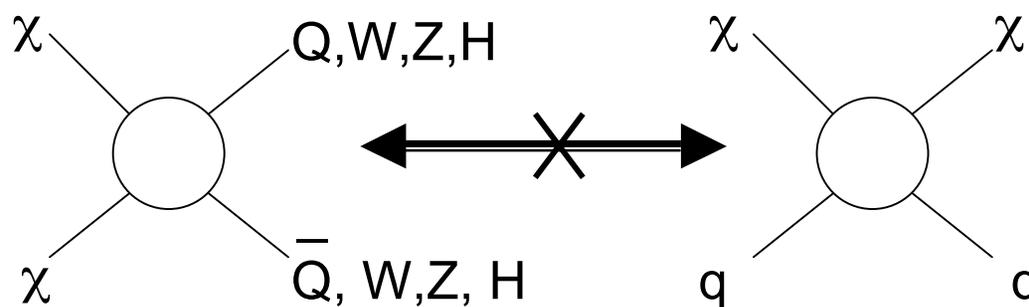
Coupling Needed To Acquire Measured Abundance

This candidate and many others are already excluded by the null results of CDMS and XENON!



WIMP Annihilation and Elastic Scattering

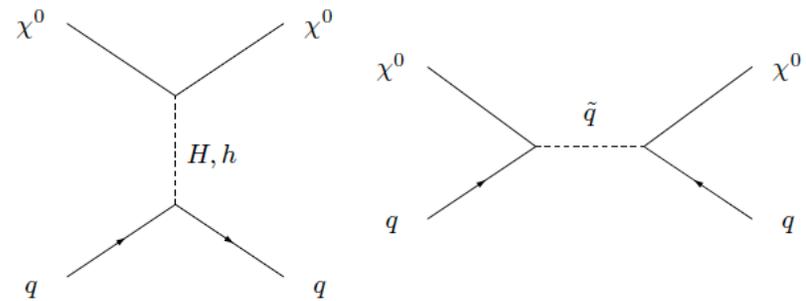
- Another Case Example: Majorana Fermion WIMP (such as a neutralino)
- Neutralino elastic scattering is suppressed by the fact that they annihilate to *heavy* quarks (and leptons) and gauge/higgs bosons, none of which are present in nuclei



- Elastic scattering is further suppressed if coannihilations or resonances play an important role in the early universe

The Direct Detection of Neutralino Dark Matter

- Neutralinos can coherently scatter with nuclei through scalar higgs and squark exchange



- Benchmark examples:

- 1) Dominated by relatively light “heavy” higgs:

$m_H \sim 200 \text{ GeV}, \mu \sim 200 \text{ GeV} \Rightarrow \sigma_{\chi p} \sim 10^{-5} \text{ to } 10^{-7} \text{ pb}$

$m_H \sim 200 \text{ GeV}, \mu \sim 1 \text{ TeV} \Rightarrow \sigma_{\chi p} \sim 10^{-7} \text{ to } 10^{-9} \text{ pb}$

- 2) Dominated by light (SM-like) higgs:

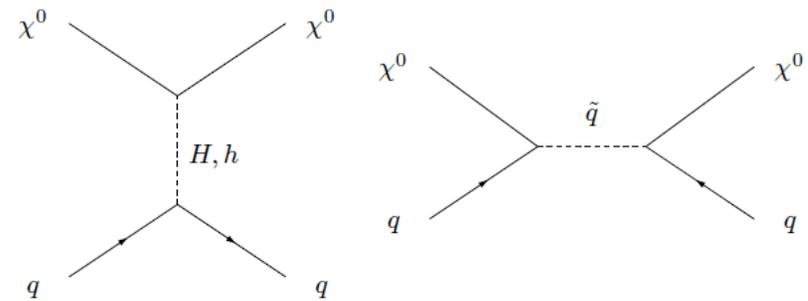
$m_H \gtrsim 500 \text{ GeV}, \mu \sim 200 \text{ GeV (1 TeV)} \Rightarrow \sigma_{\chi p} \sim 10^{-8} \text{ pb (} 10^{-10} \text{ pb)}$

- 3) Dominated by light (sub-TeV) squarks:

Potentially competitive with contributions from higgs exchange

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3) Dominated by light (sub-TeV) squarks:

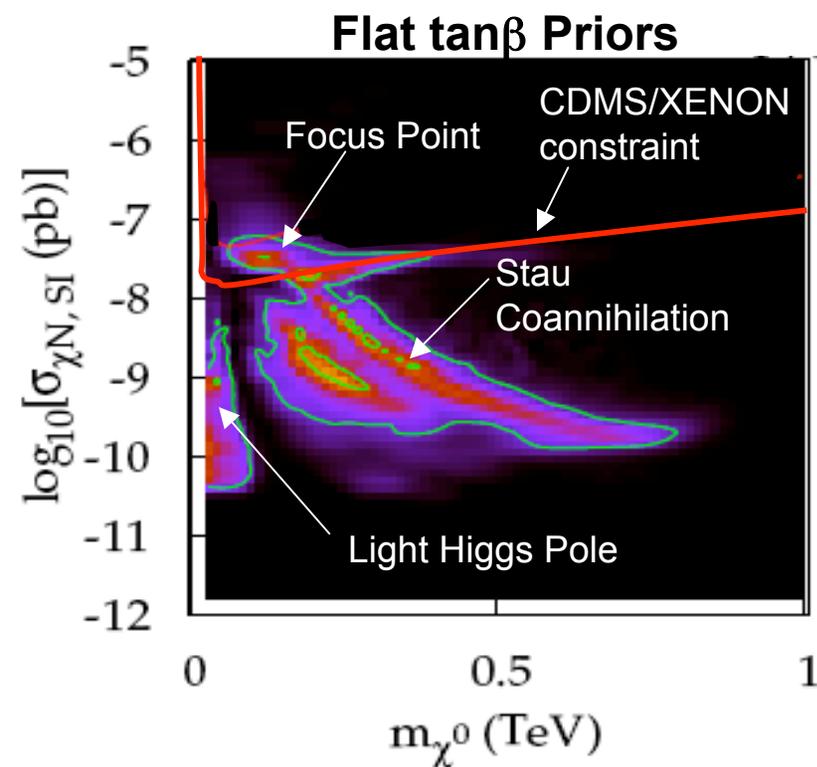
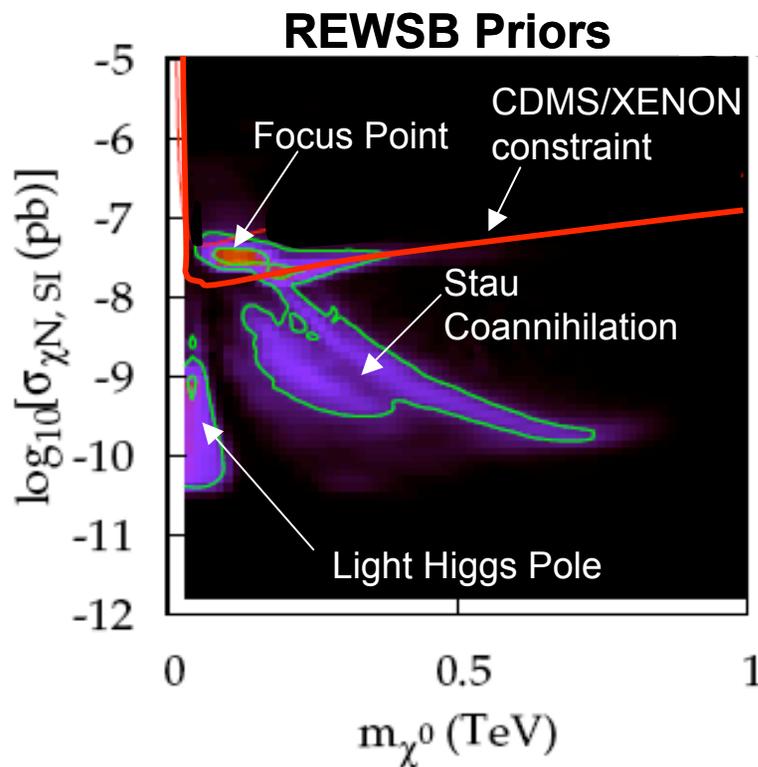
Potentially competitive with contributions from higgs exchange

~0.1-10 events per kg-day (Ge,Xe)

~0.1 events per ton-day (Ge,Xe)

The Case of Neutralino Dark Matter

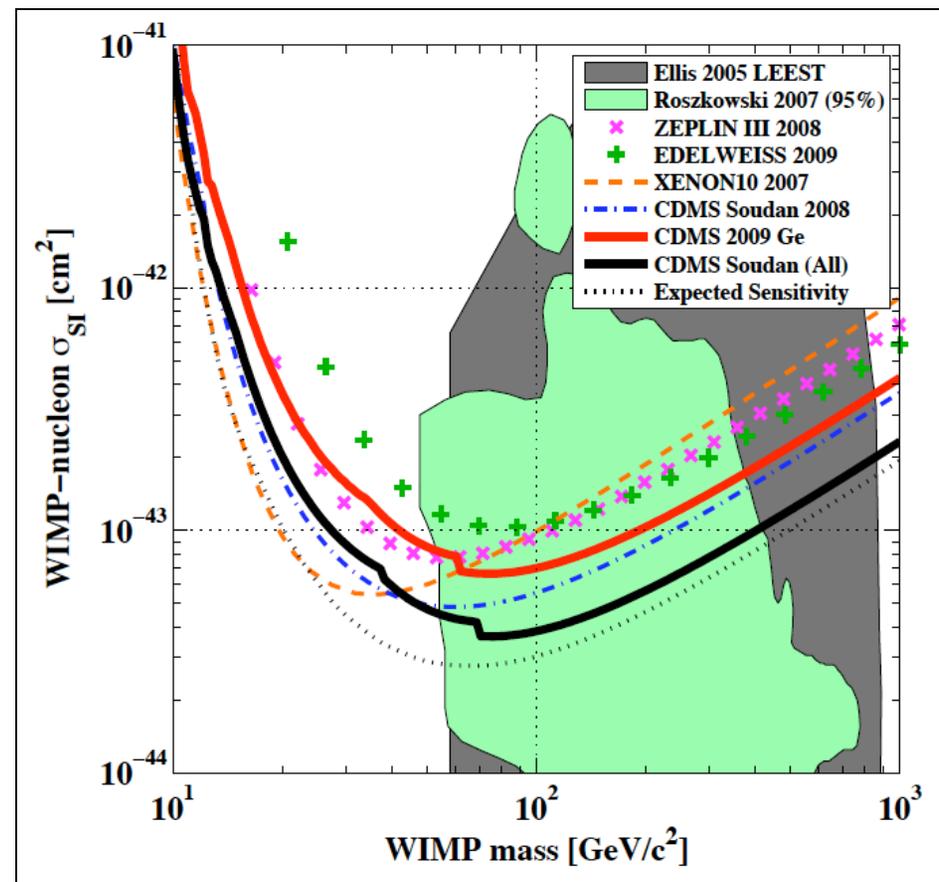
- Current direct detection experiments are only beginning to constrain a significant fraction of the likely SUSY parameter space



Direct Detection

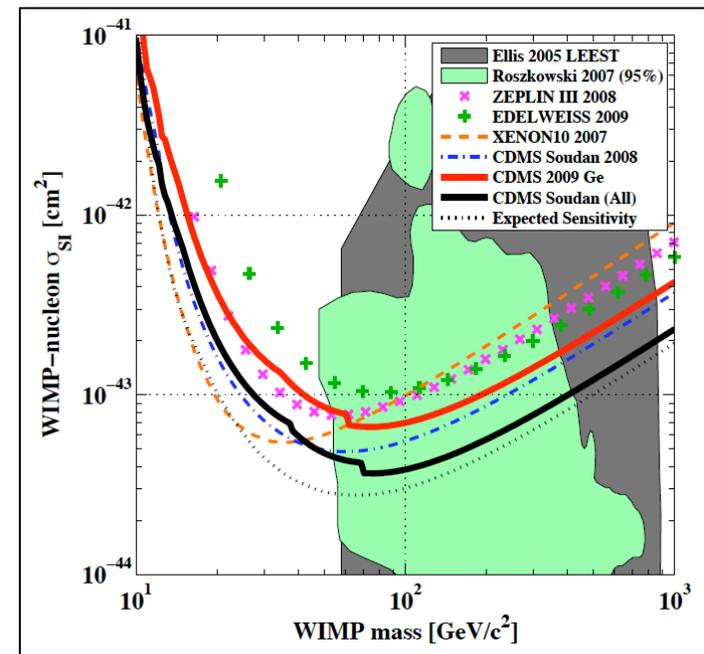
- Over the past decade, direct detection experiments have improved in sensitivity at a rate of about 1 order of magnitude every 2 years

Current Status



The Signals/Hints/Detections

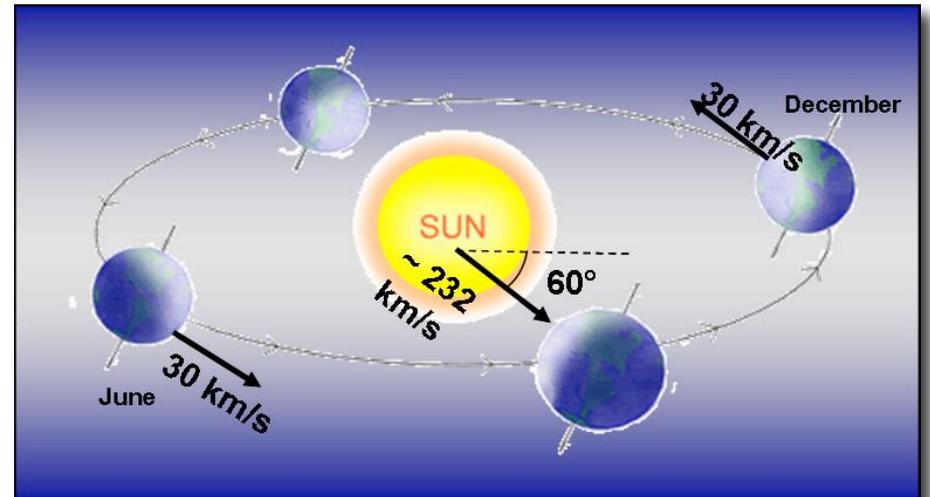
- But not all direct detection experiments have claimed null results
- For most of the past decade, DAMA has claimed to be observing dark matter scattering
- CoGeNT has recently reported a signal resembling dark matter
- CDMS recently reported 2 events



The Signals/Hints/Detections

DAMA/LIBRA

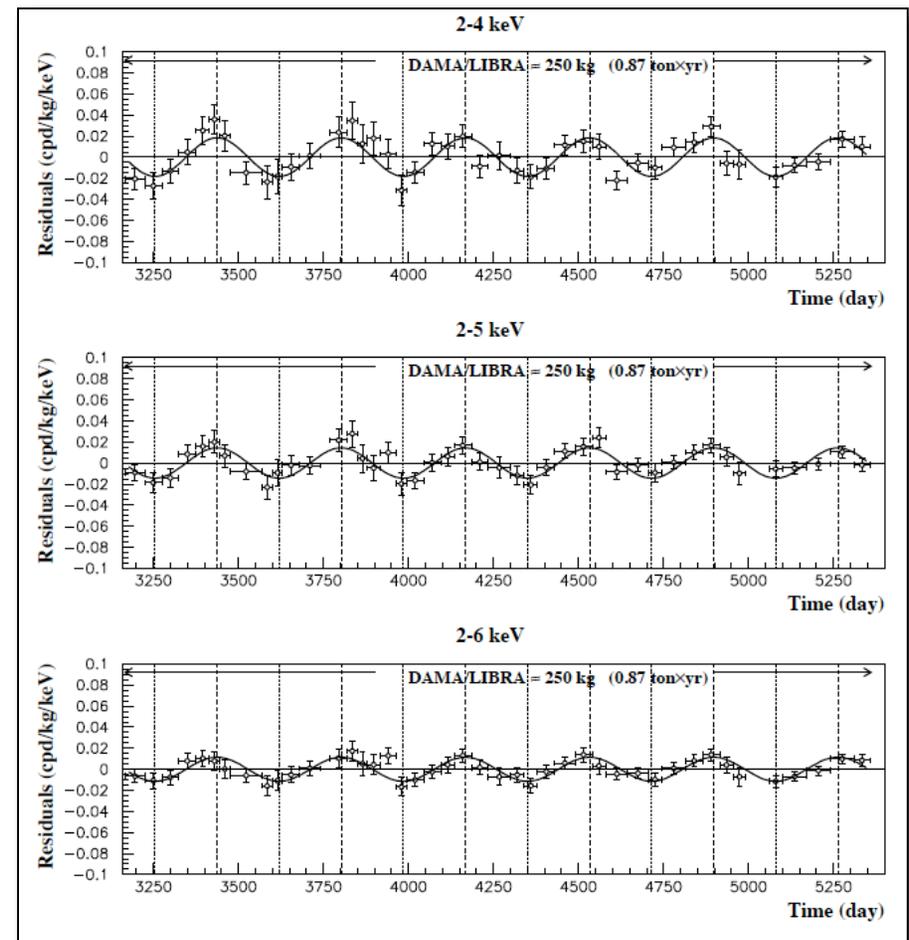
- Over the course of a year, the motion of the Earth around the Solar System is expected to induce a modulation in the dark matter scattering rate



The Signals/Hints/Detections

DAMA/LIBRA

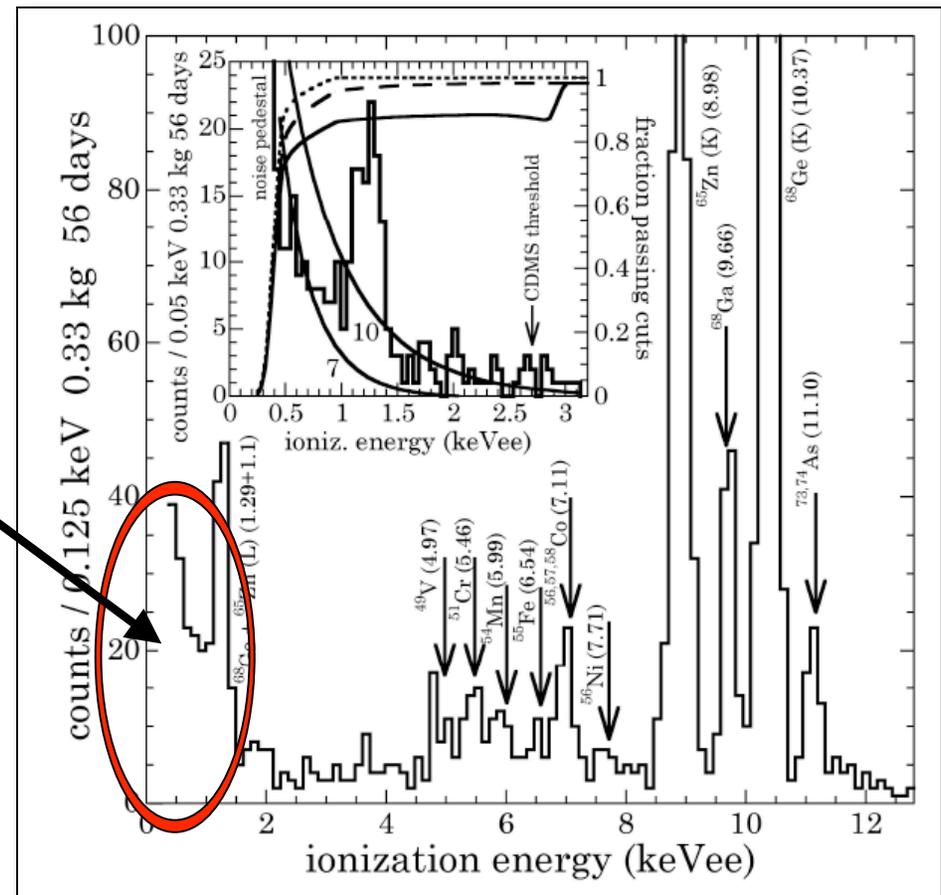
- Over the course of a year, the motion of the Earth around the Solar System is expected to induce a modulation in the dark matter scattering rate
- The DAMA collaboration reports a modulation with the right phase to be dark matter, and with high statistics (8.9σ)



The Signals/Hints/Detections

CoGeNT

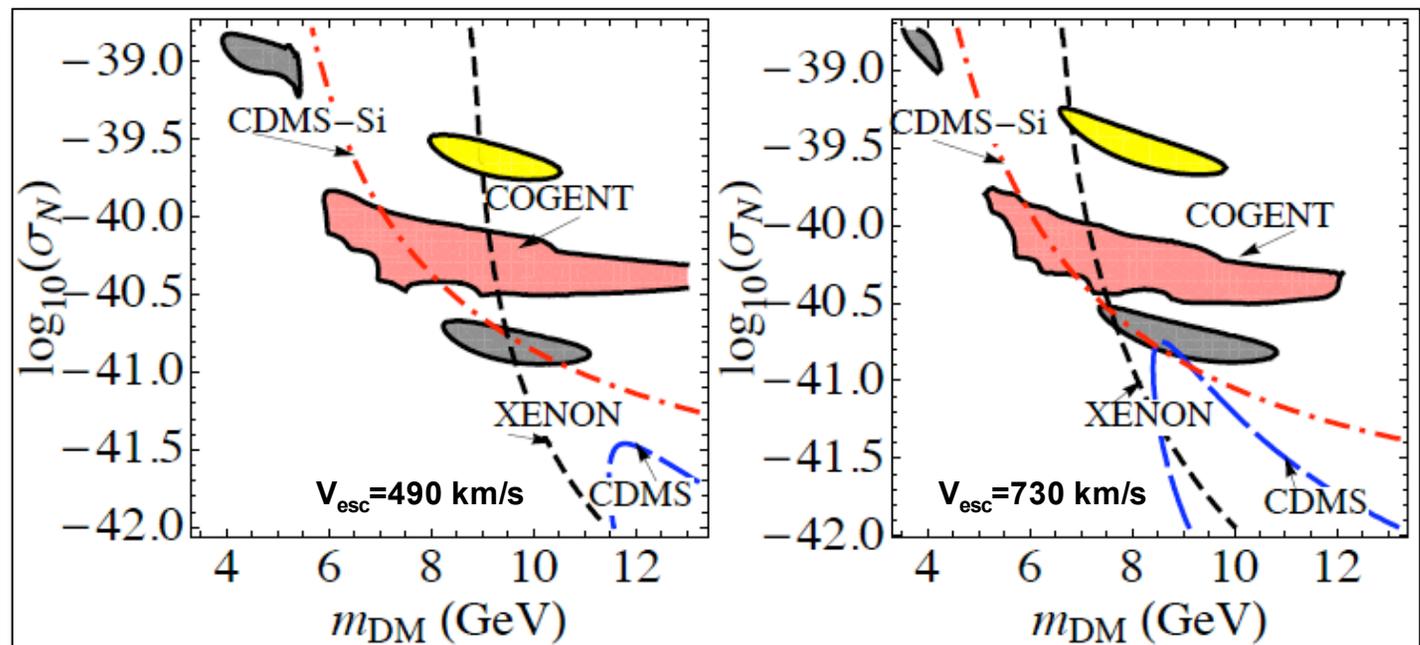
- The CoGeNT collaboration recently announced their observation of an excess of low energy events
- Although it has less exposure than other direct detection experiments, CoGeNT is particularly well suited to look for low energy events (low mass WIMPs)



CoGeNT Collaboration, arXiv:1002.4703

The Signals/Hints/Detections

- Intriguingly, the CoGeNT and DAMA signals, if interpreted as dark matter, point to a similar region of parameter space
- Depending on the velocity distribution used, and on how one treats channeling, regions can be found in which both DAMA and CoGeNT can be explained by the same ~ 7 GeV dark matter particle

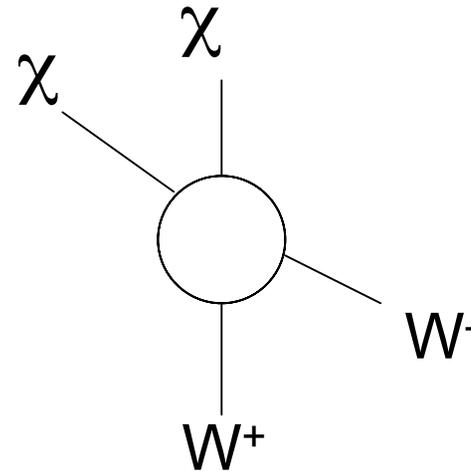


Fitzpatrick, Hooper,
Zurek,
arXiv:1003.0014

The Indirect Detection of Dark Matter

1. WIMP Annihilation

Typical final states include heavy fermions, gauge or Higgs bosons



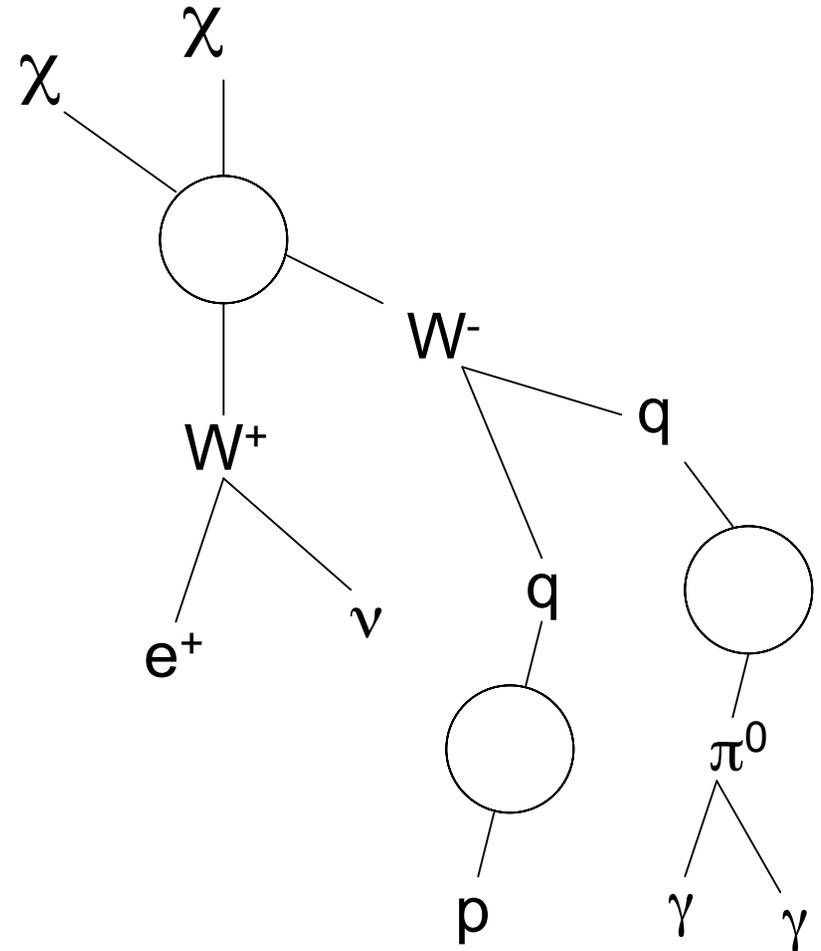
The Indirect Detection of Dark Matter

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2. Fragmentation/Decay

Annihilation products decay and/or fragment into combinations of electrons, protons, deuterium, neutrinos and gamma-rays



The Indirect Detection of Dark Matter

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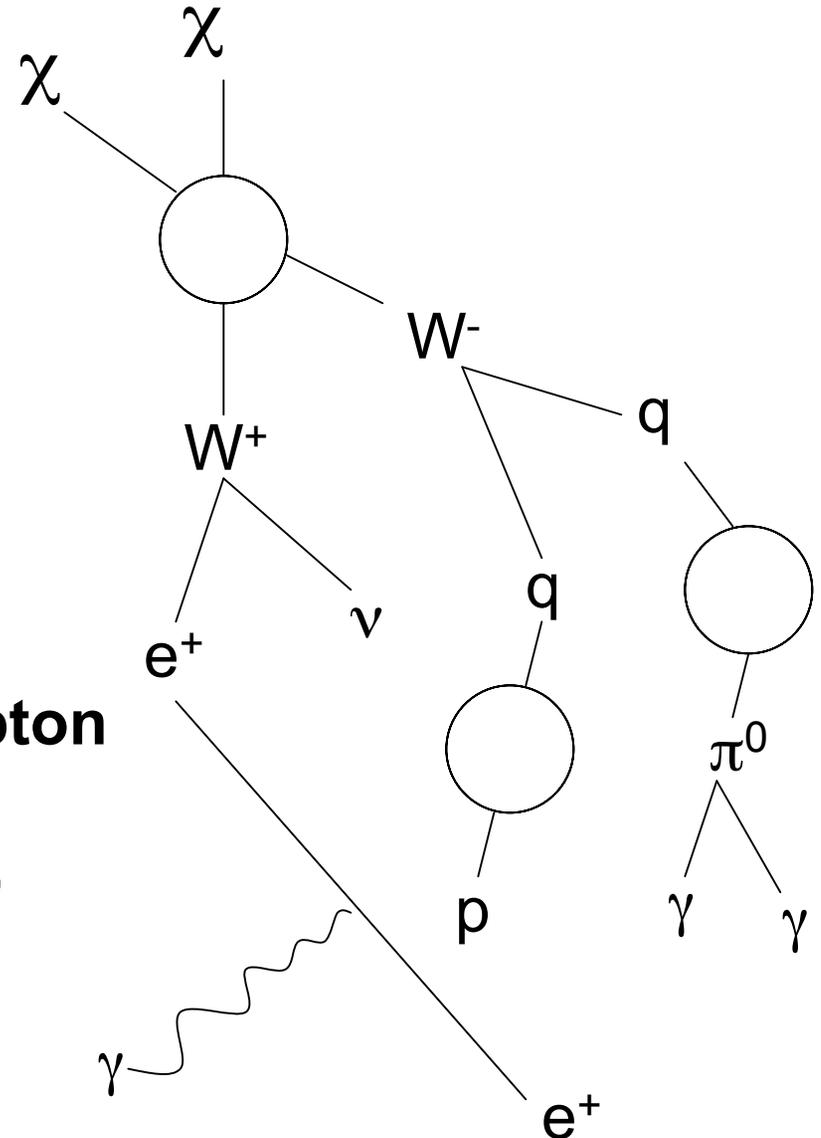
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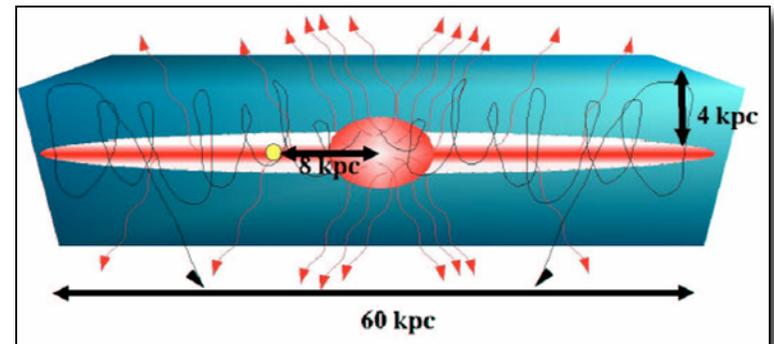
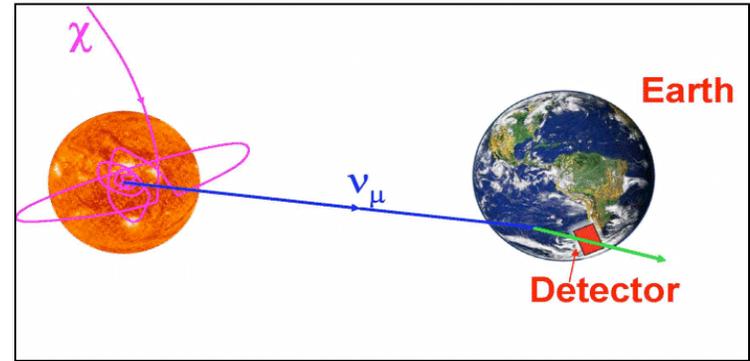
3. Synchrotron and Inverse Compton

Relativistic electrons up-scatter starlight/CMB to MeV-GeV energies, and emit synchrotron photons via interactions with magnetic fields



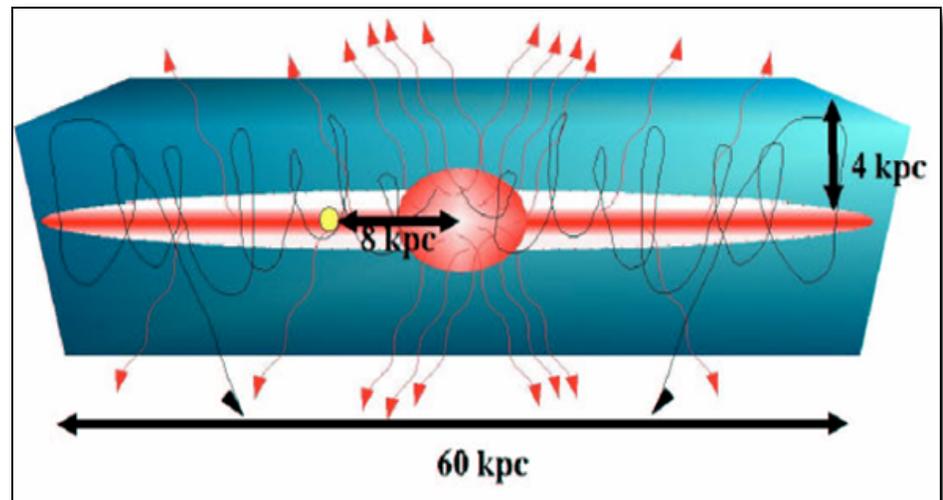
The Indirect Detection of Dark Matter

- **Neutrinos** from annihilations in the core of the Sun
- **Gamma Rays** from annihilations in the galactic halo, near the galactic center, in dwarf galaxies, etc.
- **Positrons/Antiprotons** from annihilations throughout the galactic halo
- **Synchrotron and Inverse Compton** from electron/positron interactions with the magnetic fields and radiation fields of the galaxy



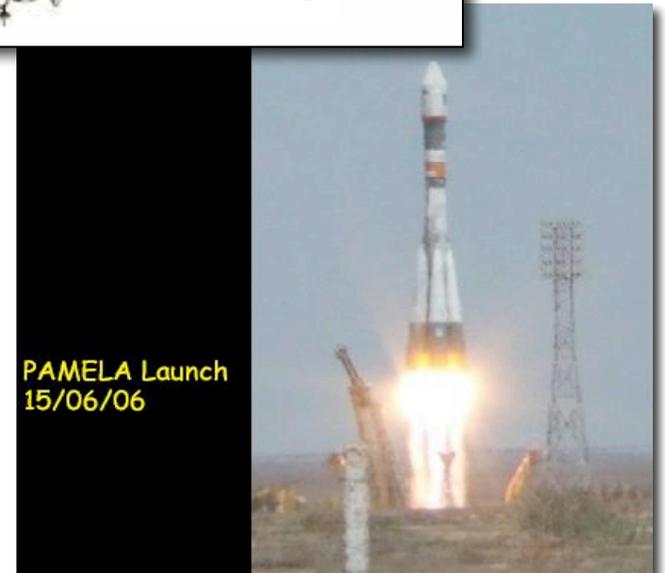
Dark Matter With Charged Cosmic Rays

- WIMP annihilation products fragment and decay, generating equal numbers of electrons and positrons, and of protons and antiprotons
- Charged particles move under the influence of the Galactic Magnetic Field; Electrons/positrons lose energy via synchrotron and inverse Compton scattering
- Astrophysical sources are generally expected to produce far more matter than antimatter; large positron/antiproton content in the cosmic ray spectrum could provide evidence for dark matter



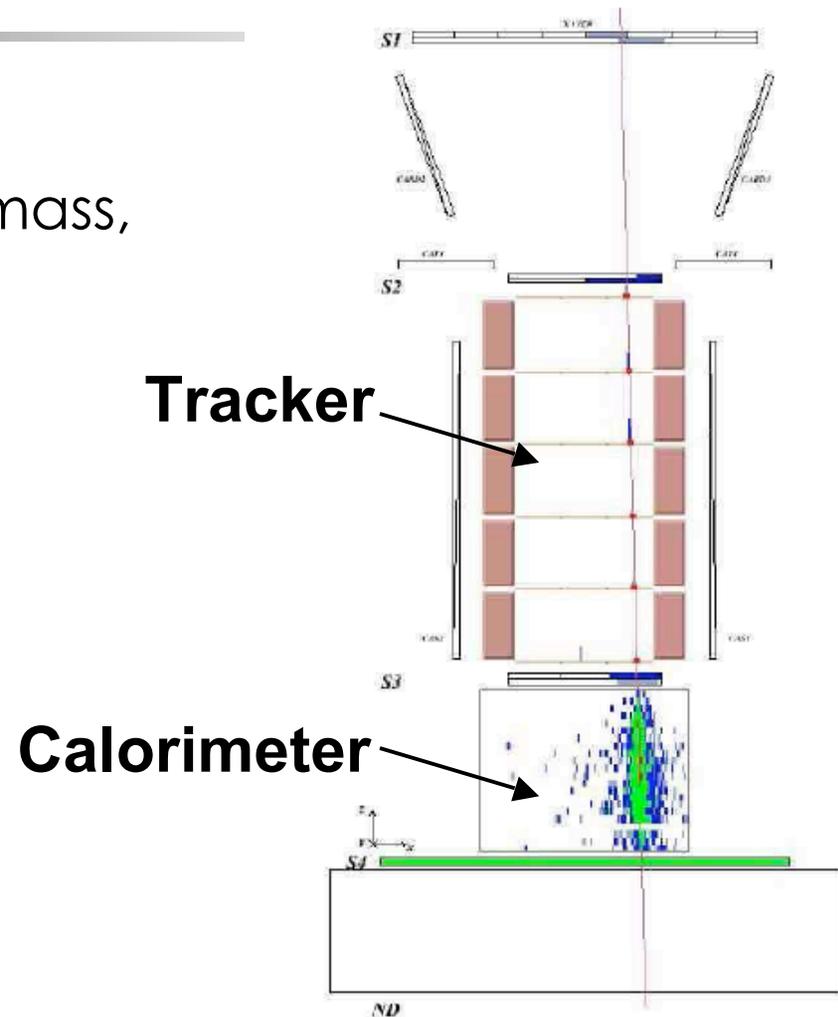
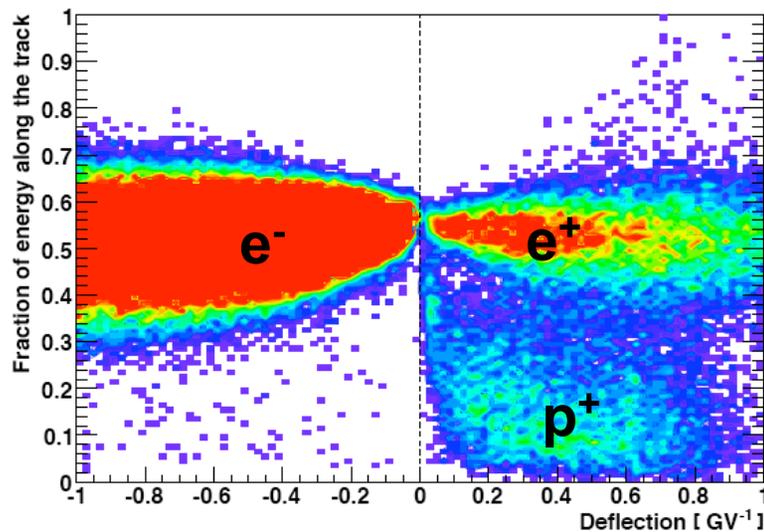
Charged Particle Astrophysics With Pamela

- Major step forward in sensitivity to GeV-TeV cosmic ray electrons, positrons, protons, antiprotons, and light nuclei
- Among other science goals, PAMELA hopes to identify or constrain dark matter annihilations in the Milky Way halo by measuring the cosmic positron and antiproton spectra

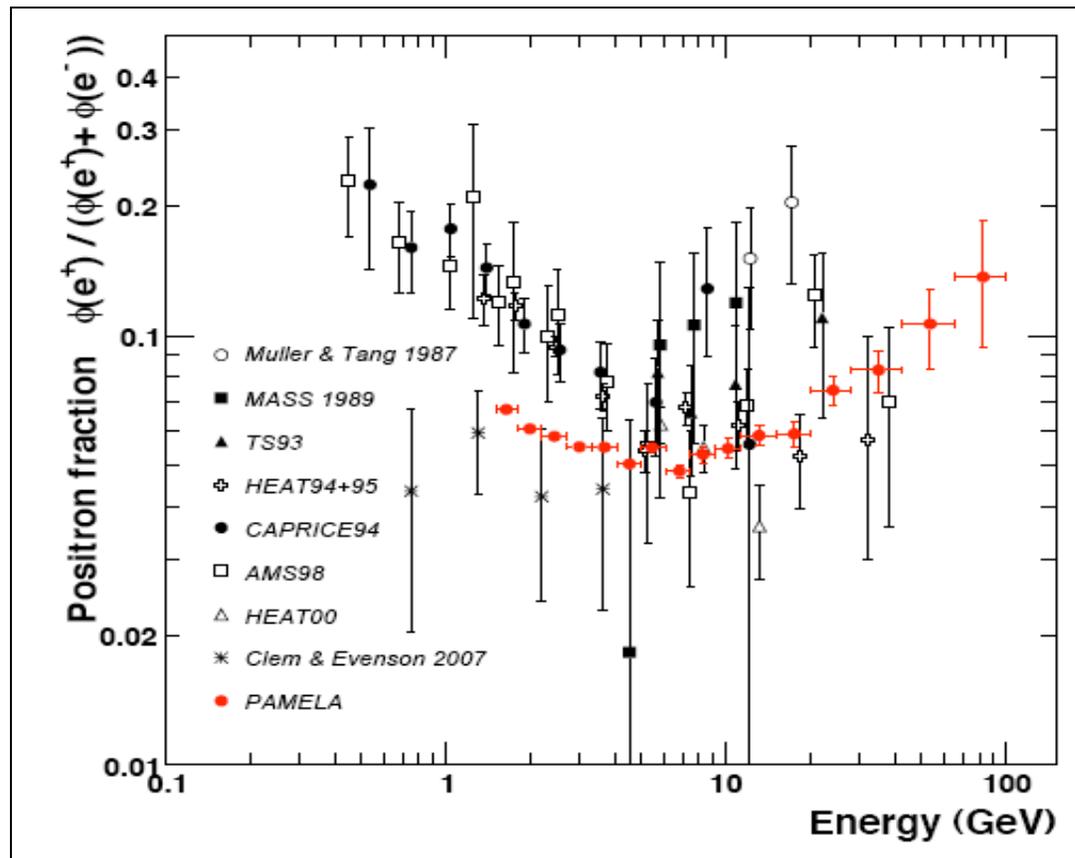


Charged Particle Astrophysics With Pamela

- Combination of tracker and calorimeter enable charge, mass, and energy determinations
- Very accurate particle ID

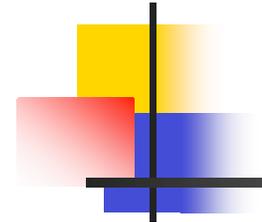


Pamela's New Positron Measurement



Pamela Collaboration, arXiv:0810.4995

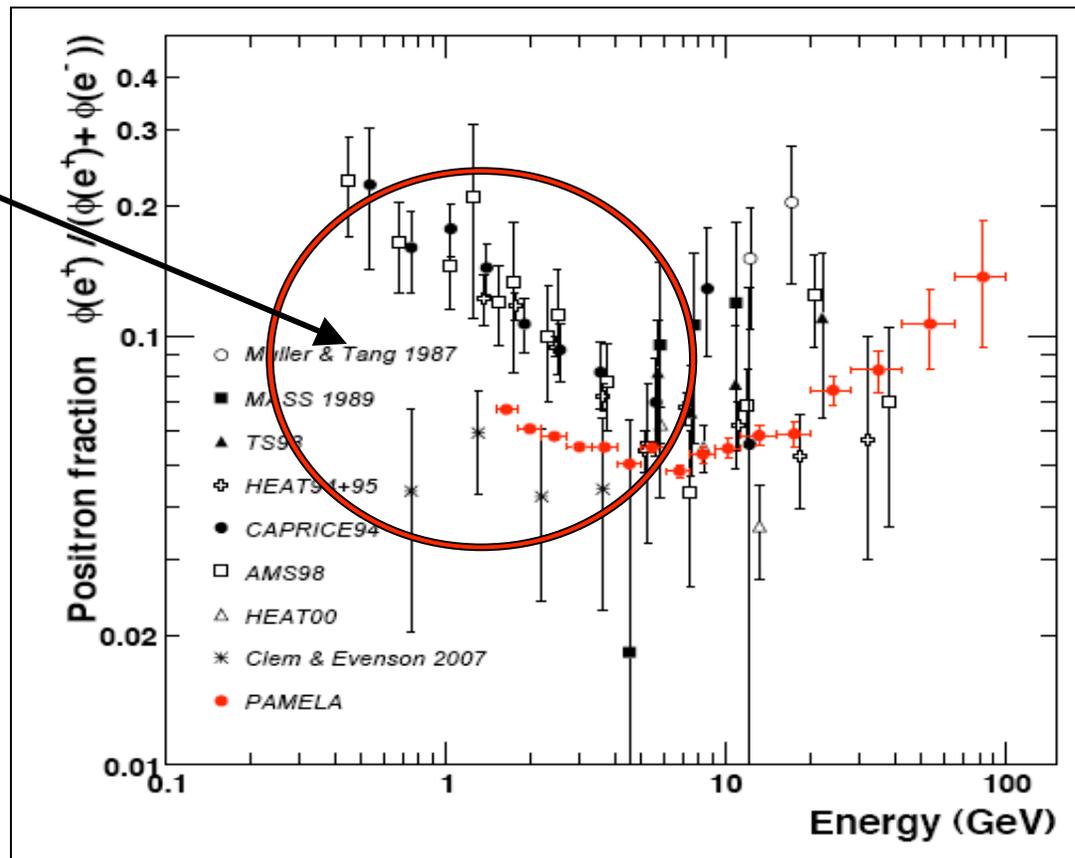
Pamela's New Positron Measurement



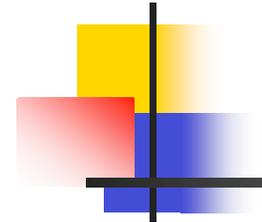
First glance:
-Is this all
screwed up?

Charge-dependent
solar modulation
important below
5-10 GeV!

***(Pamela's
sub-10 GeV
positrons appear
as they should!)***



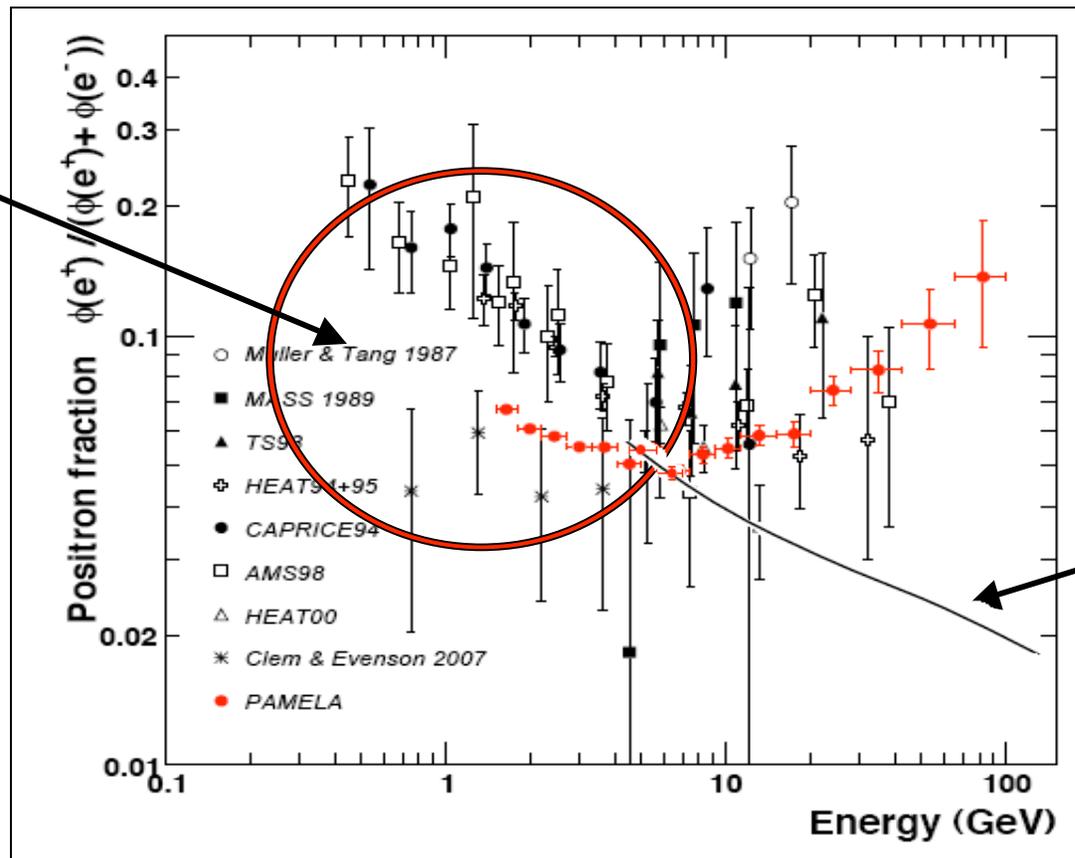
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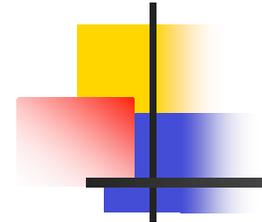
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Astrophysical
expectation
(secondary
production)

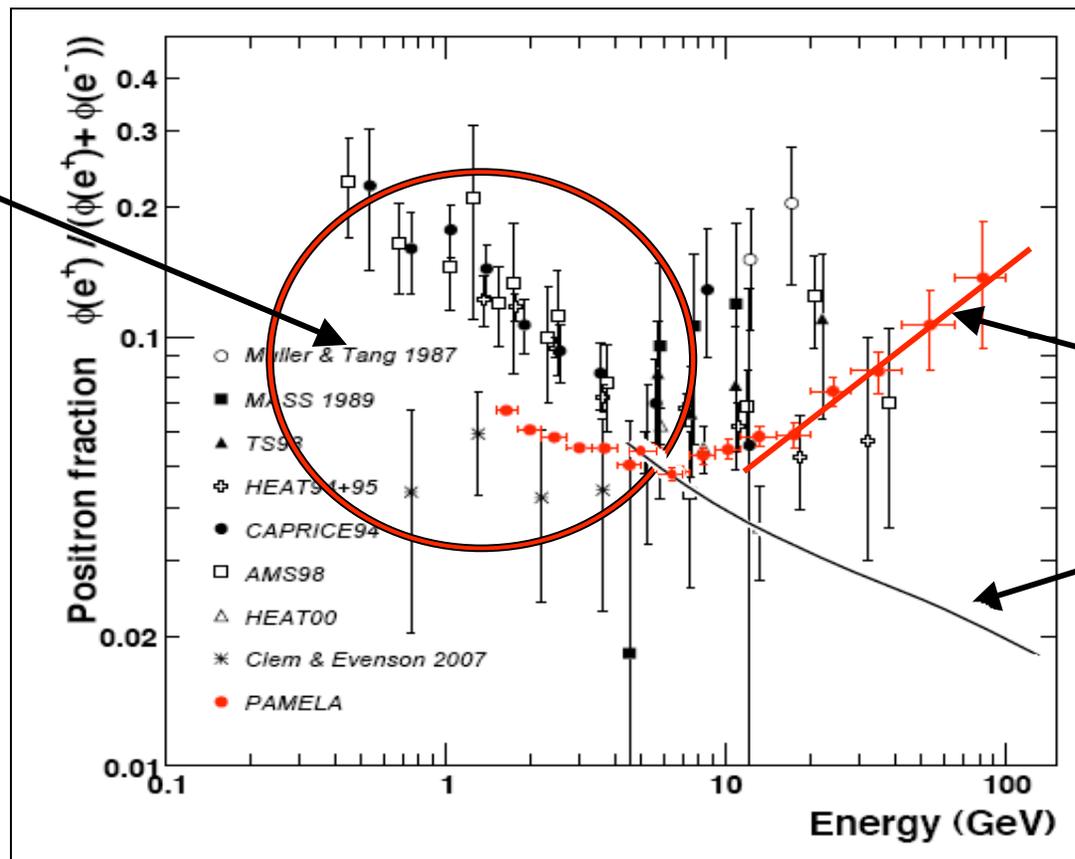
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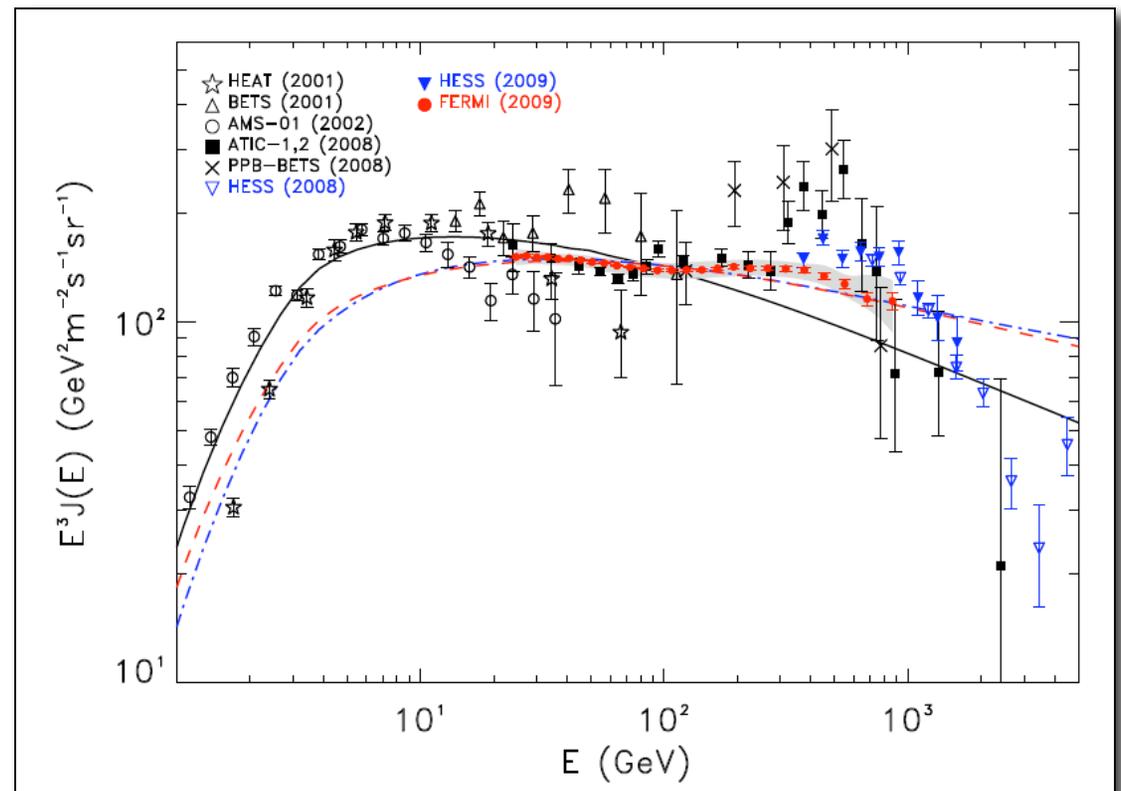


***Rapid climb
above 10 GeV
indicates the
presence of a
primary
source of
cosmic ray
positrons!***

Astrophysical
expectation
(secondary
production)

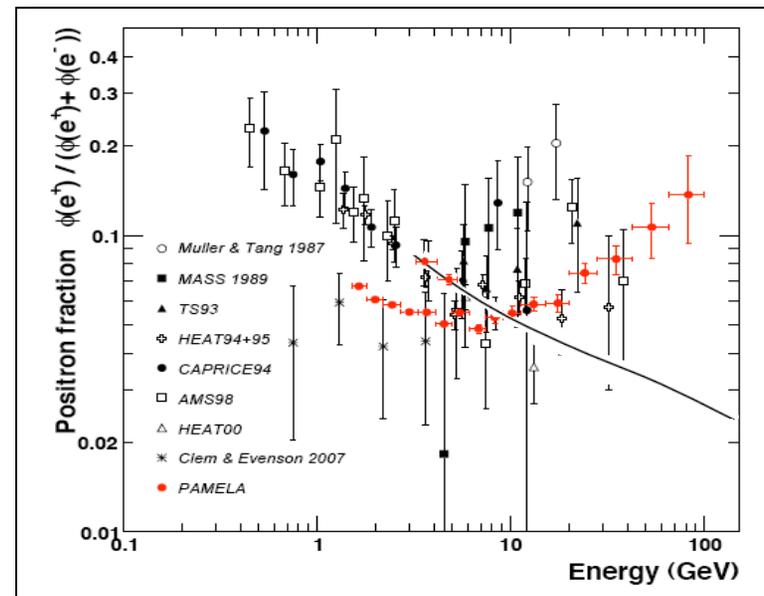
The Cosmic Ray Electron Spectrum

- Measurements from Fermi and HESS suggest a departure from simple homogeneous source models
- Not very surprising - At TeV energies, electrons sample only the surrounding ~1 kpc



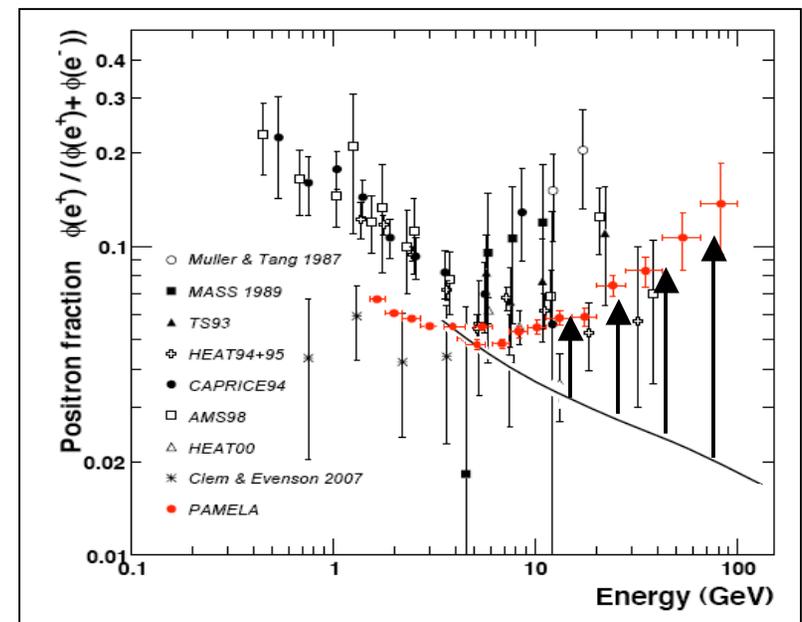
Possible Origins of the PAMELA Positron Excess

- Reacceleration of secondary positrons in or around supernova remnants
- Astrophysical primary positron sources (ie. pulsars)
- Dark matter annihilations or decays



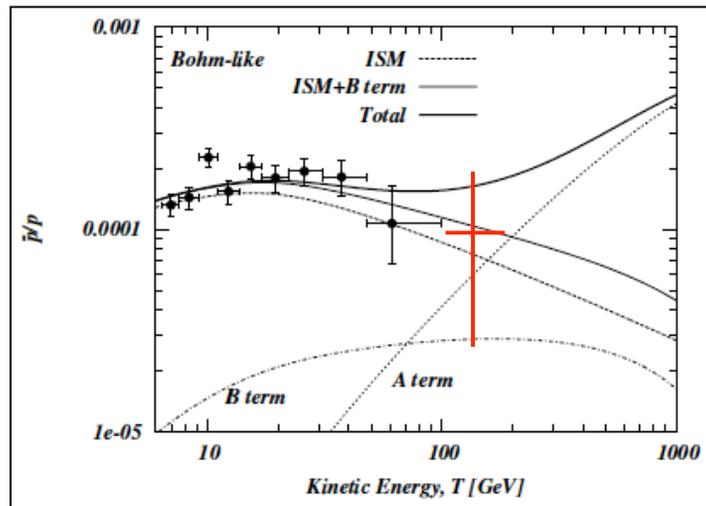
Possible Origins of the PAMELA Positron Excess

- The standard prediction for secondary positron production is calculated by combining the spectrum of cosmic ray protons, the density of targets, and the spectrum of cosmic ray electrons; Unavoidably leading to the prediction of a steadily falling positron fraction
- It has been suggested that if secondary positrons are produced *inside of* cosmic ray acceleration regions, their spectrum may be hardened, potentially causing the positron fraction to rise

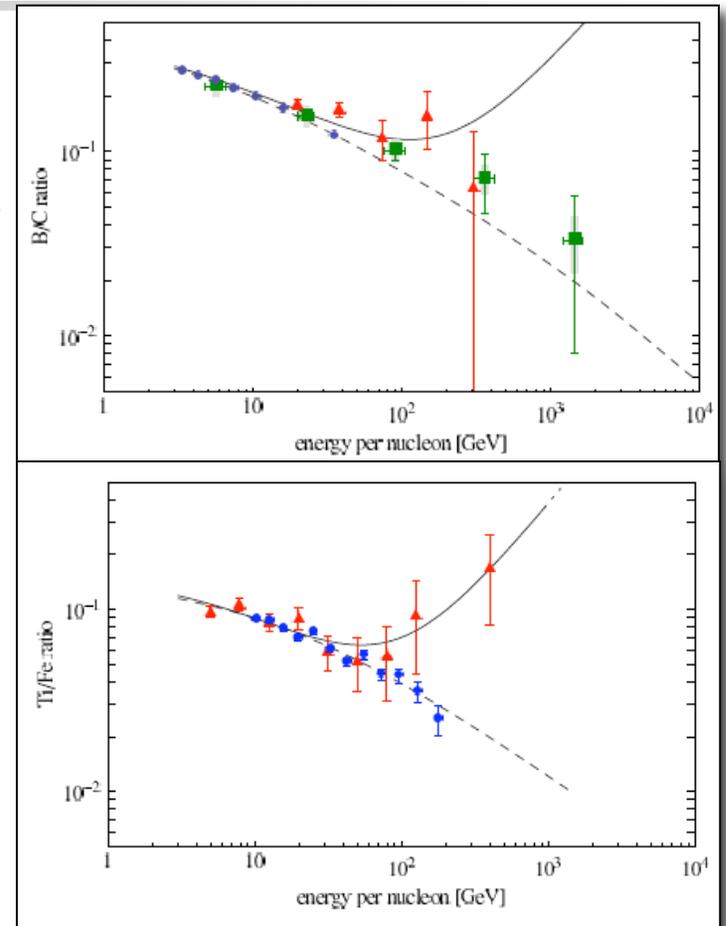


Possible Origins of the PAMELA Positron Excess

- The acceleration of positron secondaries, however, should be accompanied by the acceleration of antiproton, boron, and other secondary species
- This is not yet observed



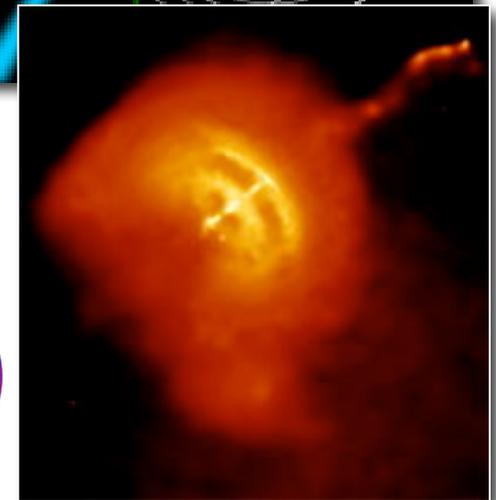
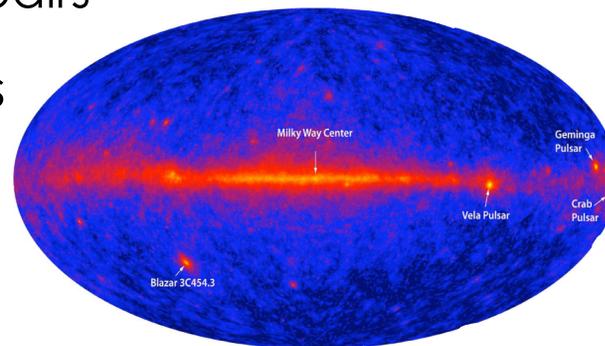
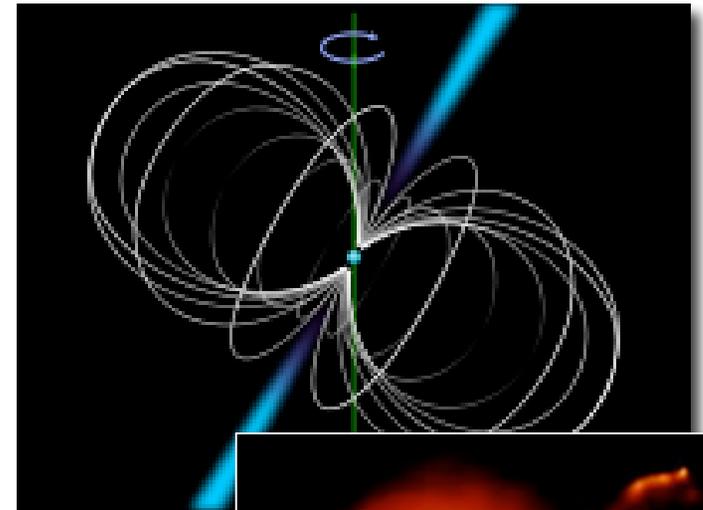
Blasi and Serpico, PRL, arXiv:0904.0871



Sarkar, Mertsch, PRL, arXiv:0905.3152;
Ahlers, Mertsch, Sarkar, PRD, arXiv:0909.4060

High-Energy Positrons and Electrons From Nearby Pulsars

- Rapidly spinning (\sim msec period) neutron stars, accelerate electrons to very high energies (power from slowing rotation - spindown)
- Energies can exceed the pair production threshold
- Very young pulsars ($\lesssim 10,000$ years) are typically surrounded by a pulsar wind nebula, which can absorb energetic pairs
- Most of the spindown power is expended in first $\sim 10^5$ years

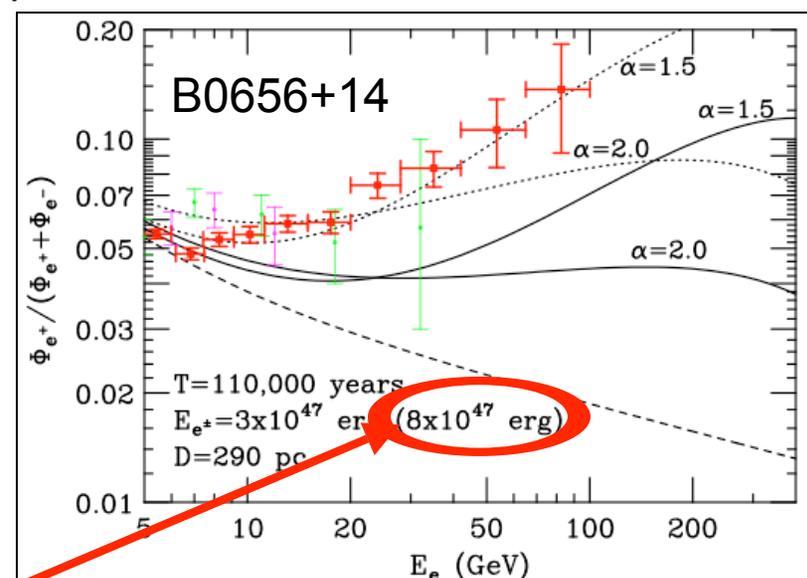
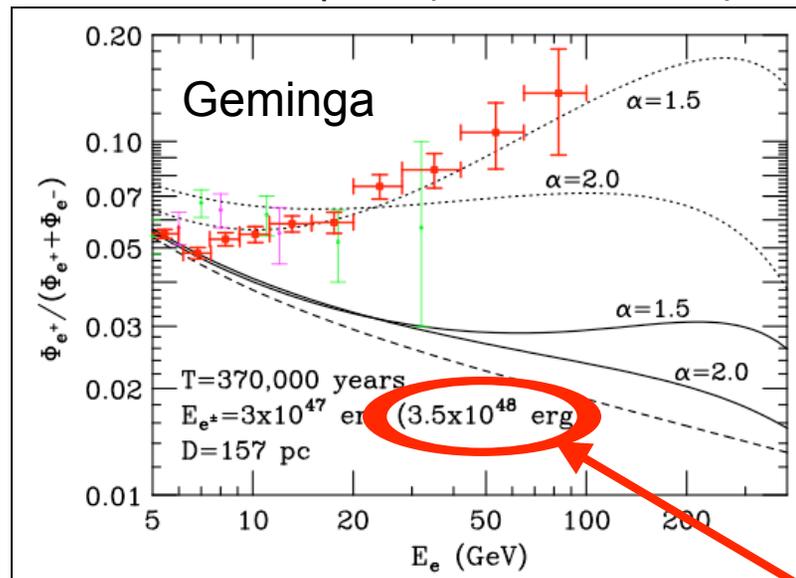


Vela Pulsar (12,000 yrs old)

High-Energy Positrons From Nearby Pulsars

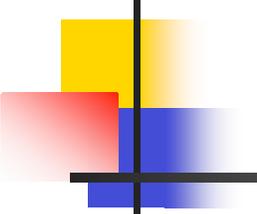
Two promising candidates:

- Geminga (157 pc away, 370,000 years old)
- B0656+14 (290 pc, 110,000 years)



Tens of percent of the total spindown energy is needed in high energy e^+e^- pairs!

Hooper, P. Blasi, P. Serpico,
JCAP, arXiv:0810.1527



Dark Matter and PAMELA

Dark matter annihilations in the halo of the Milky Way could explain the positron excess, although some obstacles exist

- **Hard positron spectrum** - Annihilation to leptons?
Local overdensity/clump?
- **Lack of excess antiprotons** - Annihilation to leptons?
Narrow diffusion region?
- **Large positron flux** - Non-thermal dark matter?
Sommerfeld enhancements?
Large degree of substructure?

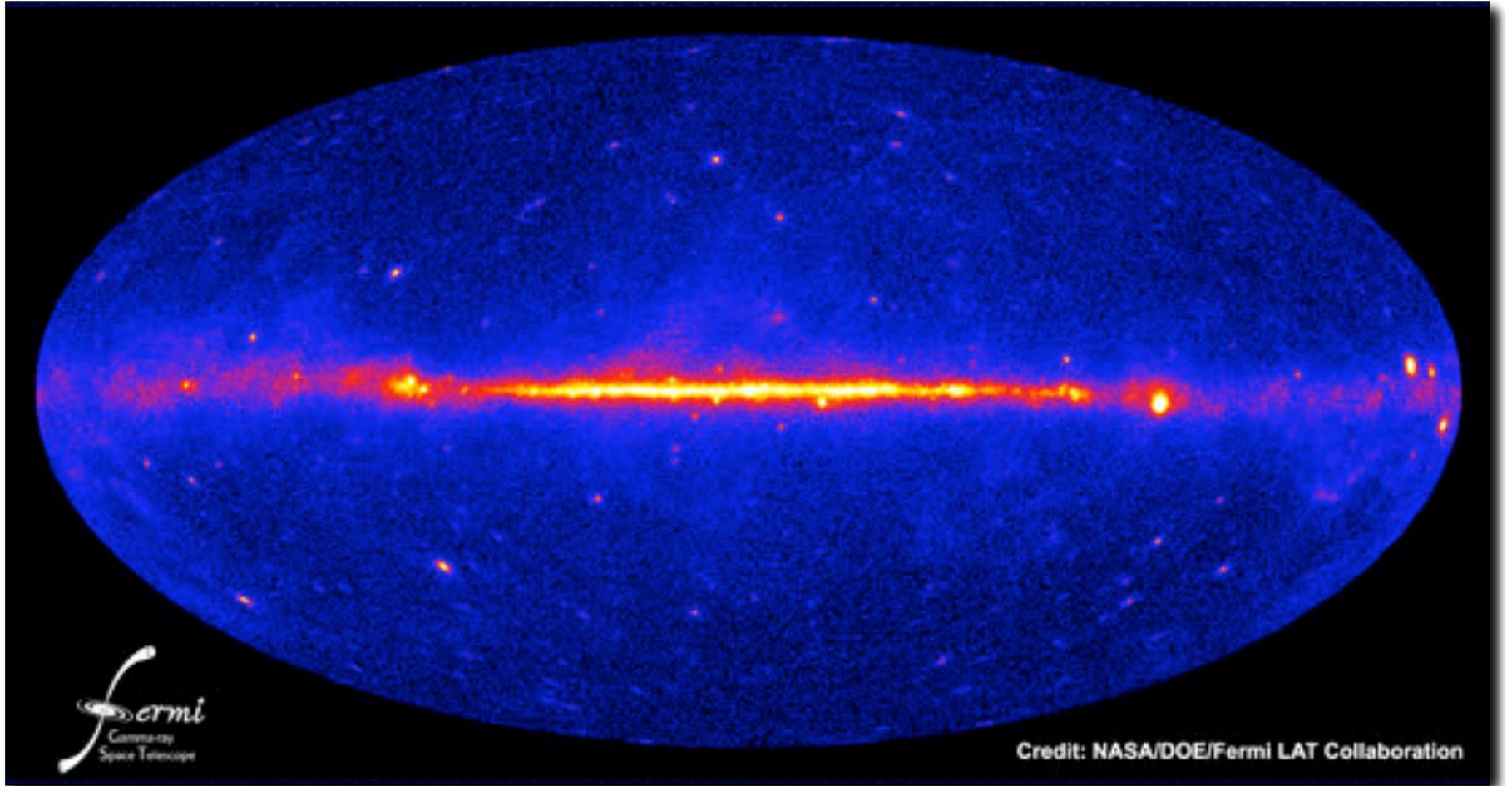
Pamela's positrons cannot be explained by Vanilla Dark Matter, but could be in more complex scenarios

An Essential Test:

Searches For Gamma Rays From Dark Matter Annihilations With Fermi

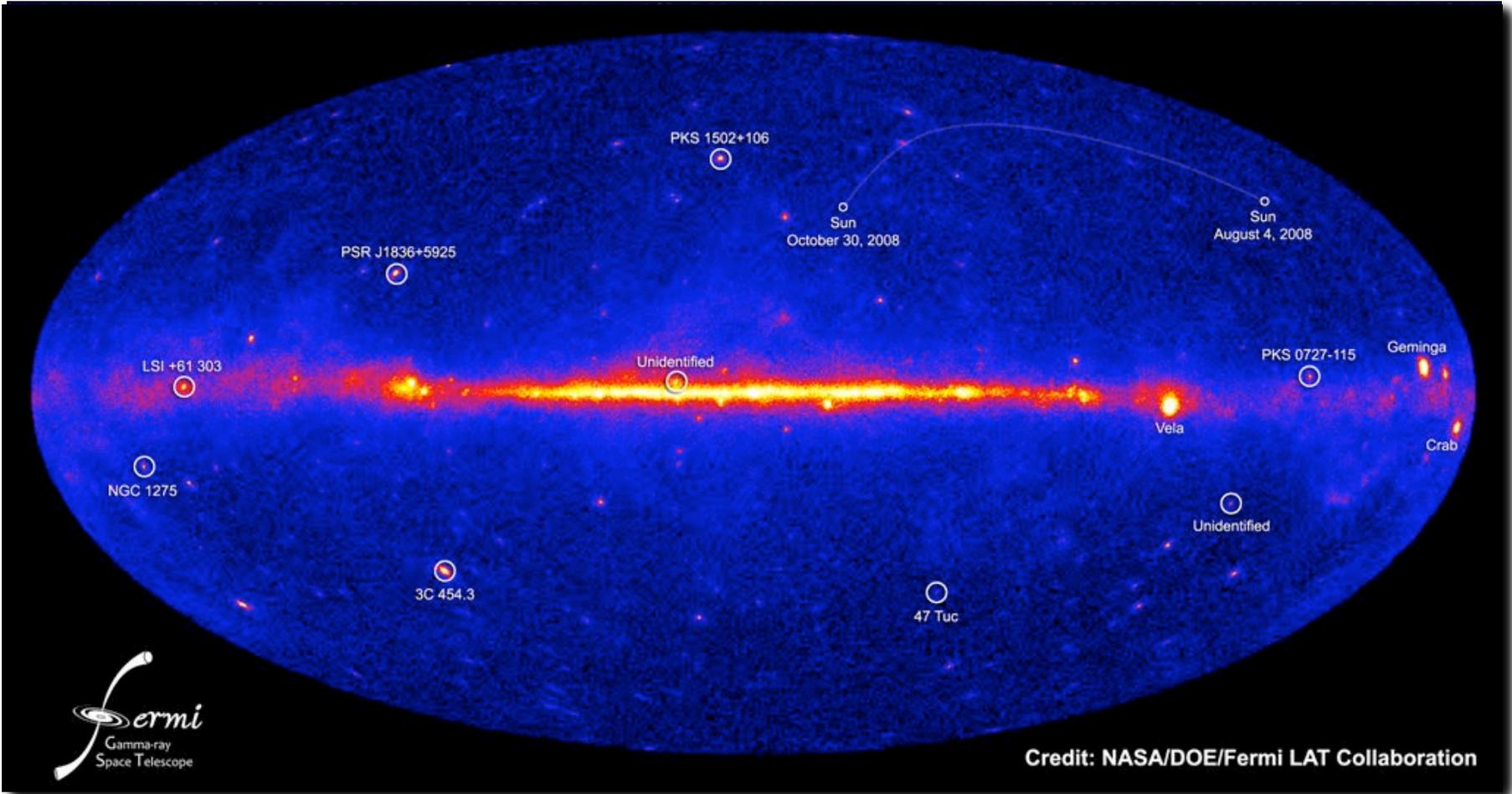
- Last year, the FERMI collaboration announced their first results!
- In August 2009, their first year data became publicly available
- Signatures of dark matter annihilation might appear clearly and quickly, or over years of exposure, or not at all, depending on the dark matter distribution, annihilation cross section, mass, and astrophysical backgrounds



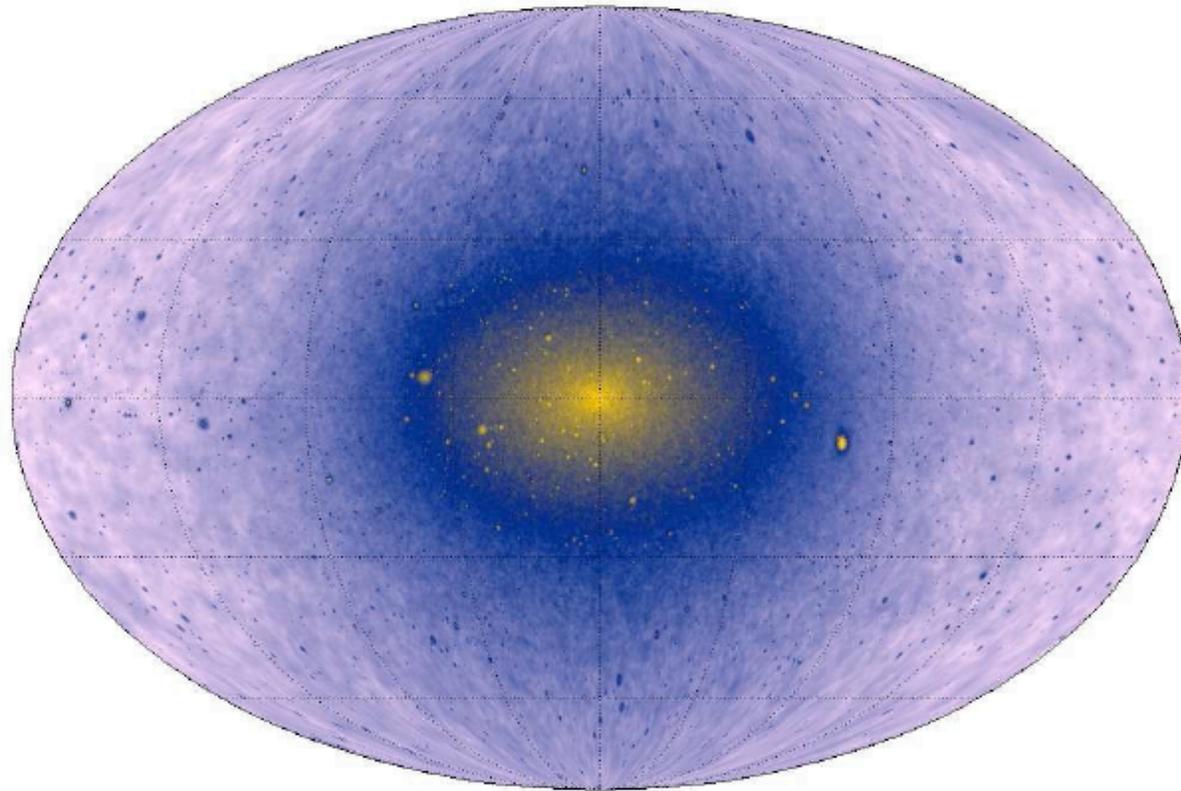


fermi
Gamma-ray
Space Telescope

Credit: NASA/DOE/Fermi LAT Collaboration



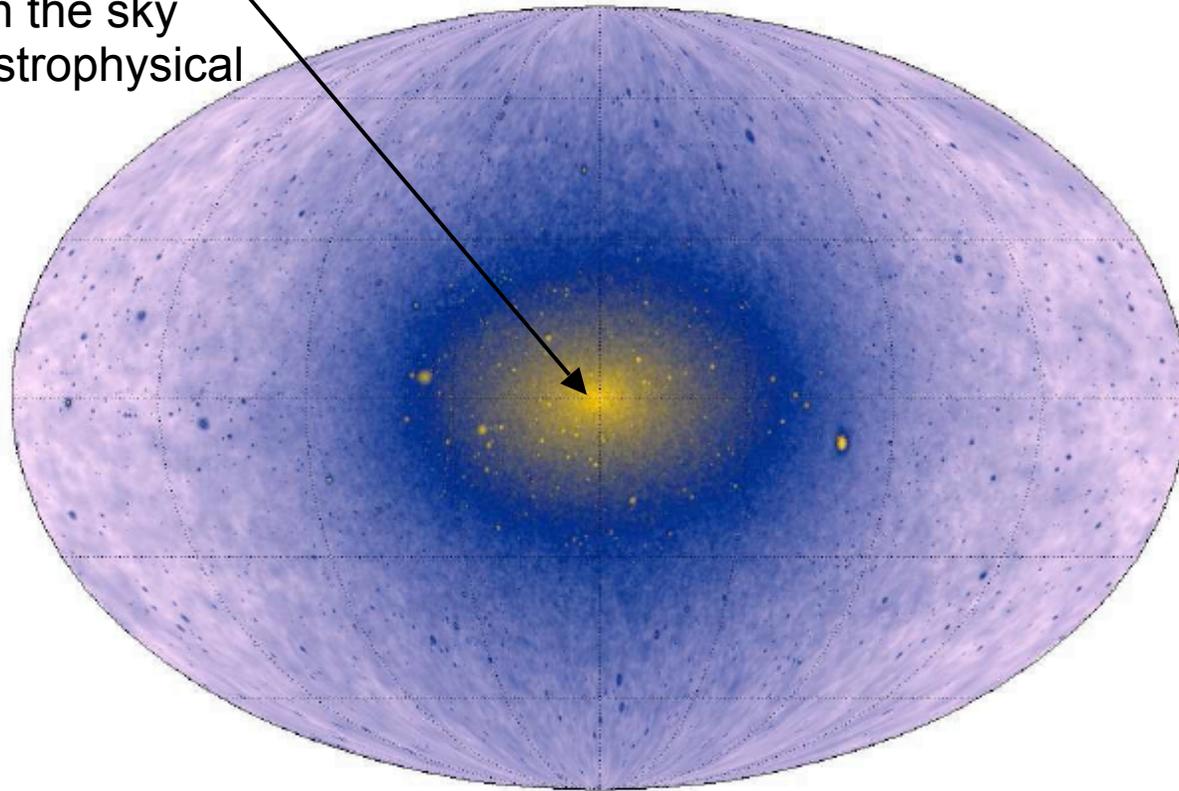
Where To Look For Dark Matter With Fermi?



Where To Look For Dark Matter With Fermi?

The Galactic Center

- Brightest spot in the sky
- Considerable astrophysical backgrounds



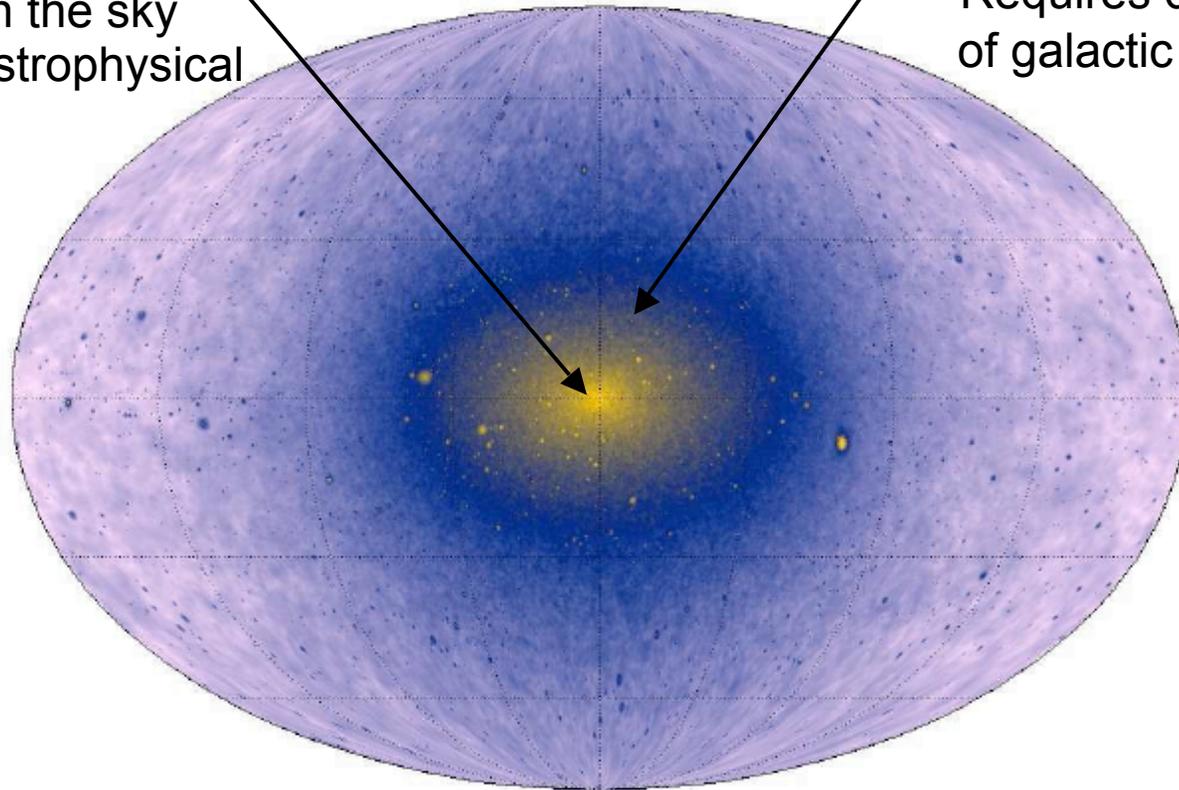
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The Galactic Halo

- High statistics
- Requires detailed model of galactic backgrounds



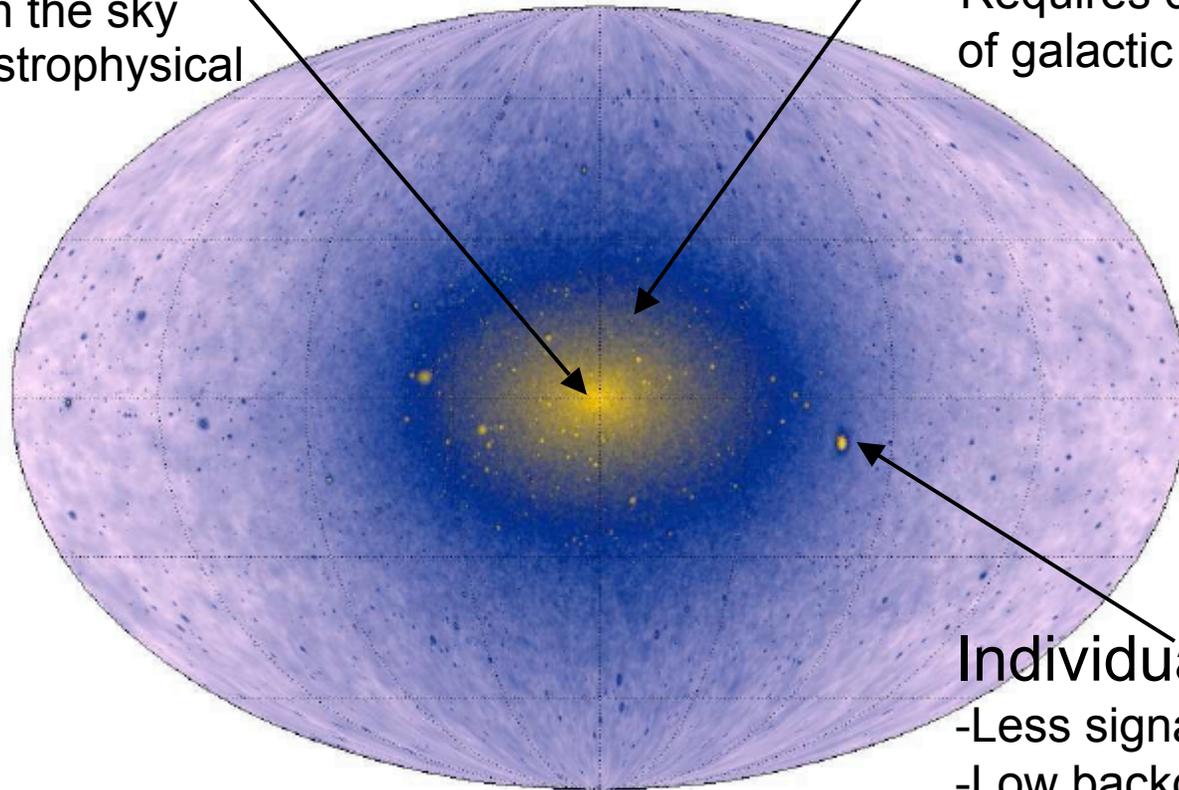
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Individual Subhalos

- Less signal
- Low backgrounds

Where To Look For Dark Matter With Fermi?

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The Galactic Halo

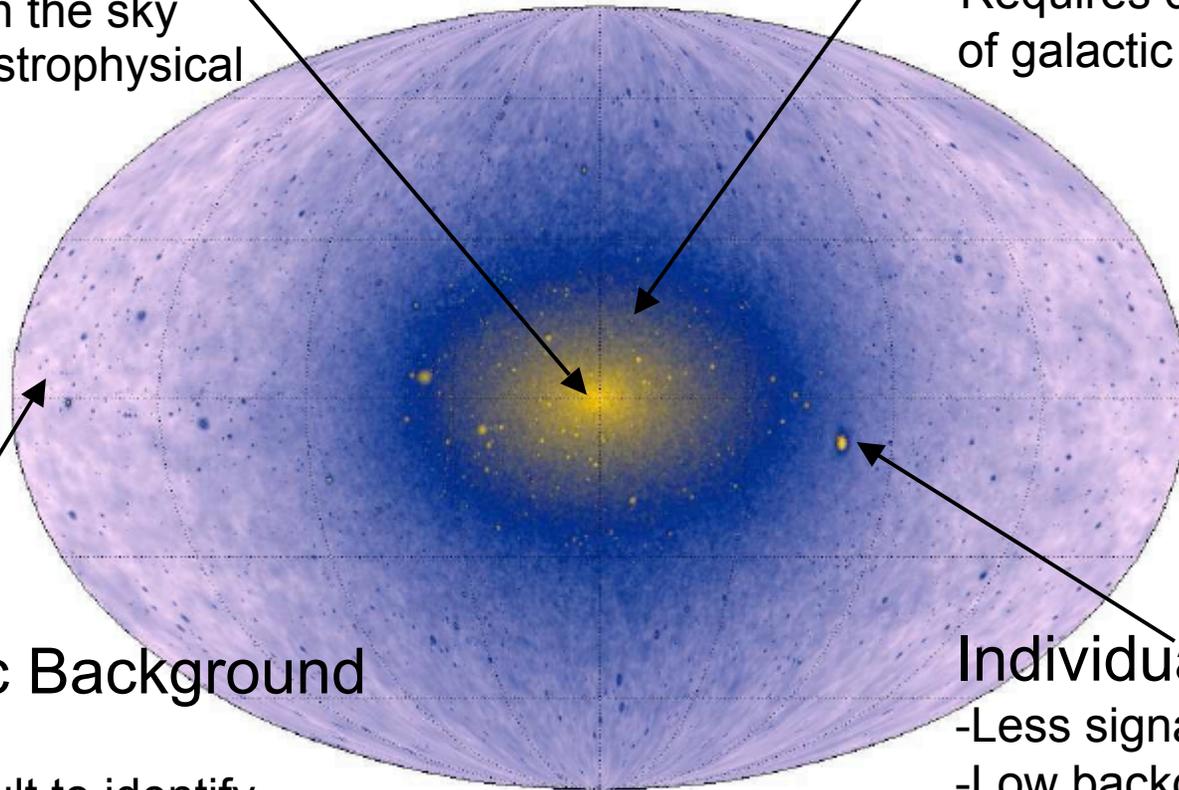
- High statistics
- Requires detailed model of galactic backgrounds

Extragalactic Background

- High statistics
- potentially difficult to identify

Individual Subhalos

- Less signal
- Low backgrounds



Some of the most interesting early dark matter results from Fermi

- 1) Galactic Diffuse Emission Measurement
- 2) The Galactic Center Region
- 3) Subhalos
- 4) Galaxy Clusters
- 5) Dwarf Spheroidal Galaxies
- 6) Line searches
- 7) Diffuse ICS (the “Fermi Haze”)
- 8) Isotropic Diffuse

~~Haze~~
Bubbles



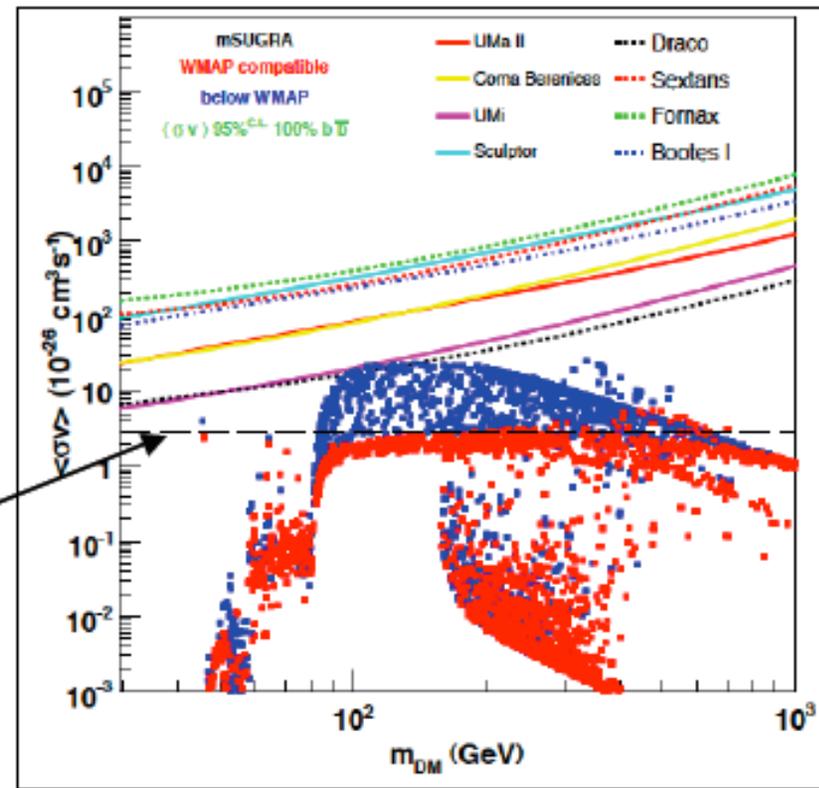
Dwarf Spheroidal Galaxies

- The FGST collaboration has recently placed some relatively stringent limits on dark matter from observations of a number of satellite galaxies (dwarf spheroidals) of the Milky Way
- The most stringent limits come from those dwarfs which are 1) dense, 2) nearby, and 3) in low background regions of the sky

Name	Distance (kpc)	year of discovery	$M_{1/2}/L_{1/2}$ ref. 8	l	b	Ref.
Ursa Major II	30± 5	2006	4000 ⁺³⁷⁰⁰ ₋₂₁₀₀	152.46	37.44	1,2
Segue 2	35	2009	650	149.4	-38.01	3
Willman 1	38± 7	2004	770 ⁺⁹³⁰ ₋₄₄₀	158.57	56.78	1
Coma Berenices	44± 4	2006	1100 ⁺⁸⁰⁰ ₋₅₀₀	241.9	83.6	1,2
Bootes II	46	2007	18000??	353.69	68.87	6,7
Bootes I	62±3	2006	1700 ⁺¹⁴⁰⁰ ₋₇₀₀	358.08	69.62	6
Ursa Minor	66± 3	1954	290 ⁺¹⁴⁰ ₋₉₀	104.95	44.80	4,5
Sculptor	79± 4	1937	18 ⁺⁶ ₋₅	287.15	-83.16	4,5
Draco	76± 5	1954	200 ⁺³⁰ ₋₂₀	86.37	34.72	4,5,9
Sextans	86± 4	1990	120 ⁺³⁰ ₋₃₅	243.4	42.2	4,5
Ursa Major I	97±4	2005	1800 ⁺¹³⁰⁰ ₋₃₀₀	159.43	54.41	6
Hercules	132±12	2006	1400 ⁺¹³⁰⁰ ₋₃₀₀	28.73	36.87	6
Fornax	138± 8	1938	8.7 ^{+2.2} _{-2.3}	237.1	-65.7	4,5
Leo IV	160±15	2006	260 ⁺¹⁰⁰⁰ ₋₂₀₀	265.44	56.51	6

Dwarf Spheroidal Galaxies

- Analysis assumed an NFW profile form, with parameters constrained by stellar velocity measurements
- Ursa Minor and Draco provide the strongest constraints, with several others within a factor of ~ 3 -10
- Constraints for $M_{\text{DM}} \sim 100$ GeV or less are within a factor of a few of the value predicted for a simple thermal relic
- Based on only 11 months of data; given the low backgrounds, this curve will come down significantly with exposure (stacked dwarf analysis in progress!)

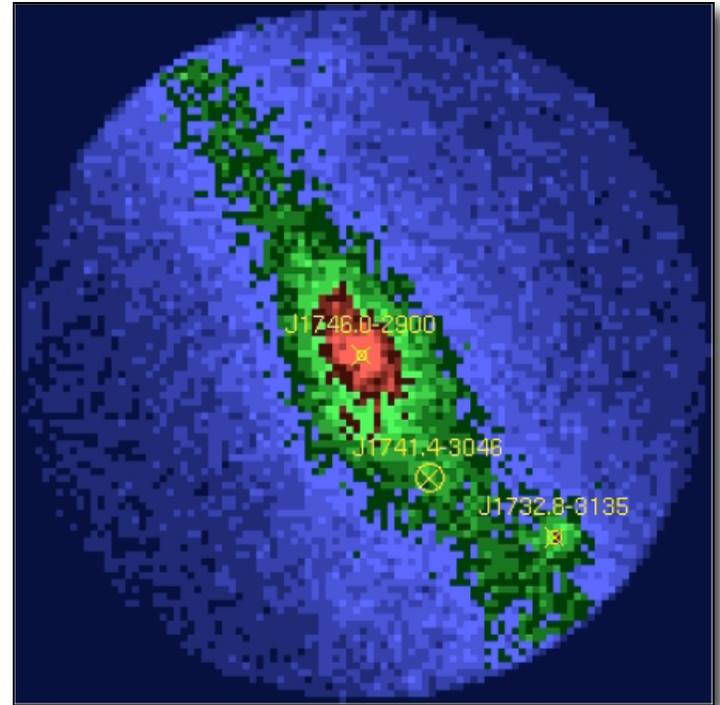


FGST Collaboration, arXiv:1001.4531

See also the analysis of Segue 1 by P. Scott *et al.*, arXiv:0909.3300

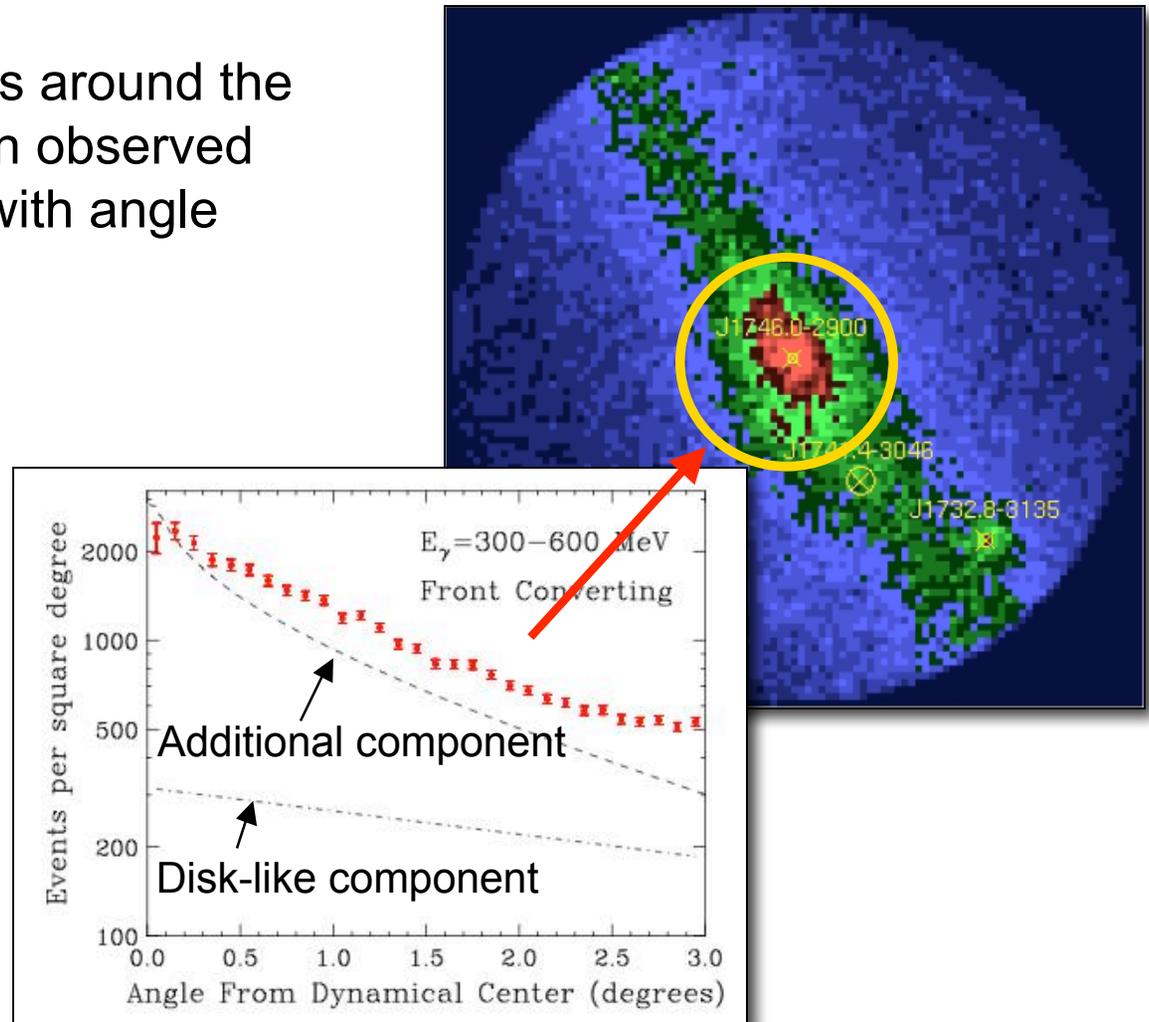
Dark Matter In The Galactic Center Region

- The region surrounding the Galactic Center is complex; backgrounds are poorly understood
- This does not necessarily make searches for dark matter in this region intractable, however



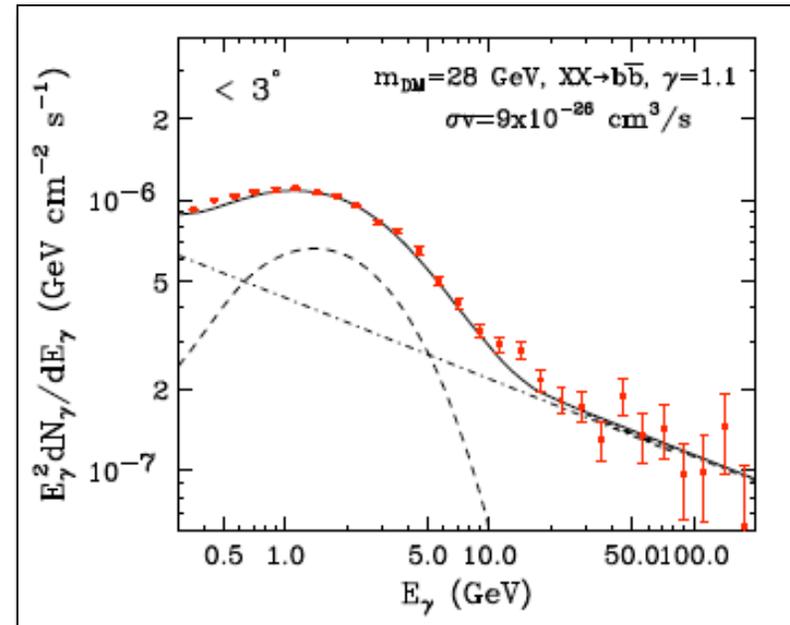
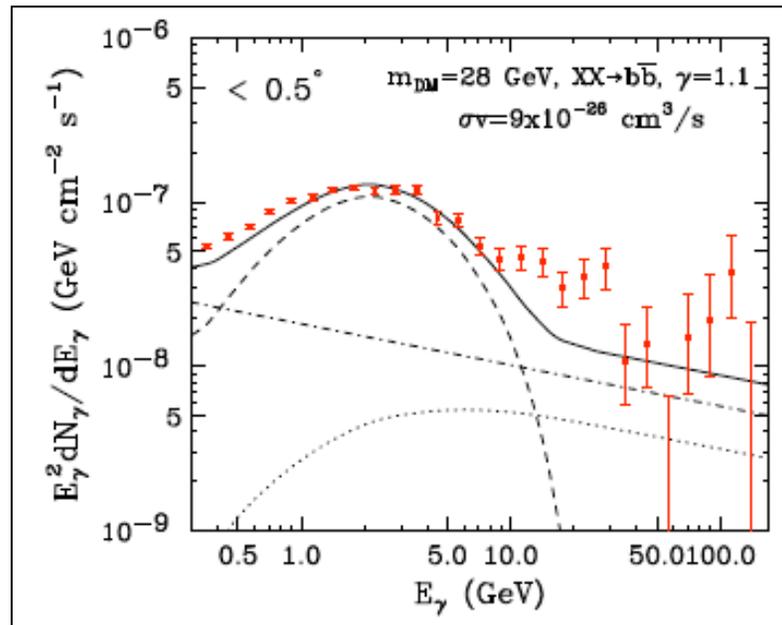
Dark Matter In The Galactic Center Region

- Within the inner few degrees around the Galactic Center, the emission observed by FGST steeply increases with angle
- If the diffuse background is modeled with the shape of the disk emission between 3° and 6° , another component is required that is more concentrated and spherically symmetric



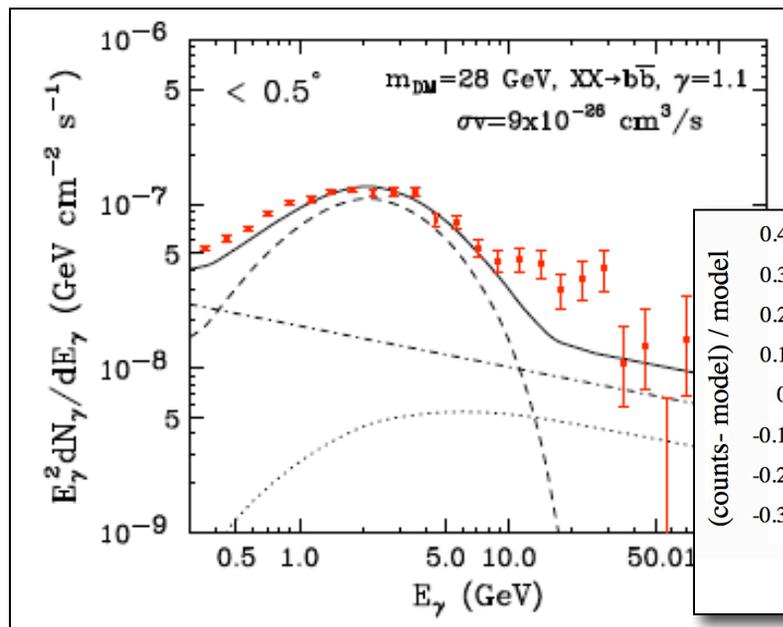
Dark Matter In The Galactic Center Region

- Intriguingly, the spectrum of the spherically symmetric component contains a “bump-like” feature at $\sim 1\text{-}5$ GeV
- Can be fit quite well by a simple 25-30 GeV dark matter particle, in a cusped distribution ($\gamma \sim 1.1$), annihilating to $b\bar{b}$ with $\sigma v \sim 9 \times 10^{-26} \text{ cm}^3/\text{s}$

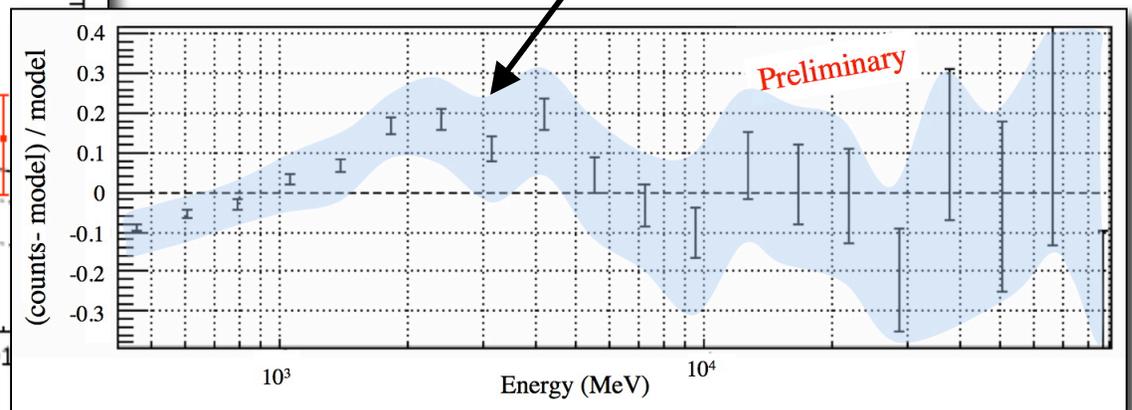


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- Recent presentations by the Fermi collaboration confirm the presence of this feature



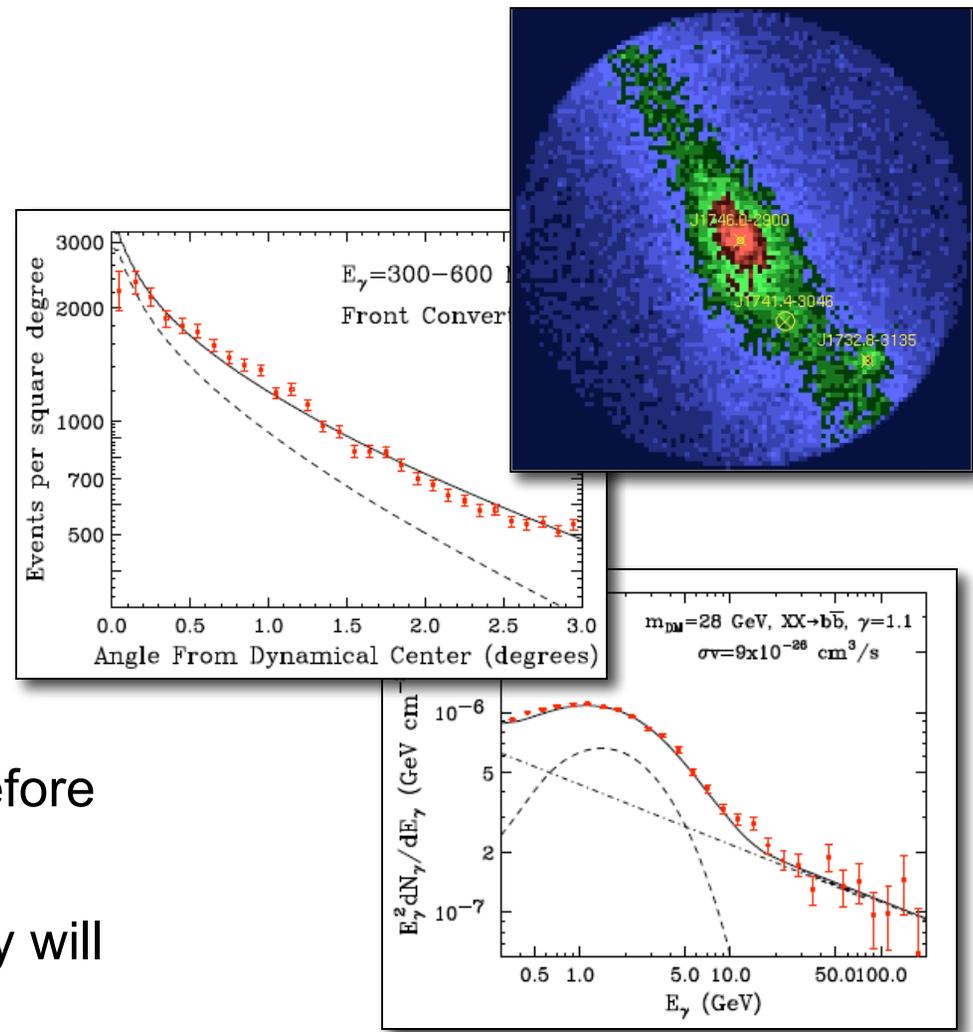
(Fermi Collaboration, Preliminary)

L. Goodenough, D. Hooper, arXiv:0910.2998

Dark Matter In The Galactic Center Region

Some words of caution:

- Although the angular distribution and spectrum observed from the inner Milky Way by FGST can be well fit by a simple annihilating dark matter scenario, an astrophysical background with a similar angular distribution and spectrum cannot be ruled out (π^0 decays have similar spectral shape, for example)
- The inner galaxy is a complex region, which must be scrutinized before any confident claims can be made
- Searches in other regions of the sky will be important to confirm or refute this interpretation



L. Goodenough, D. Hooper, arXiv:0910.2998

Dark Matter on Ice

- WIMPs can elastically scattering with nuclei in the Sun, leading them to become gravitationally captured
- After millions/billions of years, the annihilation rate is (in many models) predicted to reach equilibrium with the capture rate
- Neutrino telescopes such as IceCube are sensitive to neutrinos from WIMP annihilations in the Sun

IceTop: Air shower detector

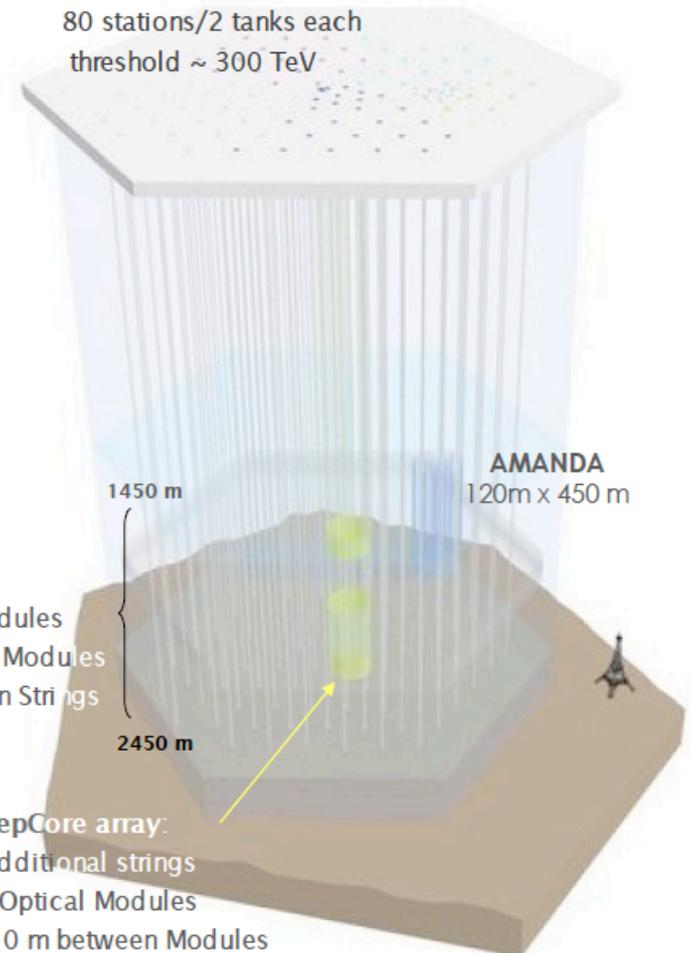
80 stations/2 tanks each
threshold ~ 300 TeV

InIce array:

80 Strings
60 Optical Modules
17 m between Modules
125 m between Strings

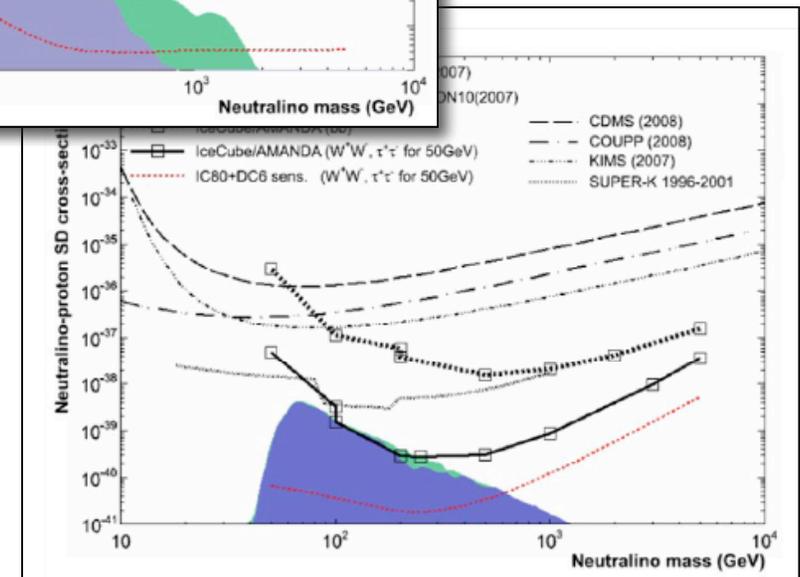
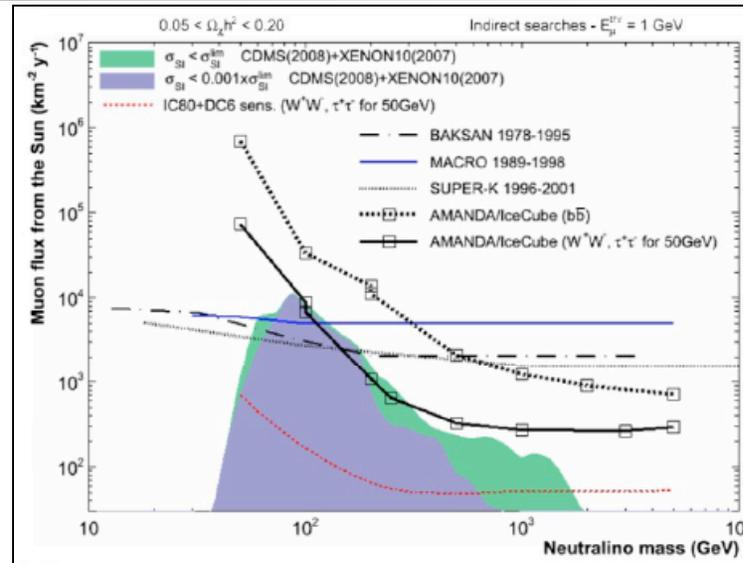
DeepCore array:

6 additional strings
60 Optical Modules
7/10 m between Modules
72 m between Strings



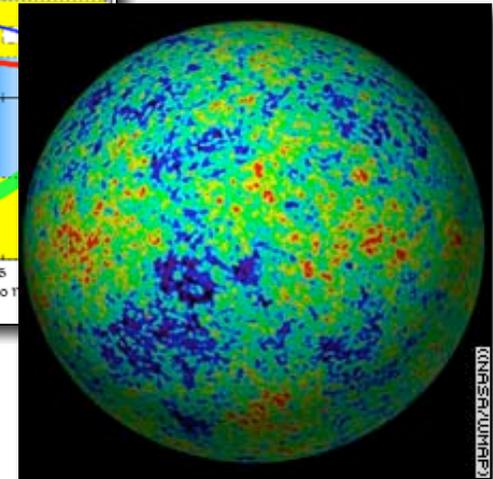
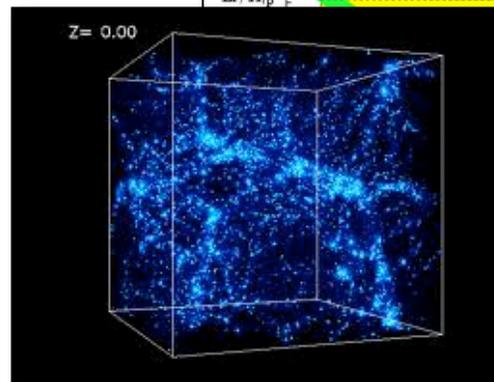
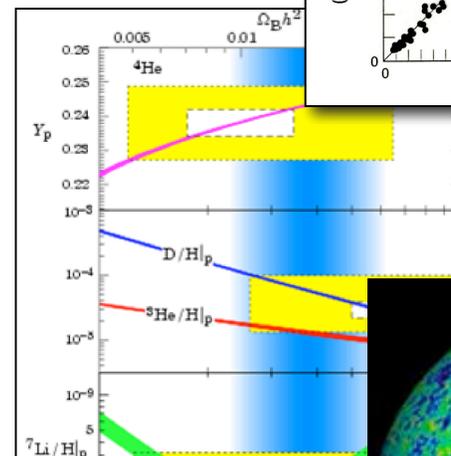
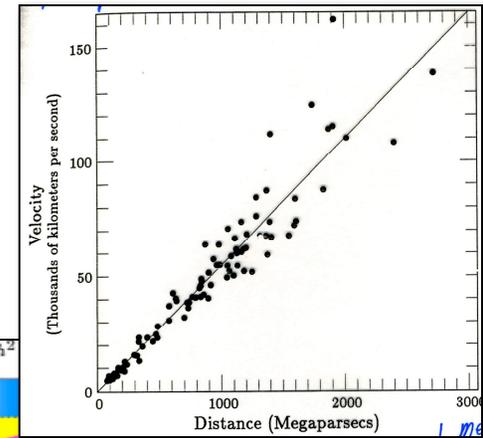
Dark Matter on Ice

- IceCube has placed constraints on the neutrino-induced muon flux from the Sun, and interpreted this in terms of a limit on the WIMP's spin-dependent elastic scattering cross section with protons
- These limits are competitive with or stronger than those from current direct detection experiments



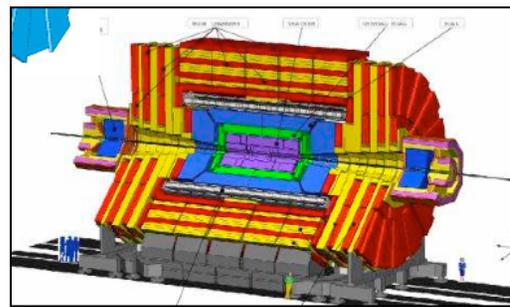
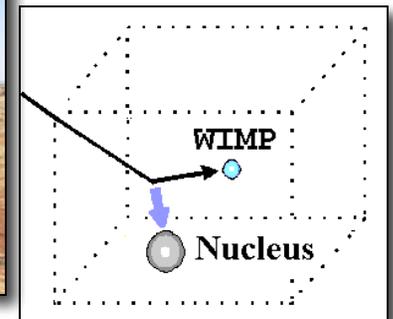
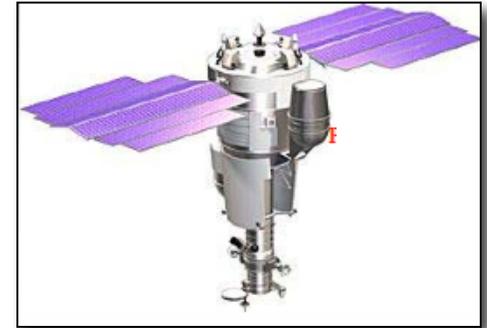
Summary

- Big Bang cosmology has been incredibly successful in explaining a wide range of observed phenomena
- Yet some aspects of our universe's history and composition remain poorly understood - dark matter, dark energy, inflation
- Weakly interacting massive particles provide a natural candidate for the dark matter, with a simple and compelling explanation for the observed dark matter abundance



Summary

- In addition to searches for dark matter at the LHC, both direct and indirect astrophysical searches are approaching the sensitivities thought to be required to observe dark matter non-gravitationally
- A number of reported signals have been interpreted as possible detections of dark matter (CoGeNT, DAMA, Fermi, Pamela)



Summary

One Year From Now

- New direct detection results from XENON 100, COUPP, and others
- Pamela positron fraction up to 200 GeV? And first data from AMS-02
- More data from Fermi, and more analysis of Fermi data
- Further input from ground based gamma ray telescopes, and observations at other wavelengths

