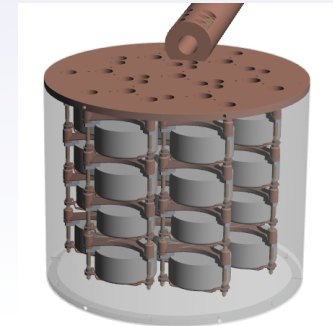
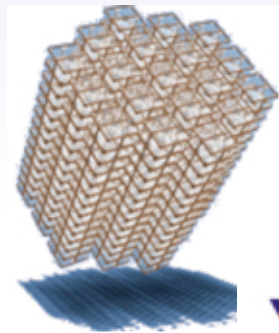
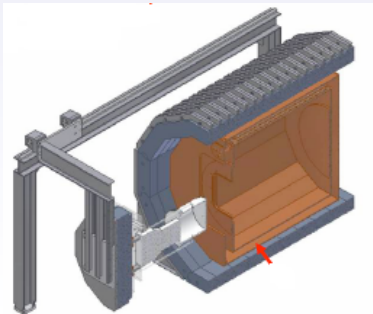
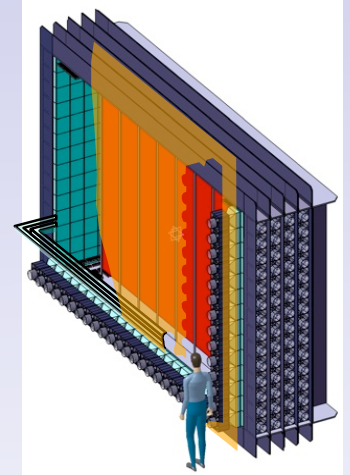
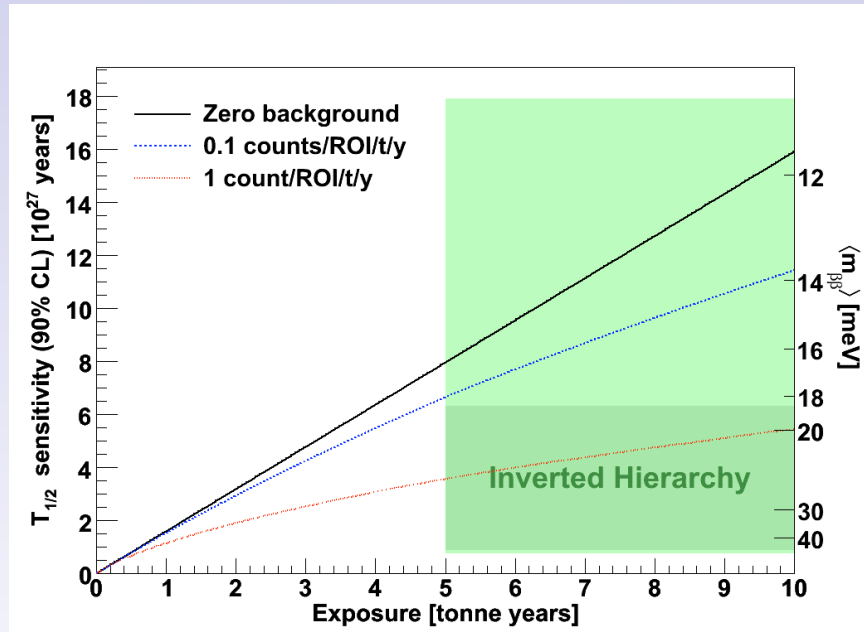


Neutrinoless Double Beta Decay II

Experimental approaches



THE UNIVERSITY
of NORTH CAROLINA
at CHAPEL HILL



J.F. Wilkerson
International Neutrino Summer School
July 13, 2009

$0\nu\beta\beta$ Decay II

Overview

Review

Early Days

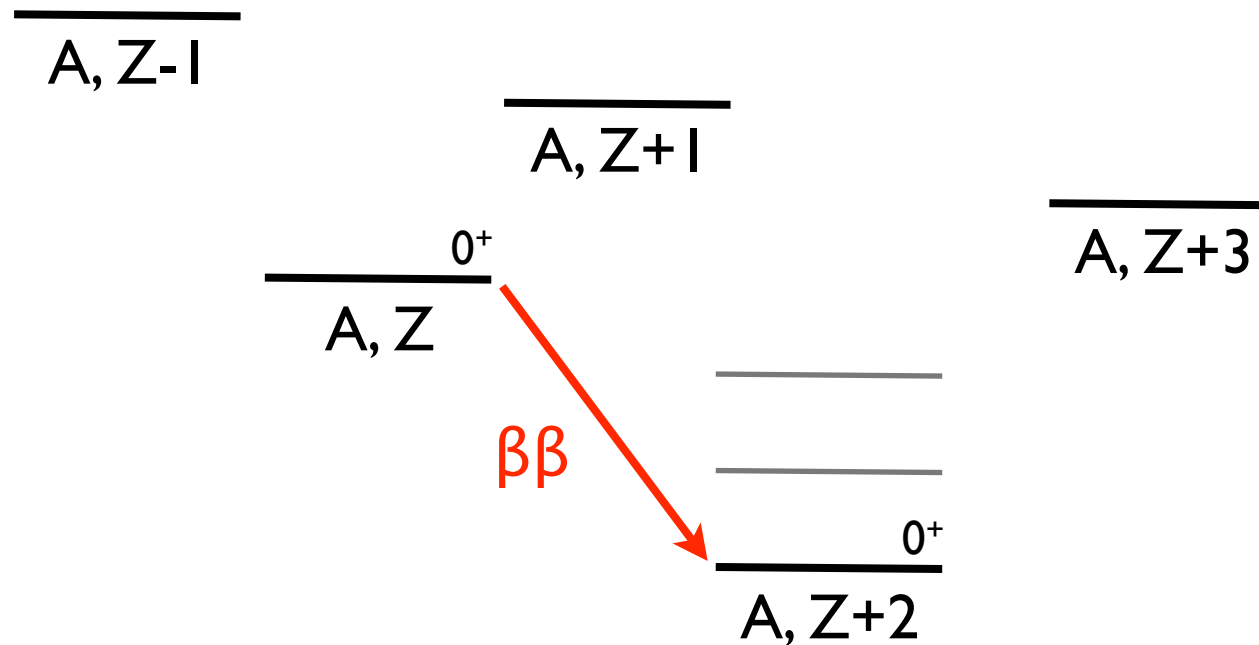
Observables

Rates

Extracting Physics

Double-Beta Decay

In a number of even-even nuclei, β -decay is energetically forbidden or strongly disfavored, while double-beta decay, from a nucleus of (A, Z) to $(A, Z+2)$, is allowed.



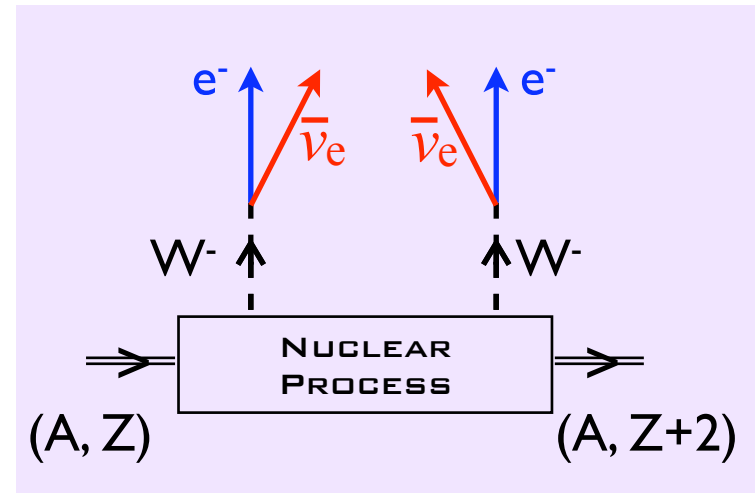
^{48}Ca , ^{76}Ge , ^{82}Se , ^{96}Zr , ^{100}Mo , ^{116}Cd , ^{128}Te , ^{130}Te , ^{136}Xe , ^{150}Nd

Double-Beta Decay Modes

2ν double-beta decay ($2\nu\beta\beta$): Nucleus $(A, Z) \rightarrow$ Nucleus $(A, Z+2) + e^- + \bar{\nu}_e + e^- + \bar{\nu}_e$



Allowed second-order
weak process
Maria Goeppert-Mayer
(1935)



0ν double-beta decay ($0\nu\beta\beta$): Nucleus $(A, Z) \rightarrow$ Nucleus $(A, Z+2) + e^- + e^-$



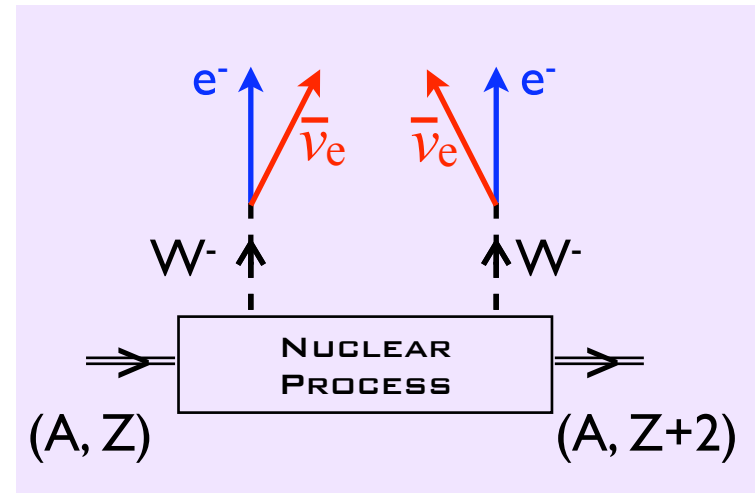
Ettore Majorana (1937)
realized symmetry properties
of Dirac's theory allowed the
possibility for electrically
neutral spin-1/2 fermions to
be their own anti-particle

Double-Beta Decay Modes

2ν double-beta decay ($2\nu\beta\beta$): Nucleus $(A, Z) \rightarrow$ Nucleus $(A, Z+2) + e^- + \bar{\nu}_e + e^- + \bar{\nu}_e$



Allowed second-order
weak process
Maria Goeppert-Mayer
(1935)

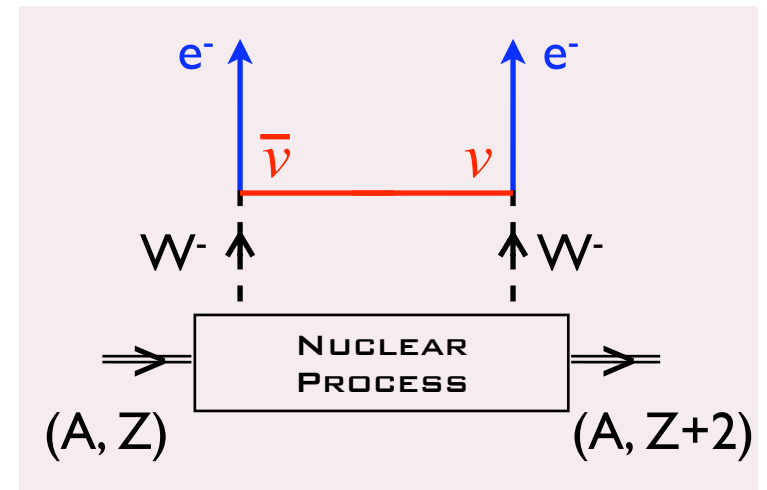


0ν double-beta decay ($0\nu\beta\beta$): Nucleus $(A, Z) \rightarrow$ Nucleus $(A, Z+2) + e^- + e^-$

Racah (1937),
Furry (1938)

$$n \rightarrow p + e^- + \bar{\nu}$$

$$\nu + n \rightarrow p + e^-$$



$\beta\beta$ Observable - electrons

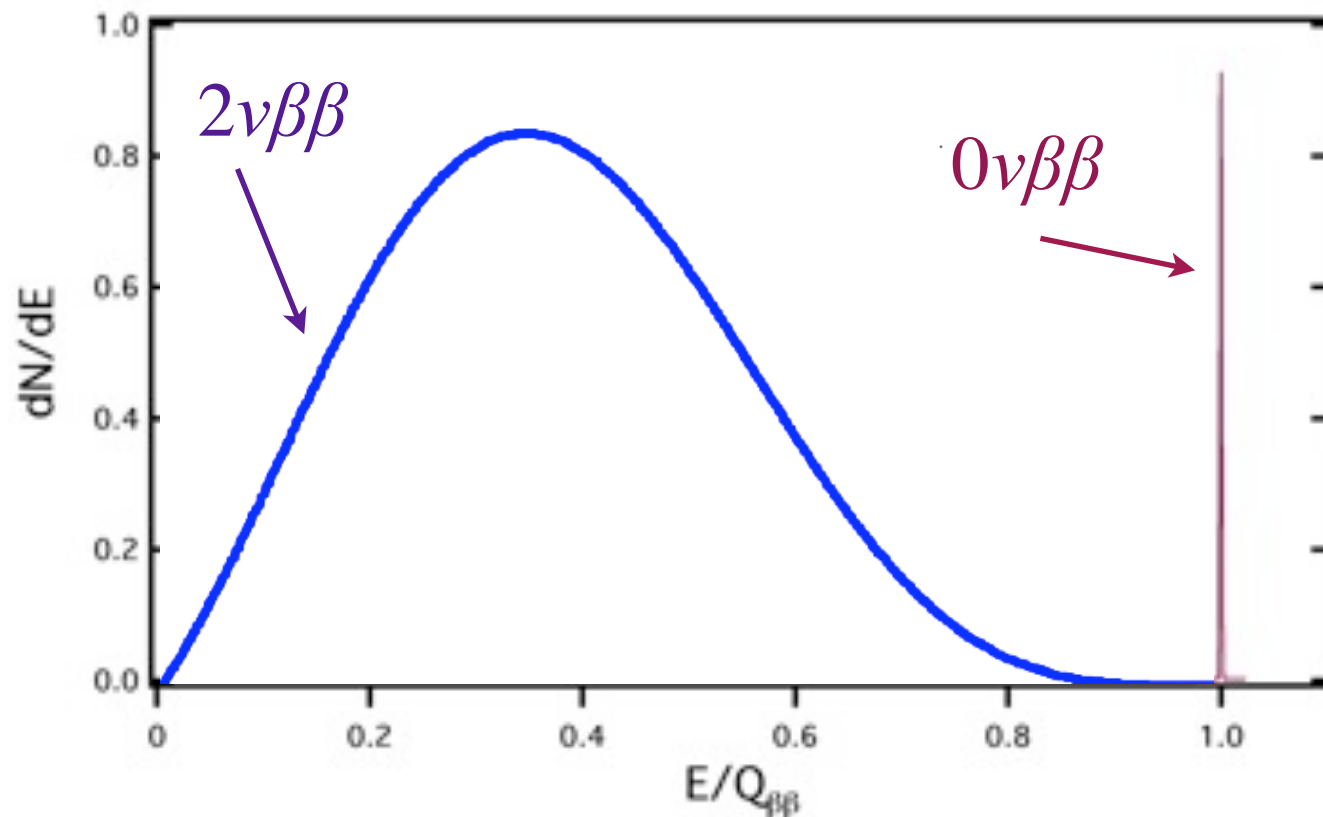
2ν double-beta decay ($2\nu\beta\beta$)

4-body decay (plus recoil nucleus)
- observe the two electrons

0ν double-beta decay ($0\nu\beta\beta$)

2-body decay (plus recoil nucleus)
- observe the two electrons

Sum of electrons energy spectrum



Early Estimates of $\beta\beta$ Decay Rates

2ν double-beta decay ($2\nu\beta\beta$)

Maria Goeppert-Mayer (1935)
using Fermi Theory

$$\left[T_{1/2}^{2\nu\beta\beta}\right]^{-1} \propto \text{Phase Space (4-body)} \propto Q^{10-12}$$

$$T_{1/2}^{2\nu\beta\beta} \approx 10^{25} \text{ years}$$

0ν double-beta decay ($0\nu\beta\beta$)

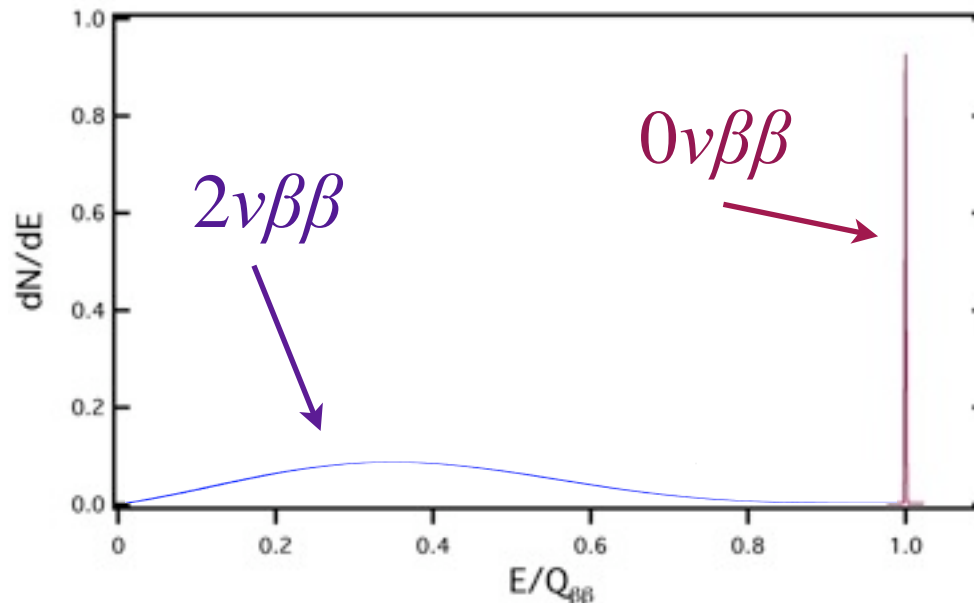
Furry (1939), assuming Parity
conserved, so no preferential handedness

$$\left[T_{1/2}^{0\nu\beta\beta}\right]^{-1} \propto \text{Phase Space (2-body)} \propto (E_\nu Q)^5$$

$$T_{1/2}^{0\nu\beta\beta} \approx 10^{19} \text{ years}$$

$0\nu\beta\beta$ mode highly favored over $2\nu\beta\beta$

If observe
 $2\nu\beta\beta \Rightarrow$
neutrinos are
Dirac

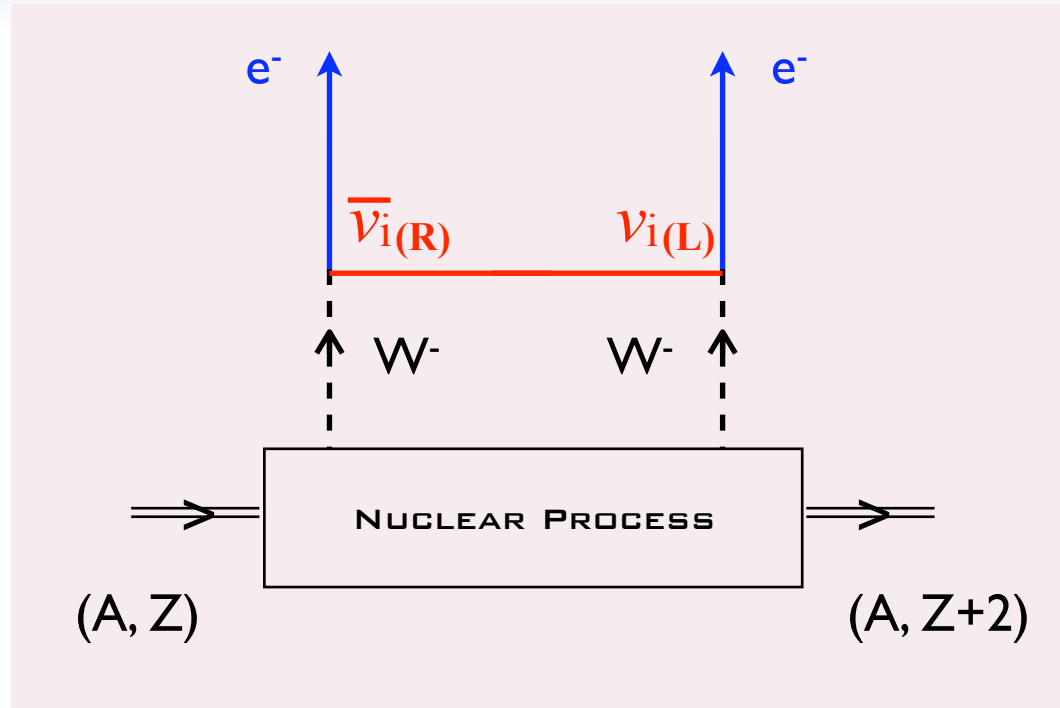


If observe
 $0\nu\beta\beta \Rightarrow$
neutrinos are
Majorana

1956-1958 Parity Violation, ν Helicity and $0\nu\beta\beta$ -decay

$$n \rightarrow p + e^- + \bar{\nu}_{(R)}$$

$$\nu_{(L)} + n \rightarrow p + e^-$$



$0\nu\beta\beta$ requires in addition to Majorana neutrinos and a lepton violating interaction, a mechanism to “flip” from right handed anti-neutrino to left-handed neutrino

The transition can occur if neutrinos have mass
“wrong-handed” helicity admixture $\sim m_i/E_{\nu_i}$

$0\nu\beta\beta$ mode highly suppressed compared to $2\nu\beta\beta$

$0\nu\beta\beta$ Decay - Current Understanding

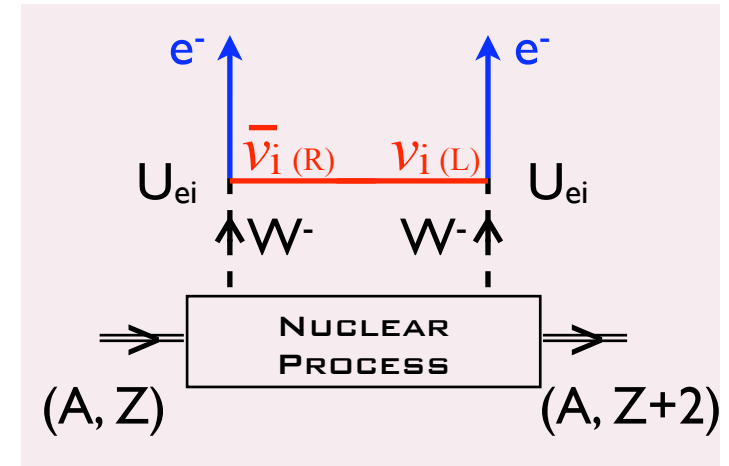
$0\nu\beta\beta$ Requires:

- neutrino to have non-zero mass
 - “wrong-handed” helicity admixture $\sim m_i/E_{\nu_i}$

Any process that allows $0\nu\beta\beta$ to occur requires Majorana neutrinos with non-zero mass.

Schechter and Valle, 1982

- Lepton number violation
 - No experimental evidence that Lepton number must be conserved
(i.e. general SM principles, such as electroweak-isospin conservation and renormalizability)



If $0\nu\beta\beta$ decay is observed $\Rightarrow$ neutrinos are Majorana particles
lepton number is violated

$\beta\beta$ Decay Rates (measured)

2ν double-beta decay ($2\nu\beta\beta$)

Observed for
 ^{48}Ca , ^{76}Ge , ^{82}Se , ^{96}Zr , ^{100}Mo ,
 ^{116}Cd , ^{128}Te , ^{130}Te , ^{150}Nd

$$T_{1/2}^{2\nu\beta\beta} \approx 10^{19} - 10^{21} \text{ years}$$

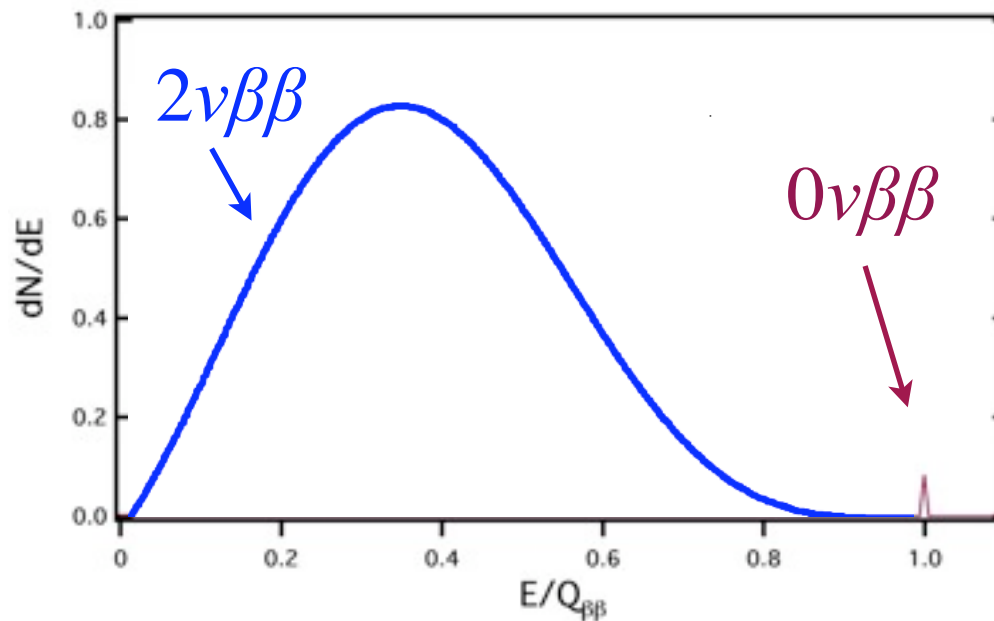
0ν double-beta decay ($0\nu\beta\beta$)

A claimed observation in ^{76}Ge
 $T_{1/2} = (1.19 + 2.99/-0.5) \times 10^{25} \text{ y}$

Limits for
 ^{48}Ca , ^{76}Ge , ^{82}Se , ^{100}Mo , ^{116}Cd , ^{128}Te , ^{130}Te , ^{136}Xe , ^{150}Nd

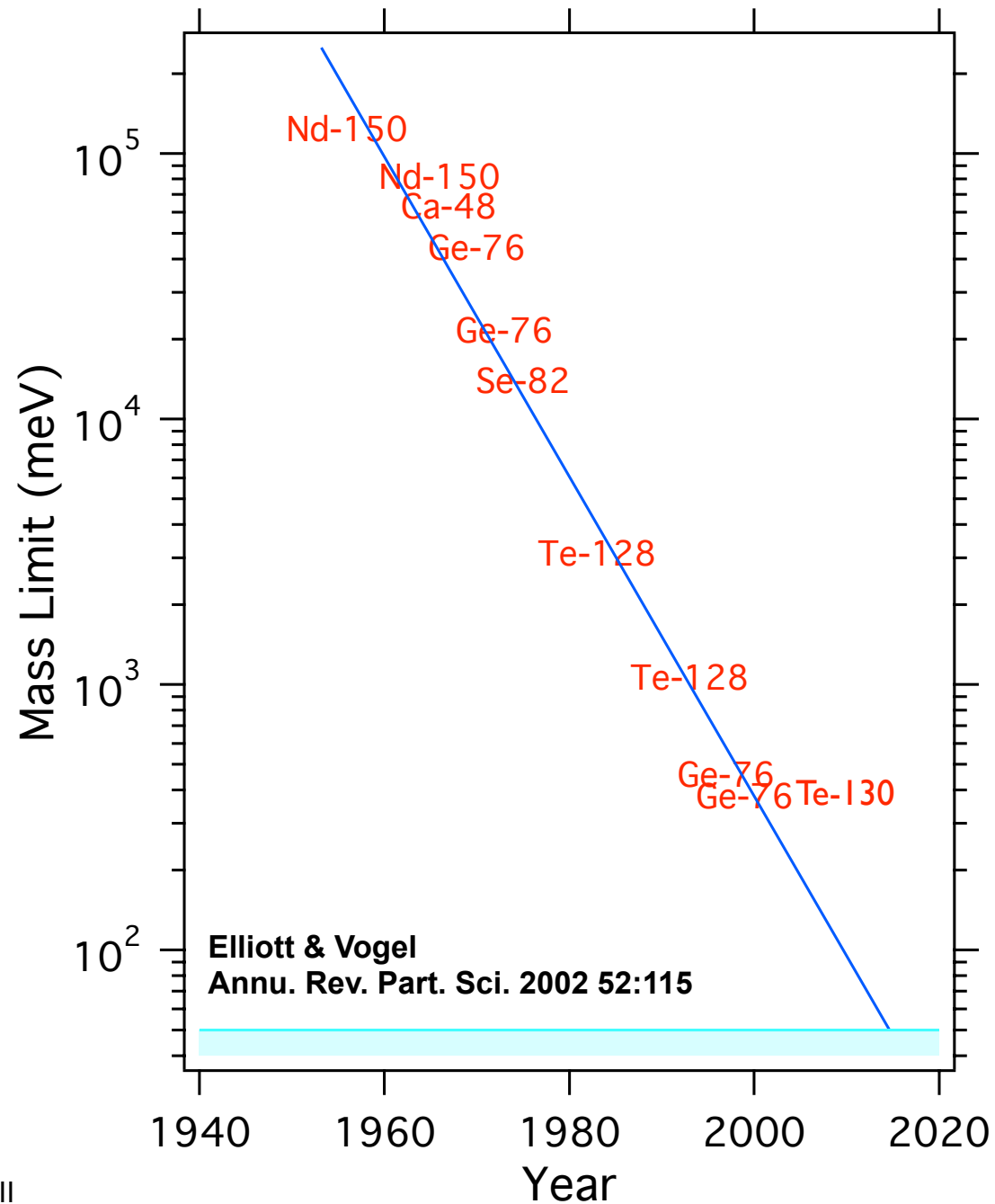
$$T_{1/2}^{0\nu\beta\beta} > 10^{25} \text{ years}$$

$0\nu\beta\beta$ mode highly suppressed compared to $2\nu\beta\beta$



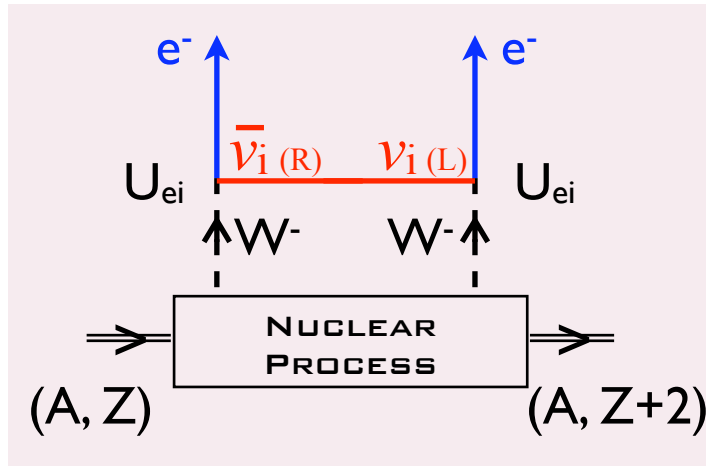
If observe
 $0\nu\beta\beta \Rightarrow$
neutrinos are
Majorana

$0\nu\beta\beta$ Progress & Best Limits



Extracting Additional Physics from $0\nu\beta\beta$

Neutrino mass & nature of lepton number violating interactions (η)



$$\left[\mathbf{T}_{1/2}^{0\nu} \right]^{-1} = G_{0\nu} \left| M_{0\nu}(\eta) \right|^2 \eta^2$$

$$\Downarrow$$

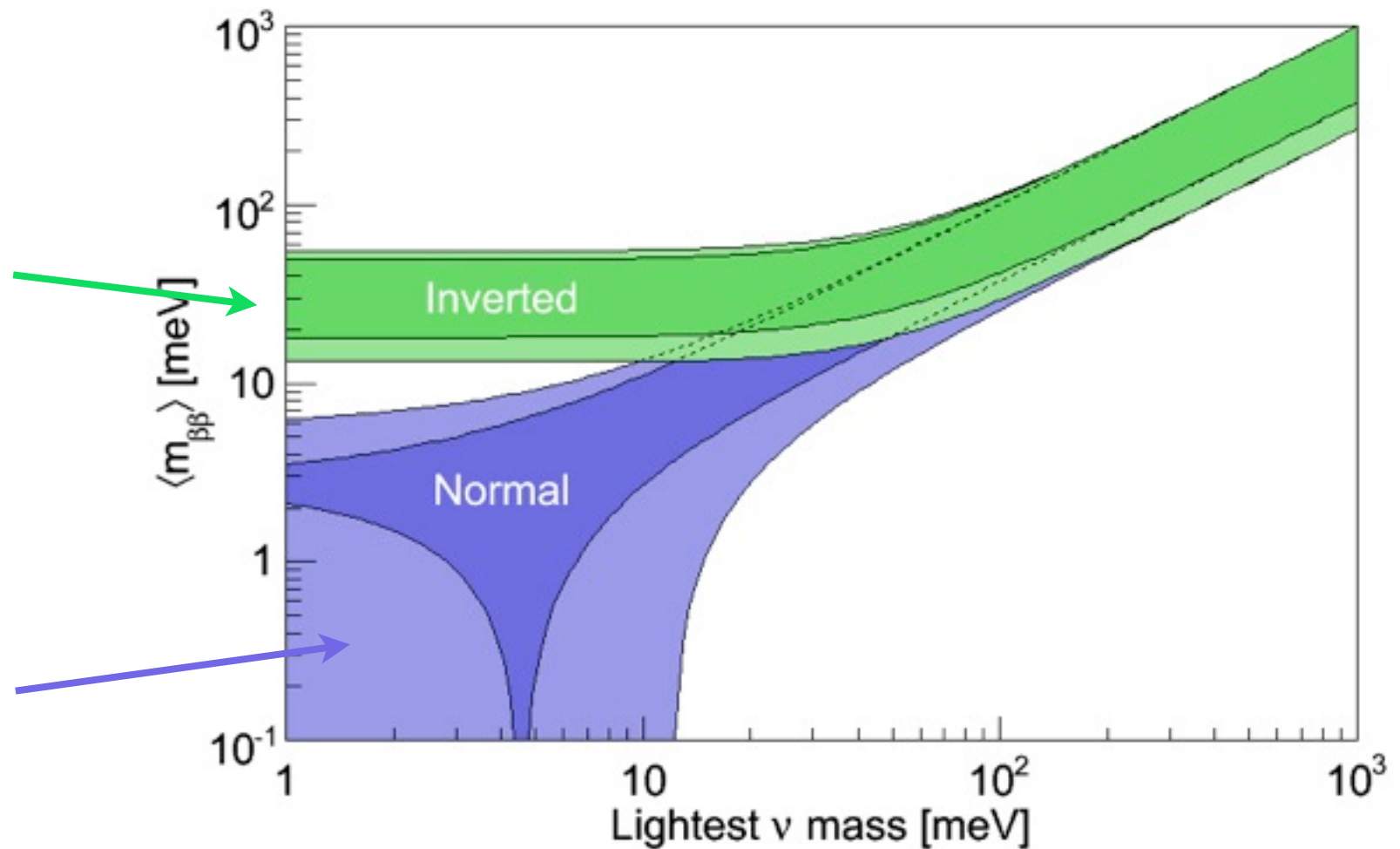
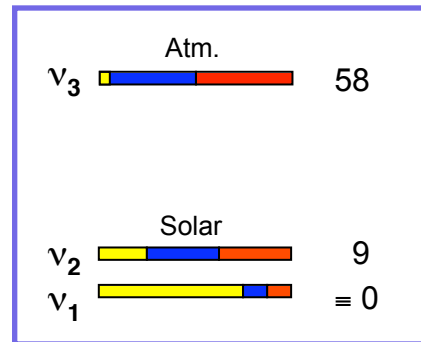
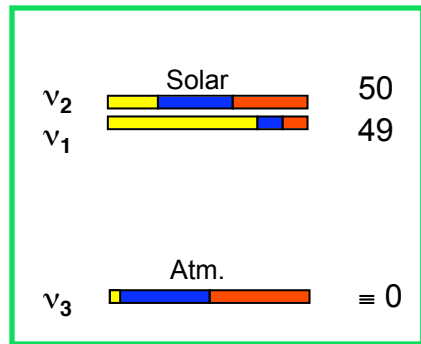
$$\left[\mathbf{T}_{1/2}^{0\nu} \right]^{-1} = G_{0\nu} \left| M_{0\nu} \right|^2 \langle m_{\beta\beta} \rangle^2$$

- Phase space, $G_{0\nu}$ is exactly calculable.
- Nuclear matrix elements (NME) via theory (Shell Model, Quasi-random phase approximation, interacting boson model).
- Effective neutrino mass, $\langle m_{\beta\beta} \rangle$, depends directly on the assumed form of lepton number violating (LNV) interactions.
 - simplest, as typically shown (e.g. above): light Majorana neutrino and “standard model” interactions (W^-)

$0\nu\beta\beta$ Decay Sensitivity to $\langle m_{\beta\beta} \rangle$

Assuming LNV mechanism is light Majorana neutrino exchange

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right| = \left| c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$

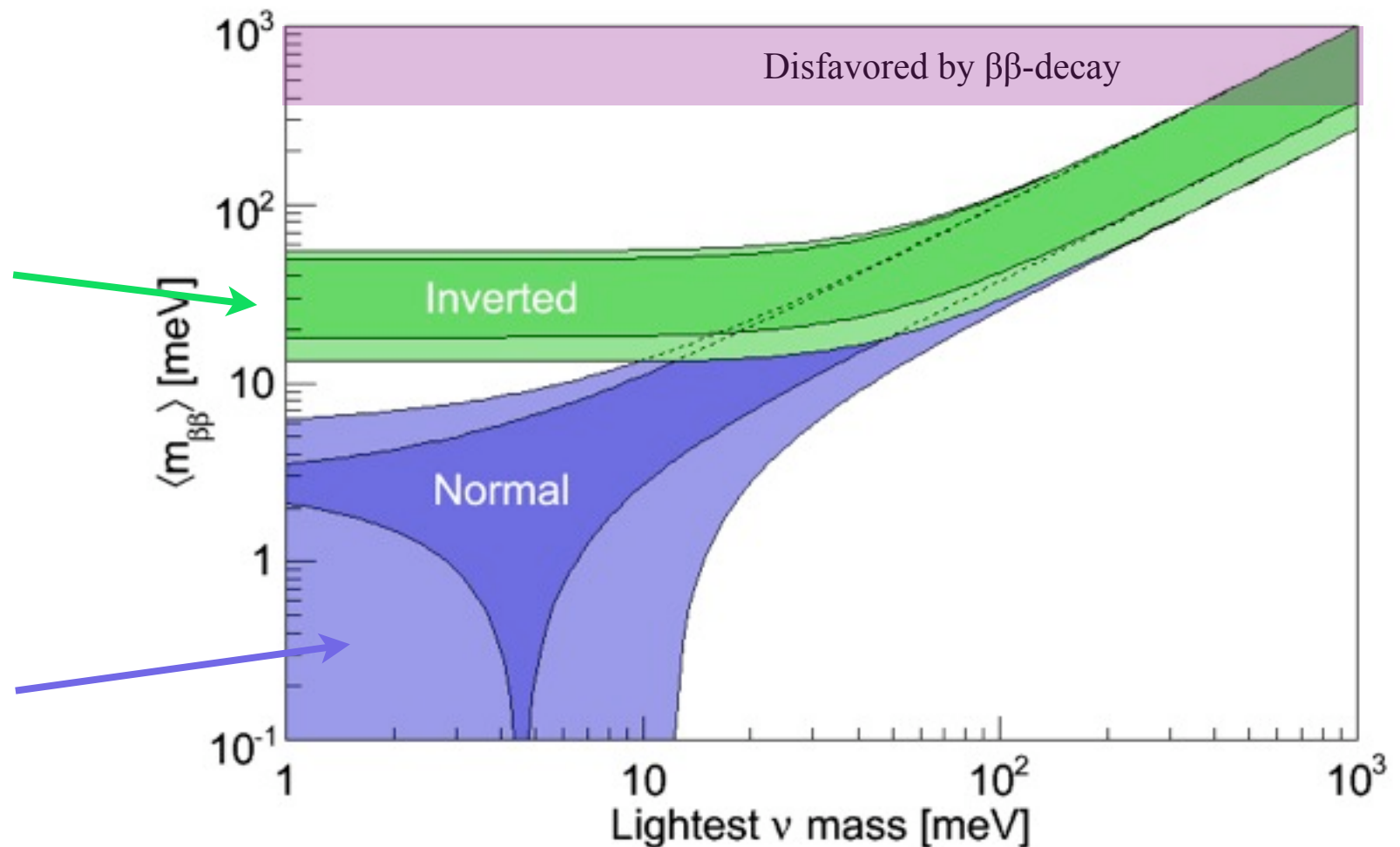
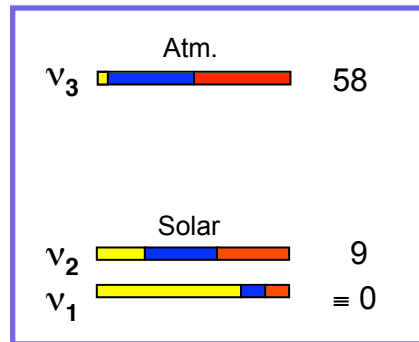
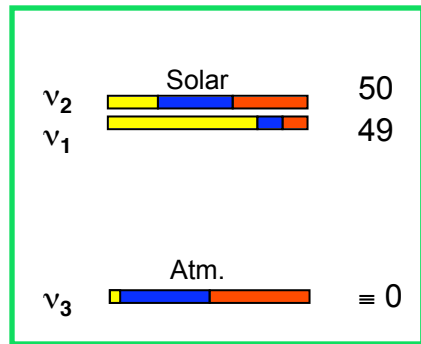


$0\nu\beta\beta$ Decay Sensitivity to $\langle m_{\beta\beta} \rangle$

Assuming LNV mechanism is light Majorana neutrino exchange

$0\nu\beta\beta$ limits for: ^{48}Ca , ^{76}Ge , ^{82}Se , ^{100}Mo , ^{116}Cd , ^{128}Te , ^{130}Te , ^{136}Xe , ^{150}Nd

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right| = \left| c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$

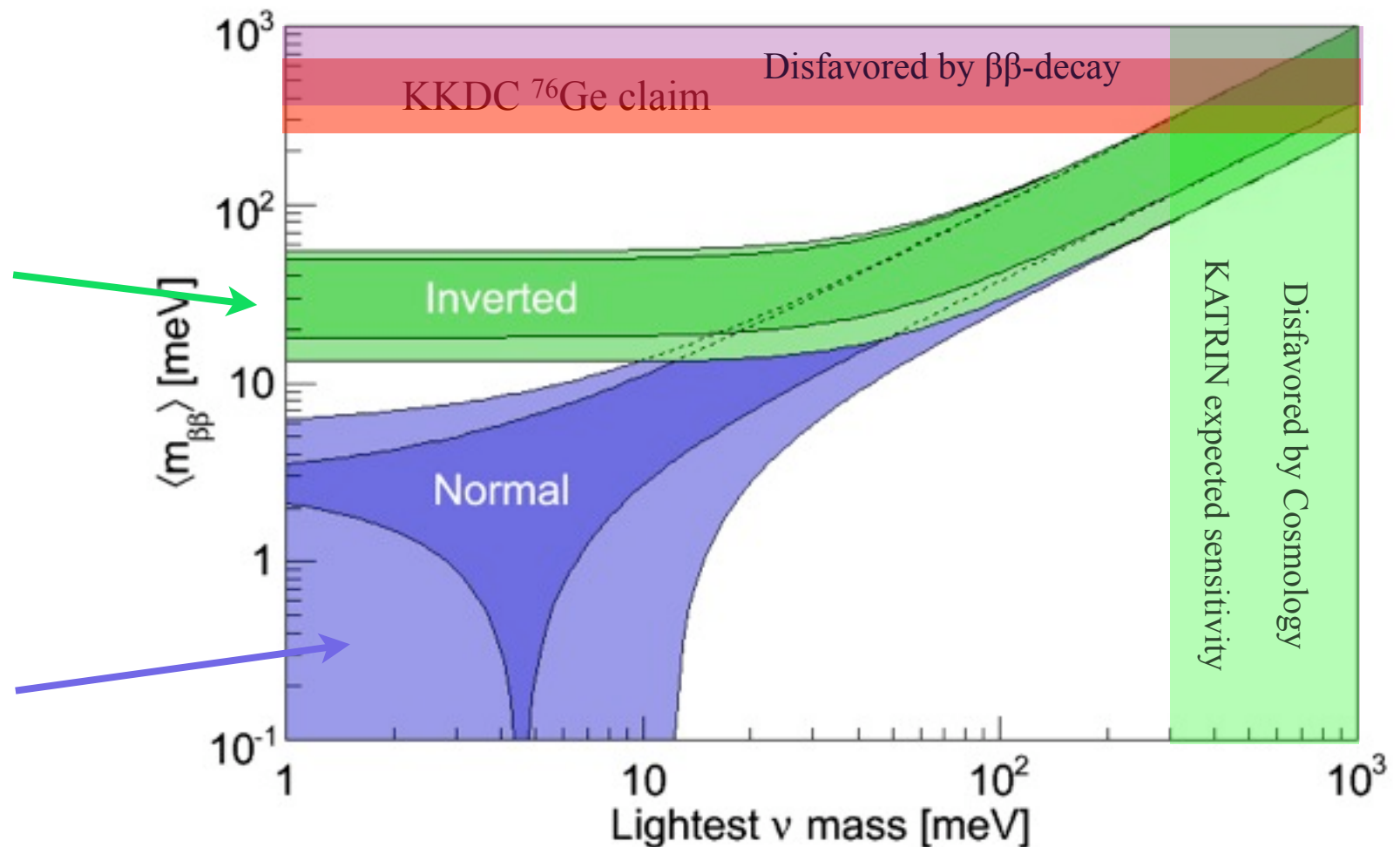
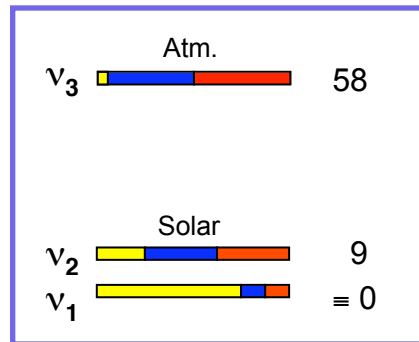
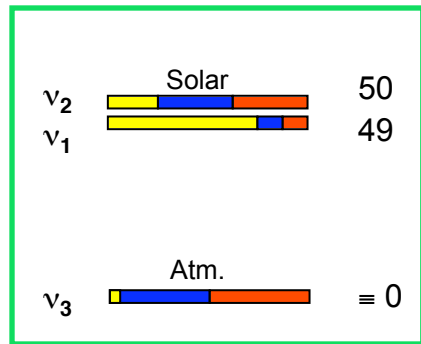


$0\nu\beta\beta$ Decay Sensitivity to $\langle m_{\beta\beta} \rangle$

Assuming LNV mechanism is light Majorana neutrino exchange

$0\nu\beta\beta$ limits for: ^{48}Ca , ^{76}Ge , ^{82}Se , ^{100}Mo , ^{116}Cd , ^{128}Te , ^{130}Te , ^{136}Xe , ^{150}Nd

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right| = \left| c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$



Searches for $0\nu\beta\beta$

$0\nu\beta\beta$ decay probes fundamental questions:

- Neutrino properties — the only practical technique to determine if neutrinos are their own anti-particles — Majorana particles.
- Lepton number violation (LNV) — might Leptogenesis be the explanation for the observed matter - antimatter asymmetry?

The observation of $0\nu\beta\beta$ would demonstrate Majorana nature of ν and LNV

If $0\nu\beta\beta$ is observed:

- Provides a promising laboratory method for determining the absolute neutrino mass scale that is complementary to other neutrino mass measurement techniques, if calculated NME are reliable.
- Measurements in a series of different isotopes can potentially help reveal the nature of the lepton number violating process(es).

Extraction of ν mass and understanding the LNV process(es) requires significant reliance on both nuclear and particle physics. (See Vogel)

$0\nu\beta\beta$ Decay II

Experimental Considerations & Sensitivity

Basics

Resolution

Backgrounds Overview

Ideal experiment

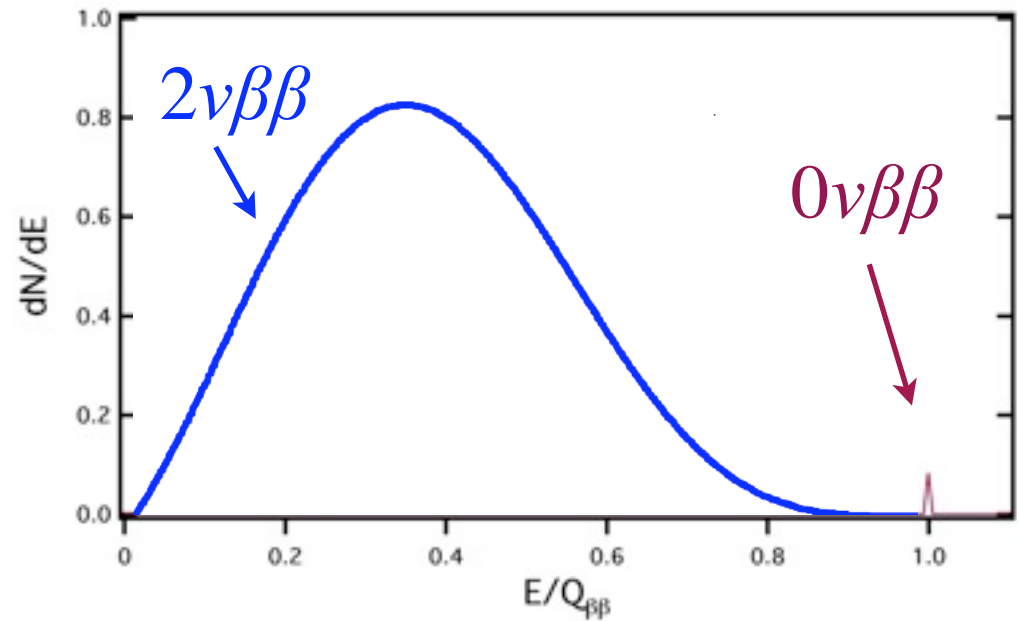
Basic Considerations

Extremely rare process
 $0\nu\beta\beta$ $T_{1/2} \sim 10^{26} - 10^{27}$ years

Large, highly efficient source mass

Best possible energy resolution

Extremely low (near-zero) backgrounds in the $0\nu\beta\beta$ peak region



Current sensitivity to $0\nu\beta\beta$ -decay

Most sensitive experiments to date (using ^{76}Ge and ^{130}Te) have attained $T_{1/2} > 10^{24} - 10^{25}$ years.

If zero backgrounds, $[T_{1/2}]^{-1} \propto M \cdot t_{\text{exp}}$

Typical masses $\sim 1\text{-}10$ kg

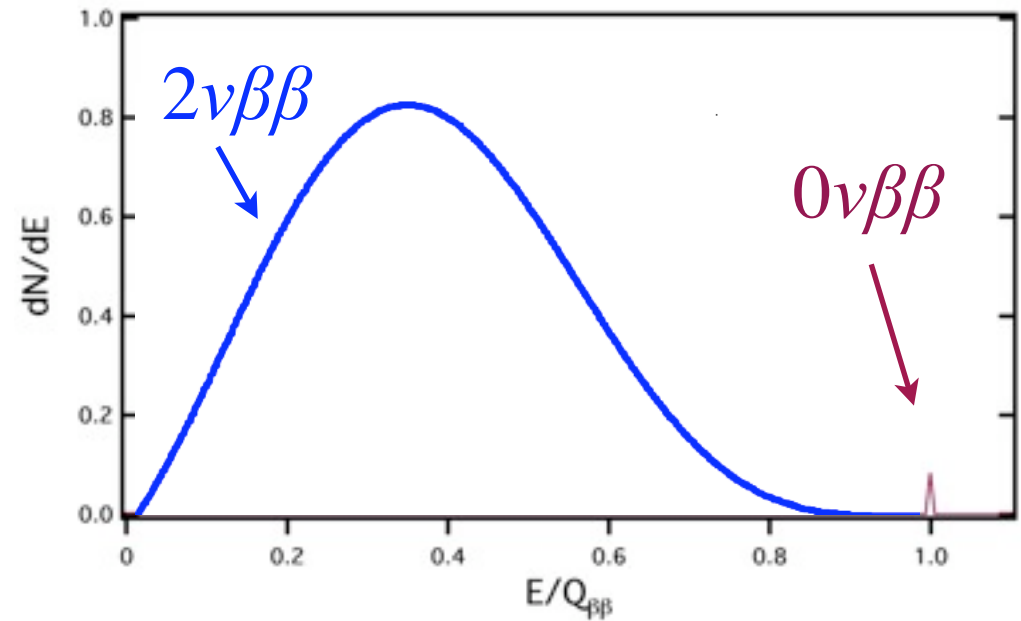
Typical exposure times of 5-8 years

Question 1: What's needed to reach sensitivities of $T_{1/2}$ on the order of $10^{25} - 10^{27}$ y?

Question 2: What would convince you that $0\nu\beta\beta$ has been discovered?

Basic Considerations

Extremely rare process
 $0\nu\beta\beta$ $T_{1/2} \sim 10^{26} - 10^{27}$ years



Large, highly efficient source mass
- detector as source

Best possible energy resolution

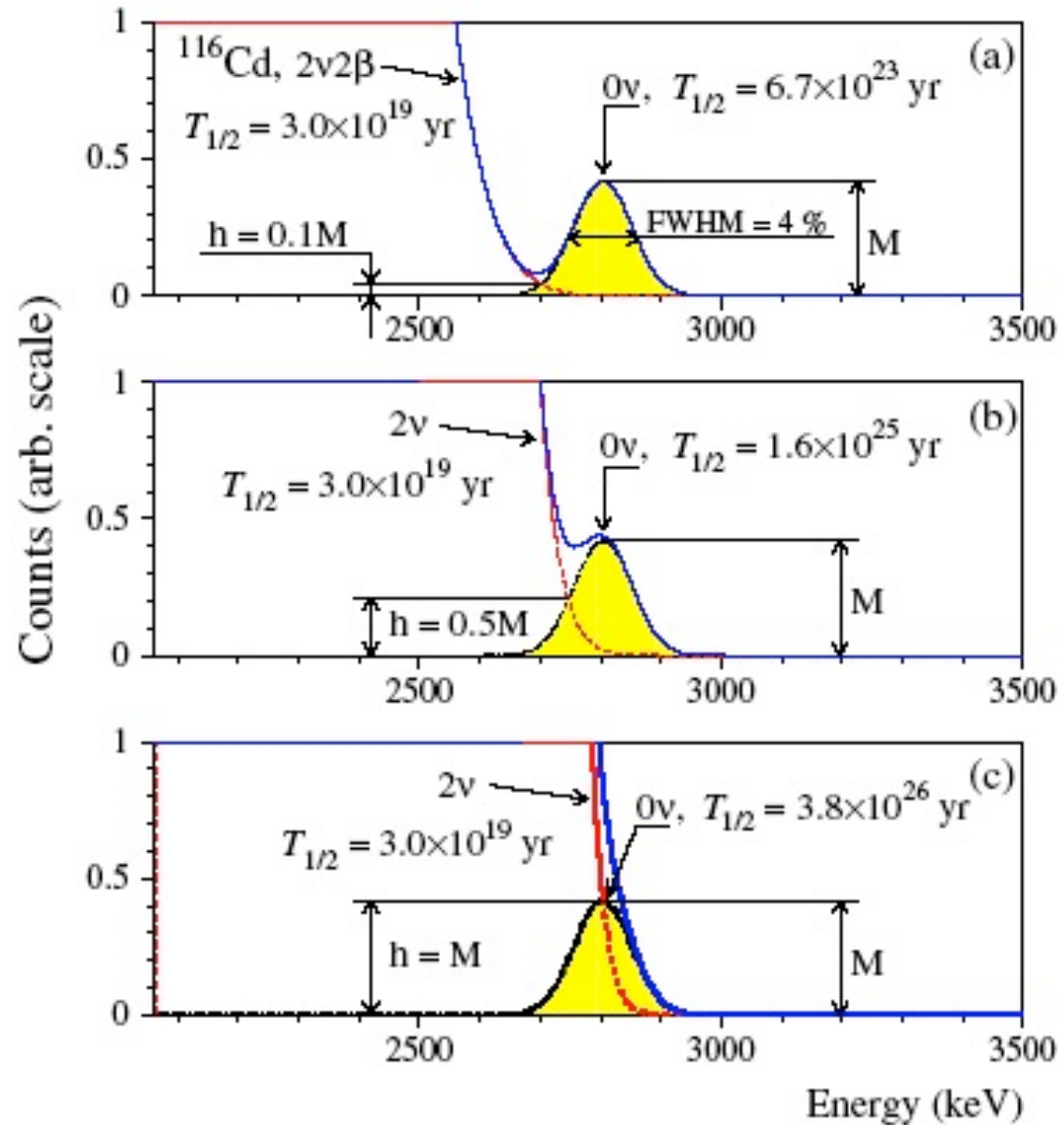
- minimize $0\nu\beta\beta$ peak ROI to maximize S/B
- separate from $0\nu\beta\beta$ from irreducible $2\nu\beta\beta$ ($\sim T_{1/2} \sim 10^{19} - 10^{21}$ years)

Extremely low (near-zero) backgrounds in the $0\nu\beta\beta$ peak region

- requires ultra-clean radiopure materials
- the ability to discriminate signal from backgrounds

Energy Resolution & Sensitivity to $0\nu\beta\beta$

For ideal case,
with only
 $2\nu\beta\beta$
background



From Zdesenko, Danevich, Tretyak, J. Phys. G 30 (2004) 971

Backgrounds & Sensitivity to $0\nu\beta\beta$

Next generation experiments are striving for backgrounds in the $0\nu\beta\beta$ region of **counts/tonne-year**.

Requires materials with sub $\mu\text{Bq/kg}$ level radioimpurities.

Very difficult to achieve this sensitivity with direct radioassays

Shielding from cosmogenic activation.

“New background regimes” -- background sources that could previously be ignored

e.g. : very weak (n,n',gamma) lines

Each experiment's susceptibility to backgrounds depends on a number of factors:

Detection technique (Solid state, TPC, bolometer, ...)

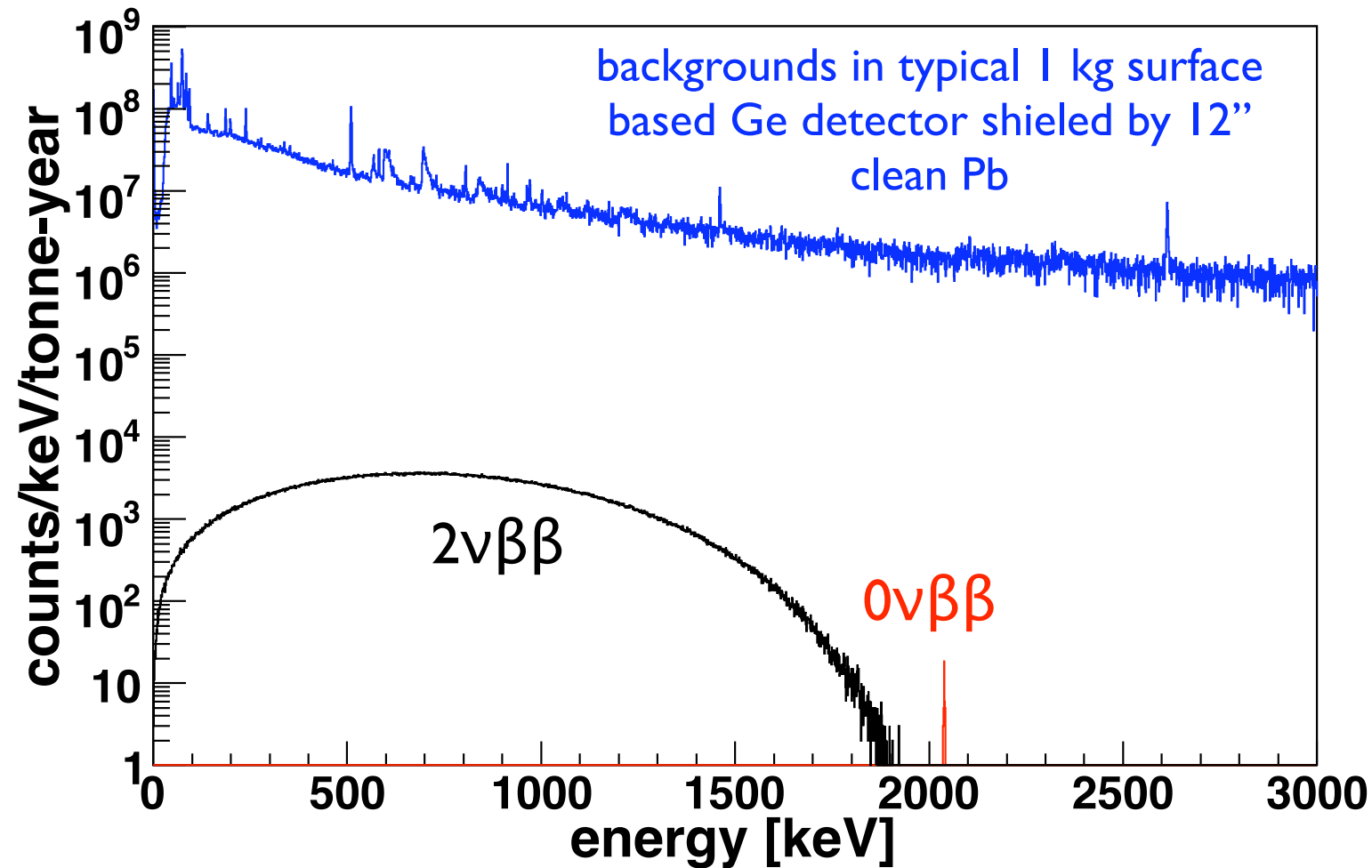
Detector resolution

Detector response function to backgrounds (γ , α , β , neutron, ...)

Construction materials and surrounding materials

Signal to background discrimination capabilities

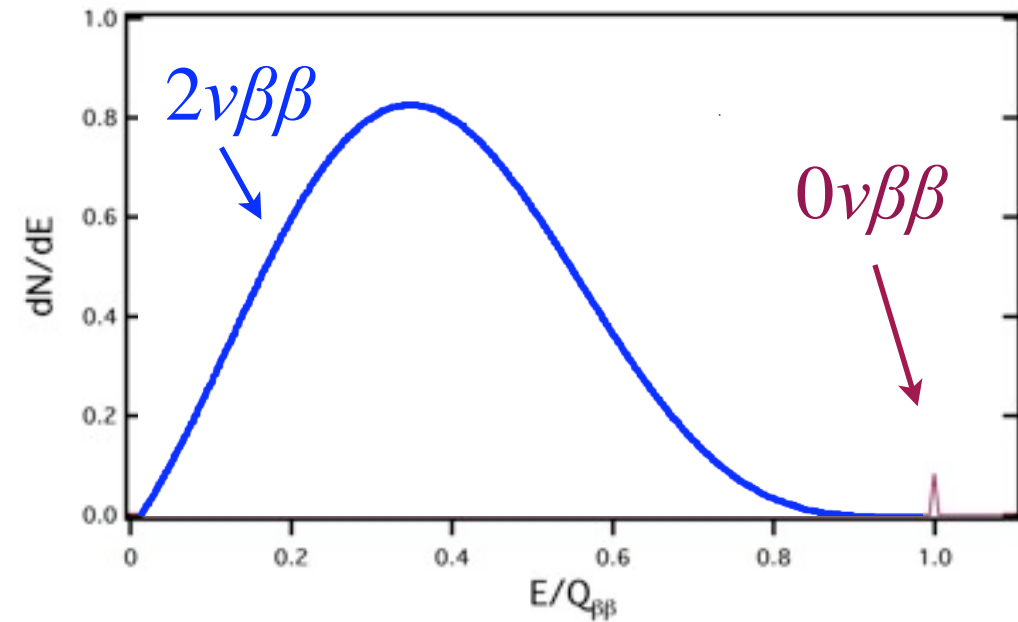
1 ct/tonne-year in context - Ge example



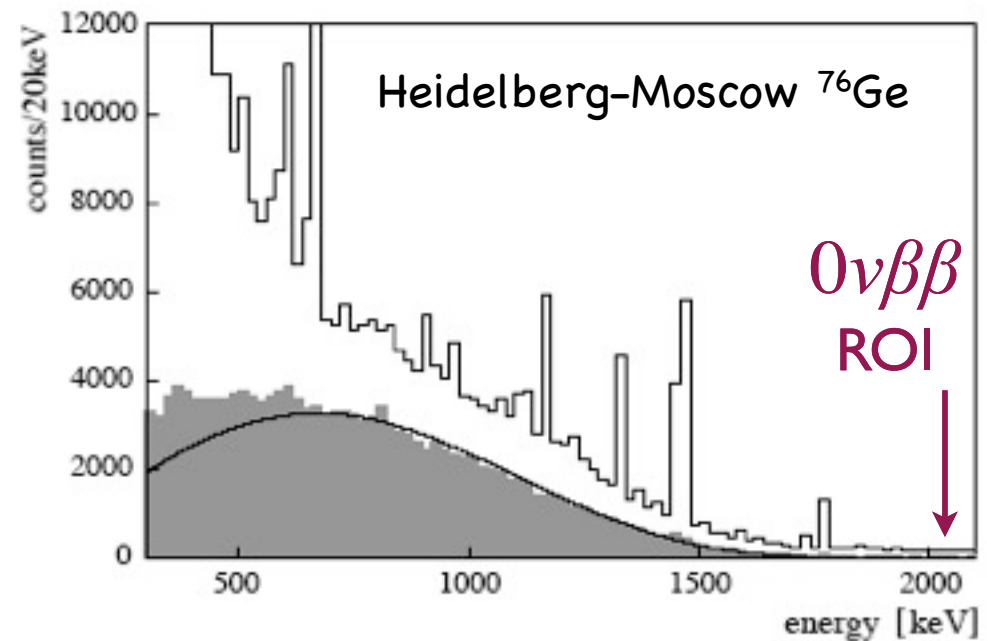
$0\nu\beta\beta$ half-life chosen to be 10x current limit

A. Schubert

Best to date - 100 cnts/tonne-year in ROI



Only $2\nu\beta\beta$

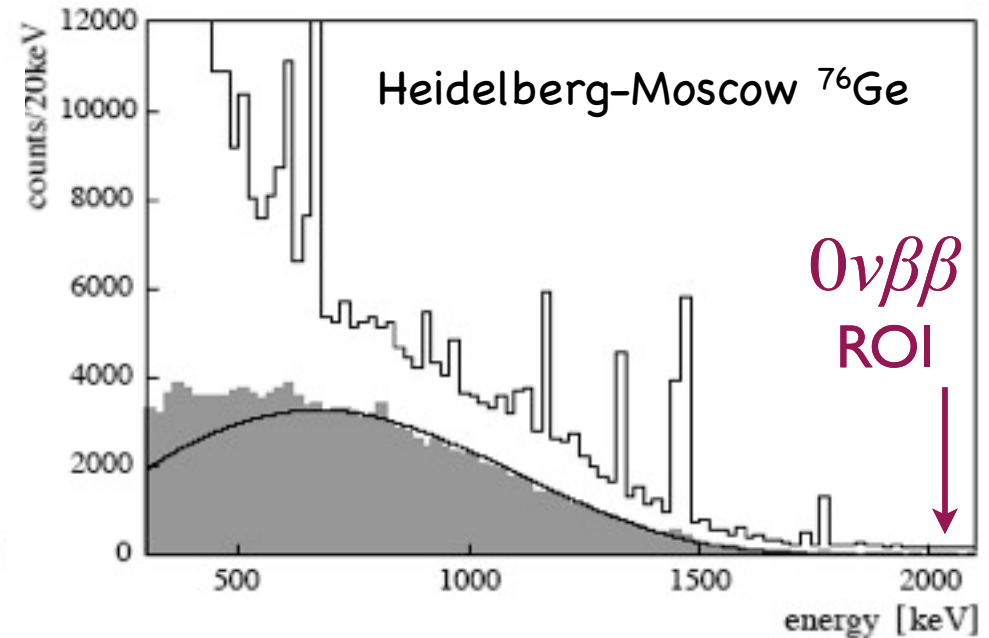
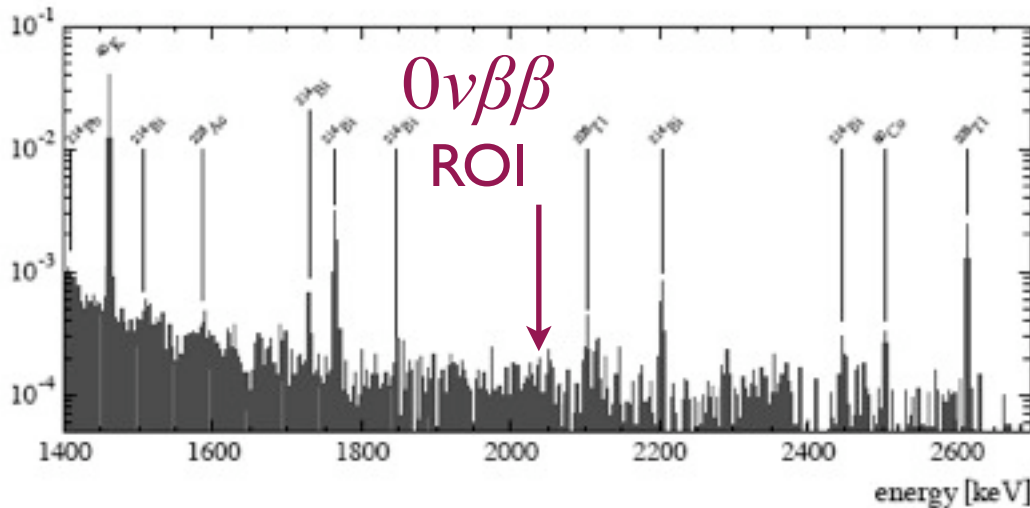
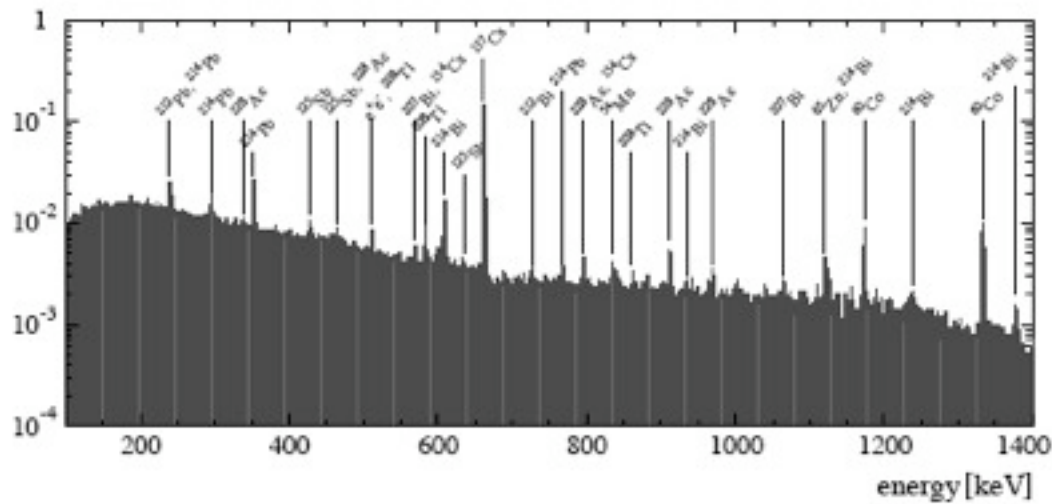


Note ^{76}Ge has best resolution : 0.16%

$2\nu\beta\beta$ (negligible)
Cosmogenic activity
Natural radioactivity
neutron induced activity

Klapdor-Kleingrothaus et al.,
Eur. Phys. J 12, 147 (2001)

Best to date - 100 cnts/tonne-year in ROI



Note ^{76}Ge has best resolution : 0.16%

$2\nu\beta\beta$ (negligible)

Cosmogenic activity

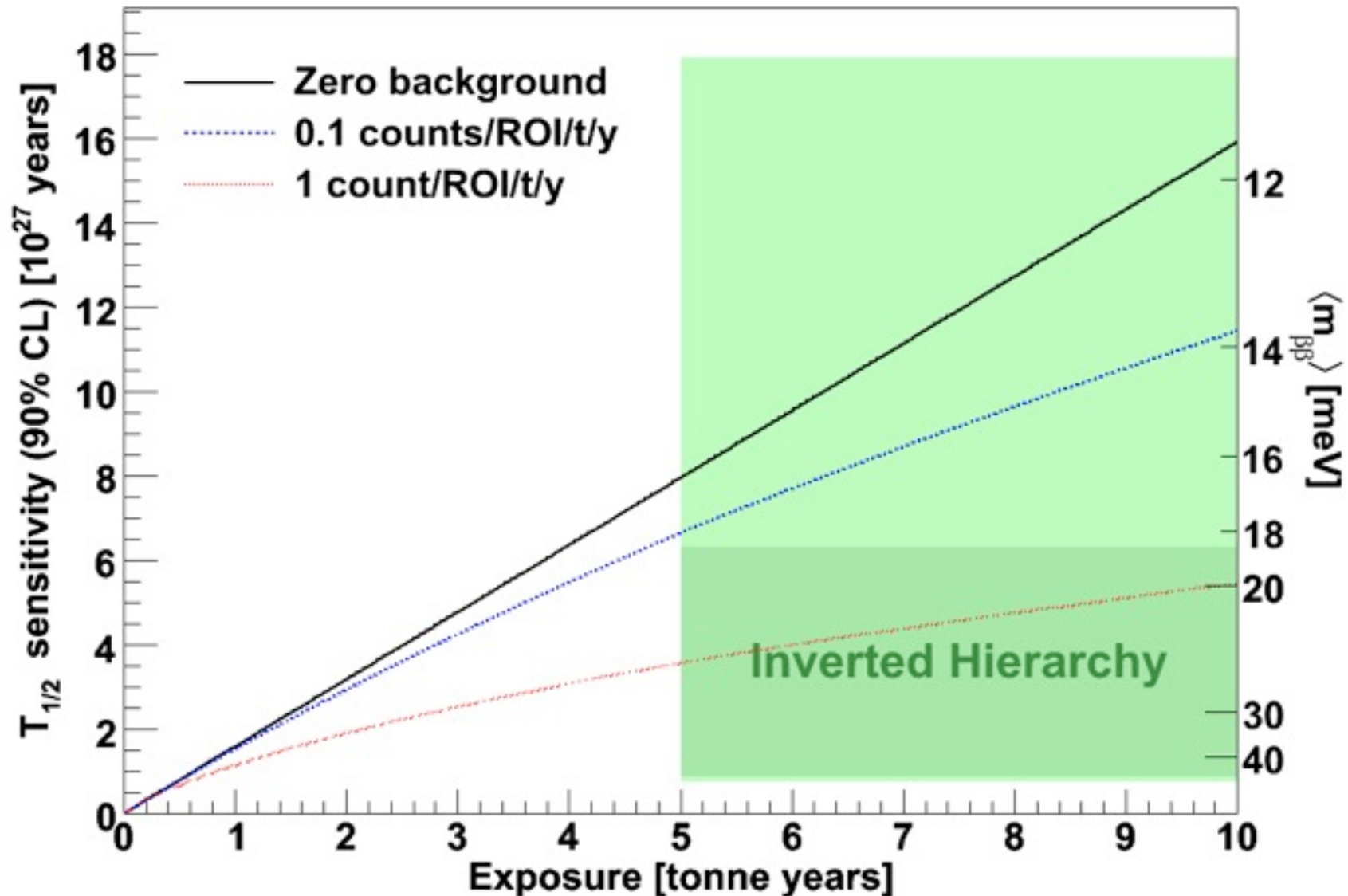
Natural radioactivity

neutron induced activity

Klapdor-Kleingrothaus et al.,
Eur. Phys. J 12, 147 (2001)

Sensitivity and backgrounds

I-tonne ^{76}Ge Example



“Ideal” Experiment

Source serves as the detector

Elemental (enriched) source to minimize active material.

Large Q value – faster $0\nu\beta\beta$ rate and also places the region of interest above many potential backgrounds.

Relatively slow $2\nu\beta\beta$ rate helps control this irreducible background.

Direct identification of the decay progeny in coincidence with the $0\nu\beta\beta$ decay eliminates all potential backgrounds except $2\nu\beta\beta$.

Full Event reconstruction, providing kinematic data such as opening angle and individual electron energy aids in the elimination of backgrounds and demonstration of signal (can possibly use $2\nu\beta\beta$)

Spatial resolution and timing information to reject background processes.

Demonstrated technology at the appropriate scale.

The nuclear theory is better understood in some isotopes than others.

The interpretation of limits or signals might be easier to interpret for some isotopes.

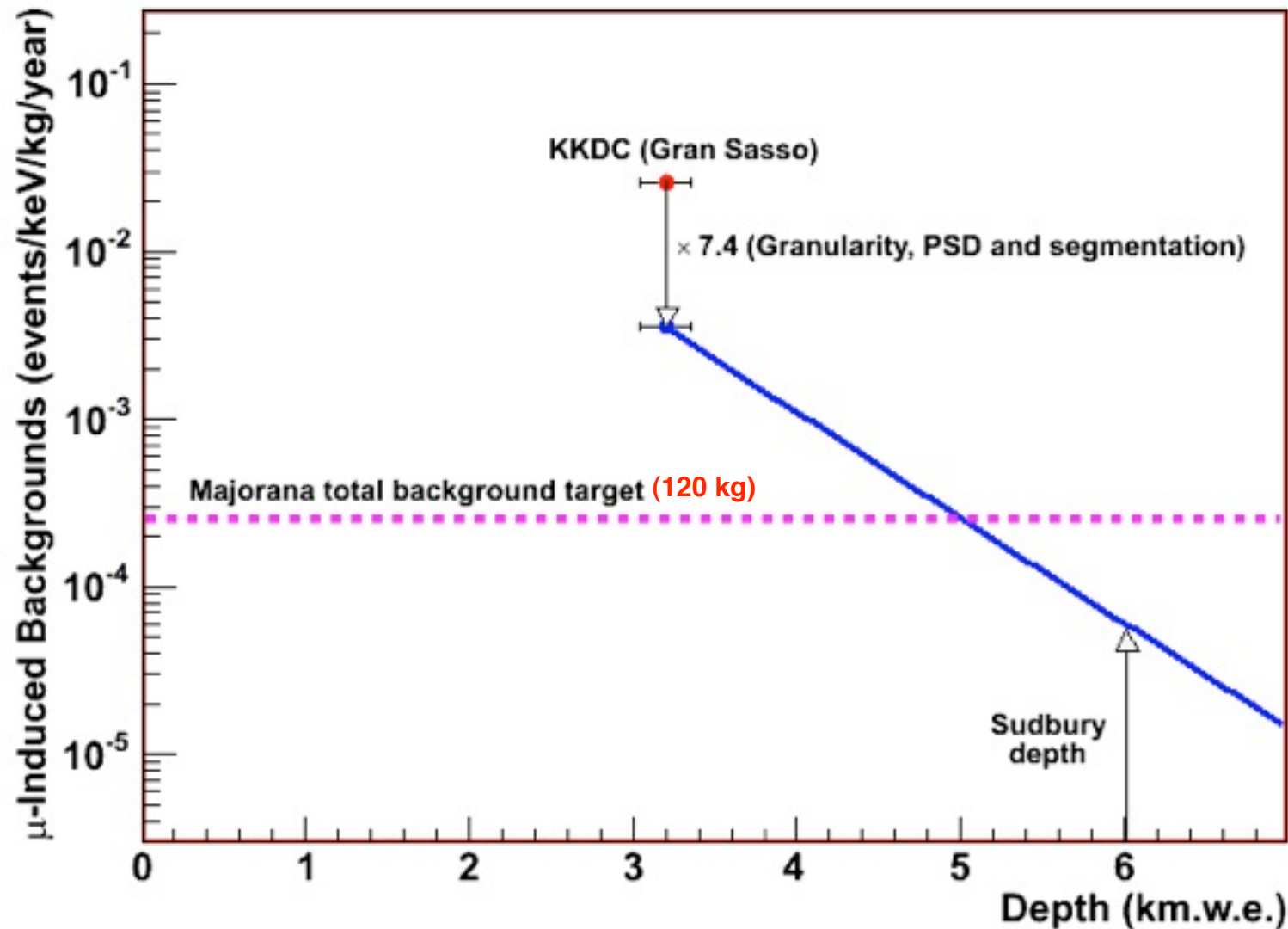
No one ideal isotope or experimental technique

$0\nu\beta\beta$ Decay II

Backgrounds & Discrimination

Cosmogenics
Natural radioactivity
Reducing Backgrounds

Cosmogenic Backgrounds (muons)



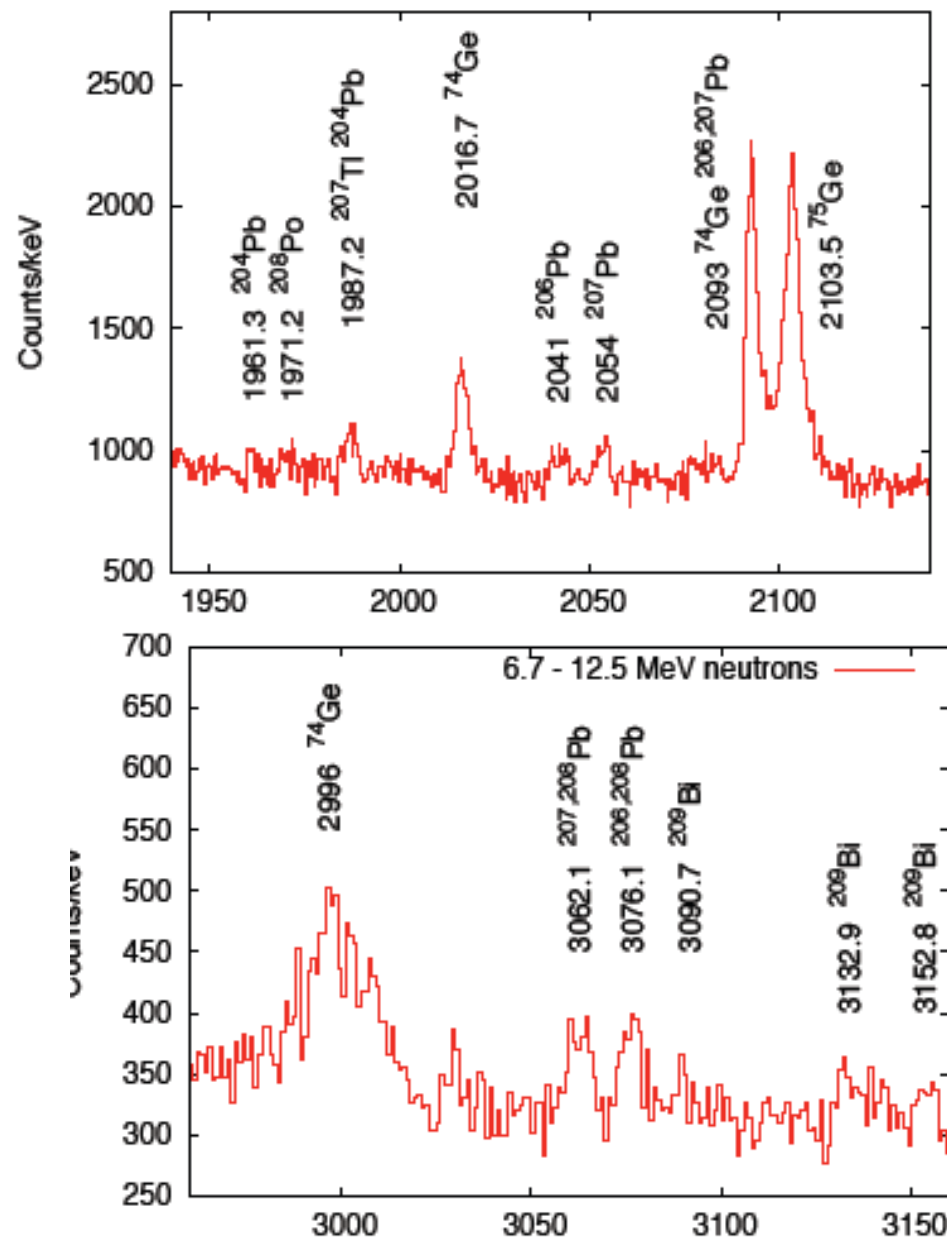
D.-M. Mei and A. Hime, Phys. Rev. D **73**, 053004 (2006).

Pb(n,n'γ) lines

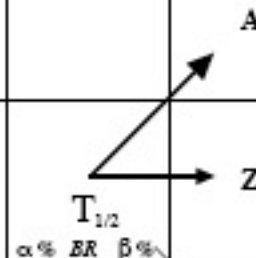
- Specific Pb gamma rays are problematic backgrounds
 - ^{206}Pb has a 2040-keV γ ray
 - ^{207}Pb has a 3062-keV γ ray
 - ^{208}Pb has a 3060-keV γ ray
- Neutron interactions in Pb can excite these levels
- The DEP of the ~ 3062 keV γ ray is a single site energy deposit at $\beta\beta$ Q-value

V. E. Guiseppe et al.

arXiv:0809.5074



Uranium (^{238}U) Decay Chain

URANIUM - RADIUM						A = 4n + 2						Th 234 Q β 0.199 70.3% 0.107 19.2% 0.106 7.6%	Th 234 24.10 d	← 4.198 79.0% 4.151 20.9%	U 238 4.468·10 ⁹ a
			Bi 214 Q β 3.272 18.2 % 1.894 7.43 % 1.542 17.8 % 1.508 17.02 % 1.425 8.18 % 1.068 5.72 %									Pa 234 Q β 0.642 19.4 % 0.502 7.0 % 0.4721 12.4 % 0.4716 33 % 0.413 8 %	Pa 234* 1.17 m 6.7 h	↘ 2.269 98.2 %	Pa234m Q β
	Pb 214 Q β 1.024 6.3% 0.729 42.2% 0.672 48.9%	Pb 214 26.8(9) m	← 6.002 99.999% 5.181 0.001% 99.980% 0.020%	Po 218 3.10(1) m	← 5.490 99.92 % 4.987 0.078%	Rn 222 3.8235(3) d	← 4.784 94.45% 4.601 5.55%	Ra 226 1600(1) a	← 4.687 76.3% 4.621 23.4%	Th 230 7.538·10 ⁴ a	← 4.775 71.38% 4.722 28.42%	U 234 2.455·10 ⁵ a			
	Tl 210 1.30(3) m	5.516 39.2% ← 5.452 53.9% 5.273 5.8%	Bi 214 19.9(4) m	↗ 6.693 90 % 6.653 6.4%	At 218 1.5 s										
	Tl 210 Q β 4.391 20% 4.210 30% 2.419 10% 2.029 10% 1.864 24% 1.609 7%	Pb 210 22.3(2) a	← 7.687 99.999% 6.902 0.010%	Po 214 164.3(20) μ s											
		Pb 210 Q β 0.064 16% 0.017 84%	Bi 210 5.013 d	Bi 210 Q β 1.162						α ← E α MeV RI%					
		Pb 206 stable	← 5.304 100 % 4.516 0.001%	Po 210 138.376 d									β Q β MeV RI%		

Thorium (^{232}Th) Decay Chain

THORIUM $A = 4n$						Ra 228 Q_{β} 0.040 10% 0.039 40% 0.026 20% 0.013 30%		Ra 228 5.75 a		Th 232 $1.405 \cdot 10^{10}$ a		
								Ac 228 Q_{β} 2.069 8% 1.731 12% 1.158 30% 1.004 6%		Ac 228 6.15 h		
Pb 212 Q_{β} 0.574 12.3 % 0.335 82.5 % 0.159 5.17%	Pb 212 10.64(1) h	← 6.778 99.998% 5.985 0.002%	Po 216 145(2) ms	← 6.288 99.886% 5.747 0.114%	Rn 220 55.6(1) s	← 5.685 94.92% 5.449 5.06%	Ra 224 3.66(4) d	← 5.423 72.2% 5.340 27.2%	Th 228 1.9116(16) a			
Tl 208 3.053(4) m	← 6.090 27.12% 6.051 69.91%	Bi 212 60.55(6) m	Q_{β} 2.254 55.46% 1.527 4.36%									
Tl 208 Q_{β} 1.803 48.7% 1.526 21.8% 1.293 24.5%	Pb 208 stable	← 8.784	Po 212 299(2) ns									

Backgrounds due to natural radioactivity

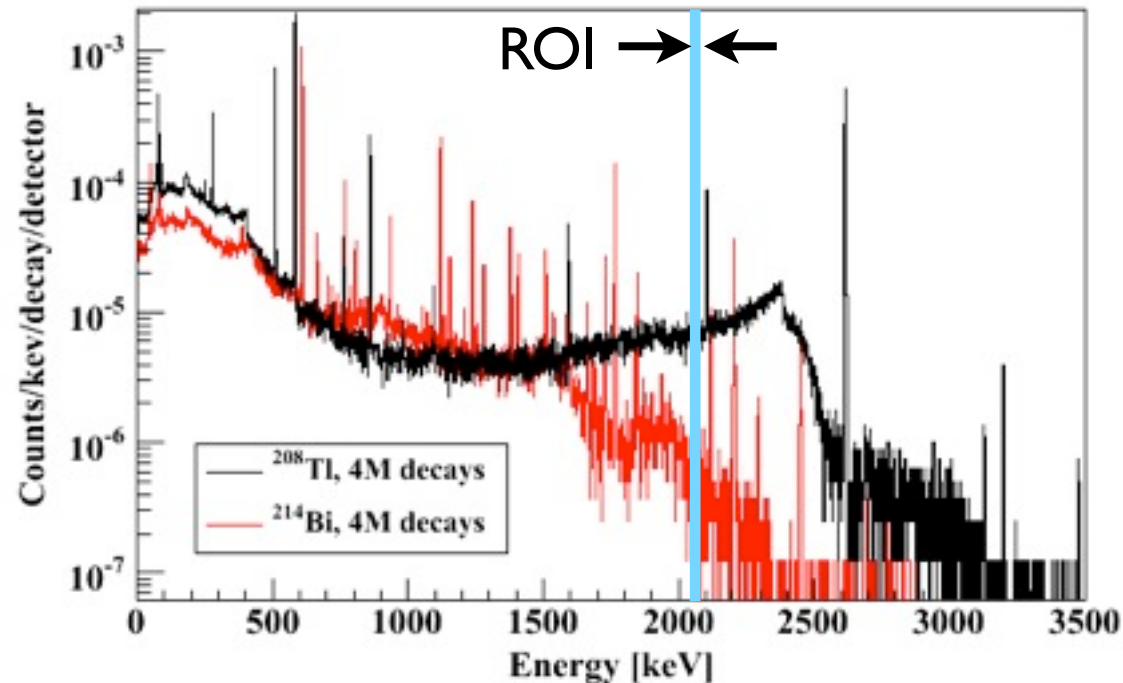
Uranium (^{238}U)

- $T_{1/2} = 4.5 \times 10^9$ years
- decays to ^{206}Pb
- ^{214}Bi produced

Thorium (^{232}Th)

- $T_{1/2} = 1.4 \times 10^{10}$ years
- decays to ^{208}Pb
- ^{208}Tl produced

simulated detector response to
 ^{208}Tl and ^{214}Bi



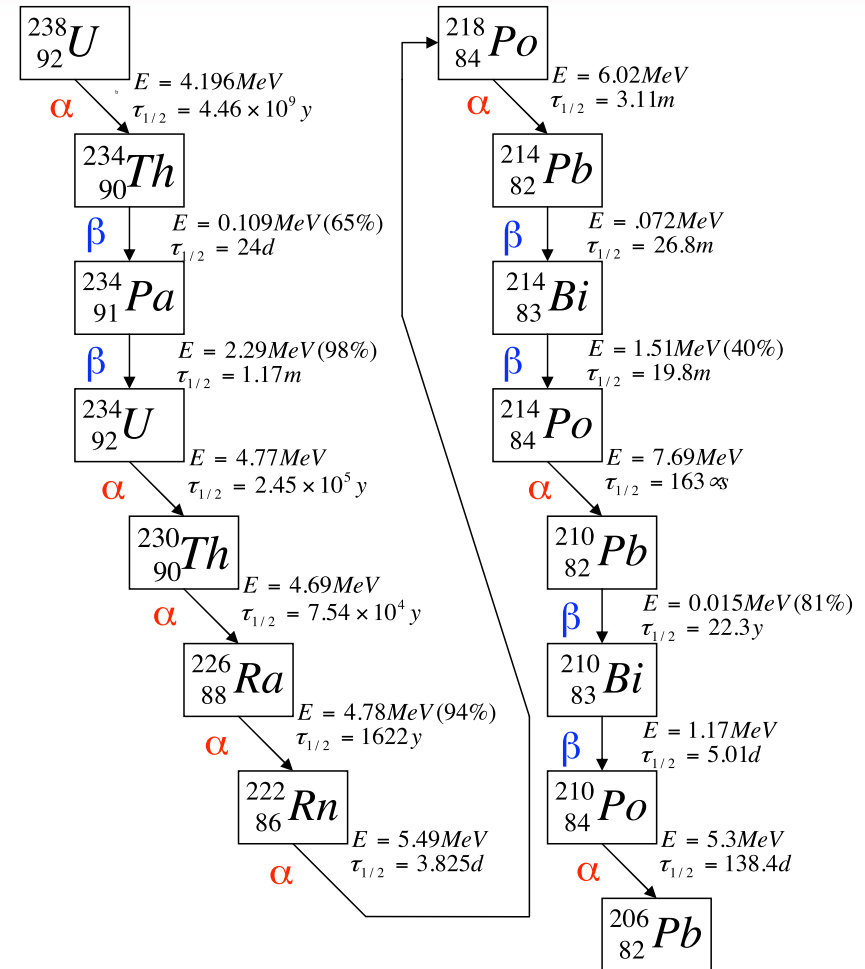
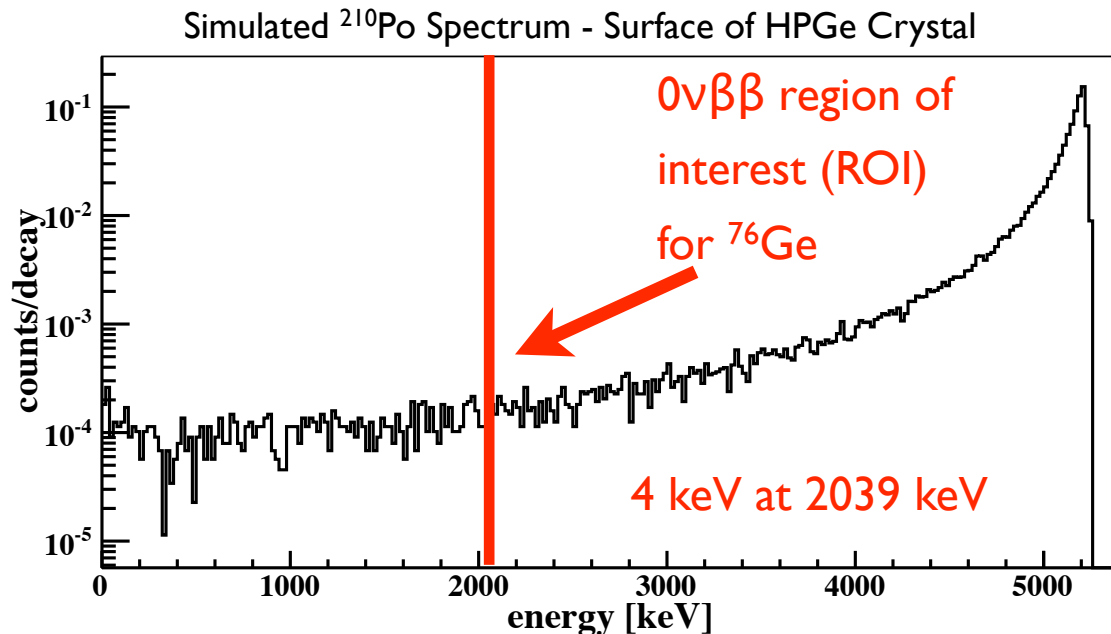
Simulations by A. Schubert

Backgrounds: Alpha Decays

Alphas emitted in decays
of ^{232}Th and ^{238}U

Kinetic energies:
4.0-8.8 MeV

Range of alphas in
germanium: 13-40 μm



Alphas only a background if

- Decays are near the detector
- The alpha only deposits a fraction of its energy in active region of germanium crystal

Simulations by R. Johnson

Reducing Backgrounds - Two Basic Strategies

- Directly reduce intrinsic, extrinsic, & cosmogenic activities
 - Select and use ultra-pure materials
 - Minimize all non “source” materials
 - Clean passive shield
 - Go deep — reduced μ 's & related induced activities
- Utilize background rejection techniques

$0\nu\beta\beta$ is a single site phenomenon, many backgrounds have multiple site interactions

- Energy resolution
- Active veto detector
- Tracking
- Energy & Angular correlations
- Ion Identification
- Granularity [multiple detectors]
- Pulse shape discrimination (PSD)
- Segmentation
- Single Site Time Correlated events (SSTC)

EXO tonne scale with Ba tagging

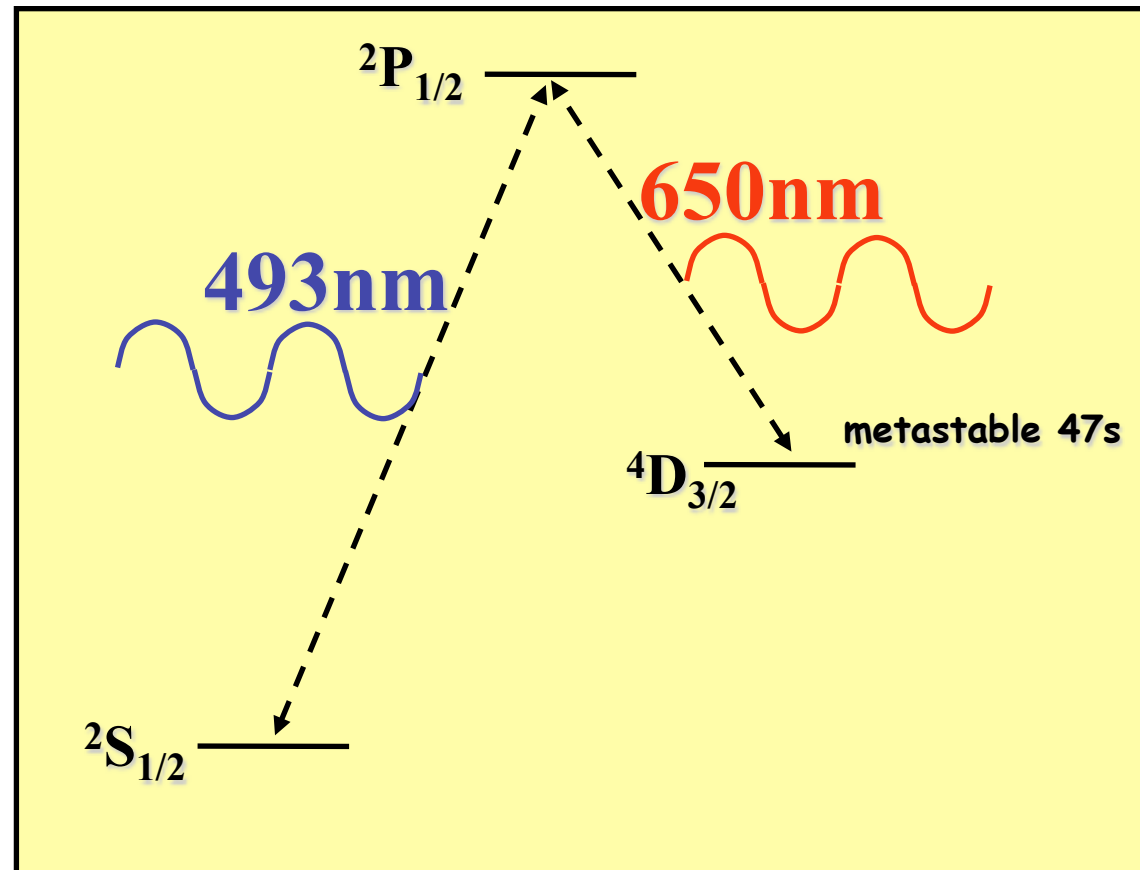


Xe offers a qualitatively new tool against background:
 $^{136}\text{Xe} \rightarrow ^{136}\text{Ba}^{++} e^- e^-$ final state can be identified
using optical spectroscopy (M.Moe PRC44 (1991) 931)

Ba^+ system best studied
(Neuhauser, Hohenstatt,
Toshek, Dehmelt 1980)
Very specific signature
“shelving”

Single ions can be detected
from a photon rate of $10^7/\text{s}$

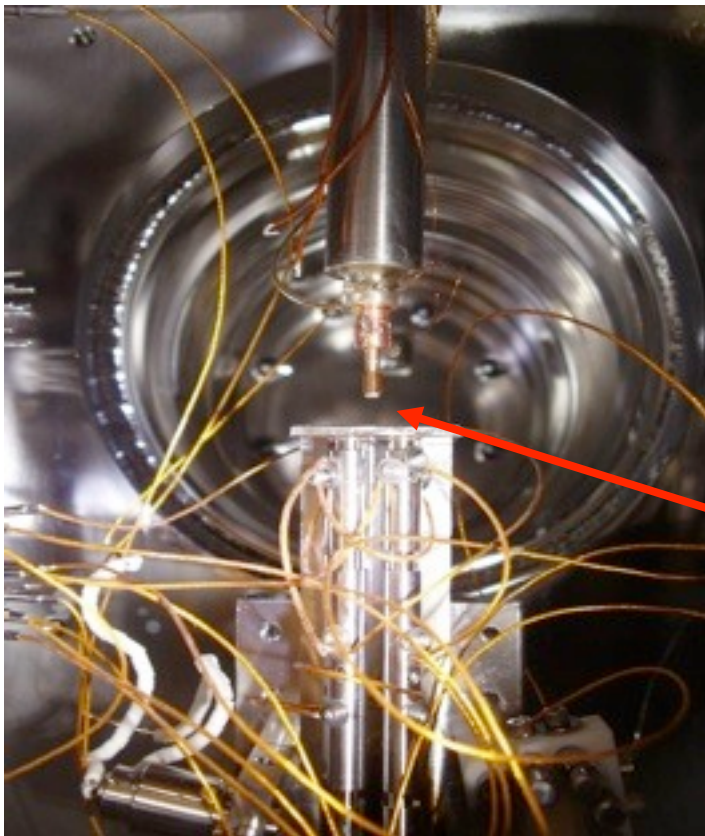
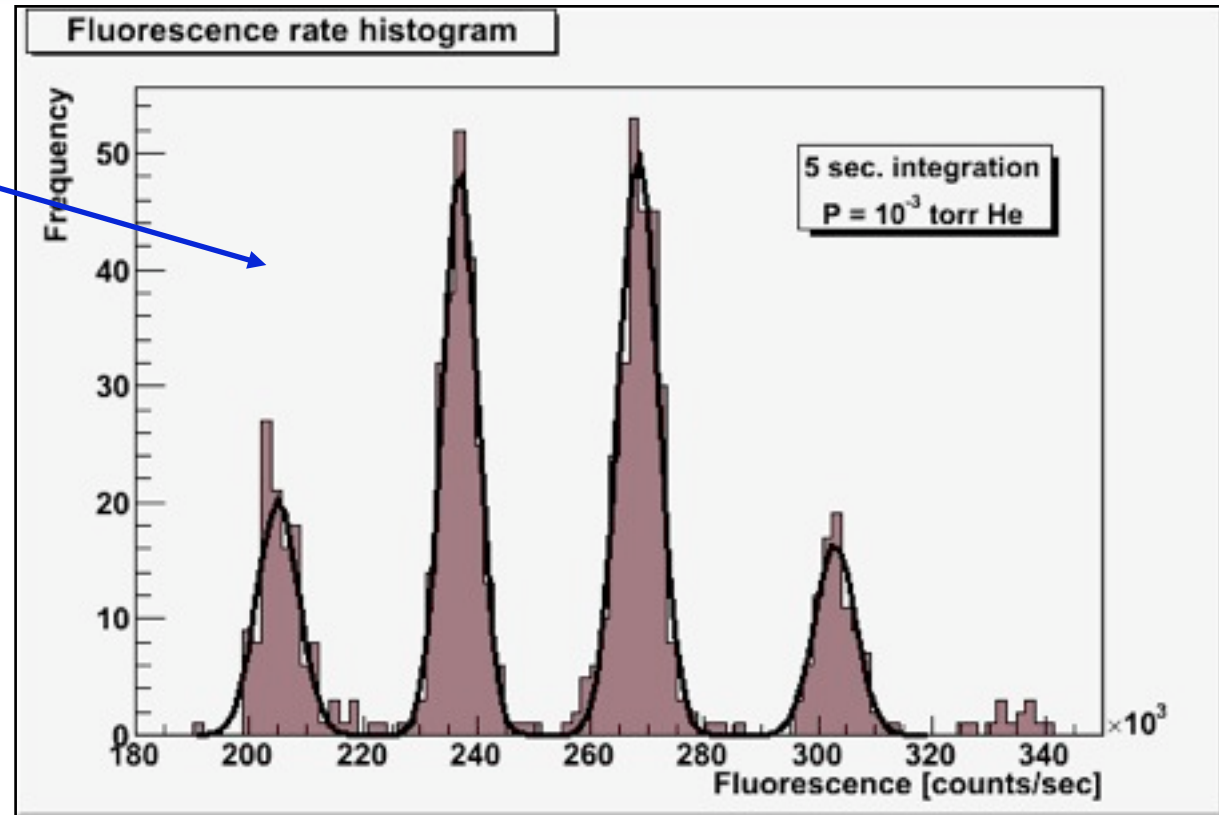
Important additional constraint
Dramatic Background reduction



EXO R&D on ion-tagging

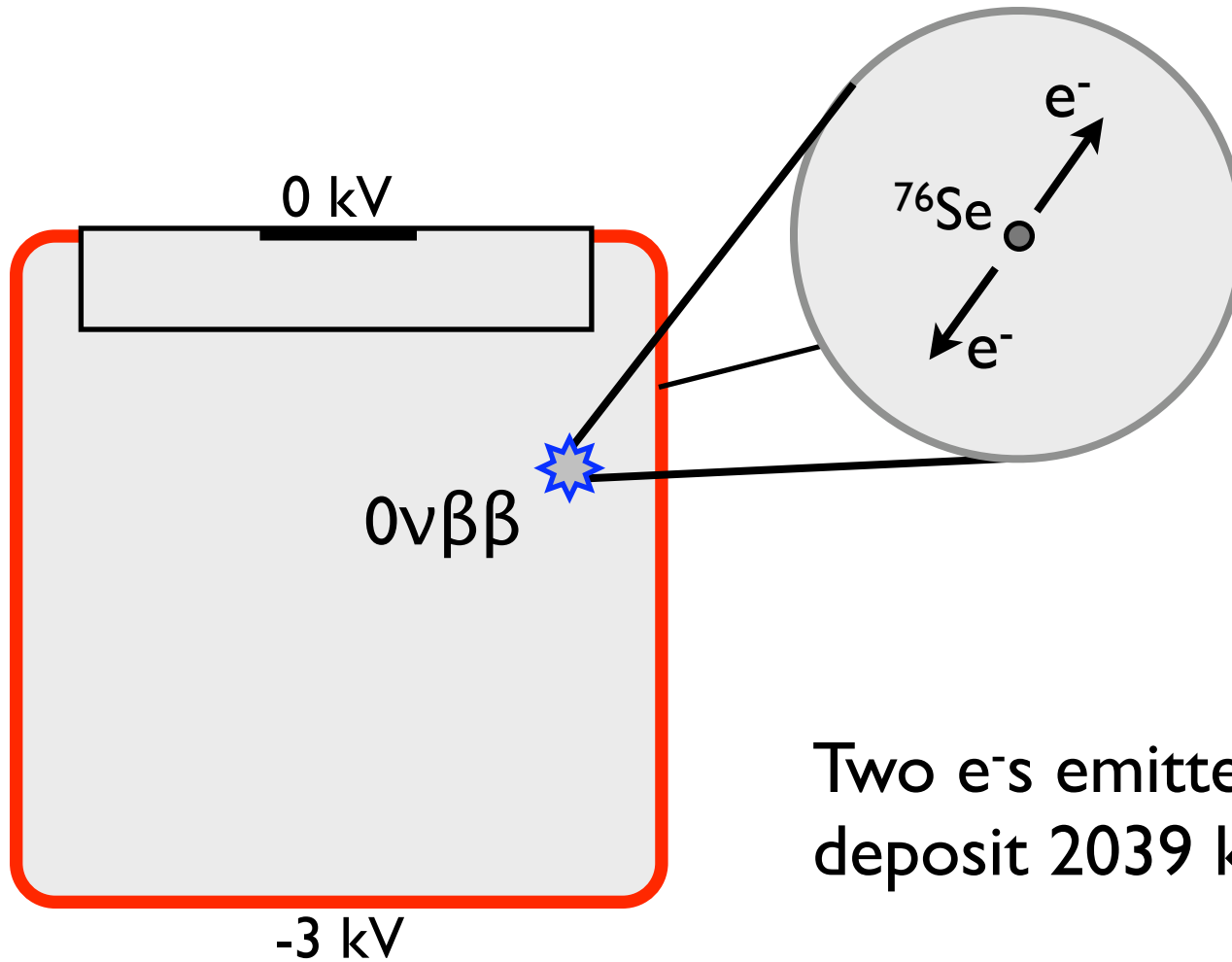


EXO linear trap can see single Ba ions in gas with large S/N ratio



Learning how to transfer single Ba ions from Xe to the ion trap

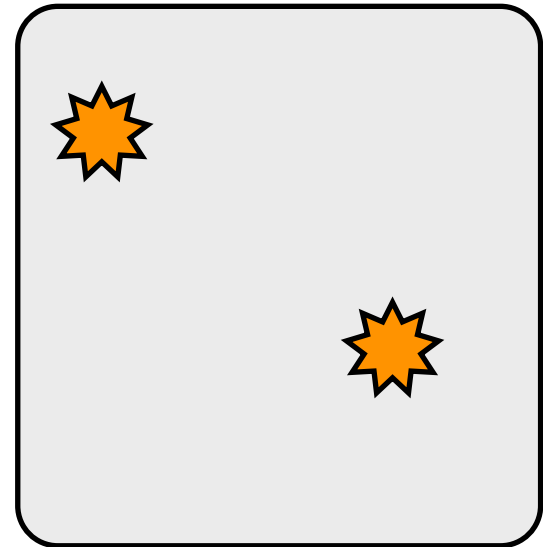
$0\nu\beta\beta$ of ^{76}Ge



Two e^- s emitted by ^{76}Ge nucleus
deposit 2039 keV in detector

a $0\nu\beta\beta$ event

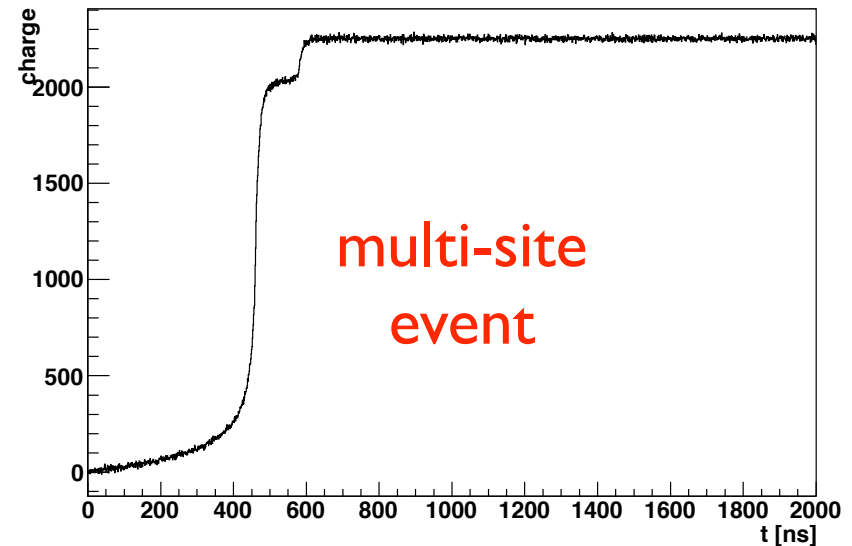
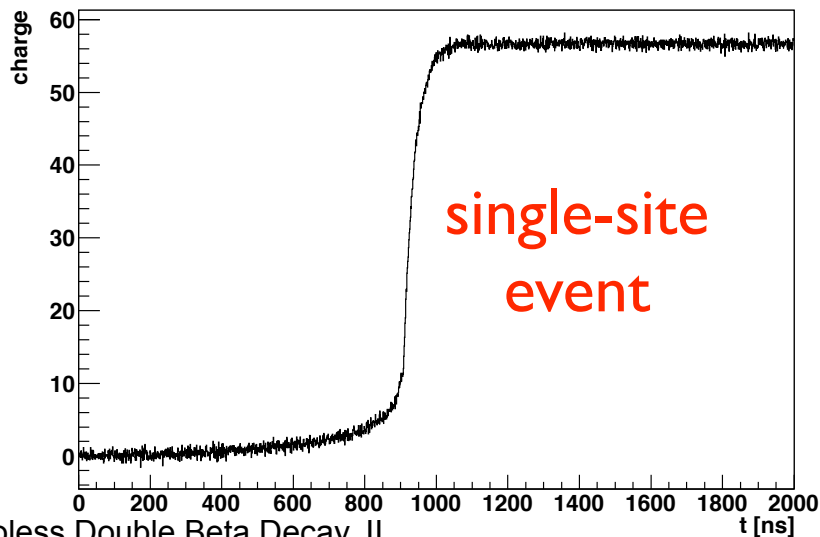
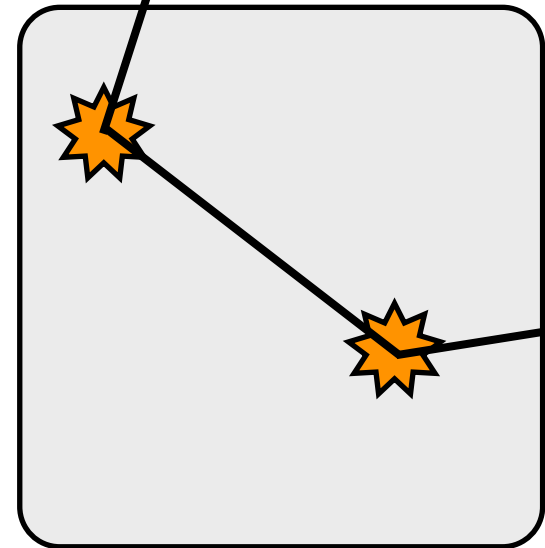
- 2039 keV
- one detector
- one site



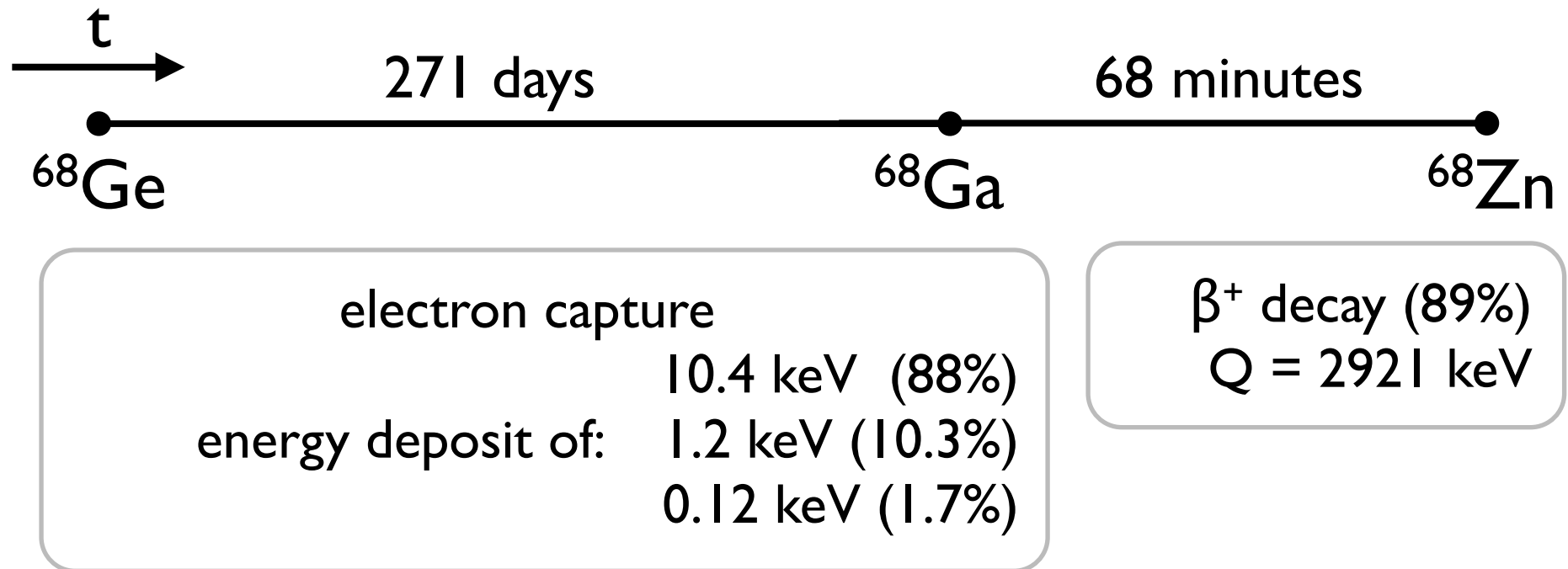
not a $0\nu\beta\beta$ event

- 2039 keV
- one detector
- one site

single-site
cut



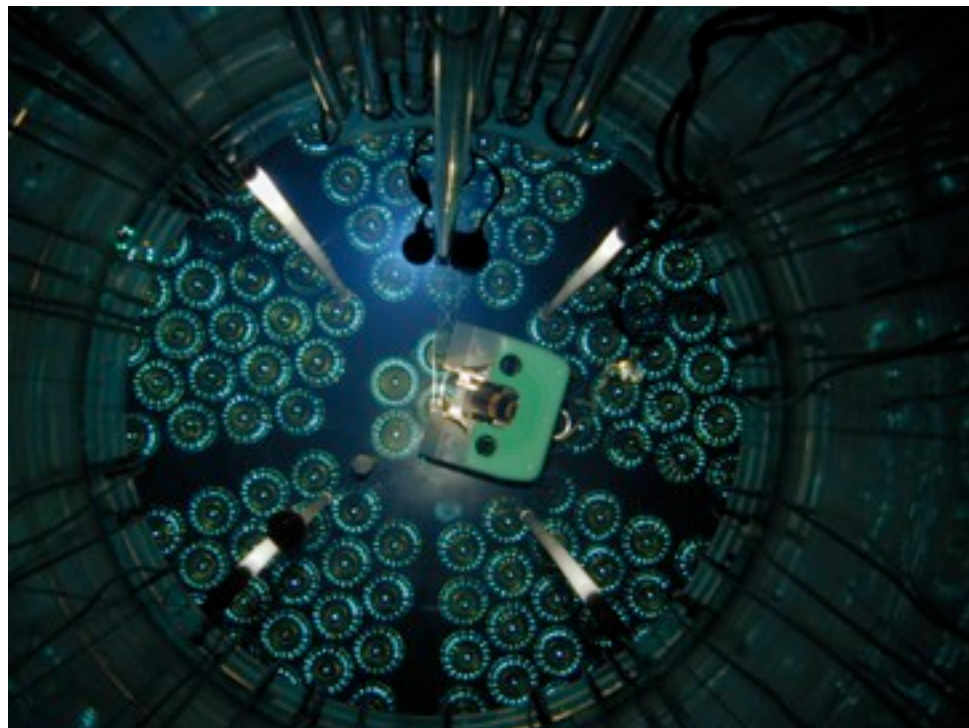
single-site time correlation



look back in time a few half lives to
identify background events

Background reduction at the larger scale

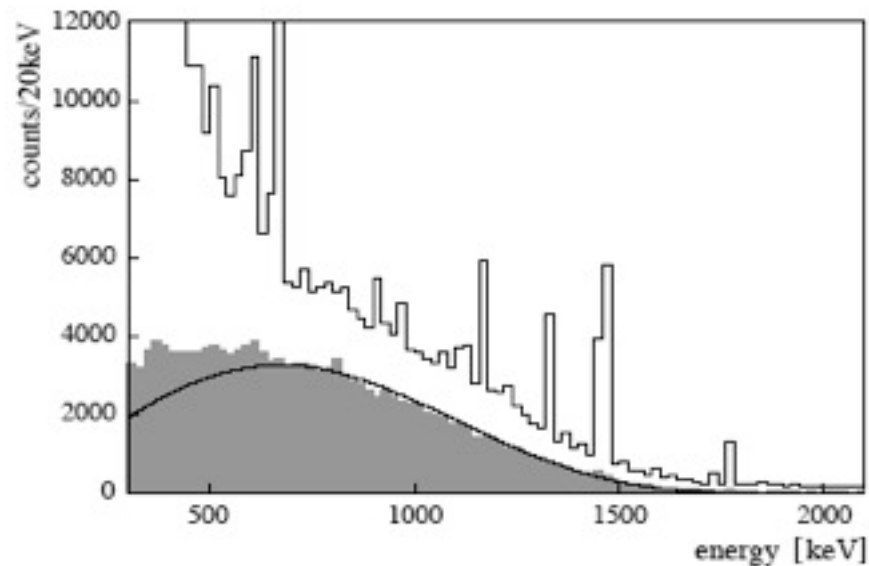
- Many groups have built $0\nu\beta\beta$ -decay experiments at the few to 10 kg level. - **Need to scale this up to the 100s of kg level.**
- Can utilize knowledge from groups that have demonstrated the construction of low-background, large-scale detectors underground: **e.g. KamLAND, SNO, SAGE, GNO, Borexino CTF**
- SNO Neutral Current Detector
Array of ^3He proportional counters
 - 450 kg of material
 - 300 detector segments
 - Activity (Stonehill, 2005)
 - $23 \pm 4 \text{ } \mu\text{Bq/kg } ^{232}\text{Th}$
 - $35 \pm 8/-10 \text{ } \mu\text{Bq/kg } ^{238}\text{U}$



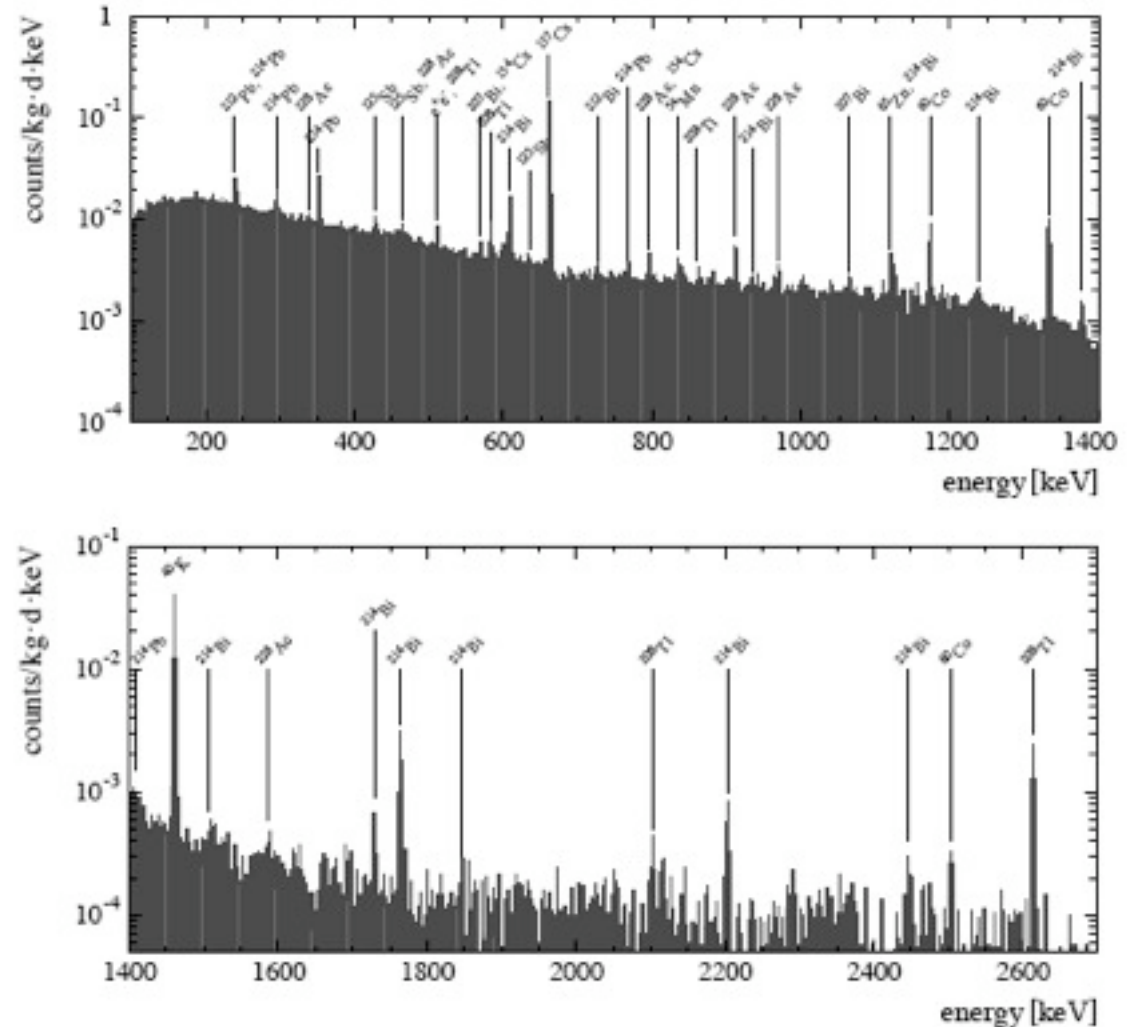
Example: Heidelberg-Moscow ^{76}Ge

Five ^{76}Ge crystals, 11 kg of total mass, ~ 70 kg-years of data.

- $T_{1/2} > 1.9 \times 10^{25} \text{ y}$ (90%CL)



Klapdor-Kleingrothaus et al.,
Eur. Phys. J 12, 147 (2001)



The KKDC Result

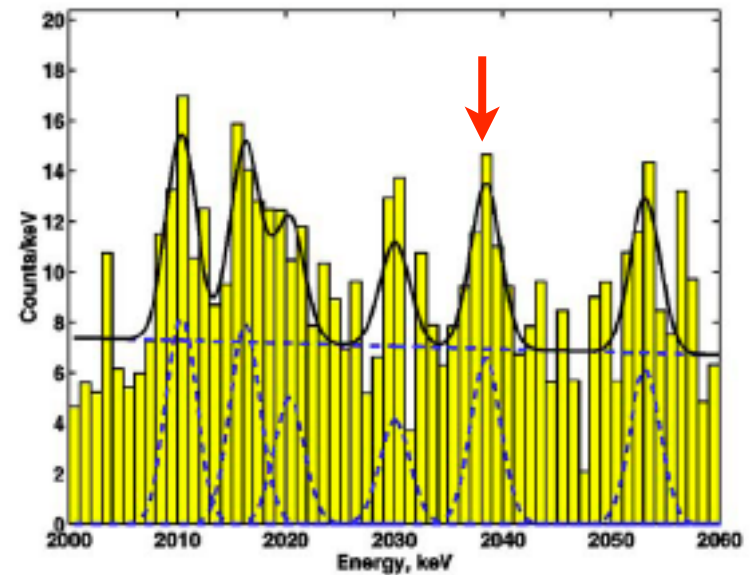
Klapdor-Kleingrothaus, Krivosheina, Dietz and Chkvorets, *Phys. Lett. B* **586** 198 (2004).

Result: Five ^{76}Ge crystals, 10.96 kg of total mass, 71 kg-years of data.

$$T_{1/2} = (1.19 + 2.99/-0.5) \times 10^{25} \text{ y}$$
$$0.24 < m_\nu < 0.58 \text{ eV} \quad (3\sigma)$$

Plotted a subset of the data for four of five crystals, 51.4 kg-years of data.

$$T_{1/2} = (1.25 + 6.05/-0.57) \times 10^{25} \text{ y}$$



The KKDC Result

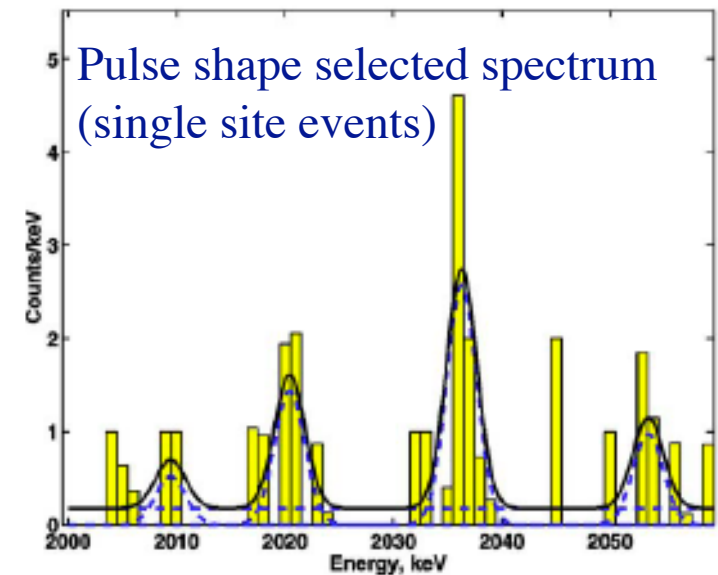
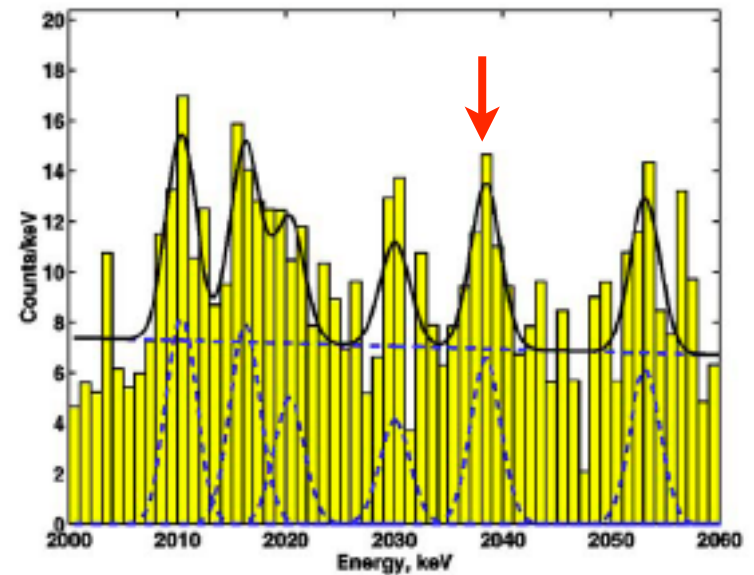
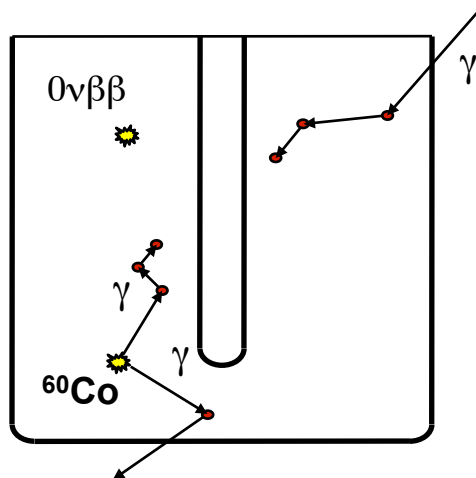
Klapdor-Kleingrothaus, Krivosheina, Dietz and Chkvorets, *Phys. Lett. B* **586** 198 (2004).

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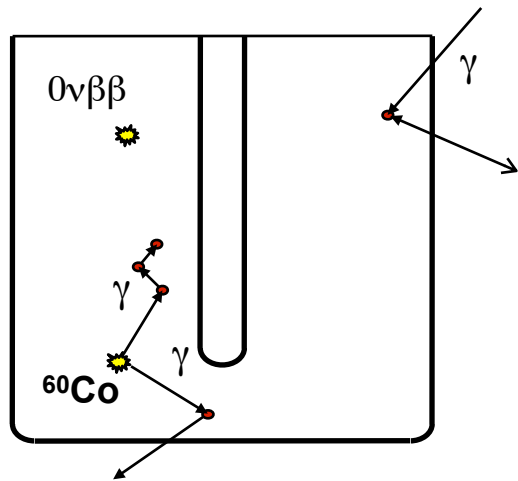
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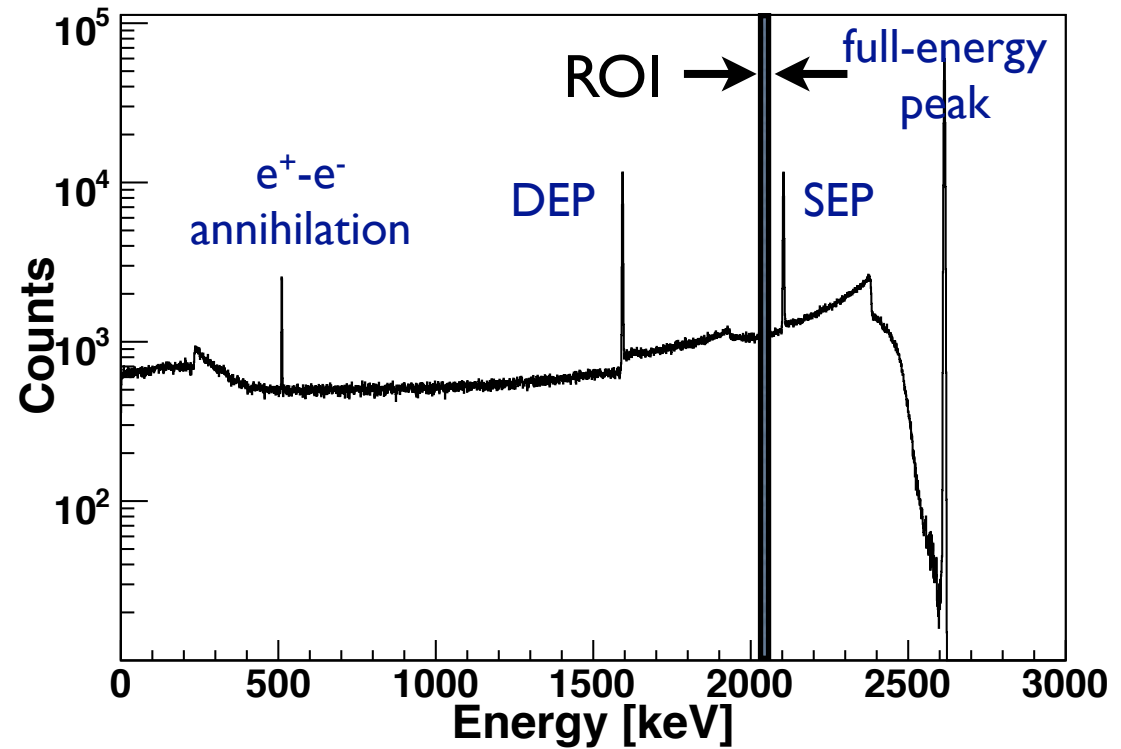


Backgrounds from ^{208}Tl

simulations and measurements show that some ^{208}Tl gammas from outside materials will interact in the detector, and then scatter back out, leaving a continuum of “point like” interactions



simulated response of Ge detector
exposed to 2615-keV γ s



80-mm diameter, 30-mm tall detector
simulated with MaGe

Simulations by A. Schubert

The KKDC Result

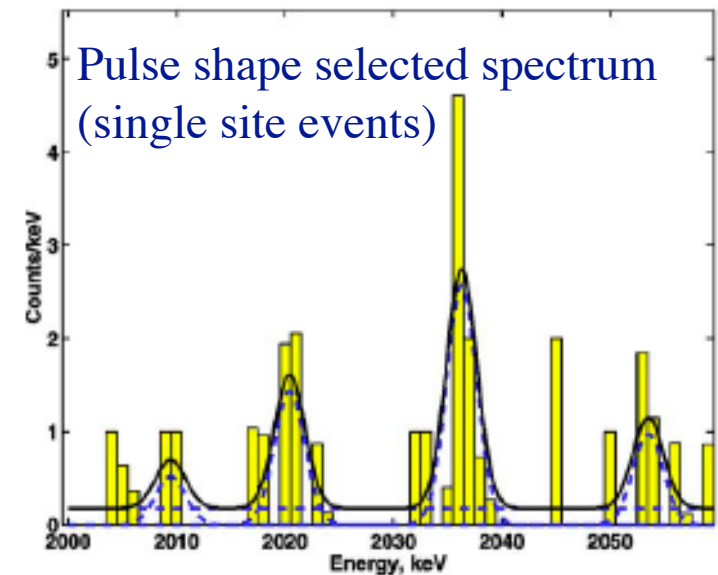
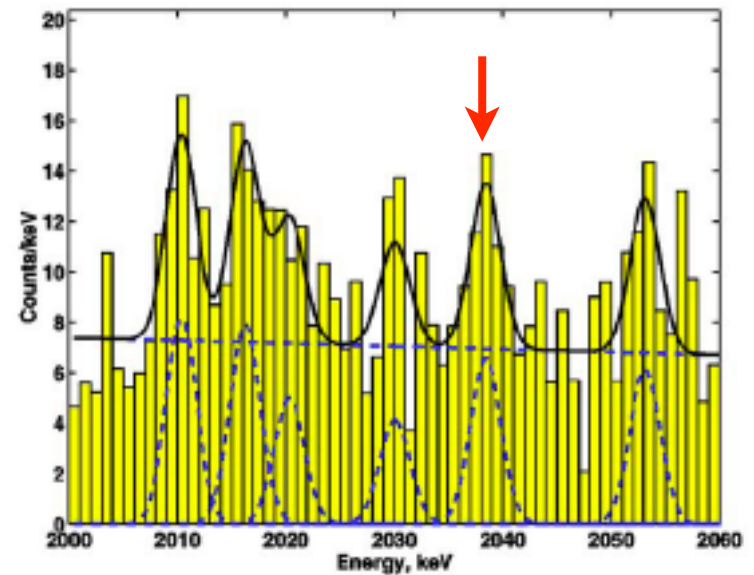
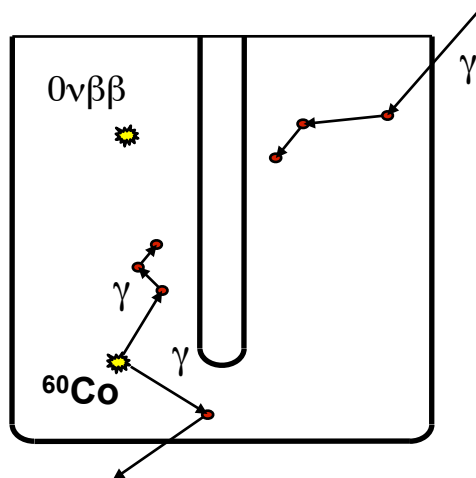
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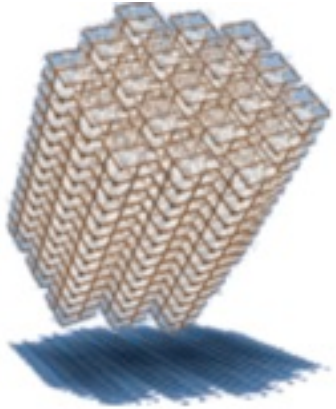
$0\nu\beta\beta$ Decay II

Next Generation Experiments

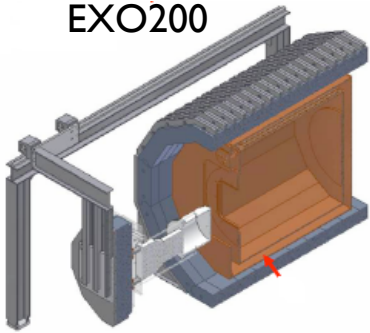
Overview of Experiments
Selected Techniques

$0\nu\beta\beta$ decay Experiments - Efforts Underway

CUORE



EXO200



NEMO

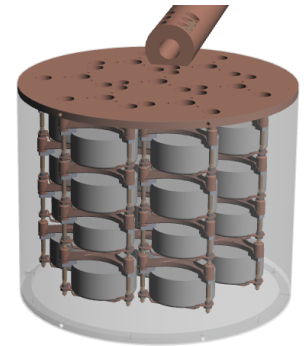


Collaboration	Isotope	Technique	Mass	Status
CAMEO	Cd-116	CdWO ₄ crystals	1 t	
CANDLES	Ca-48	60 CaF ₂ crystals in liq. scint	6 kg	Construction
CARVEL	Ca-48	⁴⁸ CaWO ₄ crystal scint.	100 kg	
COBRA	Cd-116, Te-130	CdZnTe detectors	10 kg	R&D
CUROICINO	Te-130	TeO ₂ Bolometer	11 kg	Operating
CUORE			206 kg	Construction
DCBA	Nd-150	Nd foils & tracking chambers	20 kg	R&D
EXO200	Xe-136	Xe TPC	200 kg	Construction
EXO		Xe TPC with ion ID		Future
GEM	Ge-76	Ge diodes in LN	1 t	
GERDA	Ge-76	Seg. and UnSeg. Ge in LAr	35-40 kg 1 t	Construction Future
GSO	Gd-160	Gd ₂ SiO ₅ :Ce crystal scint. in liquid scint	2t	
KamLAND	Xe-136	Xe in liq. Scint.	200 kg	Construction
MAJORANA	Ge-76	Point Contact Ge	60 kg	Construction
			1 t	Future
NEMO3	Mo-100 Se-82	Foils with tracking	6.9 kg 0.9 kg	Operating
SuperNEMO			100 kg	R&D
NEXT	Xe-136	High Pressure TPC	1t	R&D
MOON	Mo-100	Mo sheets	200 kg 1 t	R&D
SNO+	Nd-150	0.1% suspended in Scint.	56 kg	Construction
Xe	Xe-136	Xe in liq. Scint.	1.56 t	

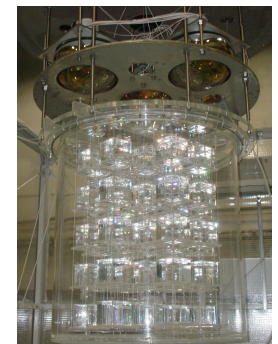
GERDA



MAJORANA



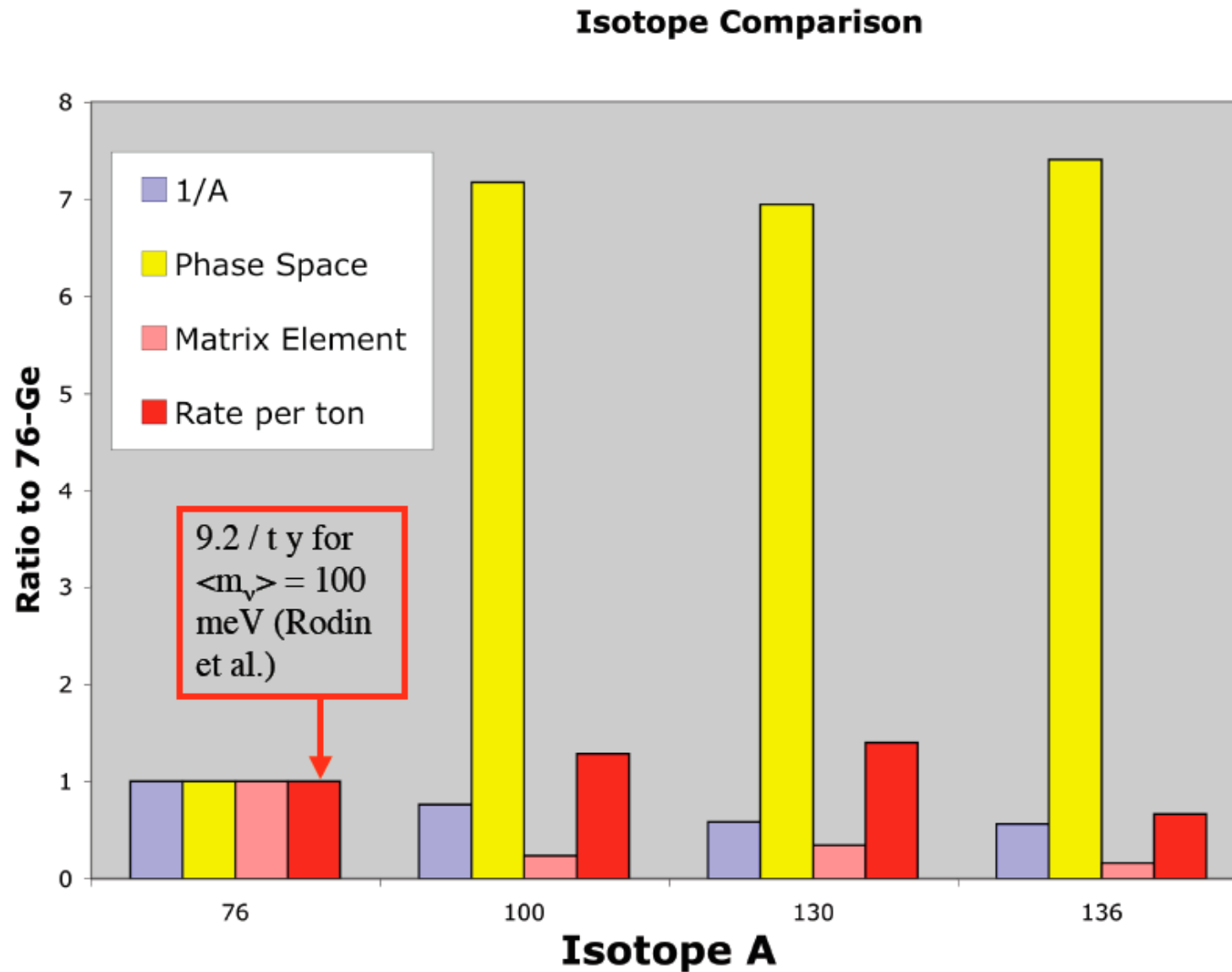
CANDLES



Operating

Construction

“Relative” Sensitivities

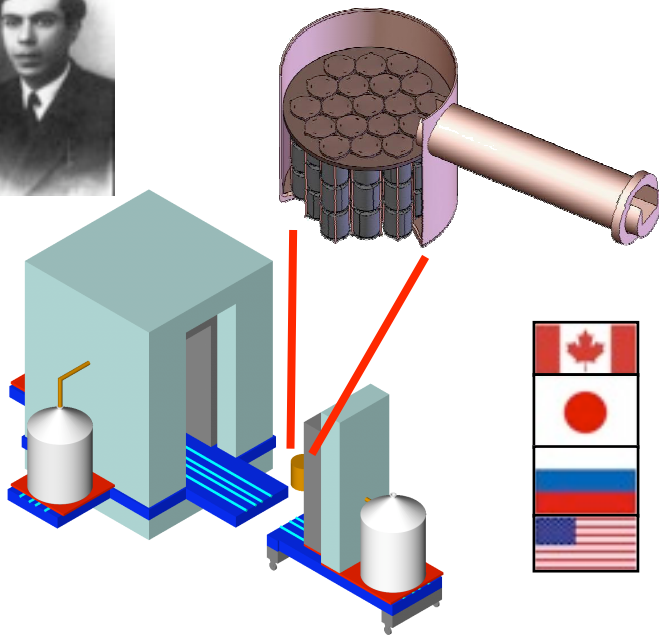


R.G.H. Robertson

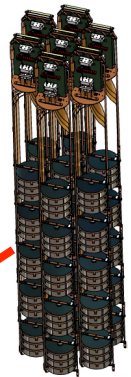
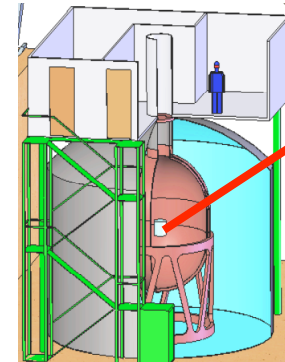
Using Rodin et al. Nucl. Matrix elements

^{76}Ge - MAJORANA and GERDA

- 86% enriched Hyper-pure ^{76}Ge crystals
- 0.16 % resolution
- best $0\nu\beta\beta$ -decay sensitivity to date



- $^{\text{enr}}\text{Ge}$ modules in electroformed Cu cryostat, Cu / Pb passive shield
- 4π plastic scintillator μ veto
- Demonstrator: 30 kg $^{\text{enr}}\text{Ge}$ /30 kg $^{\text{nat}}\text{Ge}$



- $^{\text{enr}}\text{Ge}$ array submersed in LAr
- Water Cherenkov μ veto
- Phase I: ~ 18 kg (H-M/IGEX xtals)
- Phase II: +20 kg segmented xtals

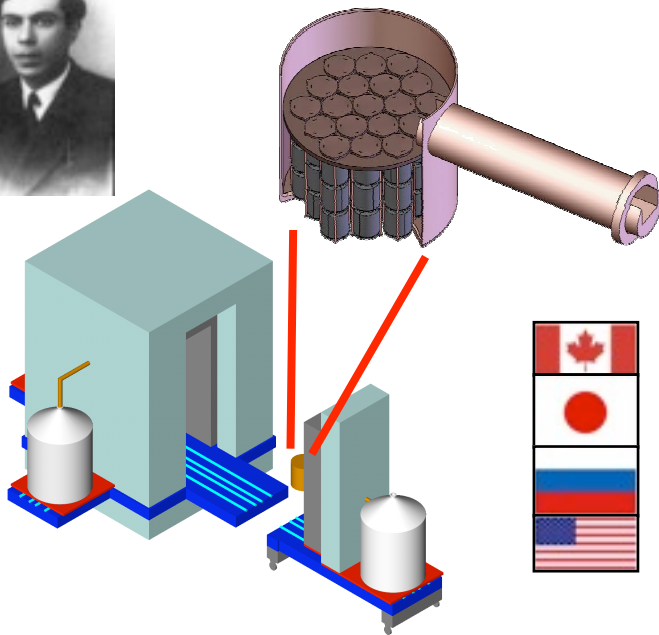
^{76}Ge - MAJORANA and GERDA

Joint Cooperative Agreement:

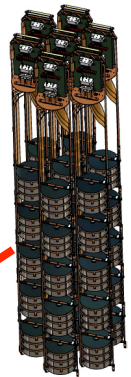
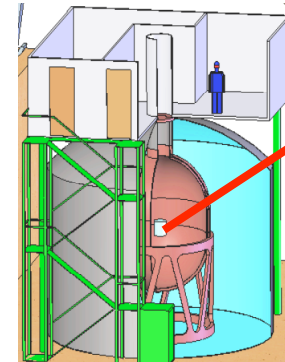
Open exchange of knowledge & technologies (e.g. MaGe, R&D)

Intention to merge for larger scale 1-tonne exp.

Select best techniques developed and tested in GERDA and MAJORANA

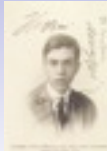


- $^{\text{enr}}\text{Ge}$ modules in electroformed Cu cryostat, Cu / Pb passive shield
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- $^{\text{enr}}\text{Ge}$ array submersed in LAr
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- Phase II: +20 kg segmented xtals

The MAJORANA ^{76}Ge Demonstrator



^{76}Ge offers an excellent combination of capabilities & sensitivities.

(Excellent energy resolution, intrinsically clean detectors, commercial technologies, best $0\nu\beta\beta$ sensitivity to date)

- 60-kg of Ge detectors

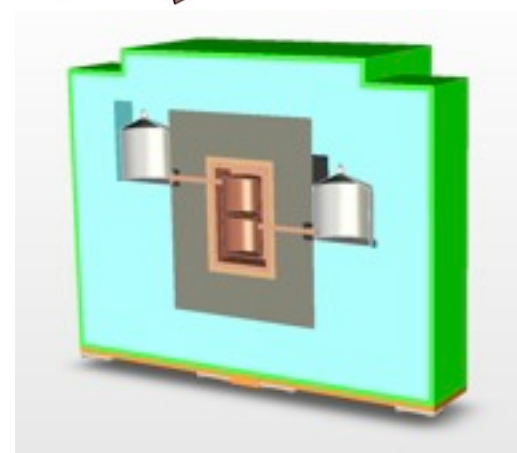
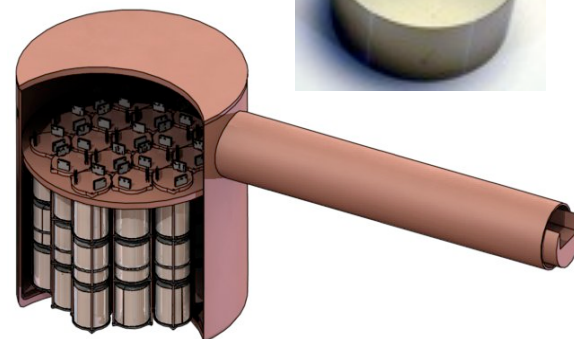
- 30-kg of 86% enriched ^{76}Ge crystals required for science goal; 60-kg for background sensitivity
- Examine detector technology options
p- and n-type, segmentation, point-contact.

- Low-background Cryostats & Shield

- ultra-clean, electroformed Cu
- naturally scalable
- Compact low-background passive Cu and Pb shield with active muon veto

- Located underground 4850' level at Sanford Lab (Homestake)

- Background Goal in the $0\nu\beta\beta$ peak region of interest (4 keV at 2039 keV) ~ **1 count/ROI/t-y** (after analysis cuts)

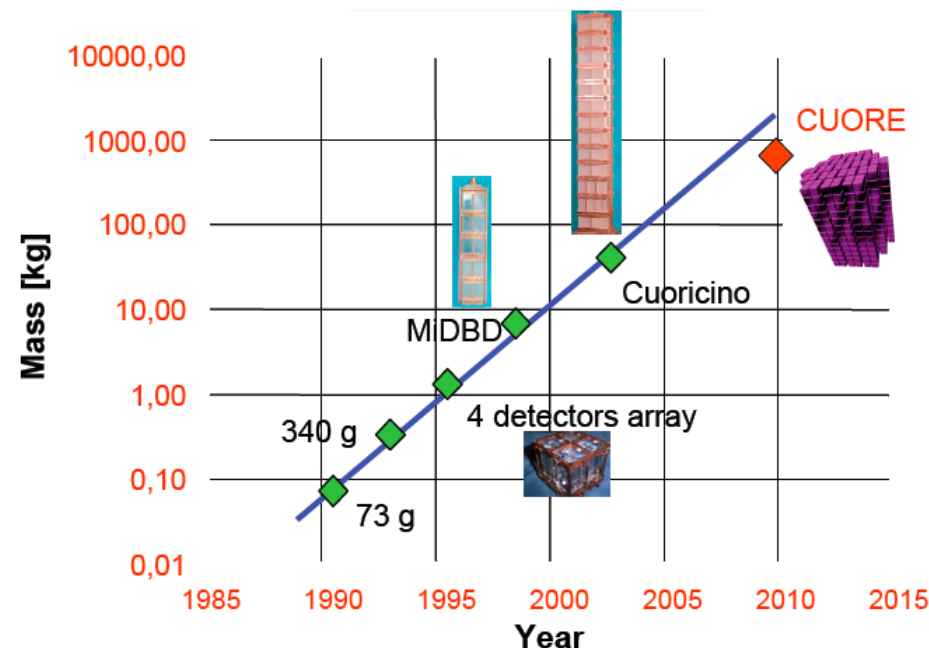
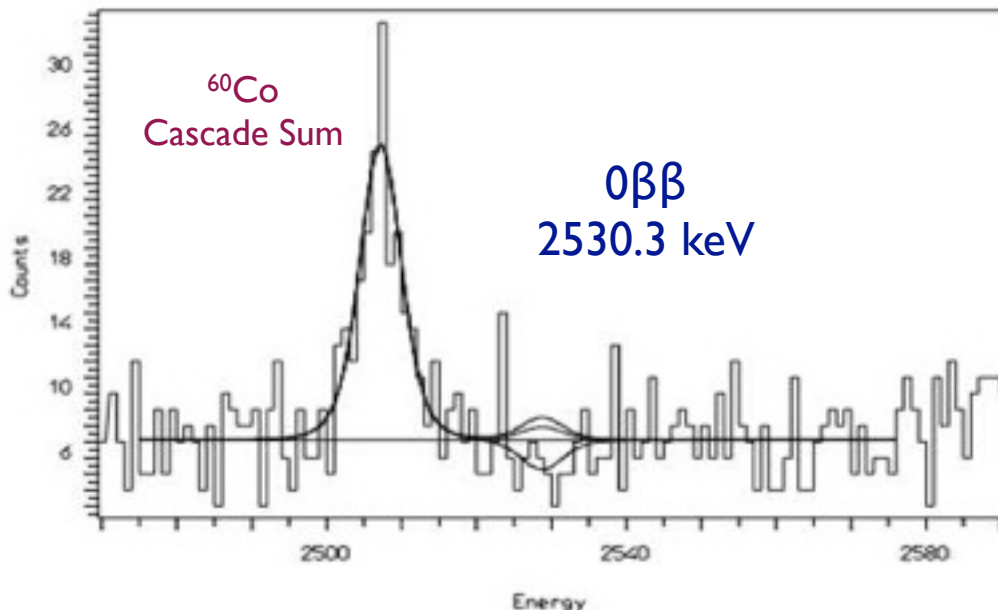
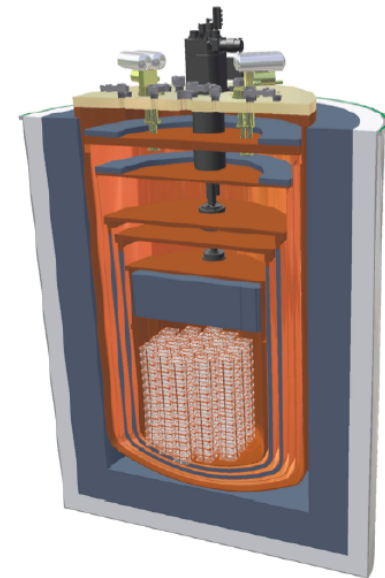


CUORE (Italy - Spain - U.S.)

- ^{130}Te (34% abundance) bolometer.
- Array of 988 TeO_2 crystals
- Expects to operate in Gran Sasso by 2012
- Builds upon success of Cuoricino

11.83 kg of ^{130}Te , April 2003-2006

$T_{1/2} > 2.94 \times 10^{24}$ y (90% CL)



SuperNEMO

Planar and modular design: ~ 100 kg of enriched isotopes ($20 \text{ modules} \times 5 \text{ kg}$)

1 module:

Source (40 mg/cm^2) $4 \times 3 \text{ m}^2$

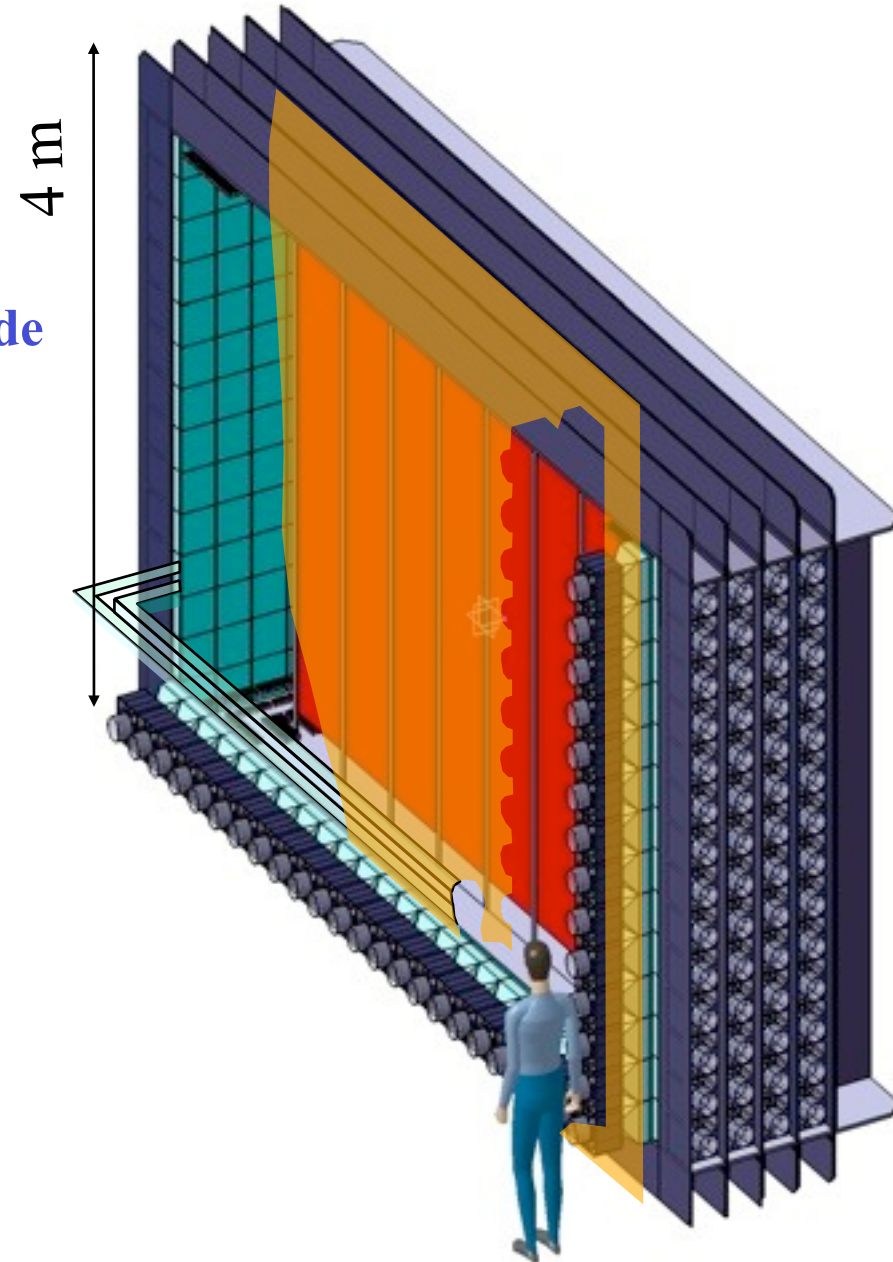
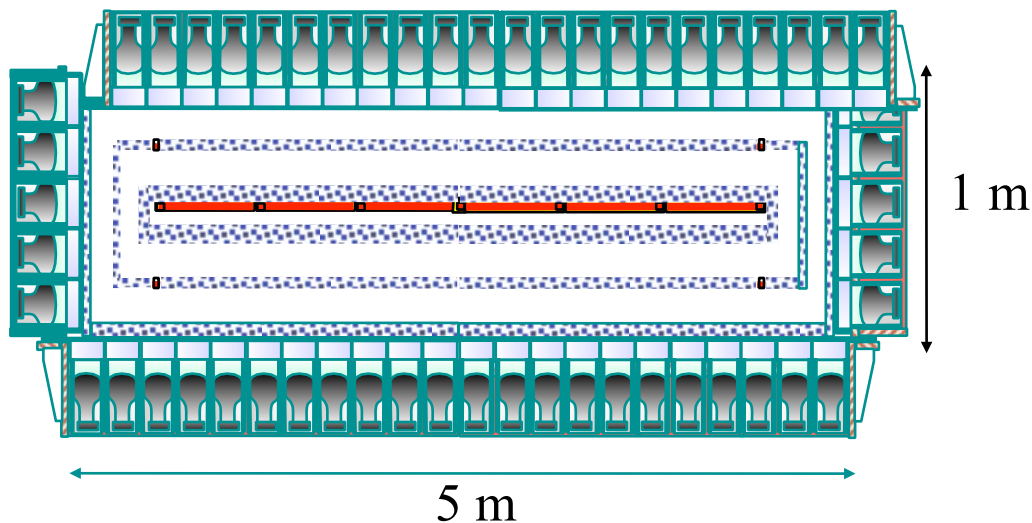
Tracking : drift chamber ~ 3000 cells in Geiger mode

Calorimeter: scintillators + PM

$\sim 1\,000$ PM if scint. blocks

~ 100 PM if scint. bars

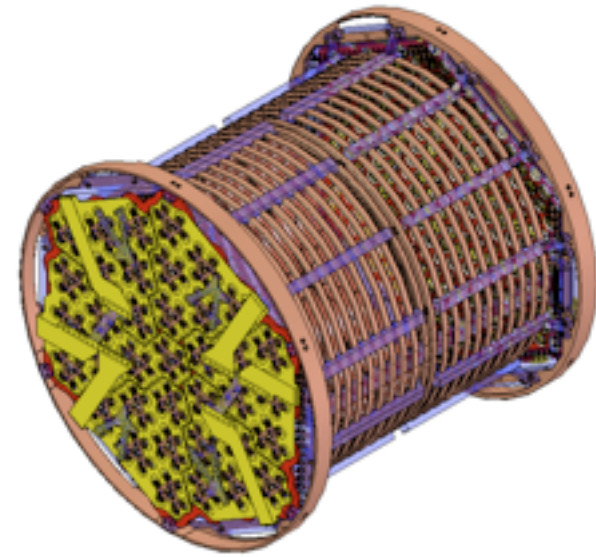
Top view



EXO-200 (U.S., Canada, Russia, Switzerland)



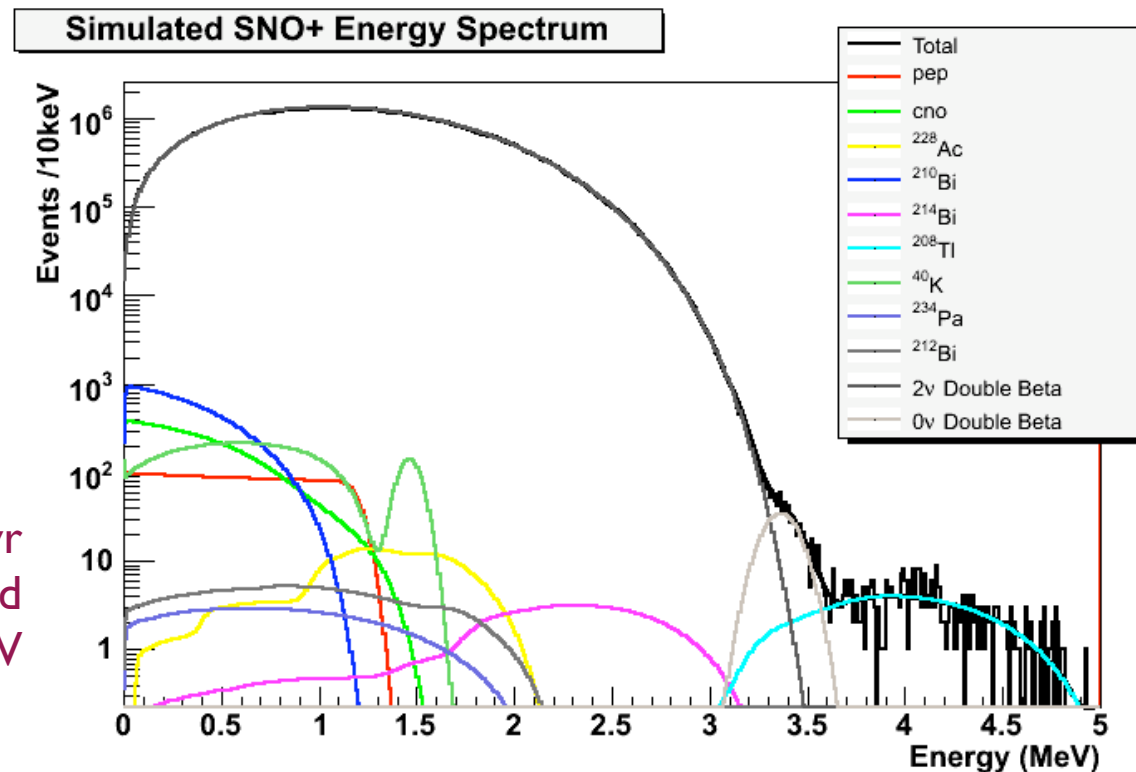
- 200 kg of 80% enriched ^{136}Xe
- Liquid time-projection chamber
- Uses charge and light collection 1.6% resolution
- Expects to operate in WIPP by end of 2009
- Aims to measure $2\nu\beta\beta$ mode (not yet observed)



Science : Neutrinoless Double Beta-decay, Geoneutrinos, Solar ν

- ^{150}Nd (5.6% abundance) 0.1% loaded in scintillator.
- Initial plan is to use 1 ton of $^{\text{nat}}\text{Nd}$ (56 kg of ^{150}Nd)
- Located at SNOLAB
- Utilizes substantial investment in SNO
- Initial success in loading into pseudocumene and also in linear alkylbenzene
- Technical challenge - available light with loaded scintillator
- Considering enrichment option

1 yr
500 kg of ^{150}Nd
 $\langle m_{ee} \rangle = 150 \text{ meV}$



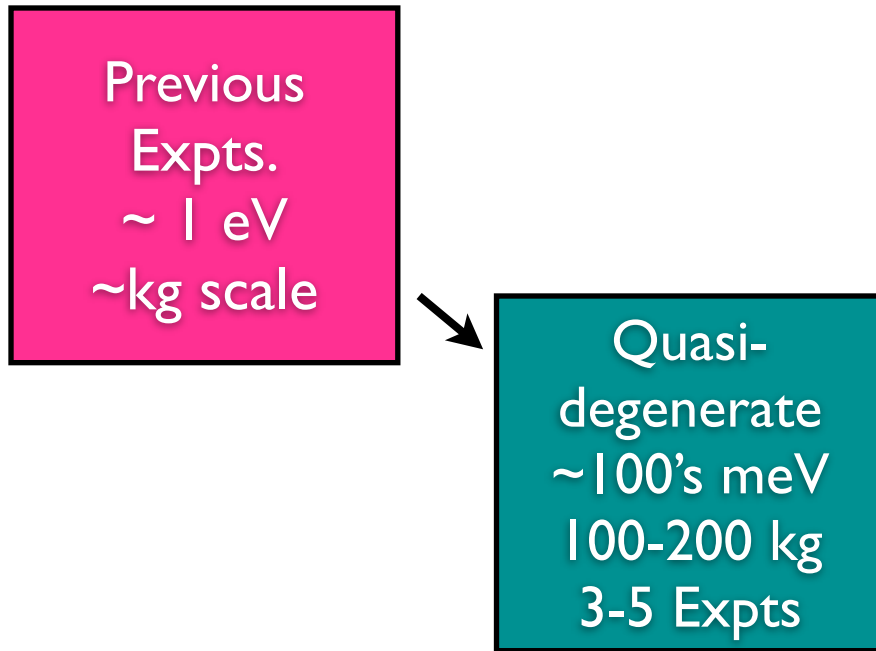
$0\nu\beta\beta$ Decay II

Future Program

Discovery of $0\nu\beta\beta$ -decay

- **Strong evidence** : a combination of
 - Correct peak energy
 - Single-site energy deposit
 - Proper detector distributions (spatial, temporal)
 - Rate scales with isotope fraction
 - Full energy spectrum understood
- **Further confirmation**: more difficult
 - Observe the two-electron nature of the event
 - Measure kinematic dist. (energy sharing, opening angle)
 - Observe the daughter
 - Observe the excited state decay
- **Irrefutable**
 - Observe the process in several isotopes, using a variety of experimental techniques

Experimental Program in $0\nu\beta\beta$

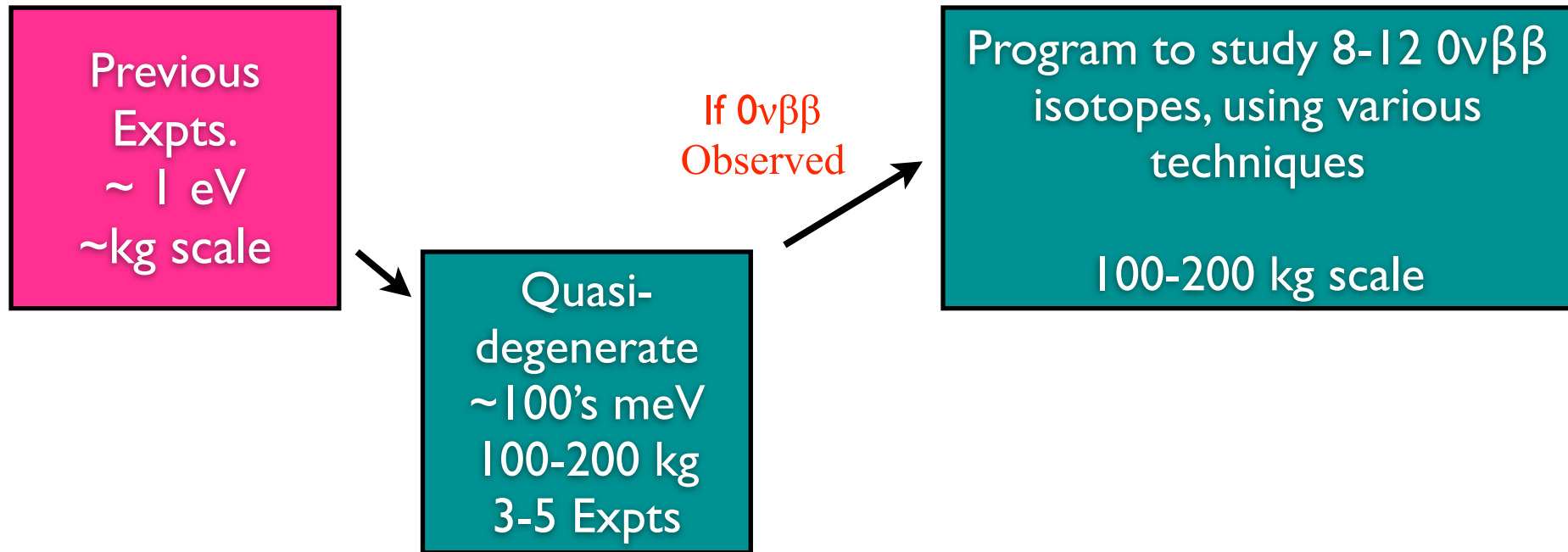


1980 - Present

2007 - 2014

2013 - 2020

Experimental Program in $0\nu\beta\beta$

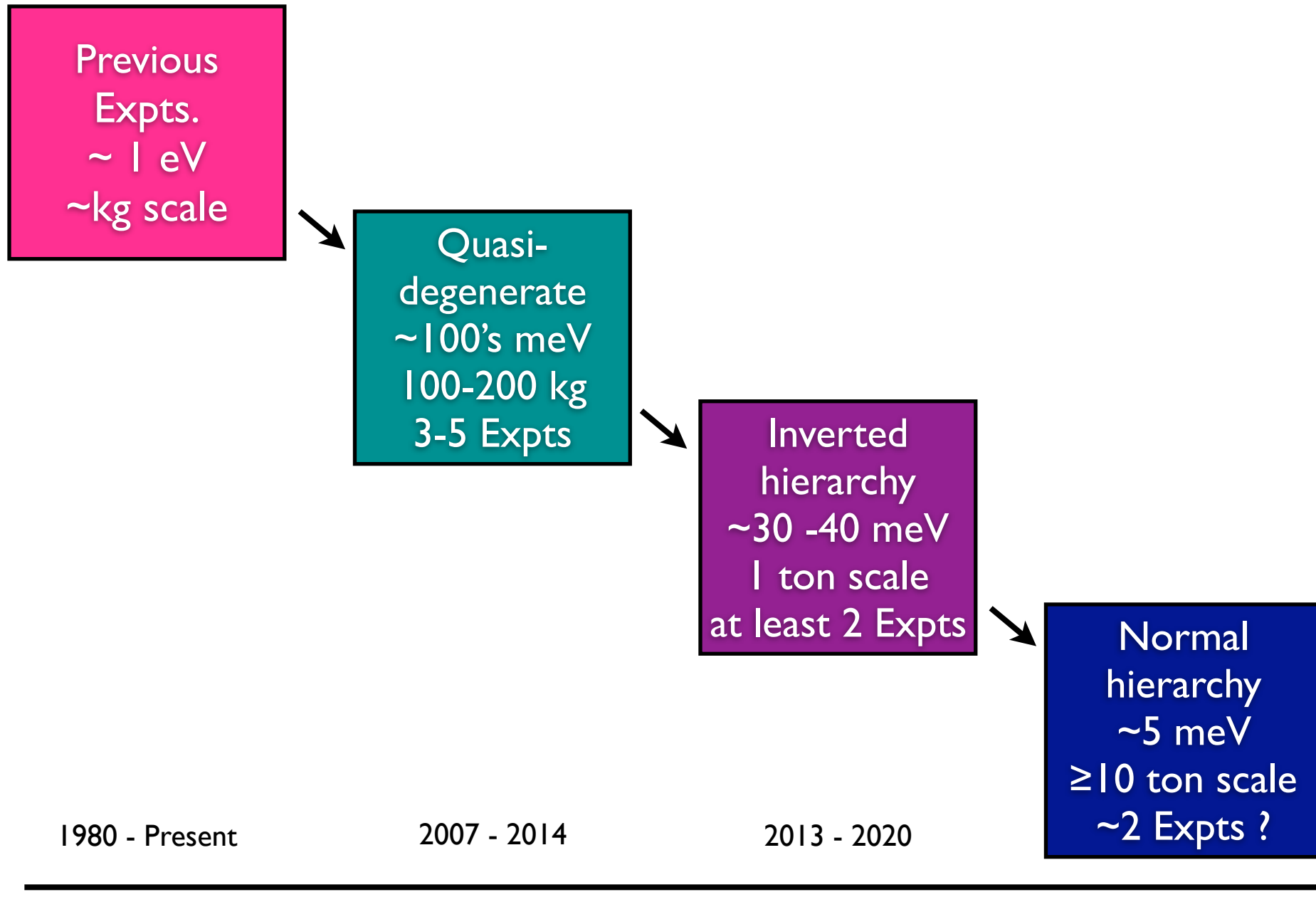


1980 - Present

2007 - 2014

2013 - 2020

Experimental Program in $0\nu\beta\beta$



Summary

- The observation of $0\nu\beta\beta$ -decay would demonstrate Lepton number violation and indicate that neutrinos are Majorana particles - **constituting a major discovery**.
 - Discovery needs to be confirmed from independent experiments using different isotopes and measurement techniques.
- If $0\nu\beta\beta$ -decay is observed then it opens an exquisitely sensitive window to search for physics beyond the Standard model.
 - Measurement of $\langle m_{\beta\beta} \rangle$ is complementary to direct and cosmological measurements of neutrino mass.
 - Measurements in different isotopes should provide insights into the underlying LNV physics process(es).

Over the past 70 years there have been dramatic changes in our understanding of the of nuclear and particle physics, yet $0\nu\beta\beta$ -decay remains relevant as we endeavor to elucidate the underlying framework of our universe.