

Extra Dimensions

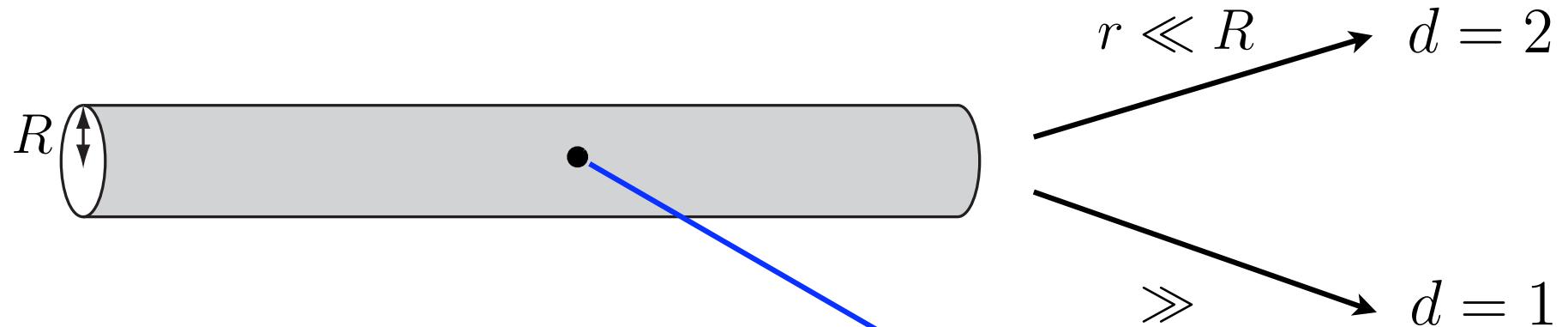
Motivation

- Non-trivial extension of spacetime symmetry
(c.f. SUSY)
- Required by string theory

Leads in many interesting directions

- Unification of spins (c.f. SUSY again)
 - Gauge-gravity unification (Kaluza 1919; Klein 1921)
 - Gauge-Higgs unification (Hosotani 1983)
- Quantum gravity at low scales (Arkani-Hamed, Dimopoulos, Dvali 1998)
- Hierarchy from gravitational warping (Randall, Sundrum 1999)
- 5D gravity \leftrightarrow 4D gauge theory (Maldacena 1999)

Extra Dimensions 101



Field theory on a cylinder:

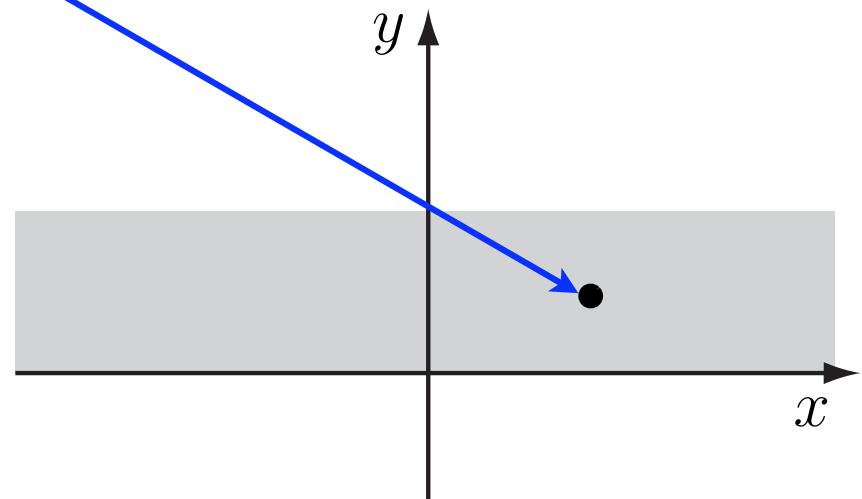
$$\Phi(x, y) = \Phi(x, y + 2\pi R)$$

4+1 dimensions just as easy:

$$\Phi(x^\mu, y) = \Phi(x^\mu, y + 2\pi R)$$

$$S = \int d^4x \int_0^{2\pi R} \left(\partial^M \Phi^\dagger \partial_M \Phi - m^2 \Phi^\dagger \Phi + \dots \right)$$

$$\dim(\Phi) = \frac{3}{2} \quad M = \text{5D Lorentz index}$$



Kaluza-Klein Decomposition

$$\Phi(x, y) = \sum_{n=0}^{\infty} \Phi^{(n)}(x) e^{iny/R}$$

$$S = \int d^4x \sum_{n=0}^{\infty} 2\pi R \left[\partial^M \Phi^{(n)\dagger} \partial_M \Phi^{(n)} - \left(m^2 + \frac{n^2}{R^2} \right) \Phi^{(n)\dagger} \Phi^{(n)} + \dots \right]$$

from integral over extra dimension

$$|\partial_y e^{iny/R}|^2 = \frac{n^2}{R^2}$$

Define $\hat{\Phi}^{(n)} = \sqrt{2\pi R} \Phi^{(n)}$

$$S = \int d^4x \sum_{n=0}^{\infty} \left[\partial^M \hat{\Phi}^{(n)\dagger} \partial_M \hat{\Phi}^{(n)} - \left(m^2 + \frac{n^2}{R^2} \right) \hat{\Phi}^{(n)\dagger} \hat{\Phi}^{(n)} + \dots \right]$$

Note $\dim(\Phi) = 1$

5D theory equivalent to infinite number of 4D fields
with increasing mass $m_n^2 = m^2 + n^2/R^2$

Interactions

$$\begin{aligned} \int dy (\Phi^\dagger \Phi)^2 &= \sum_{n_1, n_2, n_3, n_4} e^{-i(n_1 - n_2 + n_3 - n_4)y/R} \Phi^{(n_1)\dagger} \Phi^{(n_2)} \Phi^{(n_3)\dagger} \Phi^{(n_4)} \\ &= \sum_{n_1, n_2, n_3, n_4} \frac{1}{2\pi R} \underbrace{\delta_{n_1 - n_2 + n_3 - n_4, 0}}_{\text{5D momentum conservation}} \hat{\Phi}^{(n_1)\dagger} \hat{\Phi}^{(n_2)} \hat{\Phi}^{(n_3)\dagger} \hat{\Phi}^{(n_4)} \end{aligned}$$

$$\dim(\Phi^4) = 6 \quad \Rightarrow \Delta \mathcal{L}_5 \sim \frac{1}{\Lambda} (\Phi^\dagger \Phi)^2 \quad \text{irrelevant in 5D}$$

5D gauge interactions

$$D_M = \partial_M - ig_5 A_M$$

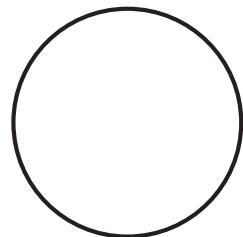
$$\dim(A_M) = \frac{3}{2} \quad \Rightarrow g_5 \sim \frac{1}{\Lambda^{1/2}} \quad \text{also irrelevant}$$

4D \rightarrow 5D \Rightarrow give up renormalizability (UV completeness)

Shape of Extra Dimensions

One extra dimension:

Compact:



or



Non-compact:

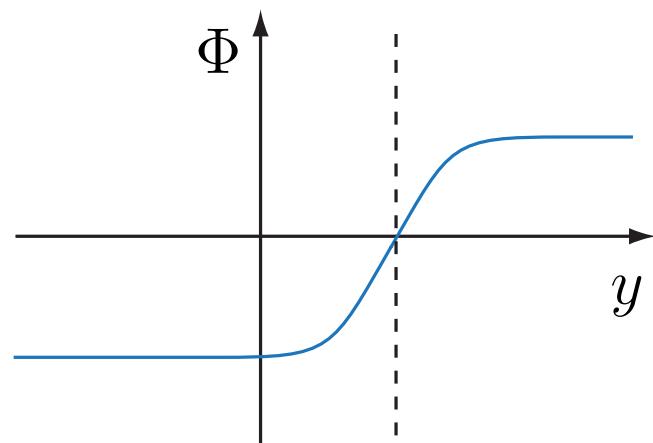


or

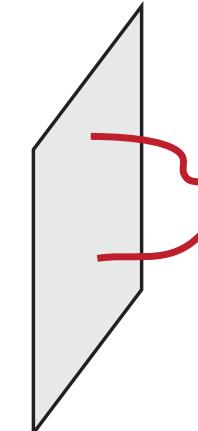


“Brane” = special point in extra dimension

Field varying in 5D

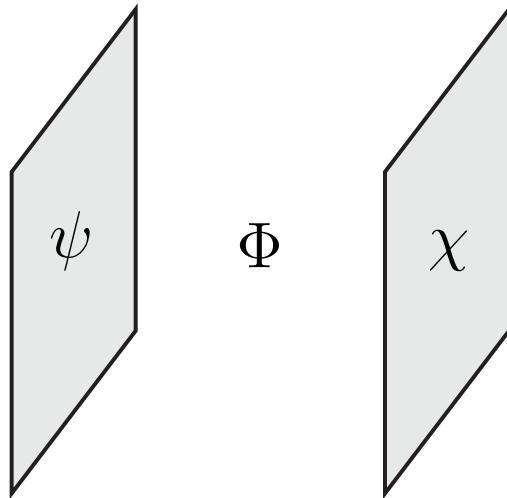


D-branes in string theory



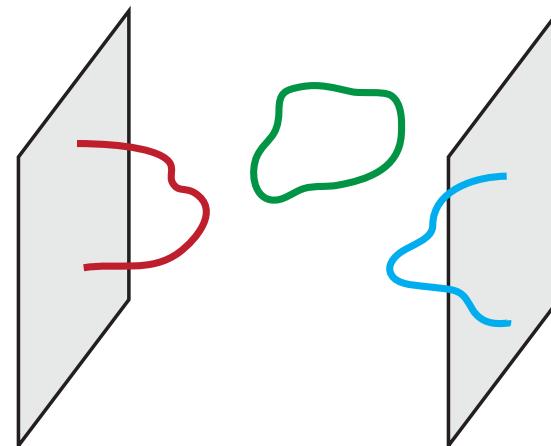
Brane Localization

Particles can be localized on branes

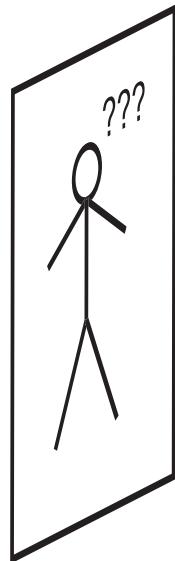


$$\begin{aligned}\mathcal{L}_{\text{eff}} = \int d^4x & \left[\int_0^L dy \mathcal{L}_{\text{bulk}}(\Phi(x, y)) \right. \\ & + \mathcal{L}_0(\Phi(x, y=0), \psi(x)) \\ & \left. + \mathcal{L}_L(\Phi(x, y=L), \chi(x)) \right]\end{aligned}$$

May occur in field theory
Required in string theory



Life on a Brane



Could we be living on a brane
in an infinite extra dimension?

a) **No:** gravity = spacetime curvature

$$V(r) \propto \frac{1}{r^{d-1}} \quad \Rightarrow d = 3 + 1$$

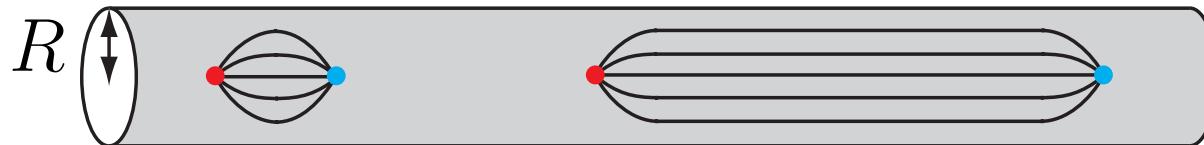
b) **Yes:** gravity can be localized on extra dimension
(Dvali, Gabadadze, Porati 2000)

⇒ interesting (dangerous) modification of gravity
at long distances

Large Extra Dimensions

(Arkani-Hamed, Dimopoulos, Dvali 1998)

Gravity in compactified extra dimensions



$$\begin{array}{ll} r \ll R & r \gg R \\ V(r) \sim \frac{1}{M_*^{d-2}} \frac{1}{r^{d-3}} & V(r) \sim \frac{1}{M_P^2} \frac{1}{r} \end{array}$$

Match at $r \sim R \Rightarrow M_P^2 \sim M_*^{d-2} R^{d-4}$

Experiment $\Rightarrow R \lesssim \mu\text{m} \sim 5 \times 10^{11} \text{ GeV}^{-1}$

Scale of quantum gravity $\sim M_* \sim \begin{cases} 10 \text{ GeV} & d = 4 + 1 \\ 5 \text{ TeV} & d = 5 + 1 \\ 5 \text{ GeV} & d = 6 + 1 \end{cases}$

No hierarchy problem!

Must still explain origin of $R \gg M_*^{-1} \dots$

Quantum Gravity at TeV

Exciting, but also dangerous

- KK gravitons $m_{\text{KK}} \sim \frac{1}{R} \sim 10^{-3}$ eV
 \Rightarrow many accessible states
- Cosmological “normalcy bound”
 $T_{\text{max}} \lesssim$ GeV
- Supernova emission, neutron star heating
 $\Rightarrow M_* \gtrsim 1000$ TeV ($d = 6$)

Asymmetric extra dimensions still allow quantum gravity at TeV

Signals

Black holes at LHC!

- Production

Black hole formed for impact parameter

$$b \lesssim R_S = \text{Schwarzschild radius} \quad (R_S)^{d-3} \sim \frac{\sqrt{\hat{s}}}{M_*^{d-2}}$$

$$\sqrt{\hat{s}} \sim \text{TeV} \Rightarrow \sigma_{\text{BH}} \sim \text{TeV}^{-2} \sim \text{nb}$$

- Decay

Hawking radiation \Rightarrow democratic decay

$\langle n \rangle \simeq 6$ particles in final state

Higgsless Models

(Csaki, Grojean, Murayama, Pilo, Terning, 2003)

Can we get W, Z masses from KK mechanism?

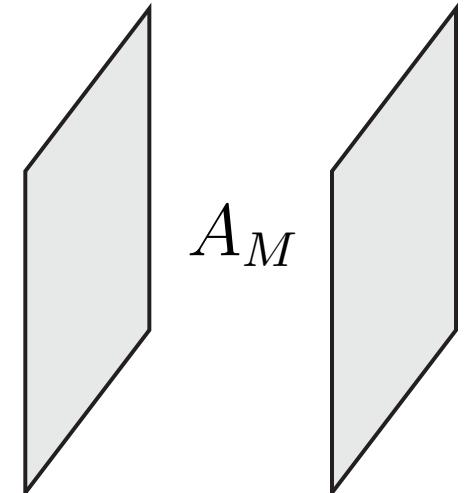
Example: $U(1)$ gauge field in 5D

$$A_M = (A_\mu, A_5) \quad \text{gauge fix } A_5 \equiv$$

$$A_\mu(x, y) = \sum_n A_\mu^{(n)}(x) f_n(y)$$

$$A_\mu^{(n)}(x) \propto e^{ip \cdot x} \Rightarrow (p^2 + \partial_y^2) f_n(y) = 0$$

Eigenvalues determine masses $\omega_n^2 = p^2 + \frac{\omega_n^2}{n}$



Boundary conditions fixed by boundary terms

$$\Delta\mathcal{L}_{\text{bdy}} = -\frac{1}{2} m A_\mu^2$$

$$\Rightarrow mA_\mu \pm F_{5\mu}|_y \quad , \quad L$$

Higgsless Models II

$$mA_\mu \pm F_{5\mu}| = 0$$

- $m \rightarrow \infty \Rightarrow \partial_5 A_\mu| = 0$

$$f_n(y) = \cos\left(\frac{n\pi y}{L}\right) \quad n = 0, 1, \dots$$

$$m_n = \frac{n\pi}{L} \quad n = 0 \text{ mode massless}$$

- $m \rightarrow \infty \Rightarrow A_\mu| = 0$

$$f_n(y) = \sin\left(\frac{n\pi y}{L}\right) \quad n = 1, 2, \dots$$

$$m_n = \frac{n\pi}{L} \quad \text{no massless 0 mode}$$

Gauge symmetry broken by boundary conditions

Higgsless Models III

Generalize to $SU(N)$

KK modes have gauge coupling $g_4 \sim \frac{g_5}{\sqrt{L}}$

Where does 5D theory break down?

$$\sim g \quad \sim \frac{g_5^4 E}{24\pi^3} \Rightarrow \Lambda_5 \sim \frac{24\pi^3}{g_5^2}$$
$$\Rightarrow \frac{m_W}{\Lambda_5} \sim \frac{L^{-1}}{24\pi^3/g_5^2} \sim \frac{g_4^2}{24\pi^3} \sim 100 \text{ TeV}$$

Compare effective standard model: $\frac{m_W}{\Lambda_5} \sim \frac{g_4}{4\pi} \sim \text{TeV}$

Partial UV completion from extra dimensions!

Higgsless Models IV

Realistic models?

- $SU(2) \times U(1) \rightarrow U(1)_{\text{EM}}$
- Custodial symmetry
- Avoid $m_W = m_Z = \frac{1}{L}$
- Avoid $m_W = 2m_W$, etc.

Requires generalized boundary conditions
 $SU(2) \times SU(2)$ bulk gauge fields

- $S \sim 10 + \dots$

Requires some tuning

Signals

- No light Higgs
- Strong WW scattering
- Many TeV resonances

Haven't we seen this before?

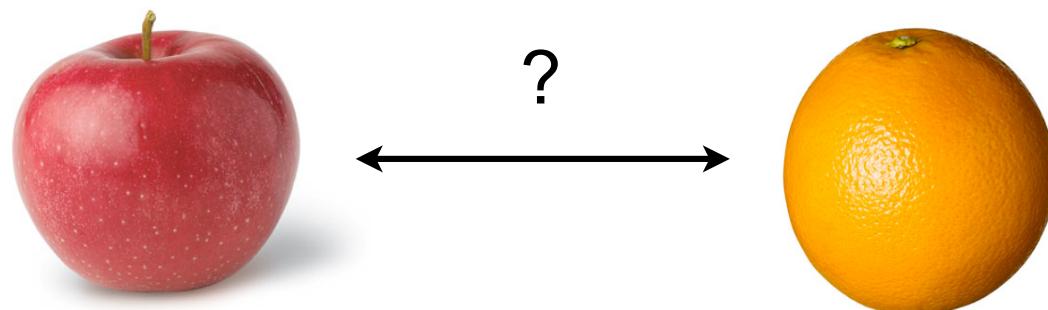
QCD? Technicolor?

Duality

Similarity between 5D and 4D theories goes deep

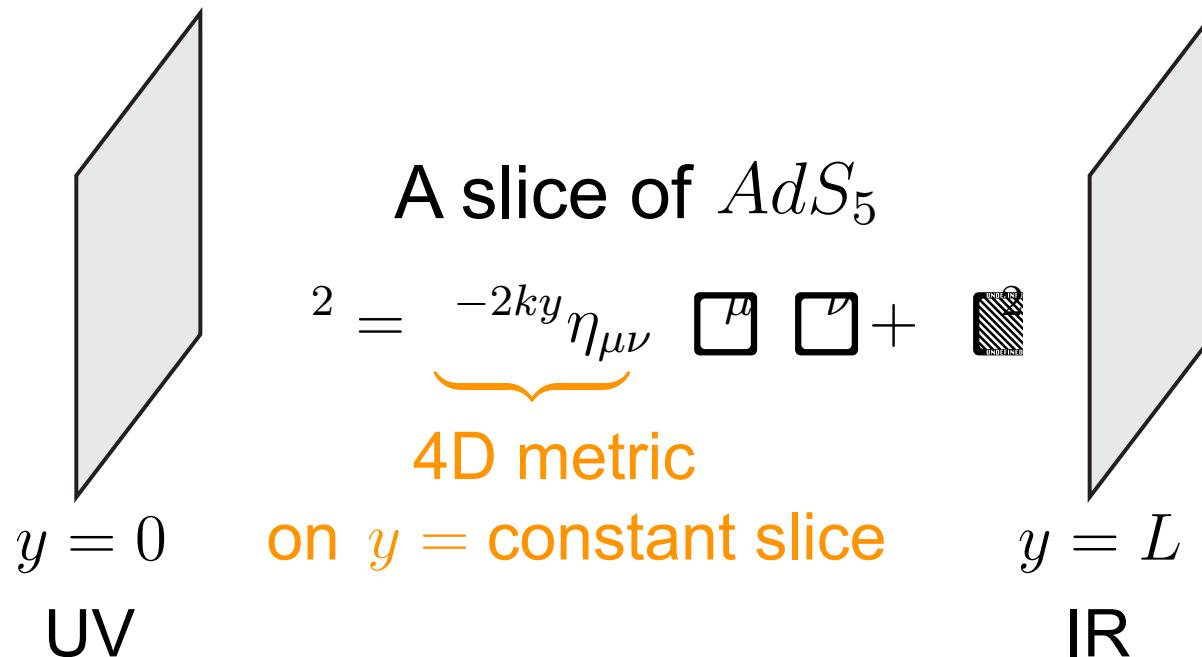
Precise correspondence in string theory:
(Maldacena 1999)

$$\begin{array}{ccc} \mathcal{N} = 4, \ d = 4 & \longleftrightarrow & \text{String theory on} \\ \text{super Yang Mills} & & AdS_5 \times S_5 \end{array}$$



Randall-Sundrum Model

(Randall, Sundrum 1999)



All fundamental parameters $\sim M_*$

Higgs localized on IR brane $\Rightarrow m_H \sim M_* e^{-kL} \ll M_*$

$k, M_* \sim M_P, kL \sim 10 \Rightarrow m_H \sim \text{TeV}$

Hierarchy problem solved by gravitational “warping”

RS with Bulk SM Fields

Several nice features

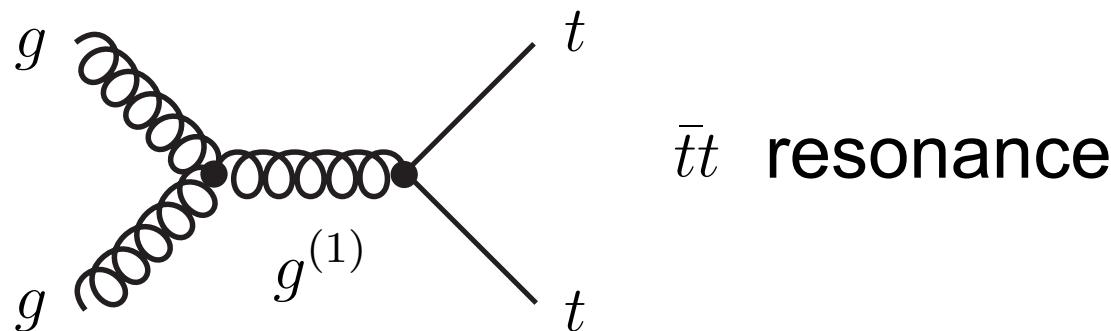
- Precision gauge coupling unification
- Fermion mass hierarchies natural

Also some problems

- Precision electroweak and flavor require some tuning
 $\sim 10\% \Rightarrow m_{KK} \gtrsim 5 \text{ TeV}$

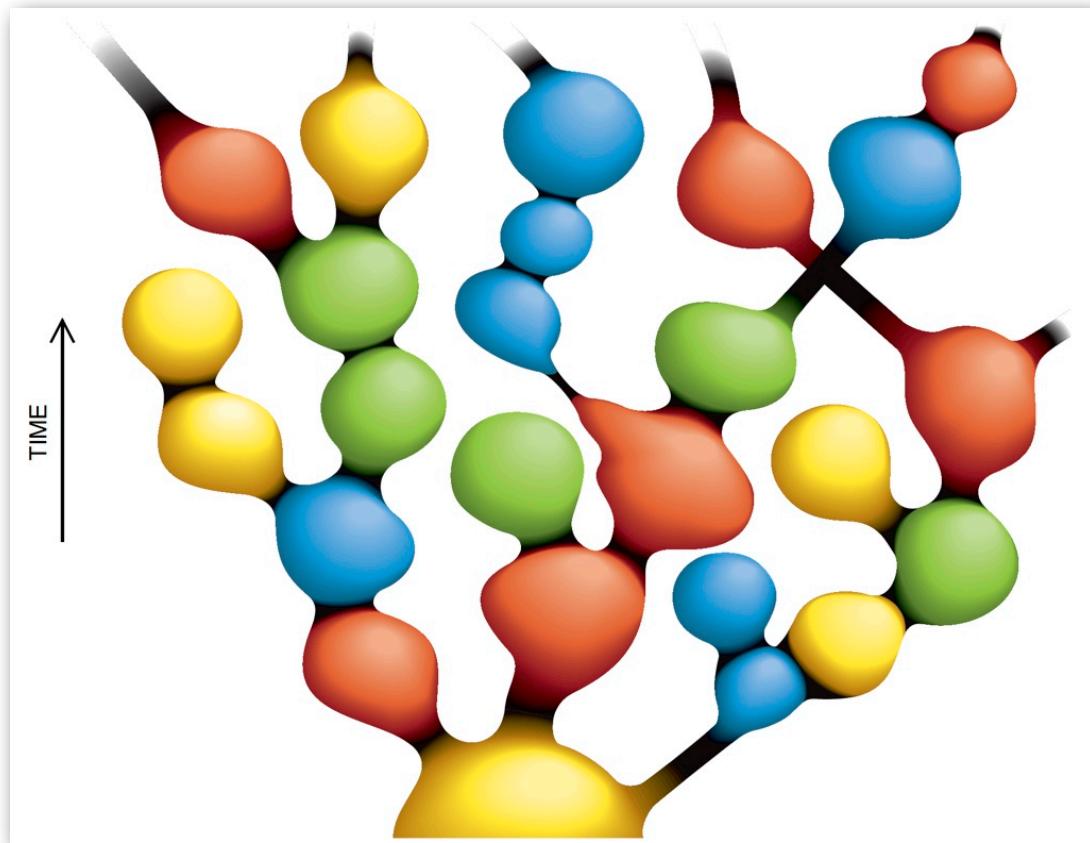
Signals

- Top quark localized near IR brane
⇒ “composite” (couples strongly to KK states)



- Other KK resonances
 $W, Z, \text{graviton}$

Multiverse



To the blackboard...